

UNITARITY CONSTRAINTS ON 3 HADRON DYNAMICS

Seminar @ UMD – 10/25/18

Maxim Mai
The George Washington University



Deutsche
Forschungsgemeinschaft
DFG

RESULTS

→ unitary 3-body scattering amplitude:

[EPJA 53 \(2017\)](#)

$$T_{3 \rightarrow 3} = v \left(\tau \frac{B}{1 + B\tau} \tau - \tau \right) v$$

RESULTS

→ unitary 3-body scattering amplitude:

EPJA53 (2017)

$$T_{3 \rightarrow 3} = v \left(\tau \frac{B}{1 + B\tau} \tau - \tau \right) v$$



→ relativistic 3-body quantization condition

EPJA53 (2017), PRD97 (2018)

$$\det \left(B_{uu'}^{\Gamma ss'} + \frac{2E_s L^3}{\vartheta(s)} \tau_s^{-1} \delta_{ss'} \delta_{uu'} \right) = 0$$

RESULTS

→ unitary 3-body scattering amplitude:

EPJA53 (2017)



→ relativistic 3-body quantization condition

EPJA53 (2017), PRD97 (2018)

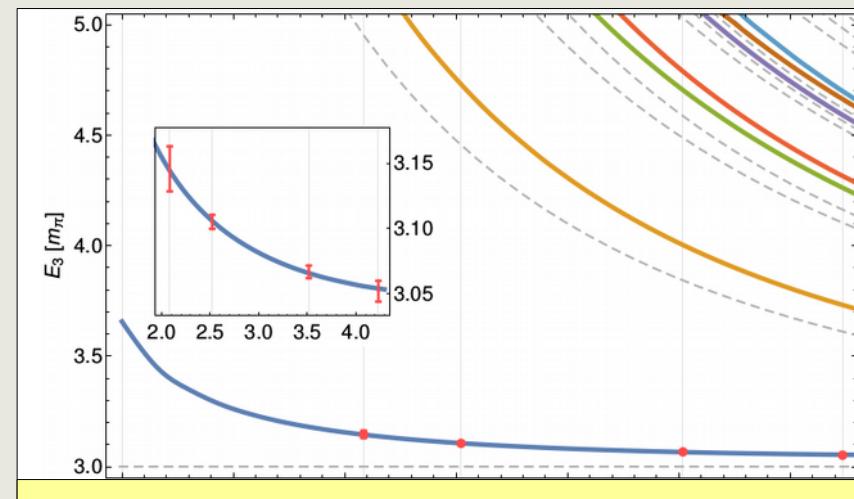
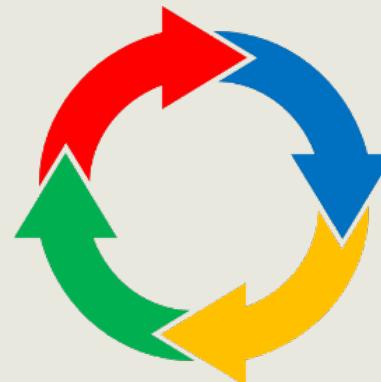
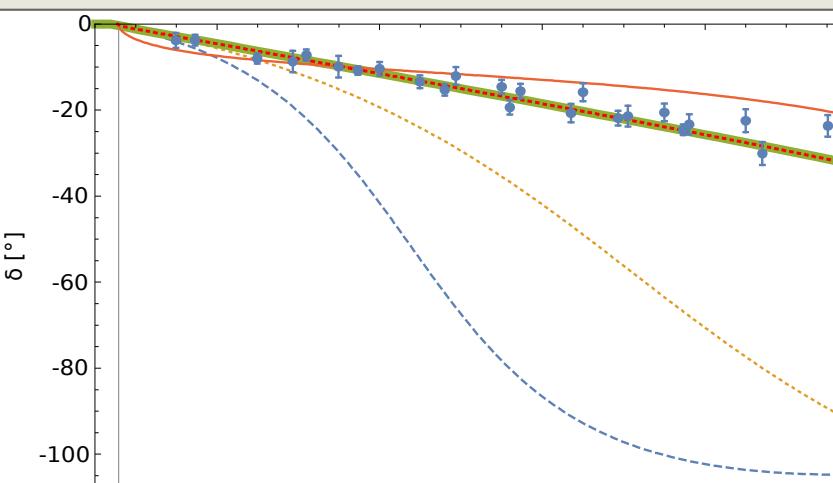


→ application to a physical system

$$T_{3 \rightarrow 3} = v \left(\tau \frac{B}{1 + B\tau} \tau - \tau \right) v$$

$$\det \left(\mathbf{B}_{\mathbf{u}\mathbf{u}'}^{\Gamma_{\mathbf{s}\mathbf{s}'}} + \frac{2\mathbf{E}_s \mathbf{L}^3}{\vartheta(s)} \tau_s^{-1} \delta_{\mathbf{s}\mathbf{s}'} \delta_{\mathbf{u}\mathbf{u}'} \right) = 0$$

arXiv:1807.04746

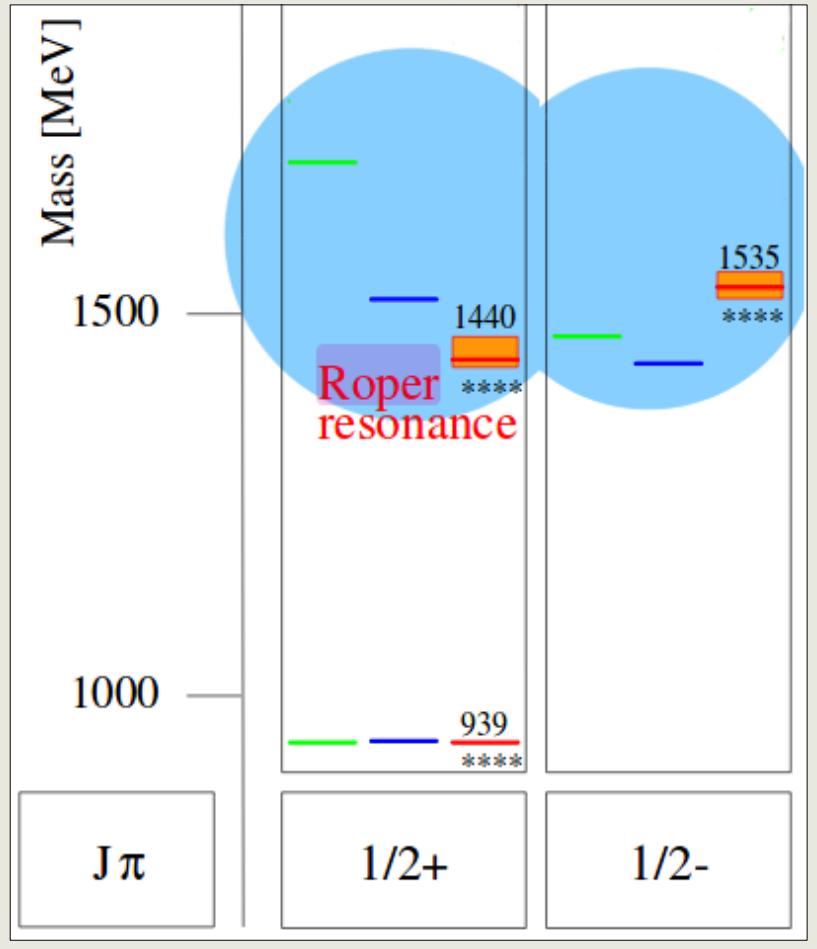


MOTIVATION

Many unsolved questions of QCD involve 3-body channels

EX. 1: Roper-puzzle

Reversed mass pattern vs. Const. Quark Model

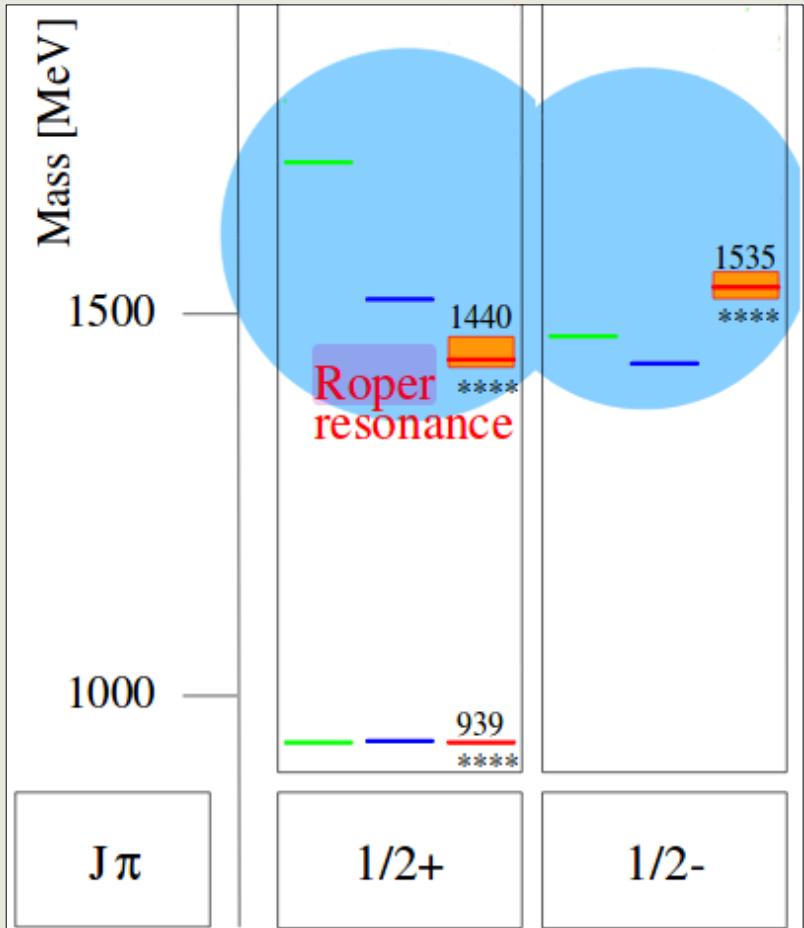


Loring et al. EPJA10 (2001)

Many unsolved questions of QCD involve 3-body channels

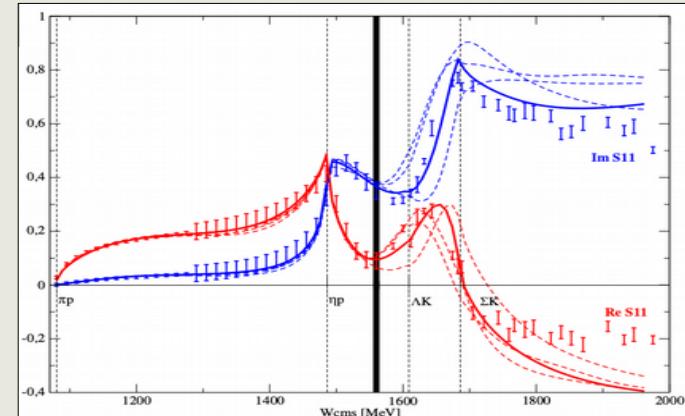
EX. 1: Roper-puzzle

Reversed mass pattern vs. Const. Quark Model



Loring et al. EPJA10 (2001)

$N(1535)1/2^-$ accessed from dyn. approaches, e.g.,

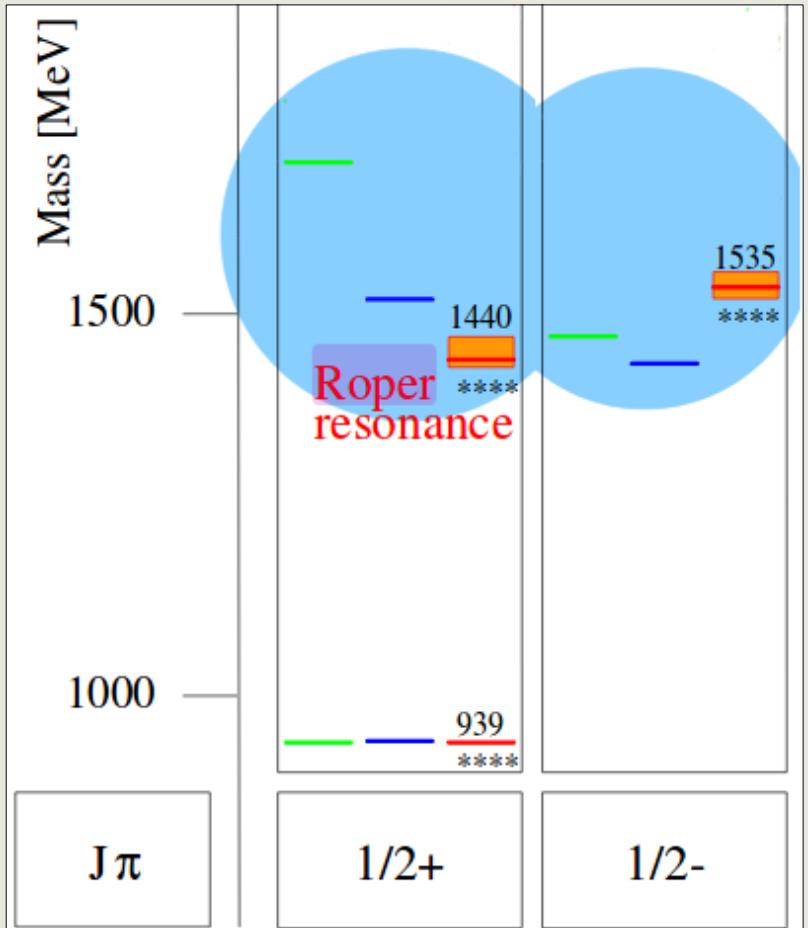


Bruns, MM, Meissner PLB 697 (2011)

Many unsolved questions of QCD involve 3-body channels

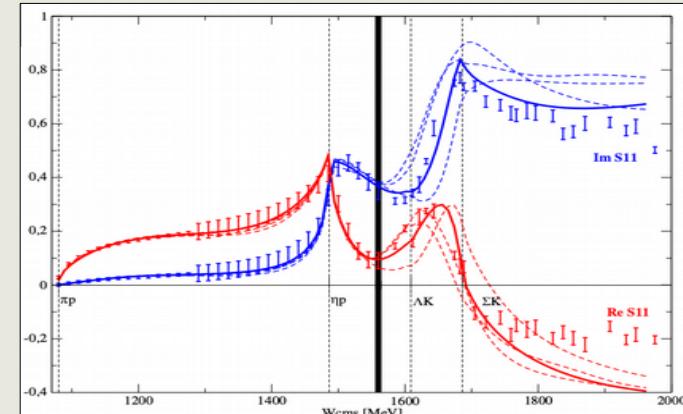
EX. 1: Roper-puzzle

Reversed mass pattern vs. Const. Quark Model



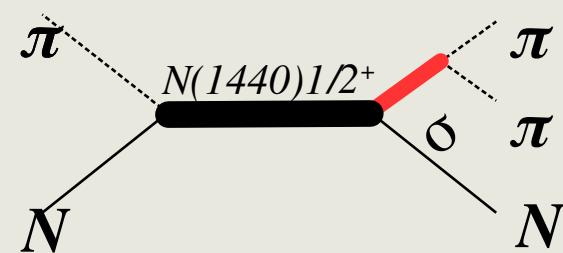
Loring et al. EPJA10 (2001)

$N(1535)1/2^-$ accessed from dyn. approaches, e.g.,



Bruns, MM, Meissner PLB 697 (2011)

...but $N(1440)1/2^+$ has large branching ratio to $\pi\pi N$



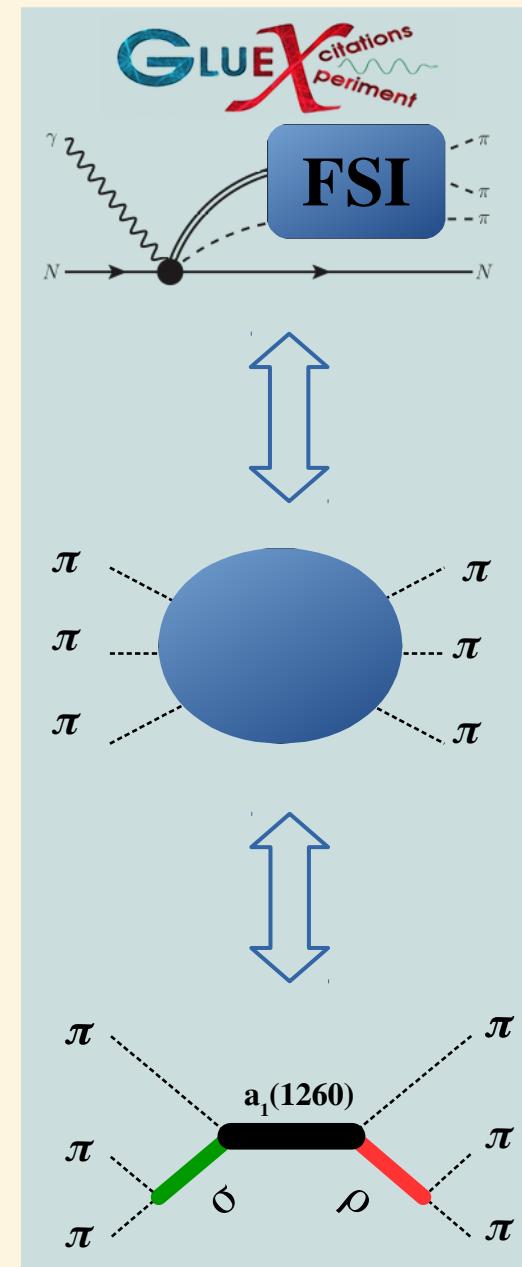
- 1st LQCD calculation w. meson-baryon operators

Lang et al. PRD 95 (2017)

Many unsolved questions of QCD involve 3-body channels

EX. 2: $a_1(1260)$

- important for the search for *spin-exotics* (*GlueX*, *COMPASS*, *BESIII*)
 - indicator for gluonic degrees of freedom
 - cannot decay in 2 but in 3 pions, such as the $a_1(1260)$



Many unsolved questions of QCD involve 3-body channels

EX. 2: $a_1(1260)$

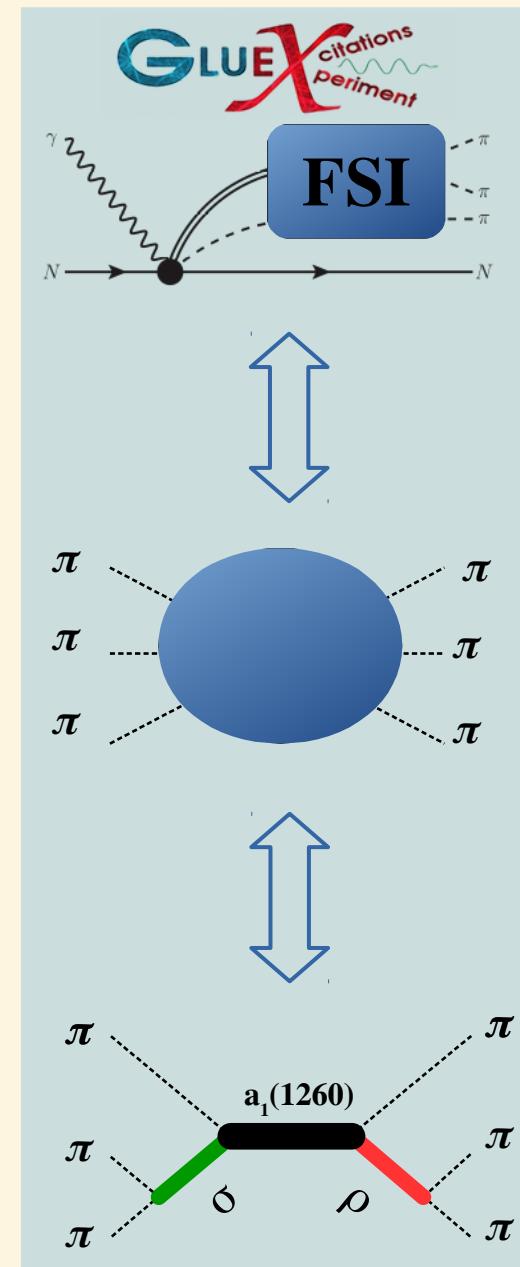
- important for the search for *spin-exotics* (*GlueX*, *COMPASS*, *BESIII*)
 - indicator for gluonic degrees of freedom
 - cannot decay in 2 but in 3 pions, such as the $a_1(1260)$
- finite-volume LQCD spectrum

Lang et al. JHEP 1404

further studies will follow...

