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# Hand in 3

## FYSC20

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**Problem:** A sphere of radius  $R$  carries a “frozen-in” polarisation  $\mathbf{P}(\mathbf{r}) = kr^2\hat{\mathbf{r}}$ , where  $k$  is a constant and  $\mathbf{r}$  is the vector from the center.

- (a) Calculate the bound charges  $\sigma_b$  and  $\rho_b$ .
- (b) Find the  $\mathbf{E}$  field inside and outside the sphere from the charge densities found at (a).
- (c) Alternatively, compute the  $\mathbf{E}$  field by first determining the  $\mathbf{D}$  field using Gauss’s law and then use that  $\mathbf{D} = \epsilon_0\mathbf{E} + \mathbf{P}$ .

**(a) Solution:** Start by drawing a sketch of the problem.

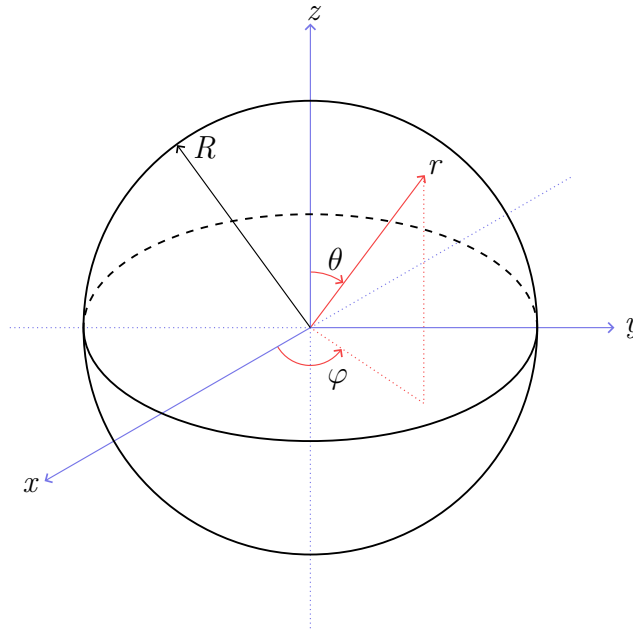


Figure 1: The problem sketch with cartesian and spherical coordinates drawn.

We can calculate the bound charge densities from the polarisation  $\mathbf{P}$ .

$$\sigma_b = \hat{\mathbf{n}} \cdot \mathbf{P}, \quad \rho_b = -\nabla \cdot \mathbf{P}. \quad (1)$$

Since we have a sphere the normal vector is  $\hat{\mathbf{n}} = \hat{\mathbf{r}}$ , the surface charge density is also only

on  $r = R$ , and  $\mathbf{P}$  only have a radial component. Thus

$$\sigma_b = \hat{\mathbf{r}} \cdot \mathbf{P} = \hat{\mathbf{r}} \cdot (kR^2 \hat{\mathbf{r}}) = kR^2 \quad (2)$$

$$\rho_b = -\nabla \cdot \mathbf{P} = -\frac{1}{r^2} \frac{\partial}{\partial r} (kr^4) = -4kr \quad (3)$$

The charge densities are only non zero on or in the sphere and zero elsewhere.

**(a) Answer:** The charge densities are

$$\sigma_b = kR^2, \quad \rho_b = -4kr. \quad (4)$$

**(b) Solution:** To calculate the  $\mathbf{E}$  field we can use Gauss's law. Adding a gaussian surface to the sketch

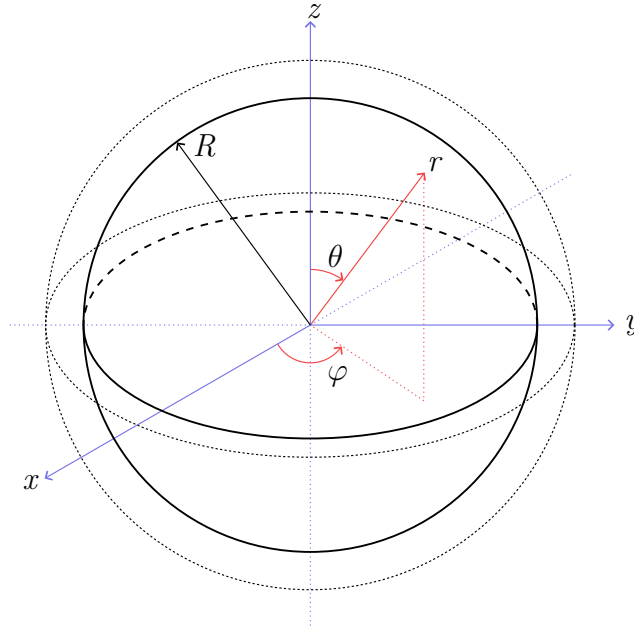


Figure 2: The sketch with and added gaussian surface.