Woss et al. JHEP 1807 [HadSpec]

$\rightarrow I=2 \pi\eta$



Many unsolved questions of QCD involve 3-body channels

EX. 2: $a_1(1260)$

- important for the search for *spin-exotics* (*GlueX*, *COMPASS*, *BESIII*)
 - indicator for gluonic degrees of freedom
 - cannot decay in 2 but in 3 pions, such as the $a_1(1260)$
- finite-volume LQCD spectrum Lang et al. JHEP 1404

further studies will follow...

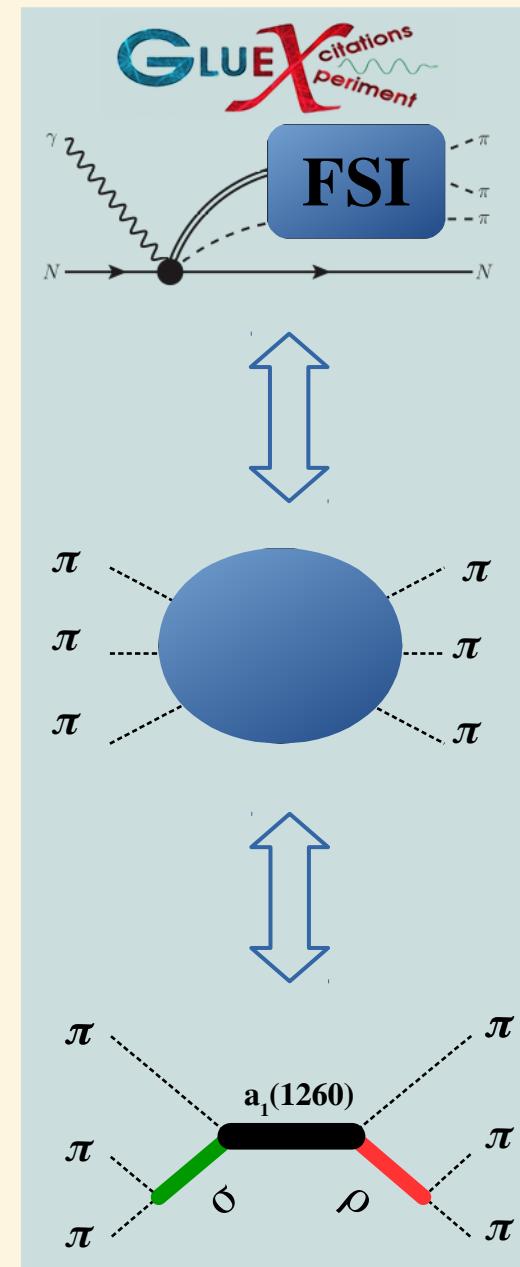
Woss et al. JHEP 1807 [HadSpec]

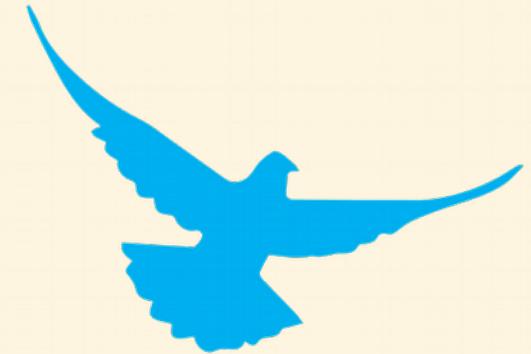
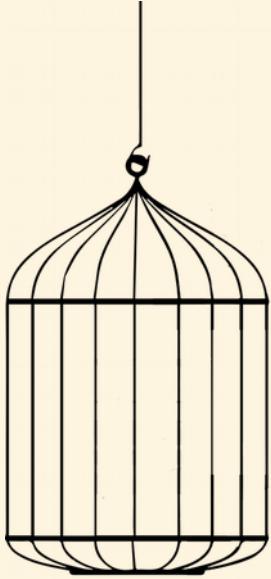
$\rightarrow I=2 \pi\eta$

EX. 3: $X(3872) \leftrightarrow D\bar{D}\pi$

Prelovsek, Leskovec PRL111 (2013)

...





3→3 SCATTERING AMPLITUDE IN INFINITE VOLUME



MM, Hu, Doring, Pilloni, Szczepaniak EPJA53 (2017)

T-MATRIX

- 3 asymptotic states (scalar particles of equal mass (m))
- Connectedness structure of matrix elements



T-MATRIX

- 3 asymptotic states (scalar particles of equal mass (m))
- Connectedness structure of matrix elements



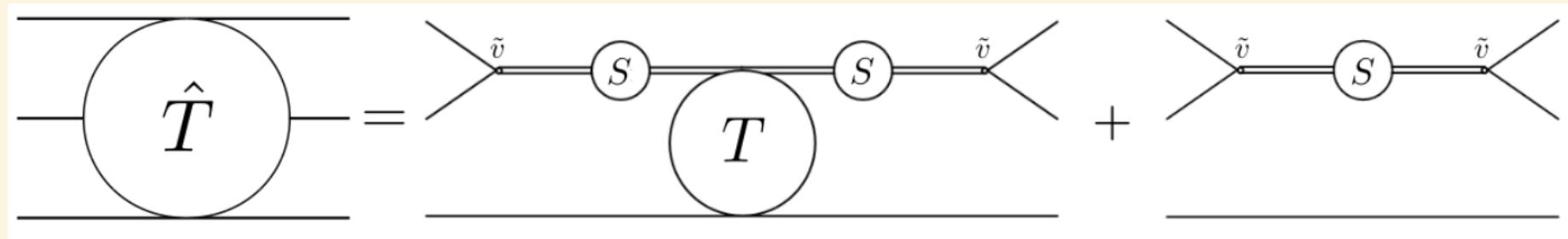
- *isobar*-parametrization of two-body amplitude

Bedaque, Griesshammer (1999)

- “*isobars*” – $S(\sigma)$ for definite QN & correct right-hand-singularities
- in general a tower of “*isobars*” for $L=0,1,2,\dots$
- coupling to asymptotic states: cut-free-function $v(q,p)$

T-MATRIX

- 3 asymptotic states (scalar particles of equal mass (m))
- Connectedness structure of matrix elements



- *isobar*-parametrization of two-body amplitude

Bedaque, Griesshammer (1999)

- “*isobars*” – $S(\sigma)$ for definite QN & correct right-hand-singularities
- in general a tower of “*isobars*” for $L=0,1,2,\dots$
- coupling to asymptotic states: cut-free-function $v(q,p)$

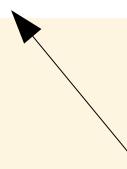
- Connected part: due to isobar-spectator interaction $\rightarrow T(q_{in}, q_{out}; s)$

- 3 unknown functions
- 8 kinematic variables

3-BODY UNITARITY

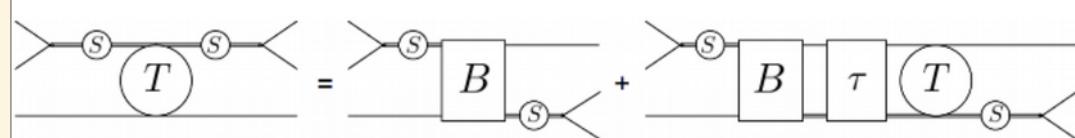
$$\langle q_1, q_2, q_3 | (\hat{T} - \hat{T}^\dagger) | p_1, p_2, p_3 \rangle = i \int_P \langle q_1, q_2, q_3 | \hat{T}^\dagger | k_1, k_2, k_3 \rangle \langle k_1, k_2, k_3 | \hat{T} | p_1, p_2, p_3 \rangle$$

phase-space integral



3-BODY UNITARITY

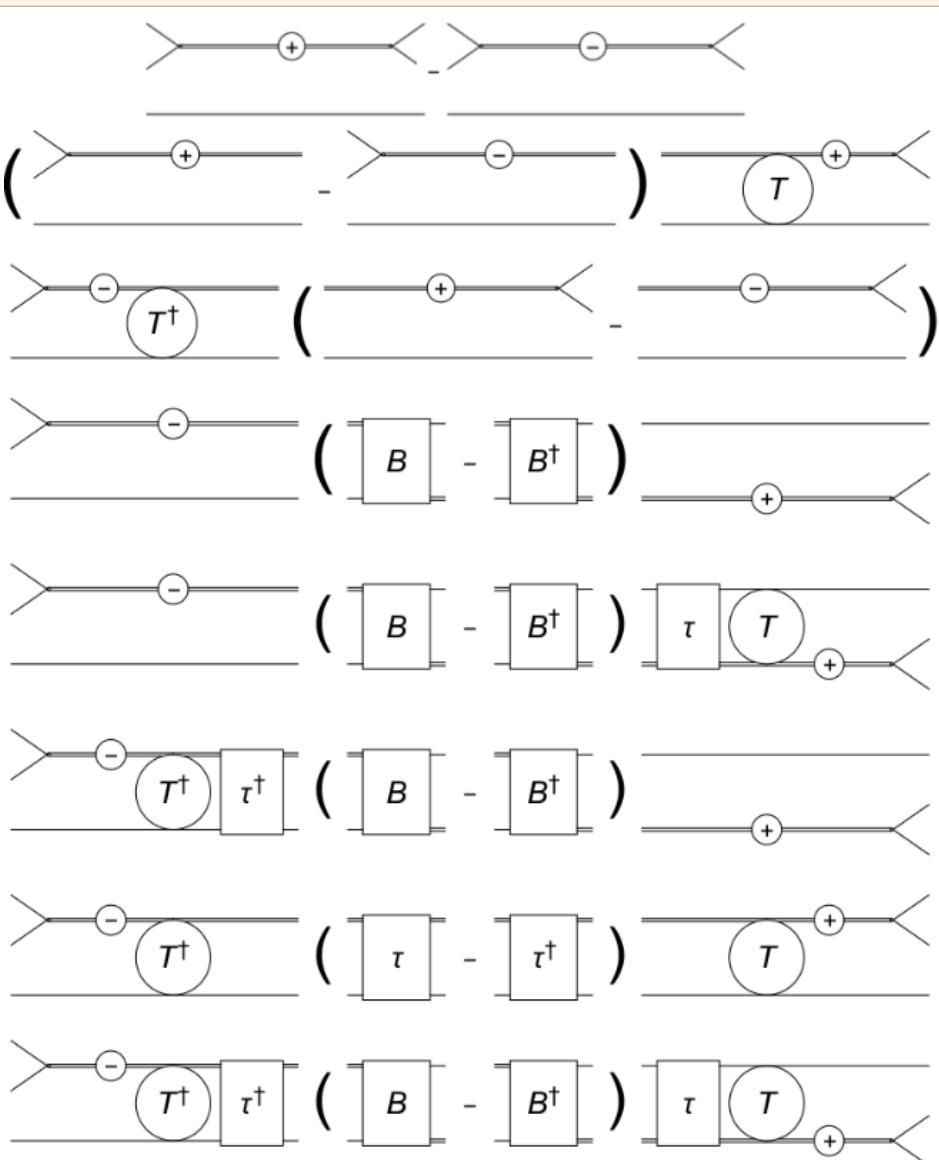
$$\langle q_1, q_2, q_3 | (\hat{T} - \hat{T}^\dagger) | p_1, p_2, p_3 \rangle = i \int_P \langle q_1, q_2, q_3 | \hat{T}^\dagger | k_1, k_2, k_3 \rangle \langle k_1, k_2, k_3 | \hat{T} | p_1, p_2, p_3 \rangle$$



General Ansatz for the isobar-spectator interaction
→ B & τ are new unknown functions

3-BODY UNITARITY

$$\langle q_1, q_2, q_3 | (\hat{T} - \hat{T}^\dagger) | p_1, p_2, p_3 \rangle = i \int_P \langle q_1, q_2, q_3 | \hat{T}^\dagger | k_1, k_2, k_3 \rangle \langle k_1, k_2, k_3 | \hat{T} | p_1, p_2, p_3 \rangle$$



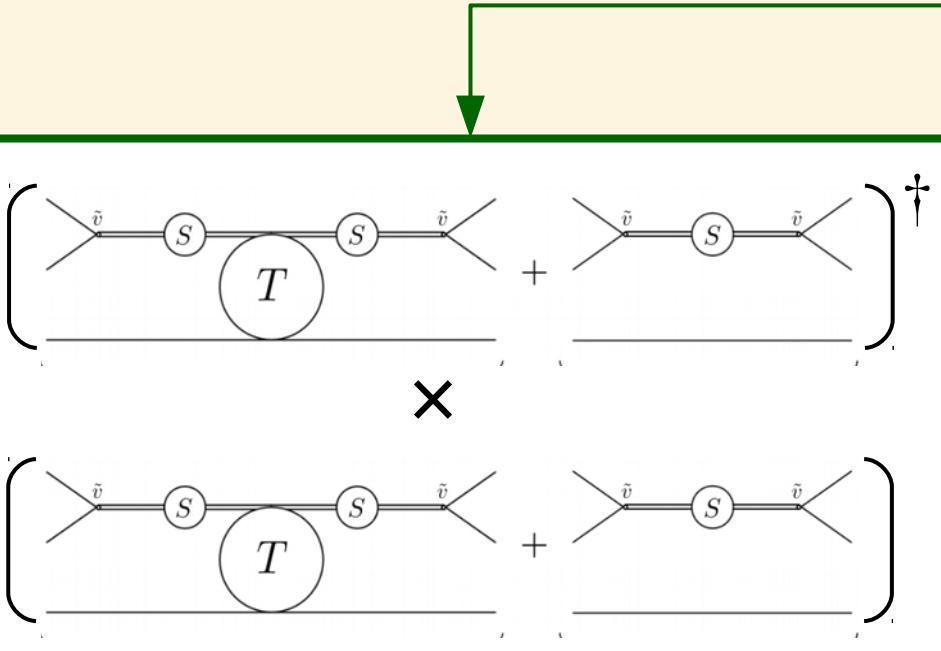
General Ansatz for the isobar-spectator interaction
 → **B & τ** are new unknown functions

8 topologies

3-BODY UNITARITY

$$\langle q_1, q_2, q_3 | (\hat{T} - \hat{T}^\dagger) | p_1, p_2, p_3 \rangle = i \int_P \langle q_1, q_2, q_3 | \hat{T}^\dagger | k_1, k_2, k_3 \rangle \langle k_1, k_2, k_3 | \hat{T} | p_1, p_2, p_3 \rangle$$

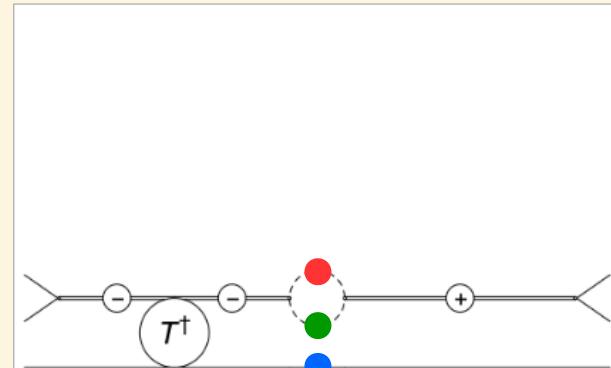
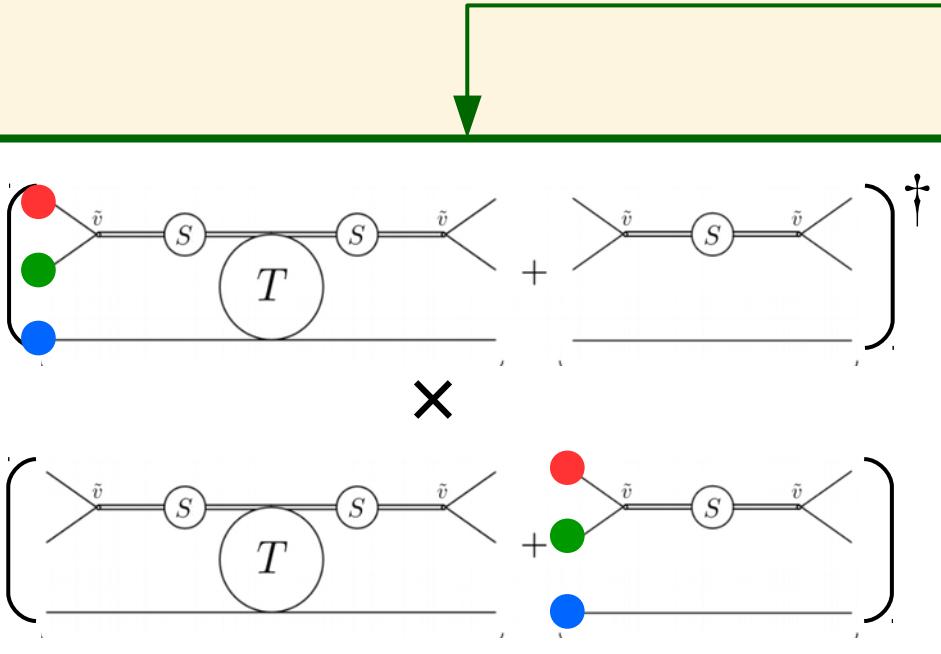
conn. + disc. parts



3-BODY UNITARITY

$$\langle q_1, q_2, q_3 | (\hat{T} - \hat{T}^\dagger) | p_1, p_2, p_3 \rangle = i \int_P \langle q_1, q_2, q_3 | \hat{T}^\dagger | k_1, k_2, k_3 \rangle \langle k_1, k_2, k_3 | \hat{T} | p_1, p_2, p_3 \rangle$$

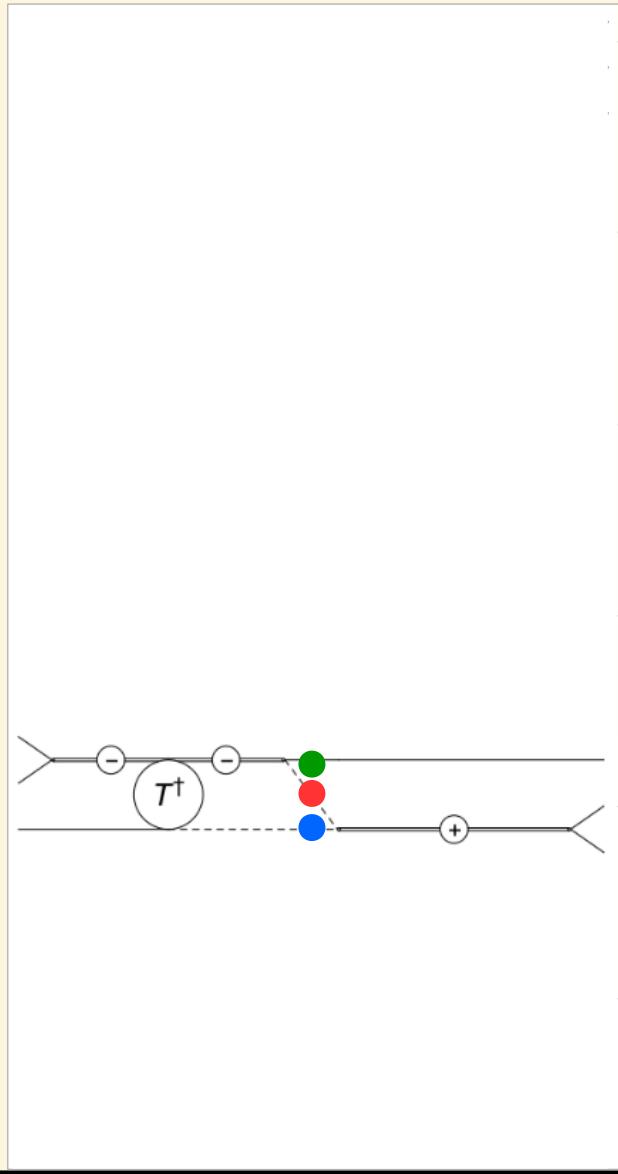
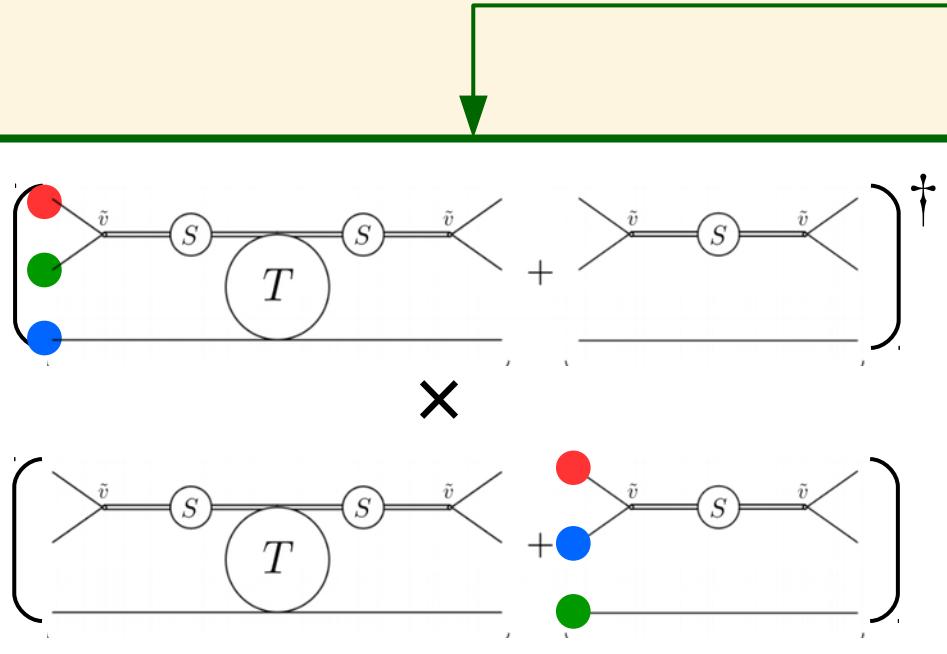
conn. + disc. parts



3-BODY UNITARITY

$$\langle q_1, q_2, q_3 | (\hat{T} - \hat{T}^\dagger) | p_1, p_2, p_3 \rangle = i \int_P \langle q_1, q_2, q_3 | \hat{T}^\dagger | k_1, k_2, k_3 \rangle \langle k_1, k_2, k_3 | \hat{T} | p_1, p_2, p_3 \rangle$$

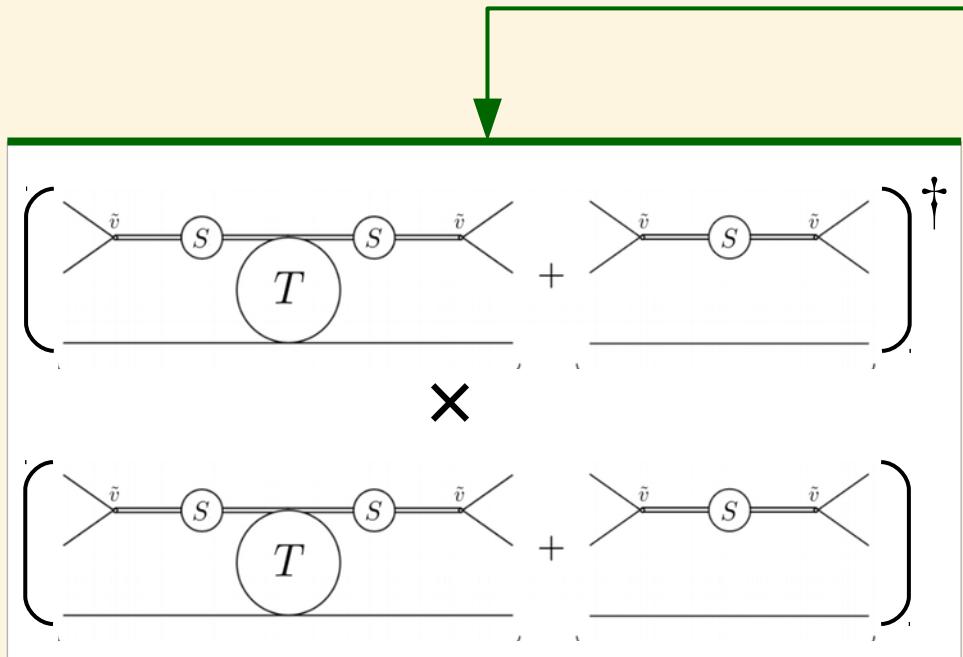
conn. + disc. parts



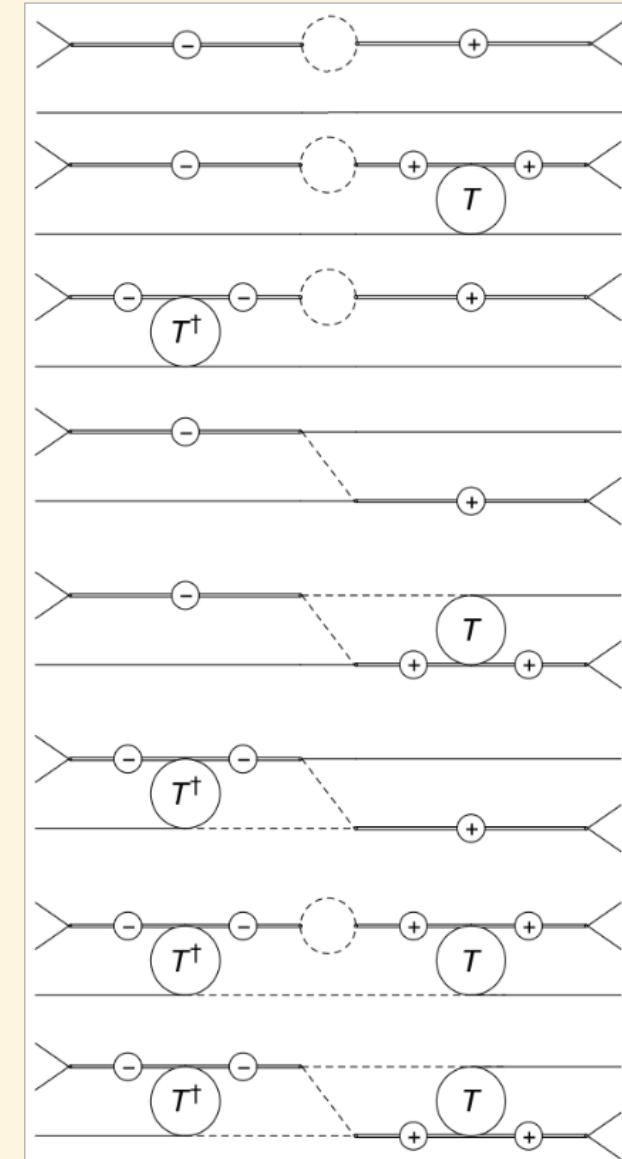
3-BODY UNITARITY

$$\langle q_1, q_2, q_3 | (\hat{T} - \hat{T}^\dagger) | p_1, p_2, p_3 \rangle = i \int_P \langle q_1, q_2, q_3 | \hat{T}^\dagger | k_1, k_2, k_3 \rangle \langle k_1, k_2, k_3 | \hat{T} | p_1, p_2, p_3 \rangle$$

conn. + disc. parts

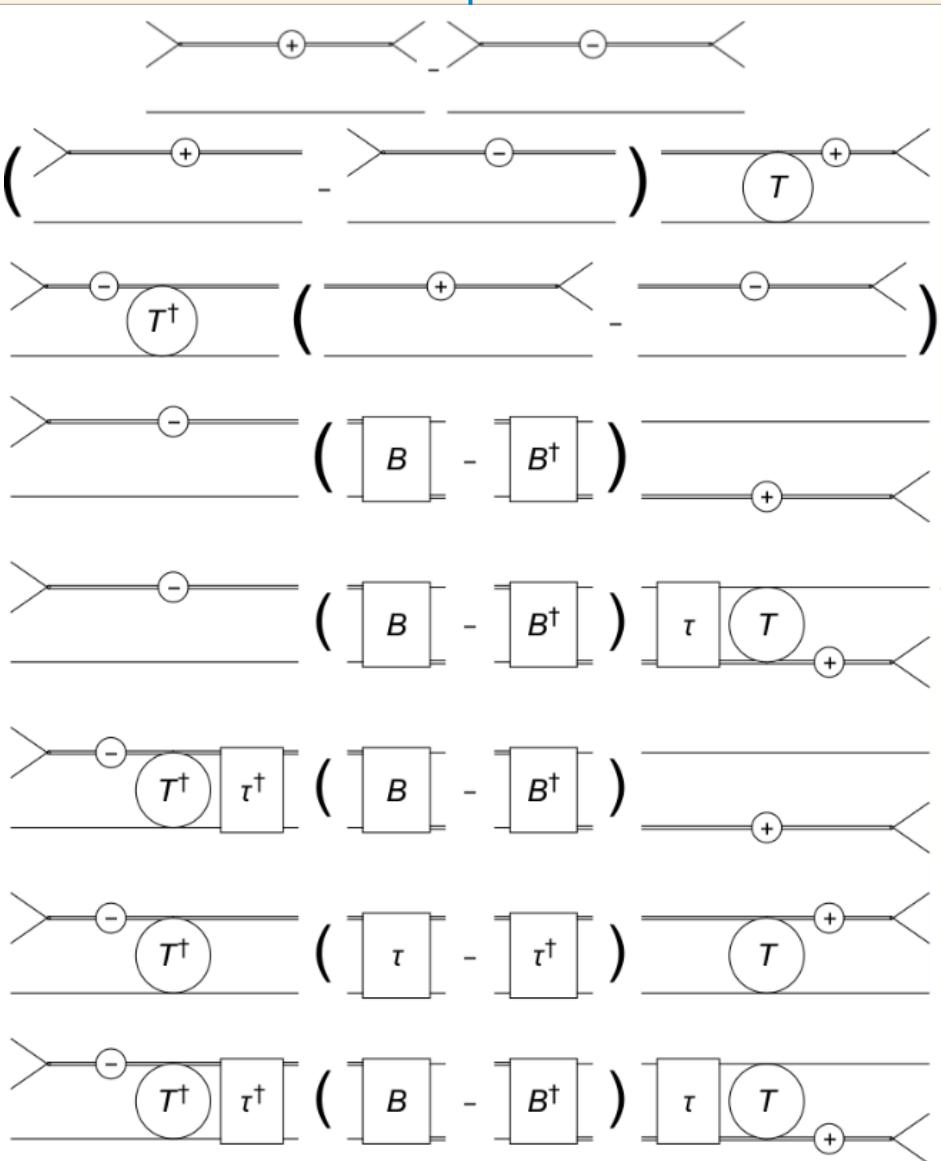


8 topologies

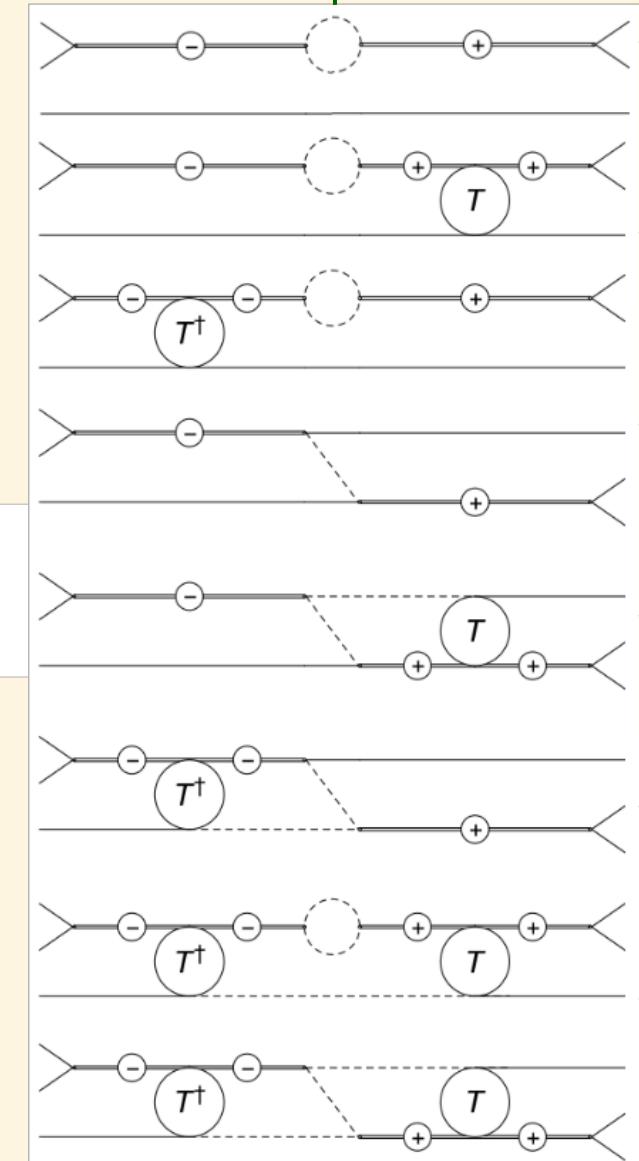


MATCHING

$$\langle q_1, q_2, q_3 | (\hat{T} - \hat{T}^\dagger) | p_1, p_2, p_3 \rangle = i \int_P \langle q_1, q_2, q_3 | \hat{T}^\dagger | k_1, k_2, k_3 \rangle \langle k_1, k_2, k_3 | \hat{T} | p_1, p_2, p_3 \rangle$$