More correctly we might write the charge densities

$$\sigma_b(\mathbf{r}) = \begin{cases} kR^2 & r = R \\ 0 & \text{elsewhere} \end{cases}, \quad \text{and} \quad \rho_b(\mathbf{r}) = \begin{cases} -4kr & 0 \leq r < R \\ 0 & \text{elsewhere} \end{cases} \quad (5)$$

Now we can define

$$\rho(\mathbf{r}) = \rho_b(\mathbf{r}) + \sigma_b(\mathbf{r}) \quad (6)$$

Due to spherical symmetry we can write

$$\mathbf{E}(\mathbf{r}) = E(r)\hat{\mathbf{r}}. \quad (7)$$

Using Gauss's law

$$\nabla \cdot \mathbf{E}(\mathbf{r}) = \frac{\rho(\mathbf{r})}{\epsilon_0} \iff E(r) \oint_{\partial V} da_r = \frac{1}{\epsilon_0} \int_V \rho(\mathbf{r}') d\tau' \quad (8)$$

where  $V$  is the volume we integrate over, that is the sphere. Thus

$$E(r) \int_0^\pi \int_0^{2\pi} r^2 \sin \theta d\theta d\varphi = \frac{1}{\epsilon_0} \int_0^\pi \int_0^{2\pi} \int_0^r (\rho_b(\mathbf{r}') + \sigma_b(\mathbf{r}')) (dr')(r' d\theta')(r' \sin \theta' d\varphi') \quad (9)$$

$$E(r)r^2 4\pi = \frac{4\pi}{\epsilon_0} \int_0^r (\rho_b(\mathbf{r}') + \sigma_b(\mathbf{r}')) r'^2 dr' \quad (10)$$

$$E(r) = \frac{1}{\epsilon_0 r^2} \int_0^r (\rho_b(\mathbf{r}') + \sigma_b(\mathbf{r}')) r'^2 dr' \quad (11)$$

Now we get two cases:  $r < R$  and  $r > R$

**Case  $r < R$ :** In this case  $\sigma_b(\mathbf{r})$  will always be zero by Eq. (5). Thus

$$E(r) = \frac{1}{\epsilon_0 r^2} \int_0^r \rho_b(\mathbf{r}') r'^2 dr' = -\frac{4k}{\epsilon_0 r^2} \int_0^r r'^3 dr' \quad (12)$$

$$E(r) = -\frac{4k}{\epsilon_0 r^2} \left[ \frac{r'^4}{4} \right]_0^r = -\frac{k}{\epsilon_0} r^2 \quad (13)$$

And thus we get that

$$\mathbf{E}(\mathbf{r})_{\text{inside}} = -\frac{k}{\epsilon_0} r^2 \hat{\mathbf{r}} \quad (14)$$

**Case  $r > R$ :** Rewriting  $\sigma_b(\mathbf{r})$  using a delta function and using that  $\rho_b(\mathbf{r})$  is zero outside the sphere we get

$$E(r) = \frac{1}{\epsilon_0 r^2} \int_0^r (\rho_b(\mathbf{r}) + \delta(r' - R)kR^2)r'^2 dr' \quad (15)$$

$$E(r) = \frac{1}{\epsilon_0 r^2} \left( -4k \int_0^R r'^3 dr' + kR^2 \int_0^r \delta(r' - R)r'^2 dr' \right) \quad (16)$$

$$E(r) = \frac{k}{\epsilon_0 r^2} \left( -4 \left[ \frac{r'^4}{4} \right]_0^R + R^4 \right) = \frac{k}{\epsilon_0 r^2} (-R^4 + R^4) = 0 \quad (17)$$

Therefore the field outside is

$$\mathbf{E}(\mathbf{r})_{\text{outside}} = \mathbf{0} \quad (18)$$

**(b) Answer:** We get the  $\mathbf{E}$  field

$$\mathbf{E}(\mathbf{r}) = \begin{cases} -\frac{k}{\epsilon_0} r^2 \hat{\mathbf{r}} & r < R \\ \mathbf{0} & r > R \end{cases} \quad (19)$$

**(c) Solution:** We use the same gaussian surface as seen in Fig. 2. We know have

$$\nabla \cdot \mathbf{D} = \rho_f \iff \oint_{\partial V} \mathbf{D} \cdot d\mathbf{a} = \int_V \rho_f d\tau. \quad (20)$$

However, we don't have any free charges so  $\rho_f = 0$ . Since we have a spherical geometry and the polarisation is spherically symmetrical we can write

$$\mathbf{D}(\mathbf{r}) = D(r) \hat{\mathbf{r}} \quad (21)$$

And thus the equation simplifies to

$$D(r)r^2 4\pi = 0 \iff D(r) = 0 \quad (22)$$

Now using that

$$\mathbf{D} = \epsilon_0 \mathbf{E} + \mathbf{P} \quad (23)$$

and since  $\mathbf{P}$  is zero outside the sphere we have two cases:  $r < R$  and  $r > R$ .

**Case  $r < R$ :** Here we have

$$\mathbf{0} = \epsilon_0 \mathbf{E}_{\text{inside}} + \mathbf{P} \iff \mathbf{E}_{\text{inside}} = -\frac{k}{\epsilon_0} r^2 \hat{\mathbf{r}} \quad (24)$$

**Case  $r > R$ :** Since  $\mathbf{P}$  is zero in this region we have

$$\mathbf{E}_{\text{outside}} = \mathbf{0} \quad (25)$$

**(c) Answer:** We see that the answer is the exact same as in (b) but with a lot less calculations. We need to be careful however, as we see in Exercise F.4, the polarisation must have spherical symmetry to be able to write  $\mathbf{D}(\mathbf{r}) = D(r)\hat{\mathbf{r}}$ .