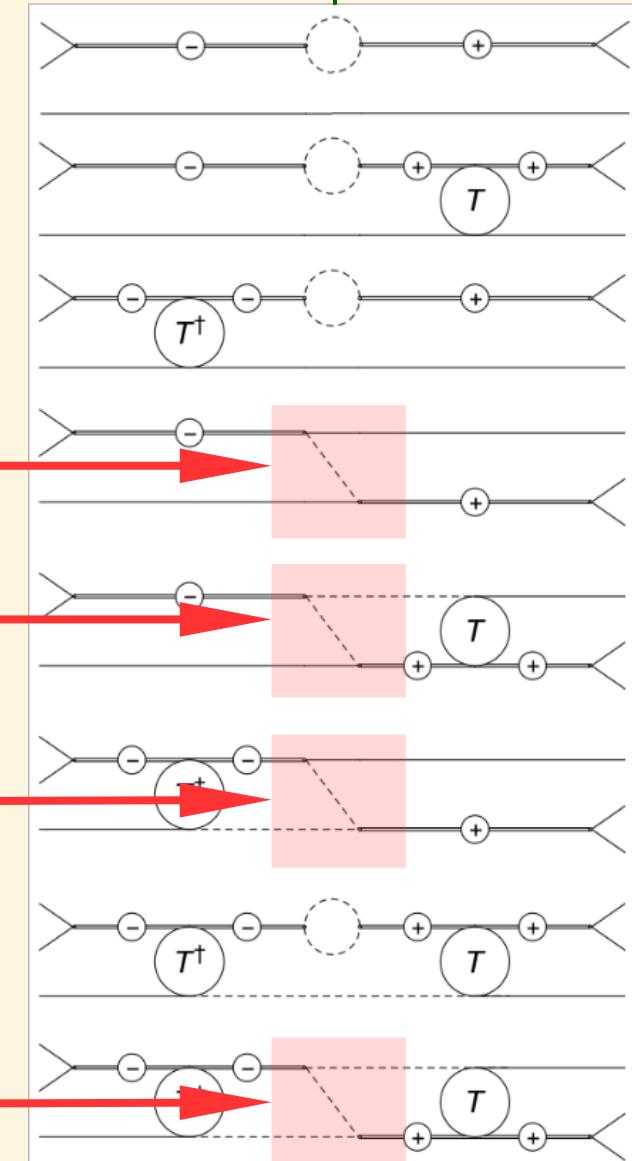
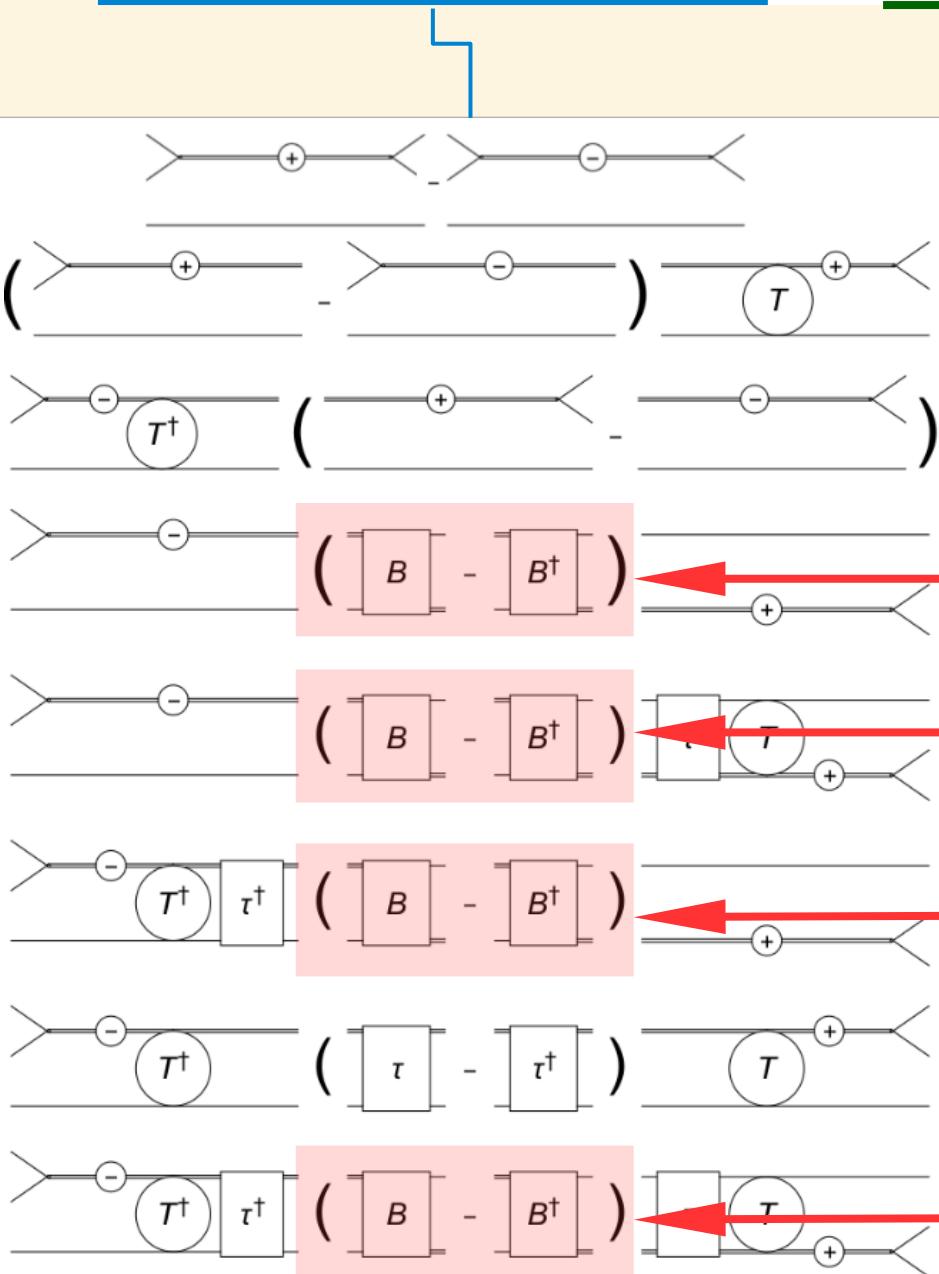


=



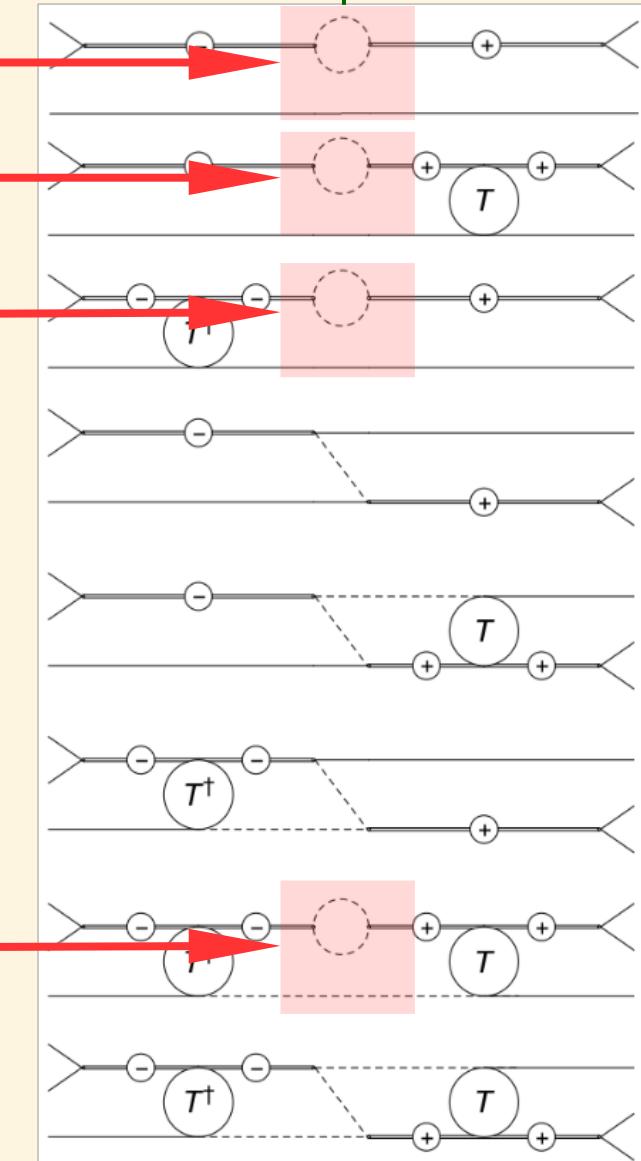
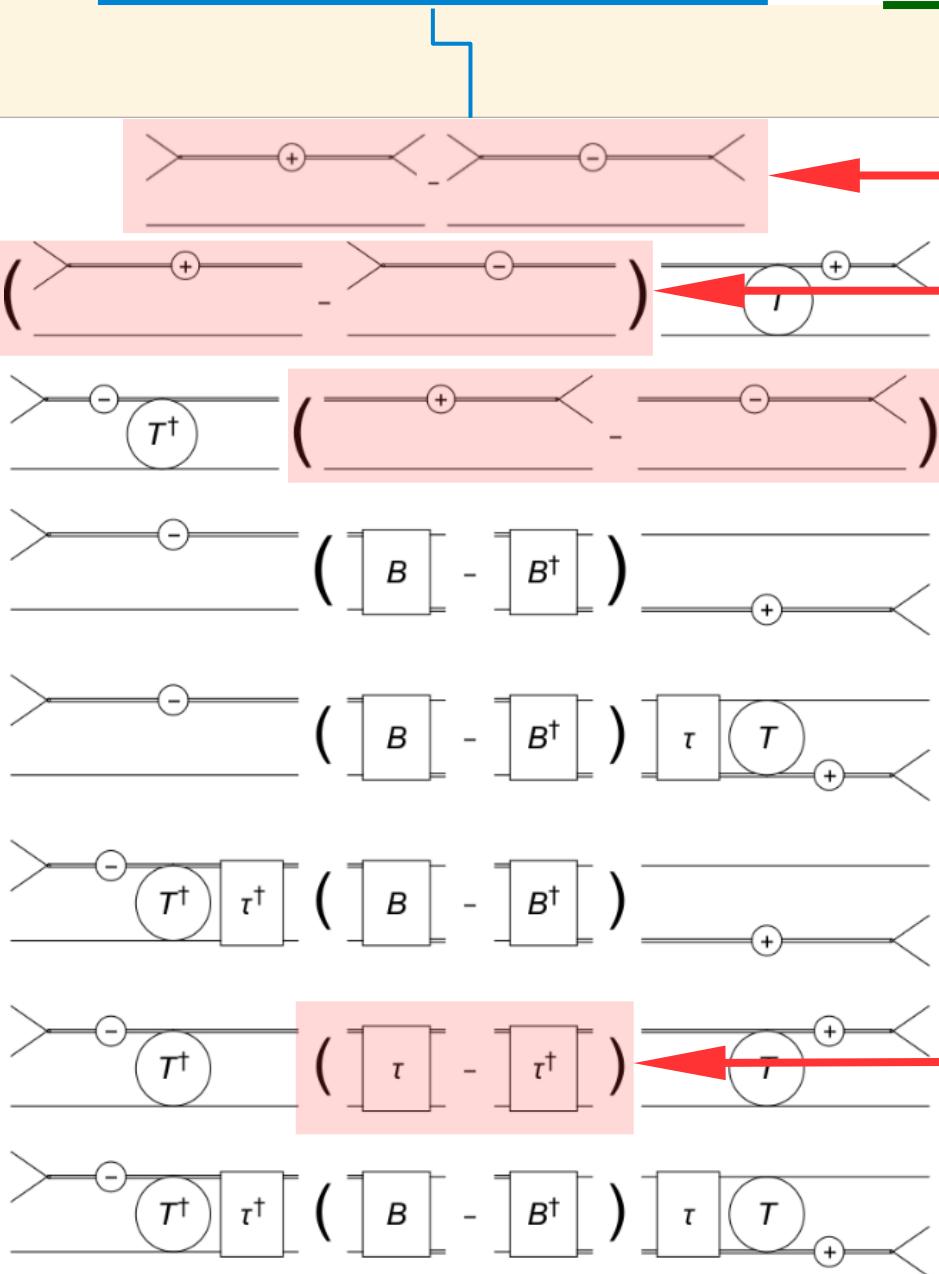
MATCHING

$$\langle q_1, q_2, q_3 | (\hat{T} - \hat{T}^\dagger) | p_1, p_2, p_3 \rangle = i \int_P \langle q_1, q_2, q_3 | \hat{T}^\dagger | k_1, k_2, k_3 \rangle \langle k_1, k_2, k_3 | \hat{T} | p_1, p_2, p_3 \rangle$$



MATCHING

$$\langle q_1, q_2, q_3 | (\hat{T} - \hat{T}^\dagger) | p_1, p_2, p_3 \rangle = i \int_P \langle q_1, q_2, q_3 | \hat{T}^\dagger | k_1, k_2, k_3 \rangle \langle k_1, k_2, k_3 | \hat{T} | p_1, p_2, p_3 \rangle$$



DISPERSION RELATION

- The only imaginary parts required by 3-body unitarity:

$$\text{Disc } B(u) = 2\pi i \frac{\delta(E_Q - \sqrt{m^2 + \mathbf{Q}^2})}{2\sqrt{m^2 + \mathbf{Q}^2}} v^2$$

$$\text{Disc } \frac{1}{S(\sigma(k))} = \frac{-i}{64\pi^2 K_{\text{cm}}} \int d^3\bar{\mathbf{K}} \frac{\delta(|\bar{\mathbf{K}}| - K_{\text{cm}})}{\sqrt{(\bar{\mathbf{K}})^2 + m^2}} v^2$$

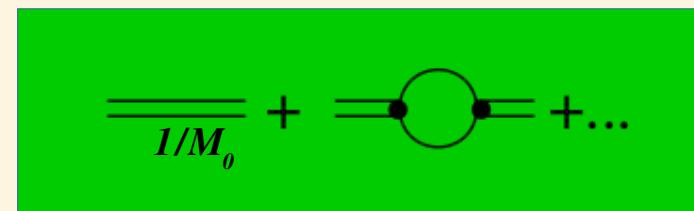
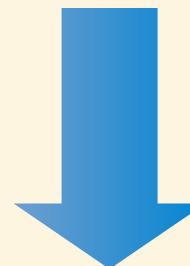
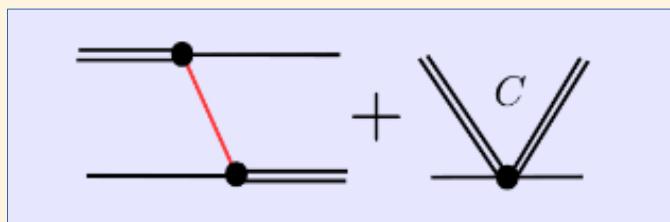
DISPERSION RELATION

- The only imaginary parts required by 3-body unitarity:

$$\text{Disc } B(u) = 2\pi i \frac{\delta(E_Q - \sqrt{m^2 + \mathbf{Q}^2})}{2\sqrt{m^2 + \mathbf{Q}^2}} v^2$$

$$\text{Disc } \frac{1}{S(\sigma(k))} = \frac{-i}{64\pi^2 K_{\text{cm}}} \int d^3\bar{\mathbf{K}} \frac{\delta(|\bar{\mathbf{K}}| - K_{\text{cm}})}{\sqrt{(\bar{\mathbf{K}})^2 + m^2}} v^2$$

- Dispersion relation



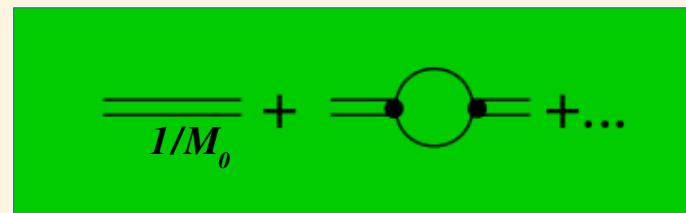
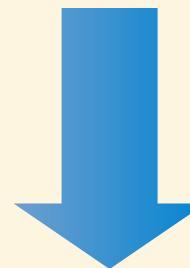
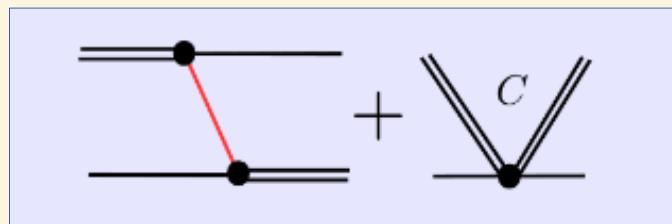
DISPERSION RELATION

- The only imaginary parts required by 3-body unitarity:

$$\text{Disc } B(u) = 2\pi i \frac{\delta(E_Q - \sqrt{m^2 + \mathbf{Q}^2})}{2\sqrt{m^2 + \mathbf{Q}^2}} v^2$$

$$\text{Disc } \frac{1}{S(\sigma(k))} = \frac{-i}{64\pi^2 K_{\text{cm}}} \int d^3\bar{\mathbf{K}} \frac{\delta(|\bar{\mathbf{K}}| - K_{\text{cm}})}{\sqrt{(\bar{\mathbf{K}})^2 + m^2}} v^2$$

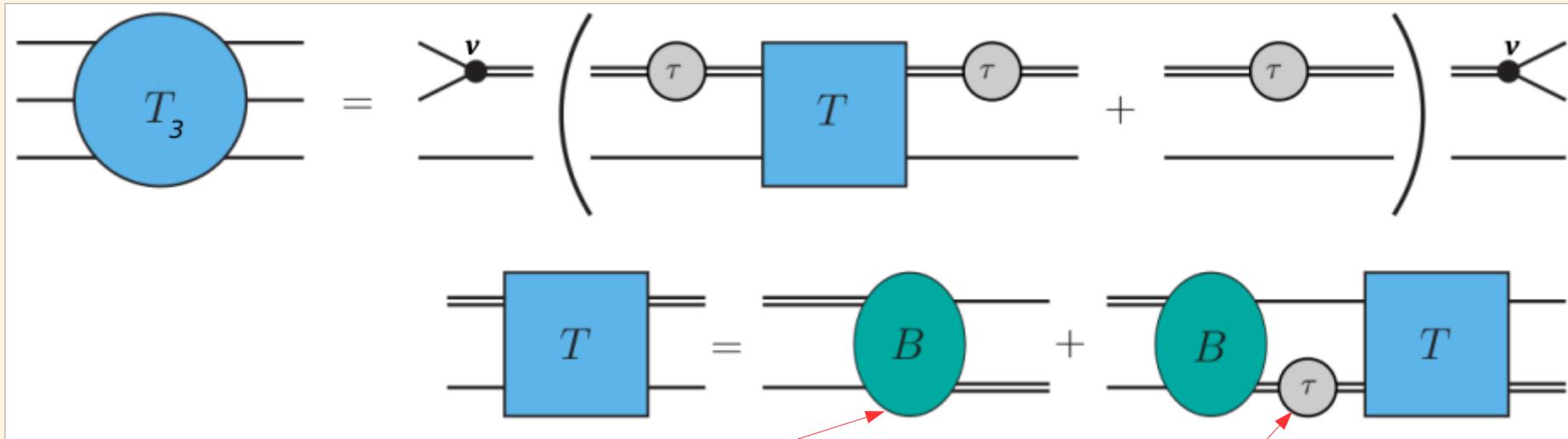
- Dispersion relation



- One-meson-exchange diagram emerges automatically \leftrightarrow **UNITARITY**
- Real constants M_0 & C are free \rightarrow input from data required

SCATTERING EQUATION

- 3-dimensional relativistic integral equation



e.g. $v=\lambda$

- un-subtracted dispersion relation

$$B(s) = \frac{-\lambda^2}{2\sqrt{m^2 + \mathbf{Q}^2} \left(E_Q - \sqrt{m^2 + \mathbf{Q}^2} + i\epsilon \right)} + C$$

- one- π exchange in TOPT \rightarrow **RESULT !**

e.g. $v=\lambda$

- twice subtracted dispersion relation in invariant mass - $\sigma(k)$

$$\frac{1}{\tau} = \sigma(k) - M_0^2 - \int \frac{d^3\ell}{(2\pi)^3} \frac{\lambda^2}{2E_\ell(\sigma(k)-4E_\ell^2+i\epsilon)}$$

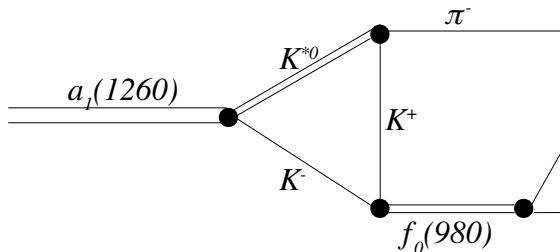
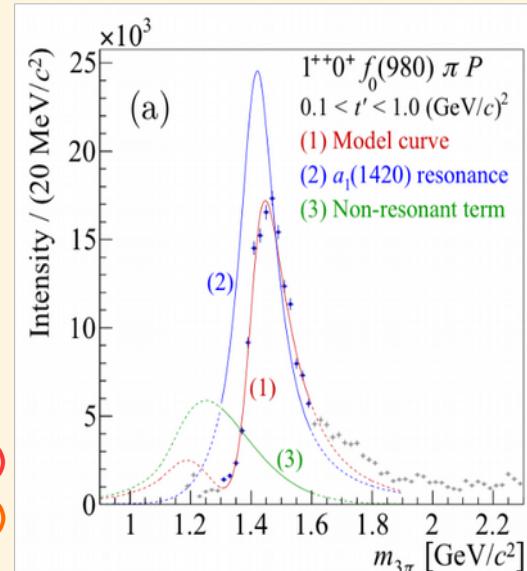
- in the rest-frame of isobar (*Lorentz invariance!*)

RECENT DEVELOPMENTS

- Interesting application: $a_1(1420)$
 - observed in COMPASS@CERN in $f_0(980)\pi$ final state
 - one explanation:
log-like behavior of the “triangle-diagram”

Mikhasenko, Ketzer, Sarantsev (2015)

Aceti, Dai, Oset (2016)

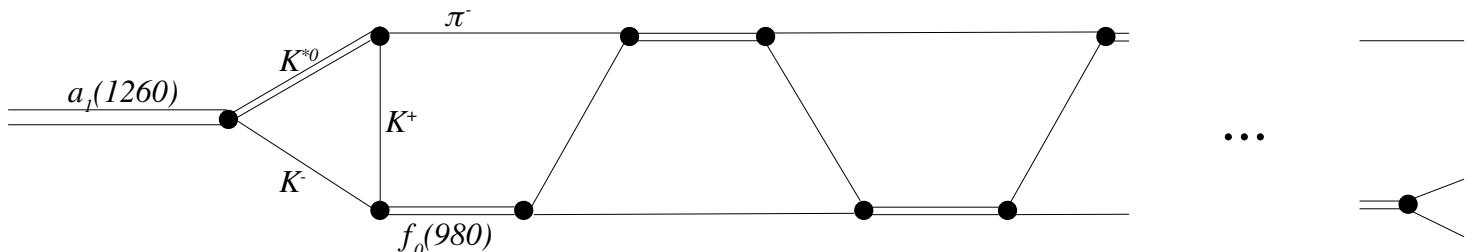
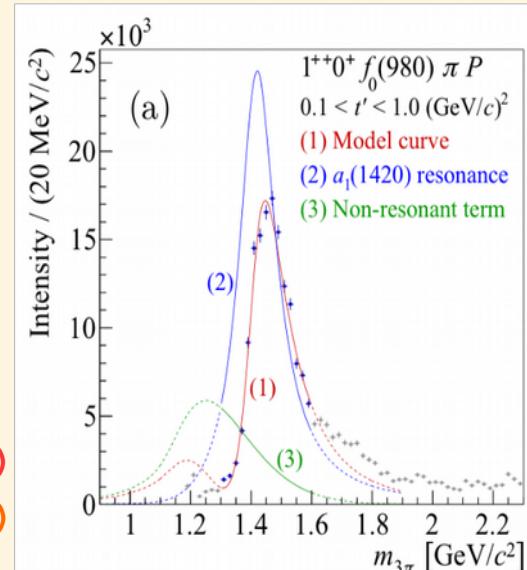


RECENT DEVELOPMENTS

- Interesting application: $a_1(1420)$
 - observed in COMPASS@CERN in $f_0(980)\pi$ final state
 - one explanation:
log-like behavior of the “triangle-diagram”

Mikhasenko, Ketzer, Sarantsev (2015)

Aceti, Dai, Oset (2016)



- Q: Does such a feature exist in full 3b-unitary FSI?

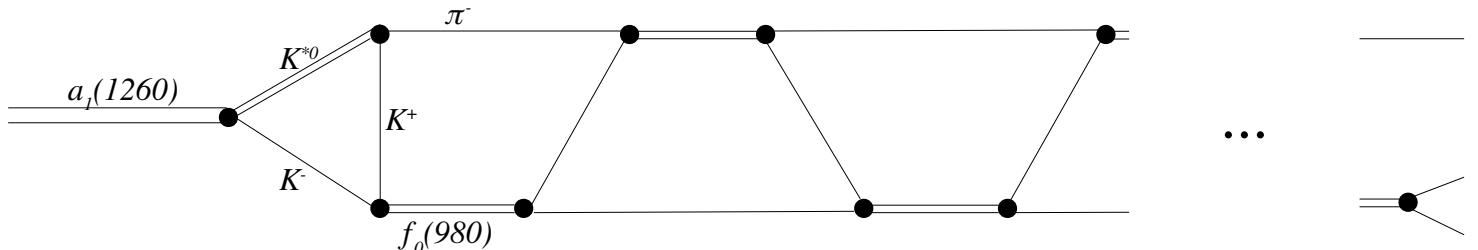
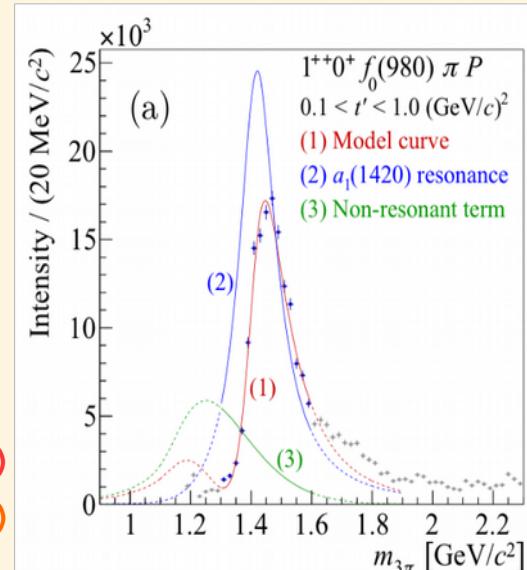
Sadisavan et al. (in progress)

RECENT DEVELOPMENTS

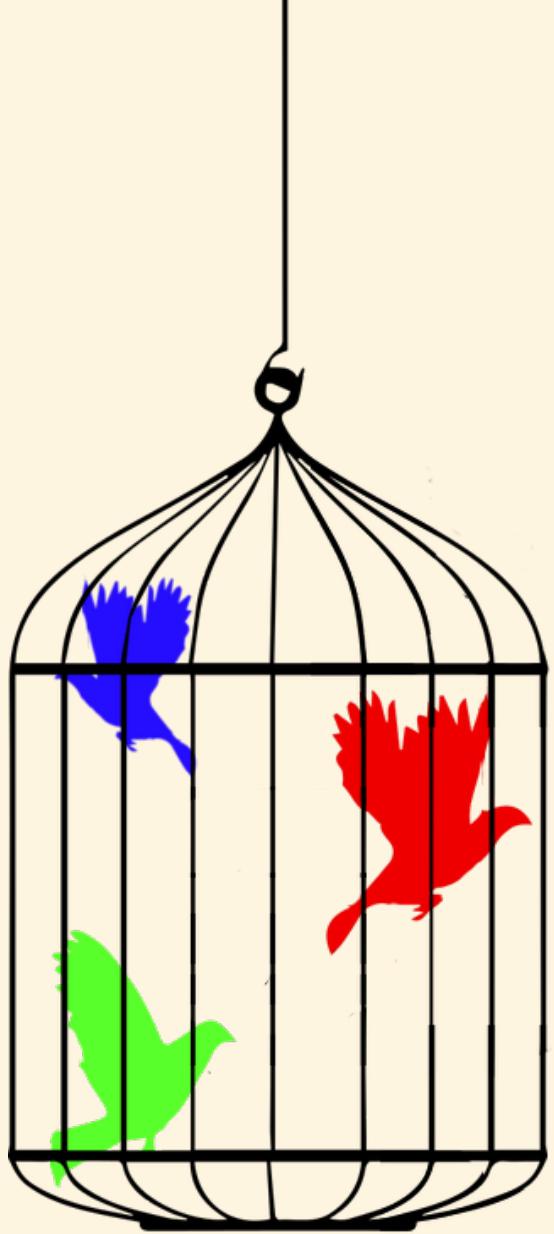
- Interesting application: $a_1(1420)$
 - observed in *COMPASS@CERN* in $f_0(980)\pi$ final state
 - one explanation:
log-like behavior of the “triangle-diagram”

Mikhasenko, Ketzer, Sarantsev (2015)

Aceti, Dai, Oset (2016)



- Q: Does such a feature exist in full 3b-unitary FSI? Sadasivan et al. (in progress)
- Recent theoretical progress Jackura et al. (2018)
 - $\text{Re}(\Delta)$ depends on the treatment of OPE – singularities in unphys. regions...



3→3 SCATTERING IN A BOX

M.M., Doring EPJA53 (2017)

LATTICE QCD

- the only systematic approach from first principles
- new interesting developments for 3-hadron systems

Lang et al. JHEP 1404

Lang et al. PRD 95 (2017)

Woss et al. JHEP 1807 [HadSpec]

...

LATTICE QCD

- the only systematic approach from first principles
- new interesting developments for 3-hadron systems

Lang et al. JHEP 1404

Lang et al. PRD 95 (2017)

Woss et al. JHEP 1807 [HadSpec]

...

- num. calculations on discretized Euclidean space-time in finite volume (at unphys. m_{π})

LATTICE QCD

- the only systematic approach from first principles
- new interesting developments for 3-hadron systems

Lang et al. JHEP 1404

Lang et al. PRD 95 (2017)

Woss et al. JHEP 1807 [HadSpec]

...

- num. calculations on discretized Euclidean space-time in finite volume (at unphys. m_{π})

continuum limit

Chiral extrapolations



LATTICE QCD

- the only systematic approach from first principles
- new interesting developments for 3-hadron systems

Lang et al. JHEP 1404

Lang et al. PRD 95 (2017)

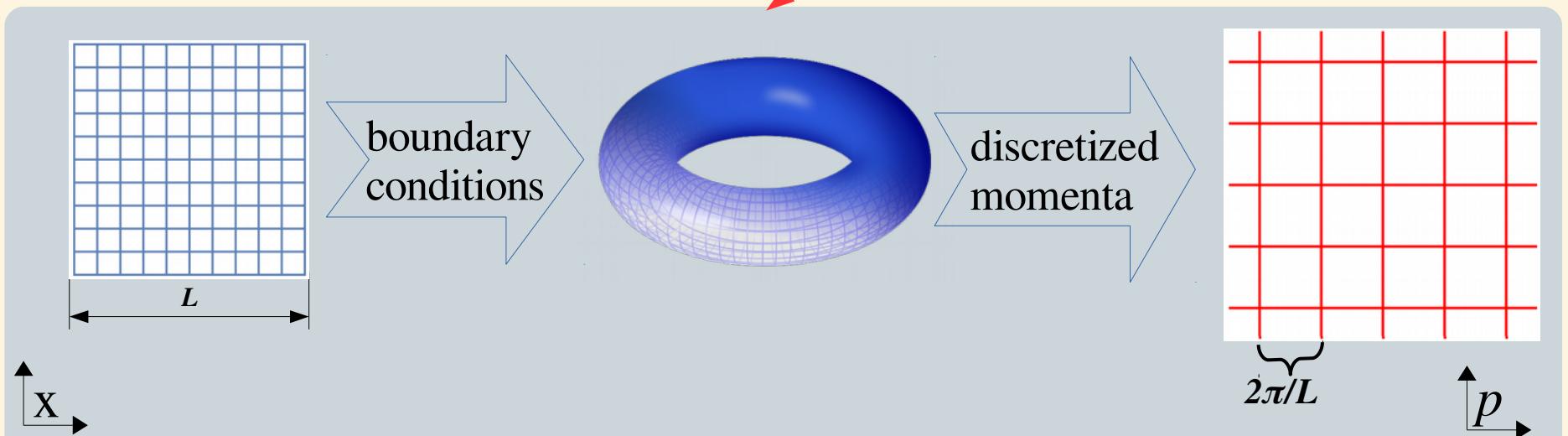
Woss et al. JHEP 1807 [HadSpec]

...

- num. calculations on discretized Euclidean space-time in finite volume (at unphys. m_π)

continuum limit

Chiral extrapolations



☞ momenta & spectra are discretized! → Quantization Condition

OVERVIEW

- **2-body case**

- well understood
- multi-channels, spin, ...

Lüscher (1986)

Gottlieb, Rummukainen, Feng, Li, Liu,
Doring, Briceno, Rusetsky, Bernard, Meissner...

OVERVIEW

- **2-body case**
 - well understood
 - multi-channels, spin, ...
- **3-body case**
 - *Lüscher*-like formalism under investigation
 - Gottlieb, Rummukainen, Feng, Li, Liu, Doring, Briceno, Rusetsky, Bernard, Meissner...*
 - Polejaeva/Rusetsky (2012)*
 - Briceño/Hansen/Sharpe (2014, 2015, 2016, 2017, 2018)*
 - Non-relativistic approaches - dimer picture & effective field theory
 - Kreuzer/Griesshammer (2012) Hammer et al. (2016, 2017) Romero/Rusetsky/Urbach et al. (2018)*

OVERVIEW

- **2-body case**
 - well understood
 - multi-channels, spin, ...
- **3-body case**
 - *Lüscher*-like formalism under investigation
 - Gottlieb, Rummukainen, Feng, Li, Liu, Doring, Briceno, Rusetsky, Bernard, Meissner...*
 - Polejaeva/Rusetsky (2012)*
 - Briceño/Hansen/Sharpe (2014, 2015, 2016, 2017, 2018)*
 - Non-relativistic approaches - dimer picture & effective field theory
 - Kreuzer/Griesshammer (2012) Hammer et al. (2016, 2017) Romero/Rusetsky/Urbach et al. (2018)*

Requirements:

- 3-body systems involve (*resonant*) two-body sub-amplitudes
- extrapolations between different energies (*problem of underdetermination*)
- all possible intermediate on-shell configurations must be identified and included to ensure all power-law finite-volume effects are taken account of

⇒ This work: *3body Quantization Condition from 3-body unitarity in isobar formulation*

MM/Doring (2017)

$$T(E) = \frac{1}{K^{-1}(E) + i\Phi(E)}$$

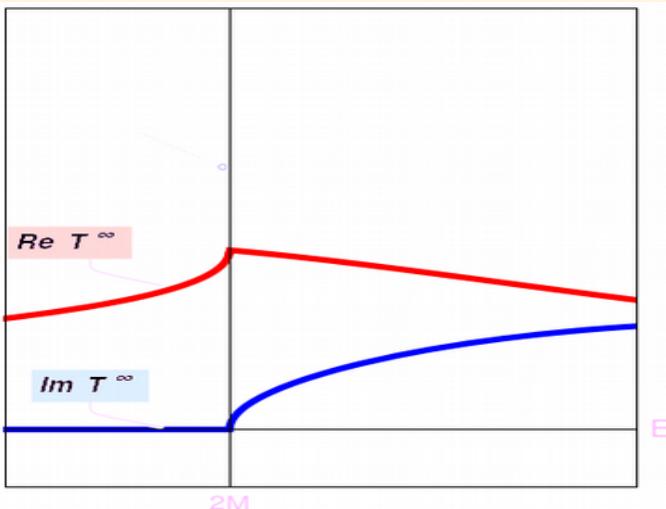
A WAY TO DERIVE 2-BODY QC

2-body Unitarity

$$T(E) = \frac{1}{K^{-1}(E) + i\Phi(E)}$$

determines imaginary parts

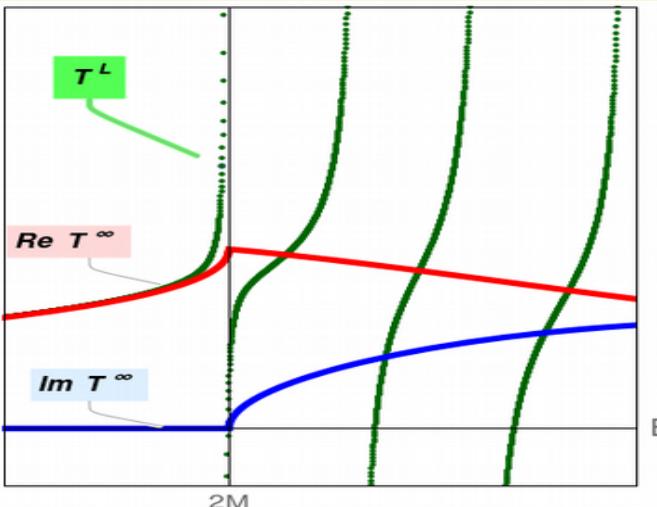
A WAY TO DERIVE 2-BODY QC



2-body Unitarity

$$T(E) = \frac{1}{K^{-1}(E) + i\Phi(E)}$$

determines imaginary parts



Discretization

Power-law fin.-vol. corrections

$$T^{FV}(E) = \frac{1}{K^{-1}(E) + \frac{2}{\sqrt{\pi L}} Z_{00}(E, L)}$$

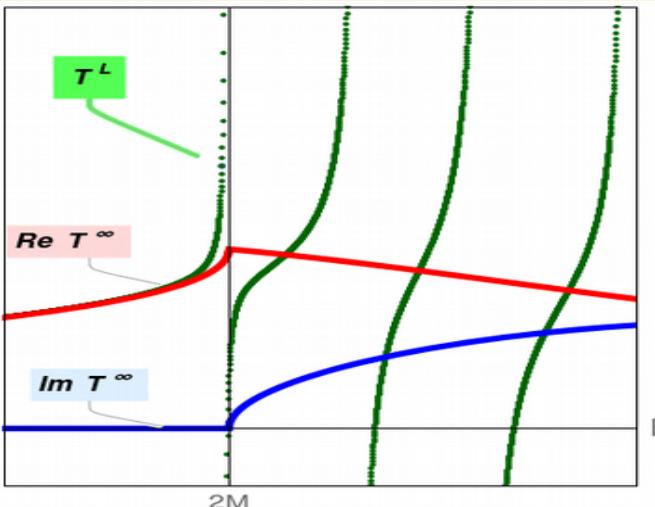
- Regular summation theorem applies for $E < 2M$
- For $E > 2M$: $T(E)$ can be singular
- LSZ formalism: energy eigenvalues \leftrightarrow pole-positions

A WAY TO DERIVE 2-BODY QC

2-body Unitarity

$$T(E) = \frac{1}{K^{-1}(E) + i\Phi(E)}$$

determines imaginary parts



Power-law fin.-vol. corrections

$$T^{FV}(E) = \frac{1}{K^{-1}(E) + \frac{2}{\sqrt{\pi L}} Z_{00}(E, L)}$$

- Regular summation theorem applies for $E < 2M$
- For $E > 2M$: $T(E)$ can be singular
- LSZ formalism: energy eigenvalues \leftrightarrow pole-positions

Quantization condition

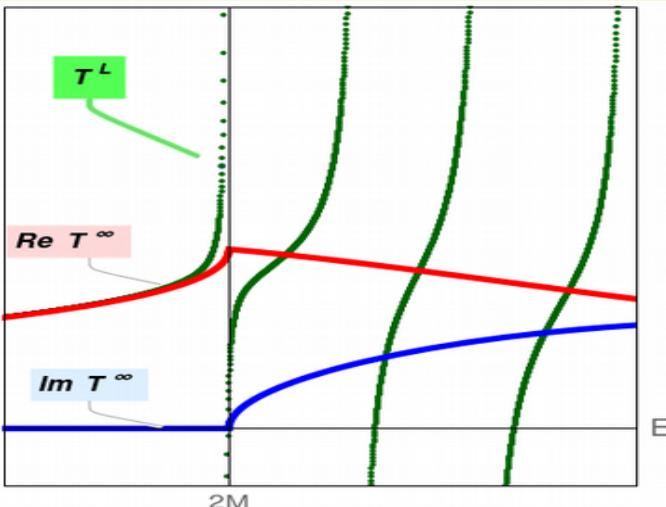
$$K^{-1}(E^*) + \frac{2}{\sqrt{\pi L}} Z_{00}(E^*, L) = 0$$

A WAY TO DERIVE 2-BODY QC

2-body Unitarity

$$T(E) = \frac{1}{K^{-1}(E) + i\Phi(E)}$$

determines imaginary parts



Power-law fin.-vol. corrections

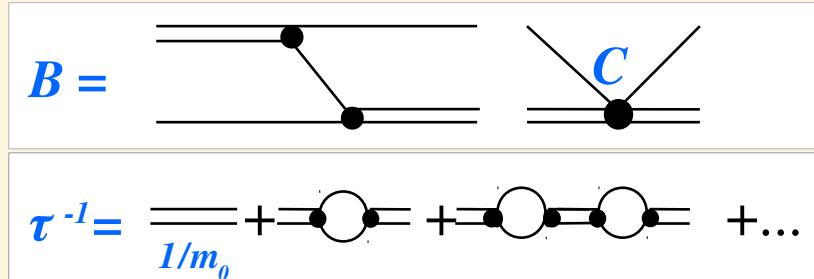
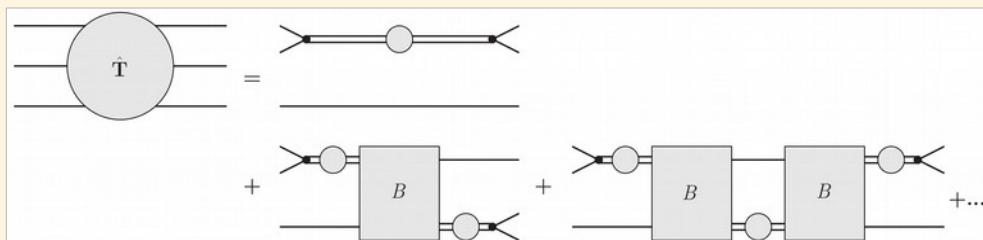
$$T^{FV}(E) = \frac{1}{K^{-1}(E) + \frac{2}{\sqrt{\pi L}} Z_{00}(E, L)}$$

- Regular summation theorem applies for $E < 2M$
- For $E > 2M$: $T(E)$ can be singular
- LSZ formalism: energy eigenvalues \leftrightarrow pole-positions

Quantization condition

$$K^{-1}(E^*) + \frac{2}{\sqrt{\pi L}} Z_{00}(E^*, L) = 0$$

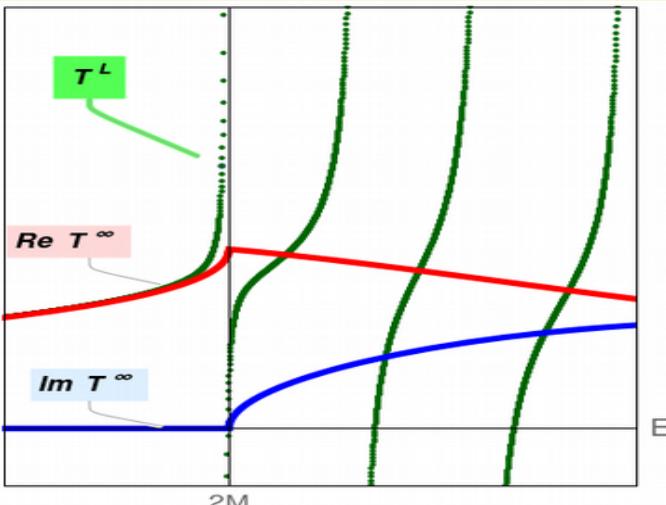
3-body Unitarity



2-body Unitarity

$$T(E) = \frac{1}{K^{-1}(E) + i\Phi(E)}$$

determines imaginary parts



Power-law fin.-vol. corrections

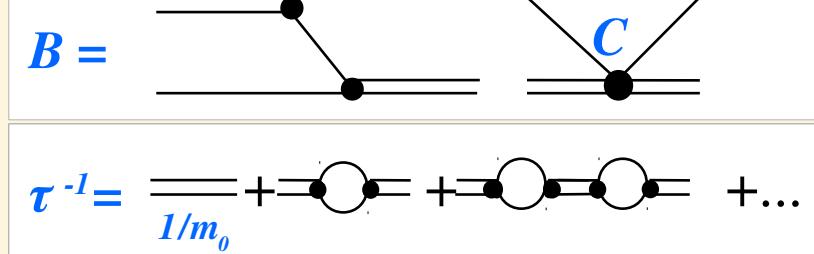
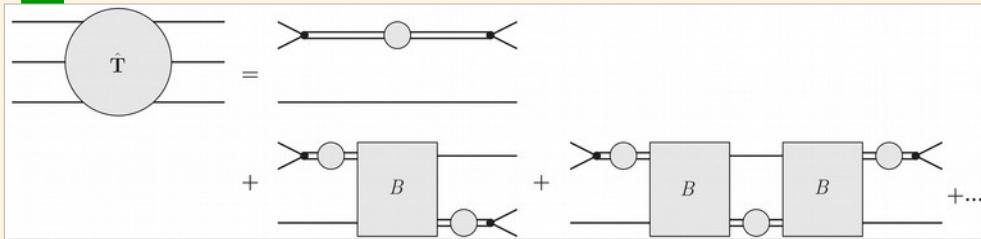
$$T^{FV}(E) = \frac{1}{K^{-1}(E) + \frac{2}{\sqrt{\pi L}} Z_{00}(E, L)}$$

- Regular summation theorem applies for $E < 2M$
- For $E > 2M$: $T(E)$ can be singular
- LSZ formalism: energy eigenvalues \leftrightarrow pole-positions

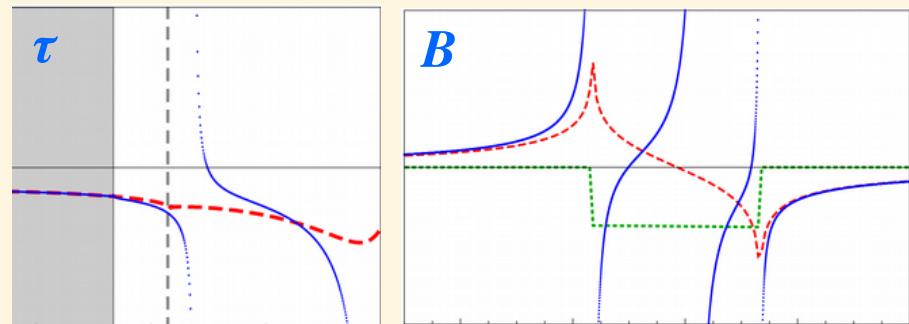
Quantization condition

$$\mathbf{K}^{-1}(\mathbf{E}^*) + \frac{2}{\sqrt{\pi L}} \mathbf{Z}_{00}(\mathbf{E}^*, \mathbf{L}) = \mathbf{0}$$

3-body Unitarity



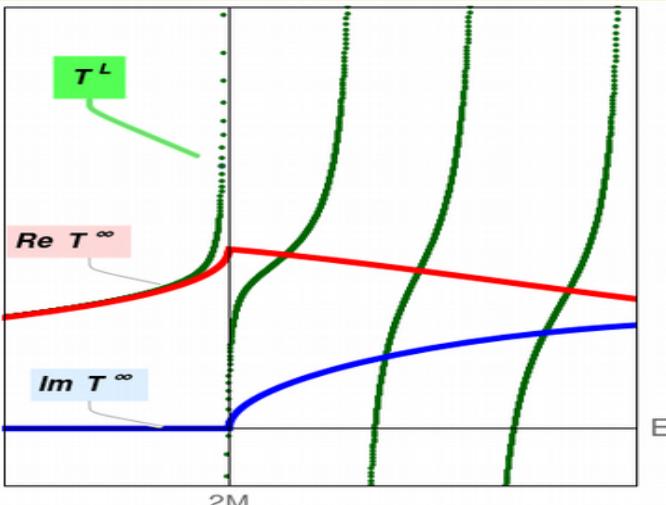
determines imaginary parts



2-body Unitarity

$$T(E) = \frac{1}{K^{-1}(E) + i\Phi(E)}$$

determines imaginary parts



Power-law fin.-vol. corrections

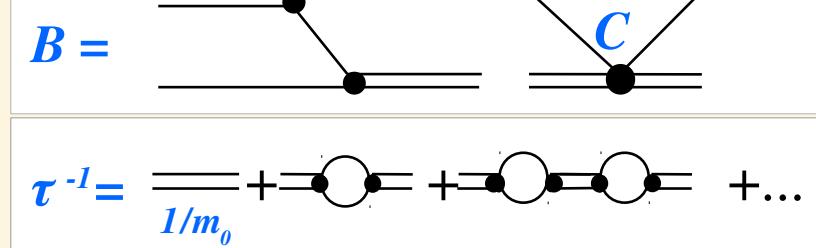
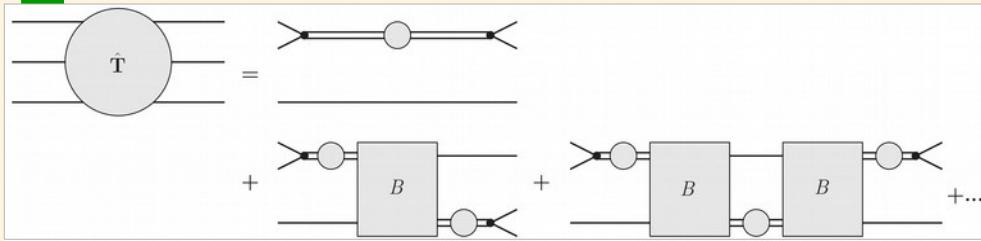
$$T^{FV}(E) = \frac{1}{K^{-1}(E) + \frac{2}{\sqrt{\pi L}} Z_{00}(E, L)}$$

- Regular summation theorem applies for $E < 2M$
- For $E > 2M$: $T(E)$ can be singular
- LSZ formalism: energy eigenvalues \leftrightarrow pole-positions

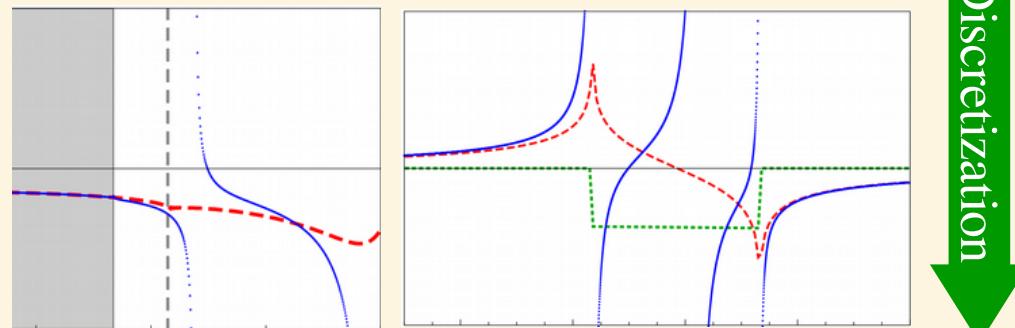
Quantization condition

$$\mathbf{K}^{-1}(\mathbf{E}^*) + \frac{2}{\sqrt{\pi L}} \mathbf{Z}_{00}(\mathbf{E}^*, \mathbf{L}) = \mathbf{0}$$

3-body Unitarity



determines imaginary parts



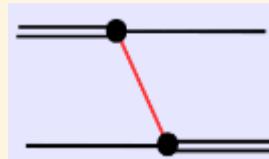
$$T_3^{FV} = v \left(\tau \frac{B}{1 + B\tau} \tau - \tau \right) v$$

Quantization condition

$$\text{Det} (\mathbf{B}(\mathbf{E}) + \tau(\mathbf{E})^{-1}) = \mathbf{0}$$

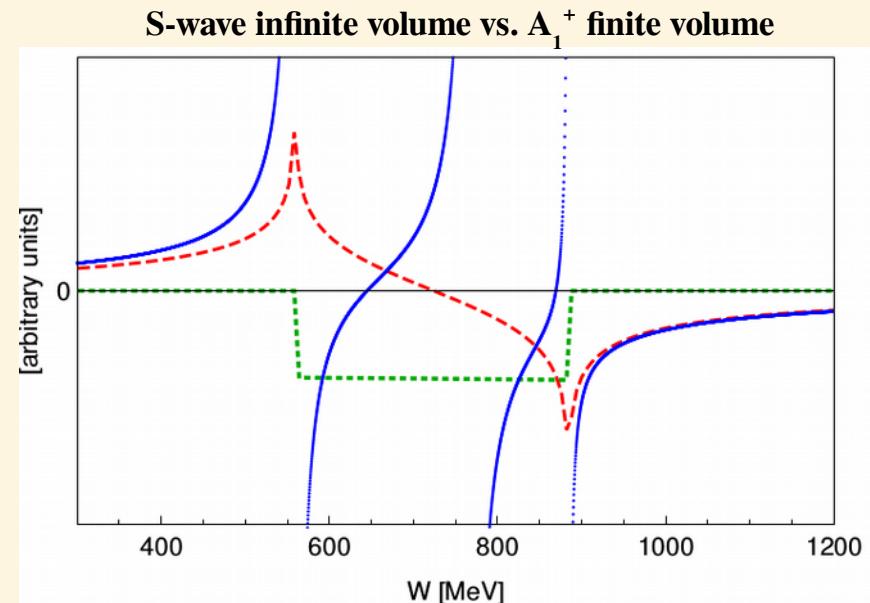
PROJECTION TO IRREPS

- High-dimensional problem
- B (OPE potential) is singular!



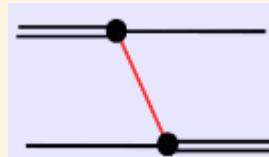
☞ Project to irreps of cubic group:

$$\Gamma \{A_p, A_2, E, T_p, T_2\}$$



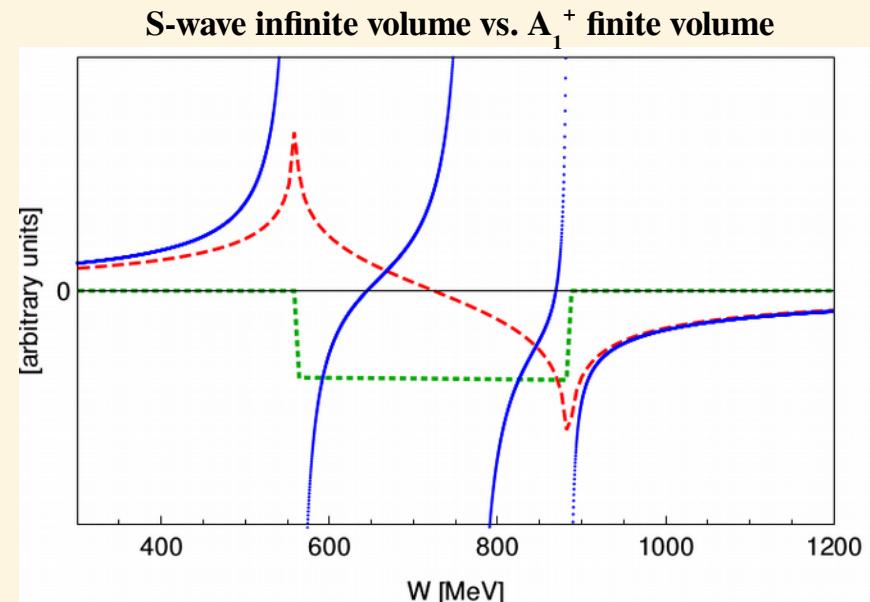
PROJECTION TO IRREPS

- High-dimensional problem
- B (OPE potential) is singular!



☞ Project to irreps of cubic group:

$$\Gamma \{A_p, A_2, E, T_p, T_2\}$$



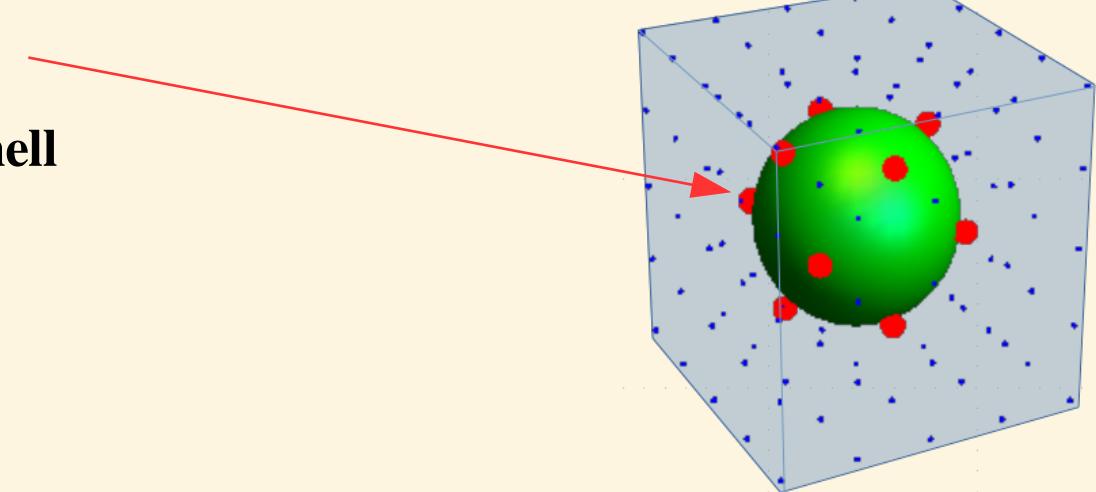
1) Separation of variables

- shells = sets of points related by O_h

2) Find the ONB of functions on each shell

$$f^s(\hat{\mathbf{p}}_j) = \sqrt{4\pi} \sum_{\Gamma\alpha} \sum_u f_u^{\Gamma\alpha s} \chi_u^{\Gamma\alpha s}(\hat{\mathbf{p}}_j)$$

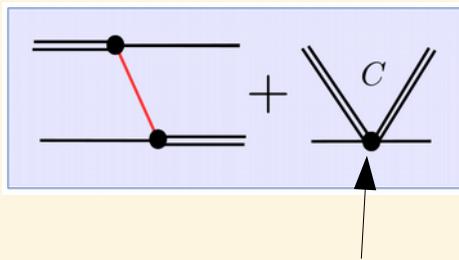
- *infinite volume analog*: PWA



Döring, Hammer, MM, Pang, Rusetsky, Wu (2018)

PROJECTED 3-BODY QC

$$\text{Det} \left(B_{uu'}^{\Gamma ss'}(\mathbf{W}^2) + \frac{2E_s L^3}{\vartheta(s)} \tau_s(\mathbf{W}^2)^{-1} \delta_{ss'} \delta_{uu'} \right) = 0$$



Fix to $3 \rightarrow 3$ data

\mathbf{W} – total energy

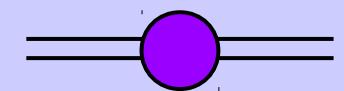
s/s' - shell index

u/u' - basis index

ϑ – multiplicity

L – lattice volume

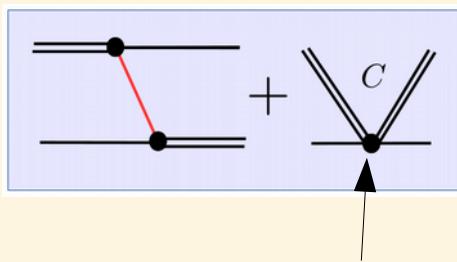
E_s – 1p. energy



Fix to $2 \rightarrow 2$ data:
 $T_{22} = v \tau v$

PROJECTED 3-BODY QC

$$\text{Det} \left(B_{uu'}^{\Gamma ss'}(\mathbf{W}^2) + \frac{2E_s L^3}{\vartheta(s)} \tau_s(\mathbf{W}^2)^{-1} \delta_{ss'} \delta_{uu'} \right) = 0$$



Fix to $3 \rightarrow 3$ data

W – total energy

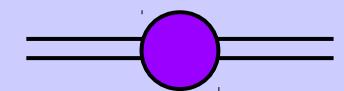
s/s' - shell index

u/u' - basis index

ϑ – multiplicity

L – lattice volume

E_s – 1p. energy

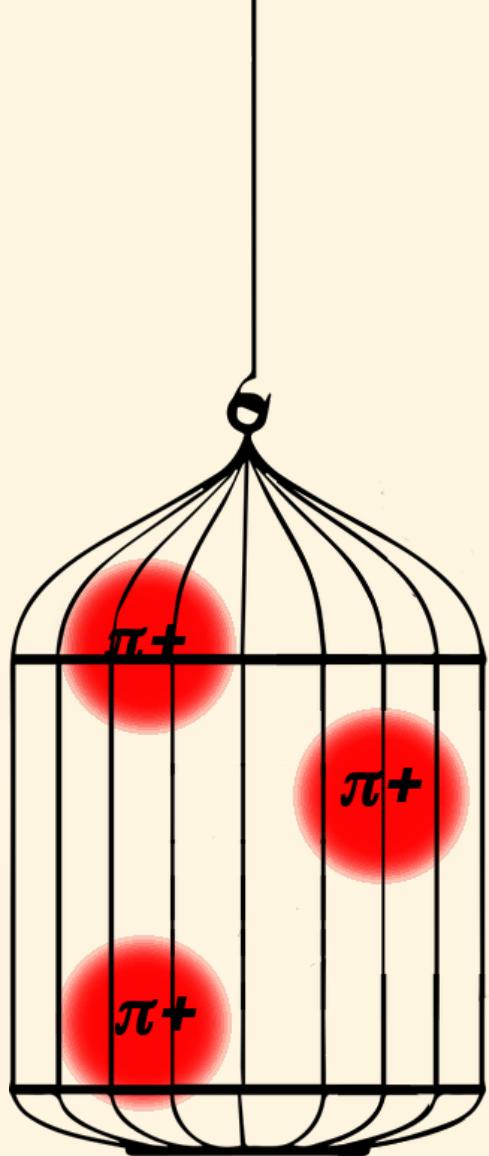


Fix to $2 \rightarrow 2$ data:
 $T_{22} = v \tau v$

- **relativistic 3-body quantization condition**
- determinant in u, s ; diagonal in Γ
- Not a Lüscher-like equation (“left”: infinite volume, “right”: finite volume)

Instead: Fix parameters to lattice eigenvalues → evaluate infinite-volume amplitude

Same work-flow as in many 2-body coupled-channel fits (e.g., Doring et al., EPJA (2012))



FINITE-VOLUME SPECTRUM OF $\pi^+\pi^+\pi^+$

MM, Doring [arXiv:1807.04746]

LQCD results from NPLQCD Detmold et al. (2008)

PHYSICAL SYSTEM

$\pi^+\pi^+$ & $\pi^+\pi^+\pi^+$ ground level @ $L=2.5\text{ fm}$ & $m_\pi=291/352/491/591\text{ MeV}$

Detmold et al. (2008)

- ☞ *Q: does the “isobar” picture hold for repulsive channel?*
 - ☞ *Bonus Q: chiral extrapolation in 3body system?*
-

PHYSICAL SYSTEM

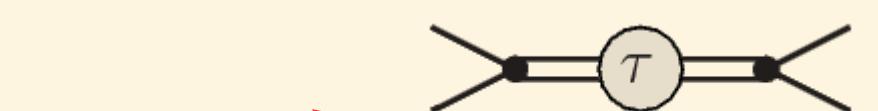
$\pi^+\pi^+$ & $\pi^+\pi^+\pi^+$ ground level @ $L=2.5\text{ fm}$ & $m_\pi=291/352/491/591\text{ MeV}$

Detmold et al. (2008)

- ☞ *Q: does the “isobar” picture hold for repulsive channel?*
- ☞ *Bonus Q: chiral extrapolation in 3body system?*

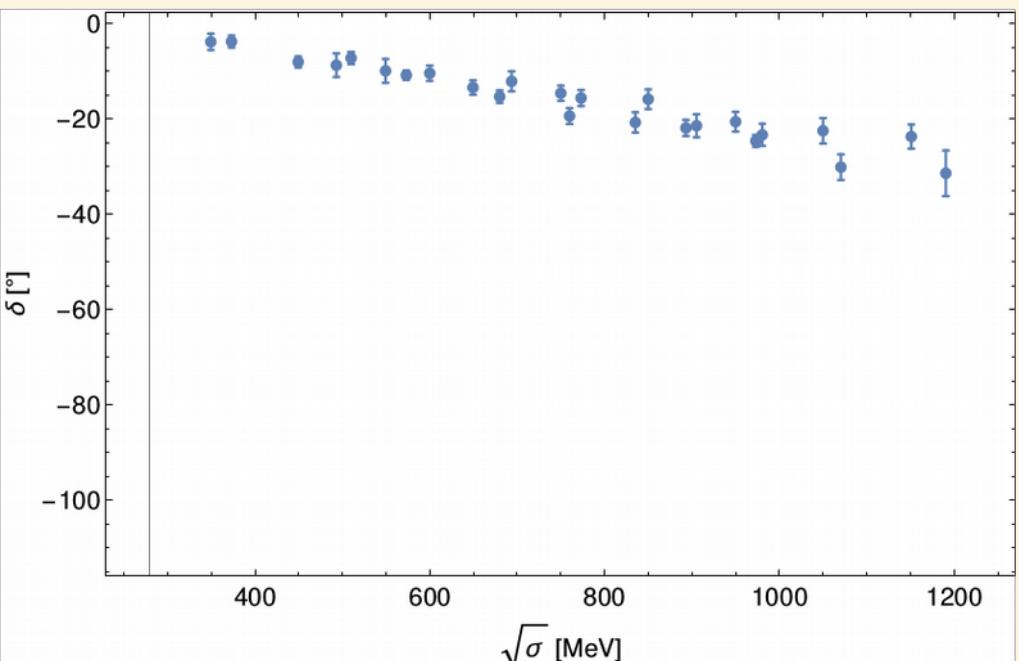
I. 2-body subchannel:

- one-channel problem: $\pi\pi$ -system in S-wave, $I=2$
- 2-body amplitude consistent with 3-body one



A Feynman diagram showing a central circular vertex labeled τ . Two incoming lines from the left and two outgoing lines to the right meet at this vertex. A red arrow points from the text "2-body amplitude consistent with 3-body one" towards this diagram.

$$T_2 = v \frac{1}{\sigma - M_0^2 - \Sigma} v$$



PHYSICAL SYSTEM

$\pi^+\pi^+$ & $\pi^+\pi^+\pi^+$ ground level @ $L=2.5\text{ fm}$ & $m_\pi=291/352/491/591\text{ MeV}$

Detmold et al. (2008)

- ☞ *Q: does the “isobar” picture hold for repulsive channel?*
- ☞ *Bonus Q: chiral extrapolation in 3body system?*

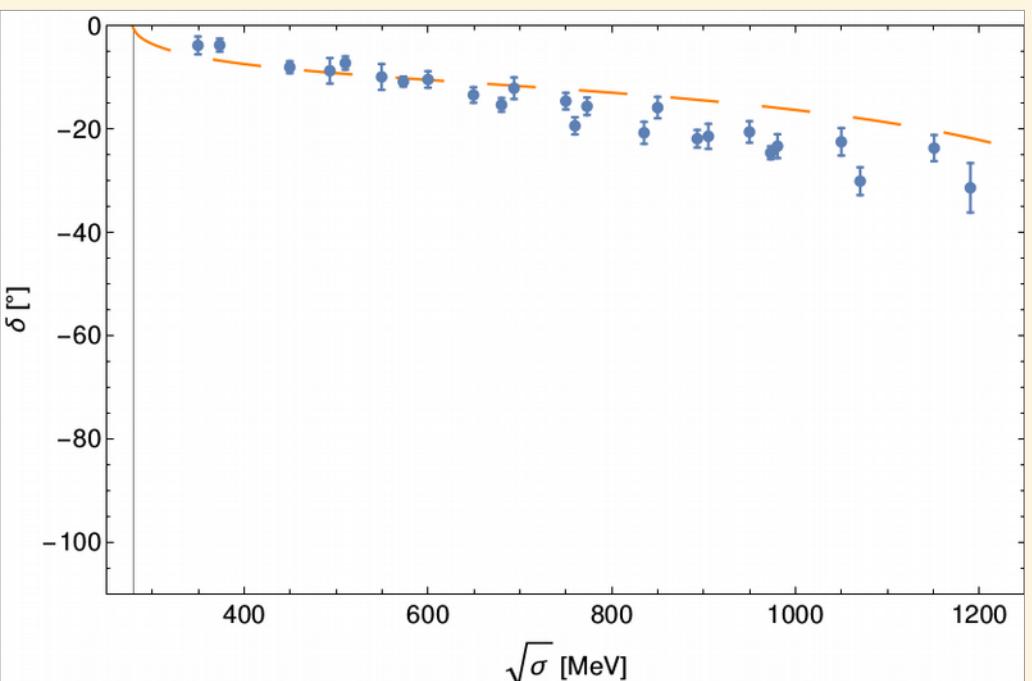
I. 2-body subchannel:

- one-channel problem: $\pi\pi$ -system in S-wave, $I=2$
- 2-body amplitude consistent with 3-body one



A Feynman diagram showing a central circular vertex labeled τ , which is connected by two horizontal lines to two external vertices. Each external vertex is connected by a diagonal line to another vertex, forming a cross-like shape.

$$T_2 = \frac{-\lambda^2/(32\pi)}{\sigma - M_0^2 - \sum_{\pm} \int \frac{d^3 k}{(2\pi)^3} \frac{\lambda^2}{4E_k \sqrt{\sigma} (\sqrt{\sigma} \pm 2E_k)}}$$



1) Fix λ, M_0 to exp. data

PHYSICAL SYSTEM

$\pi^+\pi^+$ & $\pi^+\pi^+\pi^+$ ground level @ $L=2.5\text{ fm}$ & $m_\pi=291/352/491/591\text{ MeV}$

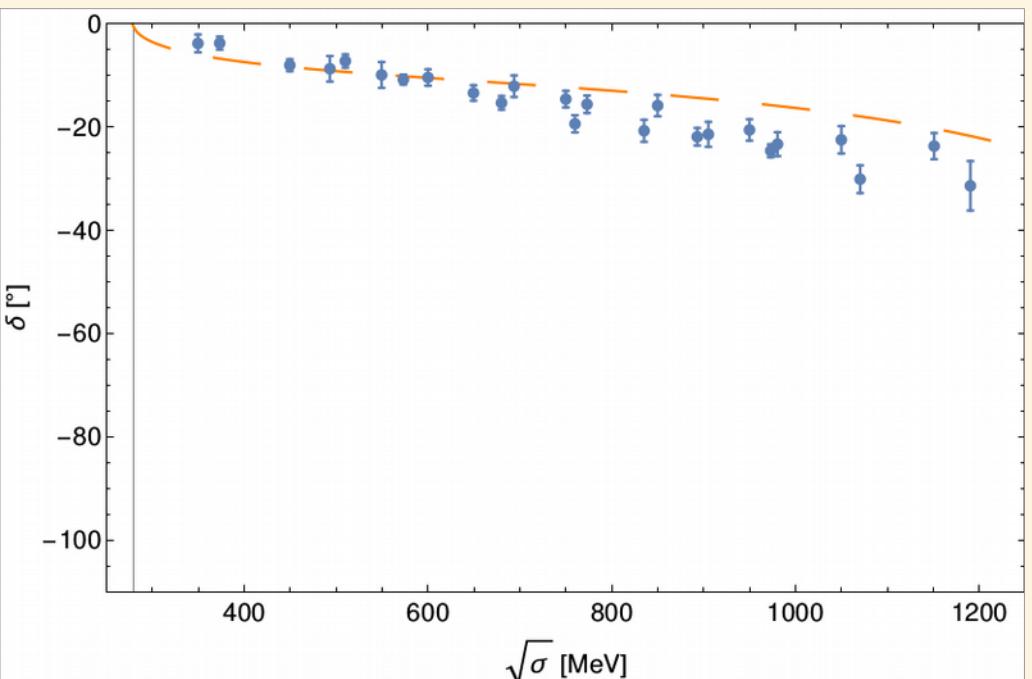
Detmold et al. (2008)

- ☞ *Q: does the “isobar” picture hold for repulsive channel?*
- ☞ *Bonus Q: chiral extrapolation in 3body system?*

I. 2-body subchannel:

- one-channel problem: $\pi\pi$ -system in S-wave, $I=2$
- 2-body amplitude consistent with 3-body one


$$T_2 = \frac{-\lambda^2/(32\pi)}{\sigma - M_0^2 - \sum_{\pm} \int \frac{d^3 k}{(2\pi)^3} \frac{\lambda^2}{4E_k \sqrt{\sigma} (\sqrt{\sigma} \pm 2E_k)}}$$



1) Fix λ, M_θ to exp. data

⌚ incorrect m_π behavior!

PHYSICAL SYSTEM

$\pi^+\pi^+$ & $\pi^+\pi^+\pi^+$ ground level @ $L=2.5\text{ fm}$ & $m_\pi=291/352/491/591\text{ MeV}$

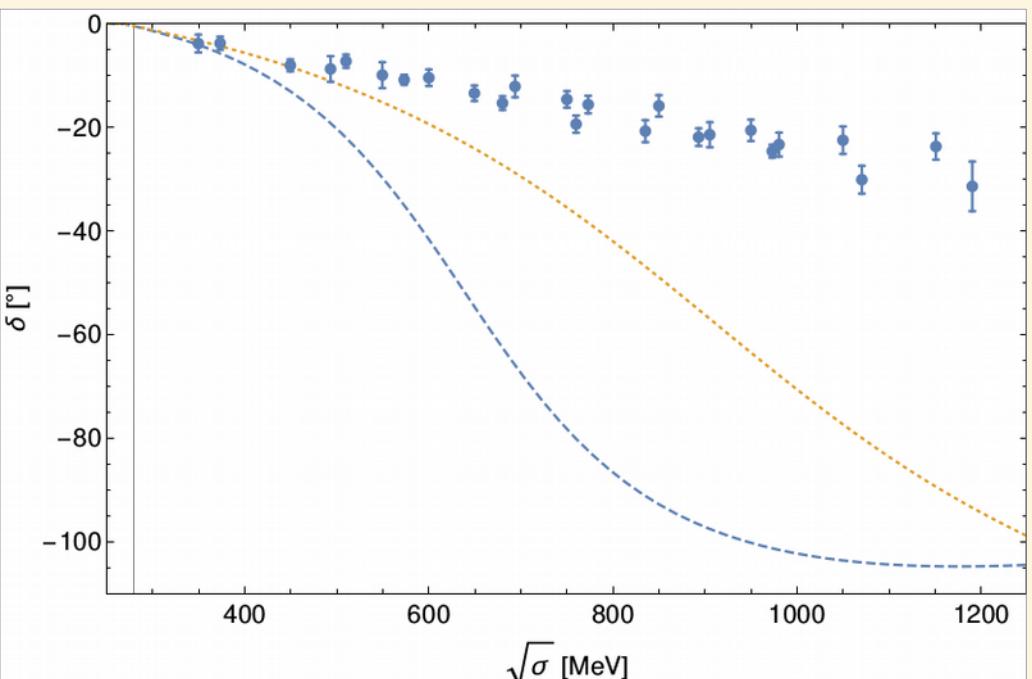
Detmold et al. (2008)

- ☞ *Q: does the “isobar” picture hold for repulsive channel?*
- ☞ *Bonus Q: chiral extrapolation in 3body system?*

I. 2-body subchannel:

- one-channel problem: $\pi\pi$ -system in S-wave, $I=2$
- 2-body amplitude consistent with 3-body one


$$T_2 = \frac{1}{K^{-1} - ip_{cms}}$$



- 1) Fix λ, M_θ to exp. data
- ⌚ incorrect m_π behavior!
- 2) Chiral NLO & K-matrix

PHYSICAL SYSTEM

$\pi^+\pi^+$ & $\pi^+\pi^+\pi^+$ ground level @ $L=2.5\text{ fm}$ & $m_\pi=291/352/491/591\text{ MeV}$

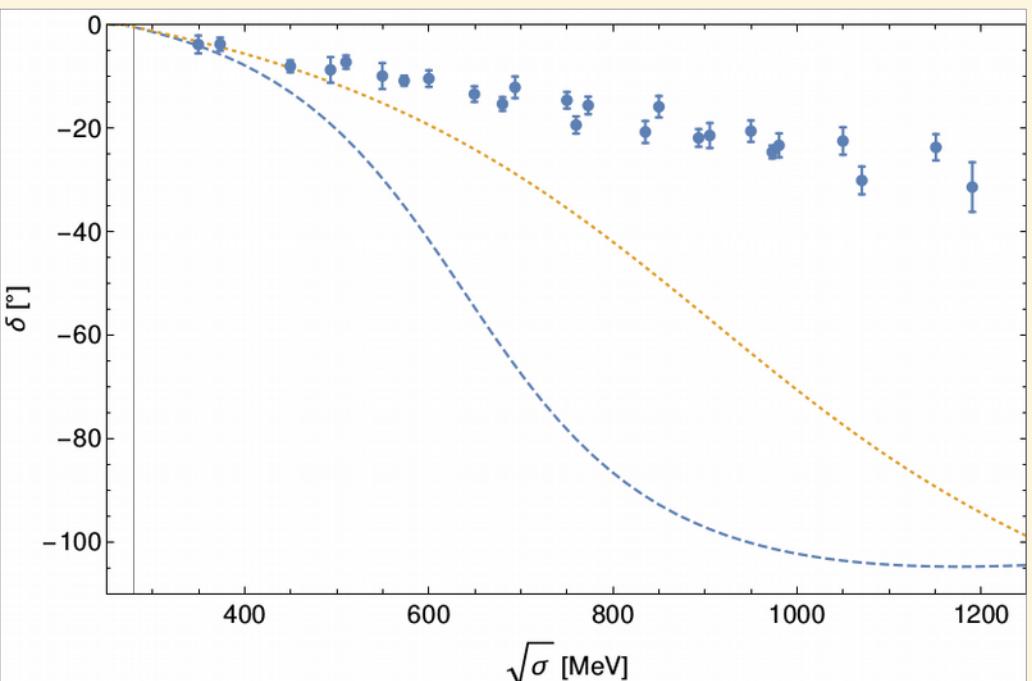
Detmold et al. (2008)

- ☞ *Q: does the “isobar” picture hold for repulsive channel?*
- ☞ *Bonus Q: chiral extrapolation in 3body system?*

I. 2-body subchannel:

- one-channel problem: $\pi\pi$ -system in S-wave, $I=2$
- 2-body amplitude consistent with 3-body one


$$T_2 = \frac{1}{K^{-1} - ip_{cms}}$$



1) Fix λ, M_θ to exp. data

⌚ incorrect m_π behavior!

2) Chiral NLO & K-matrix

⌚ works badly for high energies

PHYSICAL SYSTEM

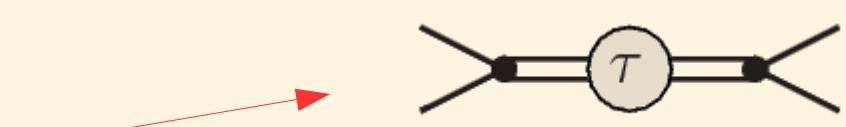
$\pi^+\pi^+$ & $\pi^+\pi^+\pi^+$ ground level @ $L=2.5\text{ fm}$ & $m_\pi=291/352/491/591\text{ MeV}$

Detmold et al. (2008)

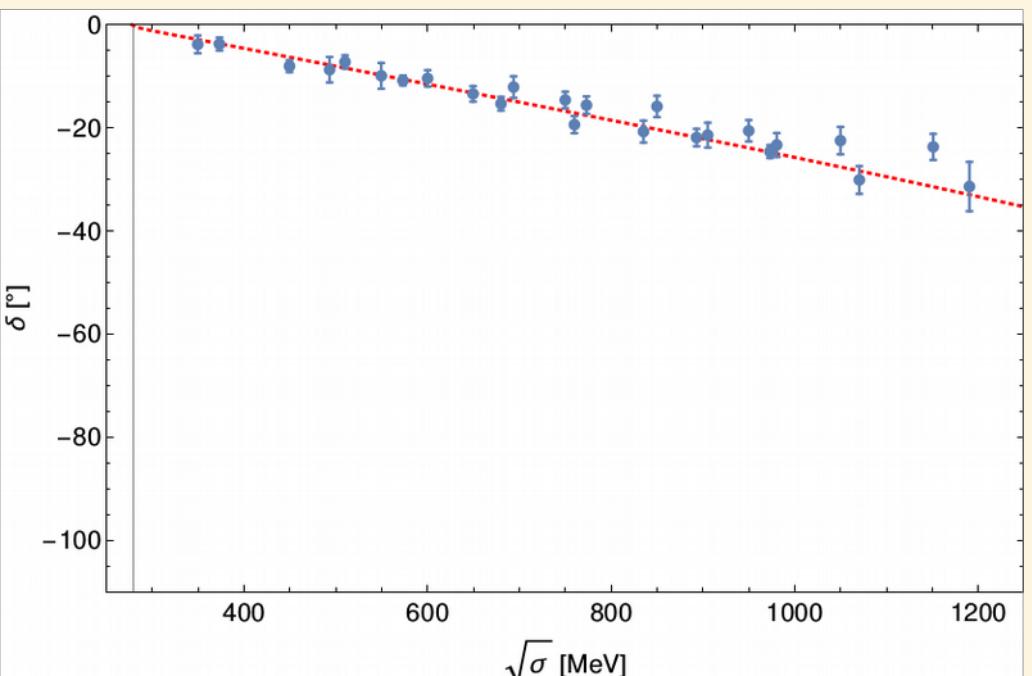
- ☞ *Q: does the “isobar” picture hold for repulsive channel?*
- ☞ *Bonus Q: chiral extrapolation in 3body system?*

I. 2-body subchannel:

- one-channel problem: $\pi\pi$ -system in S-wave, $I=2$
- 2-body amplitude consistent with 3-body one



$$T_2 = \frac{T_{\text{LO}}^2}{T_{\text{LO}} - T_{\text{NLO}}}$$



- 1) **Fix λ, M_θ to exp. data**
 - ⌚ incorrect m_π behavior!
- 2) **Chiral NLO & K-matrix**
 - ⌚ works badly for high energies
- 3) **Inverse Amplitude**
 - 😊 correct σ & m_π behavior
 - 😊 parameters known

Truong (1988)

Gasser/Leutwyler (1984)

PHYSICAL SYSTEM

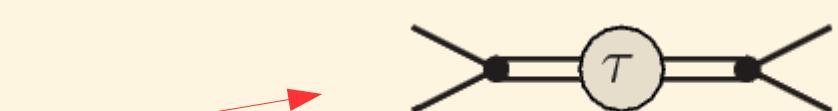
$\pi^+\pi^+$ & $\pi^+\pi^+\pi^+$ ground level @ $L=2.5\text{ fm}$ & $m_\pi=291/352/491/591\text{ MeV}$

Detmold et al. (2008)

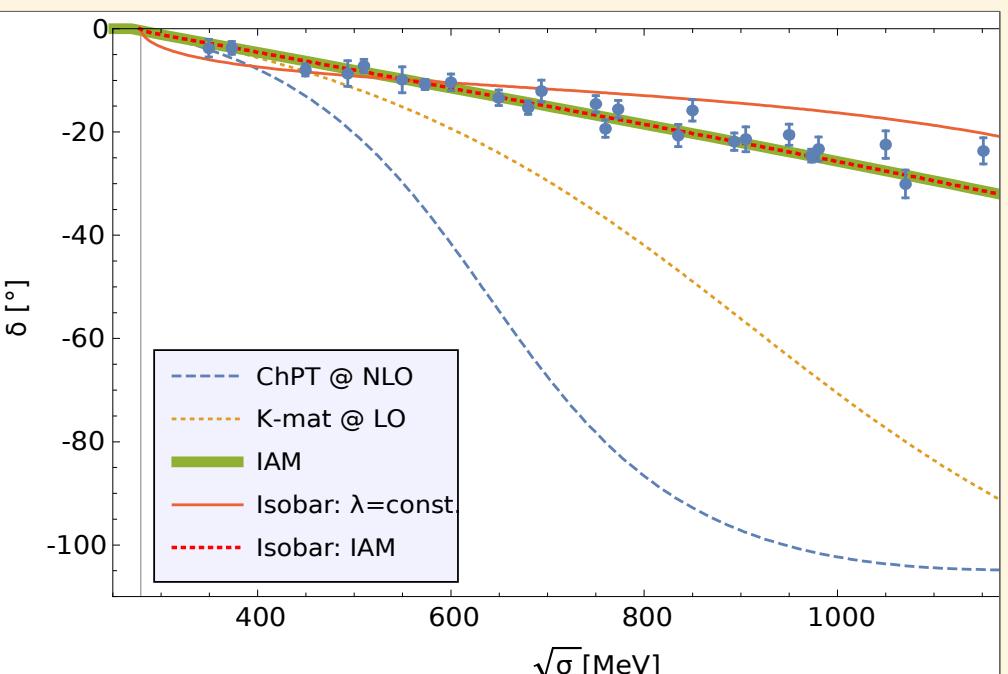
- ☞ *Q: does the “isobar” picture hold for repulsive channel?*
- ☞ *Bonus Q: chiral extrapolation in 3body system?*

I. 2-body subchannel:

- one-channel problem: $\pi\pi$ -system in S-wave, $I=2$
- 2-body amplitude consistent with 3-body one



$$T_2 = \frac{T_{\text{LO}}^2}{T_{\text{LO}} - T_{\text{NLO}}}$$



- 1) **Fix λ, M_θ to exp. data**
 - ⌚ incorrect m_π behavior!
- 2) **Chiral NLO & K-matrix**
 - ⌚ works badly for high energies
- 3) **Inverse Amplitude**
 - 😊 correct σ & m_π behavior
 - 😊 parameters known

Truong (1988)

Gasser/Leutwyler (1984)

PHYSICAL SYSTEM

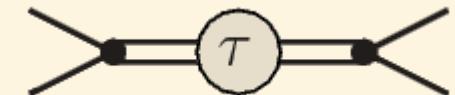
$\pi^+\pi^+$ & $\pi^+\pi^+\pi^+$ ground level @ $L=2.5\text{ fm}$ & $m_\pi=291/352/491/591\text{ MeV}$

Detmold et al. (2008)

- ☞ *Q: does the “isobar” picture hold for repulsive channel?*
- ☞ *Bonus Q: chiral extrapolation in 3body system?*

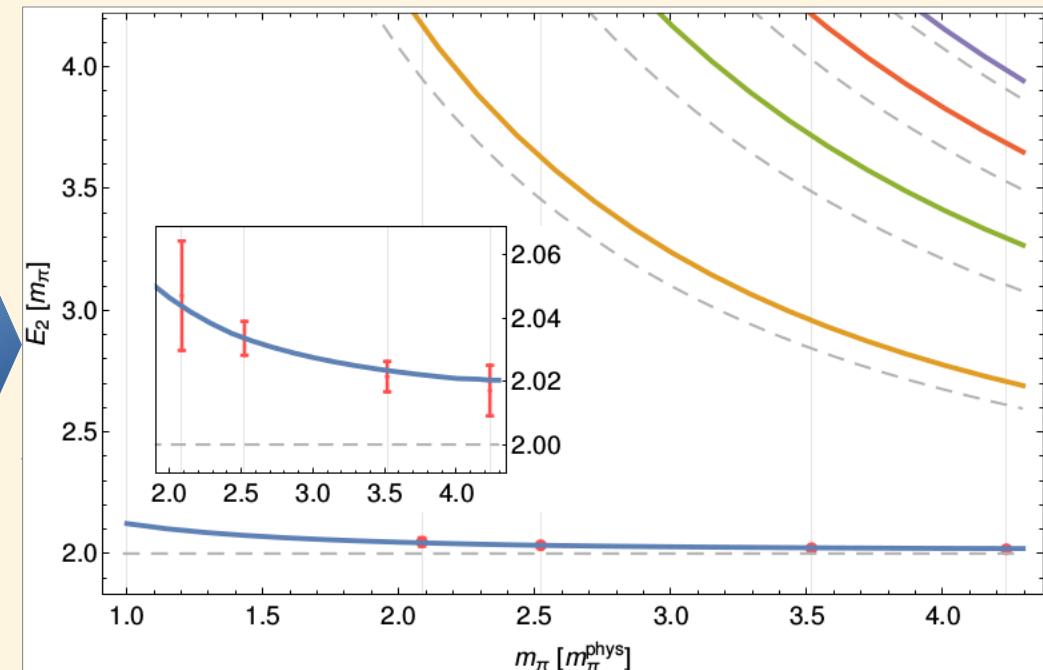
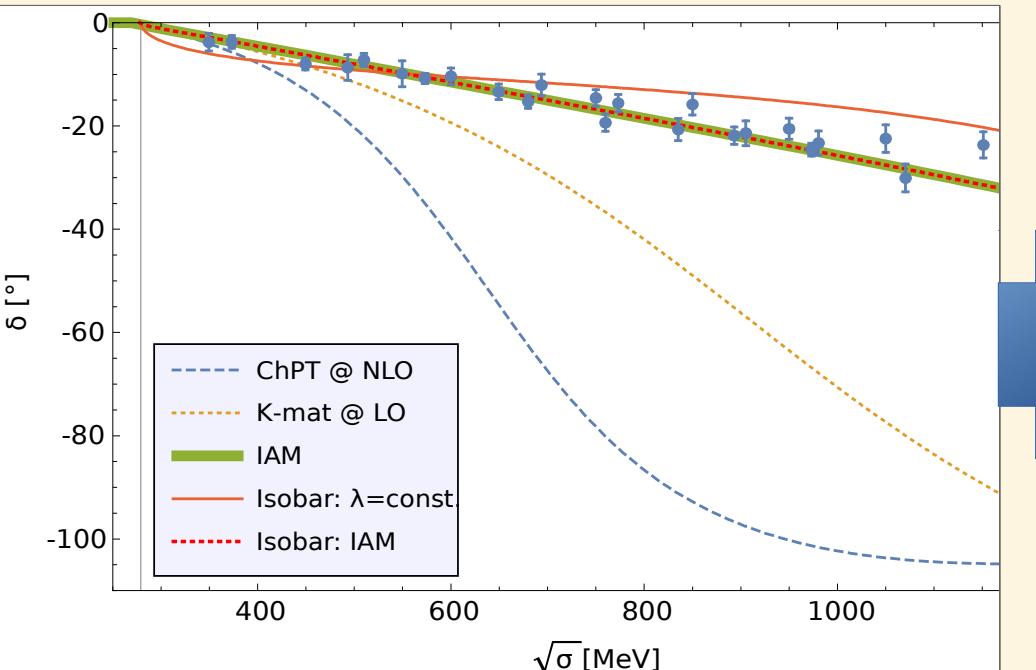
I. 2-body subchannel:

- one-channel problem: $\pi\pi$ -system in S-wave, $I=2$
- 2-body amplitude consistent with 3-body one



$$T_2 = \frac{T_{\text{LO}}^2}{T_{\text{LO}} - T_{\text{NLO}}}$$

inf-vol input for T_2 → discretize (Lüscher) → predicted fin-vol. spectrum



PHYSICAL SYSTEM

$\pi^+\pi^+$ & $\pi^+\pi^+\pi^+$ ground level @ $L=2.5\text{ fm}$ & $m_\pi=291/352/491/591\text{ MeV}$

Detmold et al. (2008)

- ☞ *Q: does the “isobar” picture hold for repulsive channel?*
- ☞ *Bonus Q: chiral extrapolation in 3body system?*

II. 3-body spectrum

Remaining unknown: $\textcolor{blue}{C}$

3-BODY QUANTIZATION CONDITION

$$\text{Det} \left(B_{uu'}^{\Gamma ss'}(\mathbf{W}^2) + \frac{2E_s L^3}{\vartheta(s)} \tau_s(\mathbf{W}^2)^{-1} \delta_{ss'} \delta_{uu'} \right) = 0$$

- genuine (momenta-dependent) 3-body “force”
- simplest case: $C_{qp} = c \delta^{(3)}(\mathbf{p}-\mathbf{q})$



PHYSICAL SYSTEM

$\pi^+\pi^+$ & $\pi^+\pi^+\pi^+$ ground level @ $L=2.5\text{ fm}$ & $m_\pi=291/352/491/591\text{ MeV}$

Detmold et al. (2008)

- ☞ *Q: does the “isobar” picture hold for repulsive channel?*
- ☞ *Bonus Q: chiral extrapolation in 3body system?*

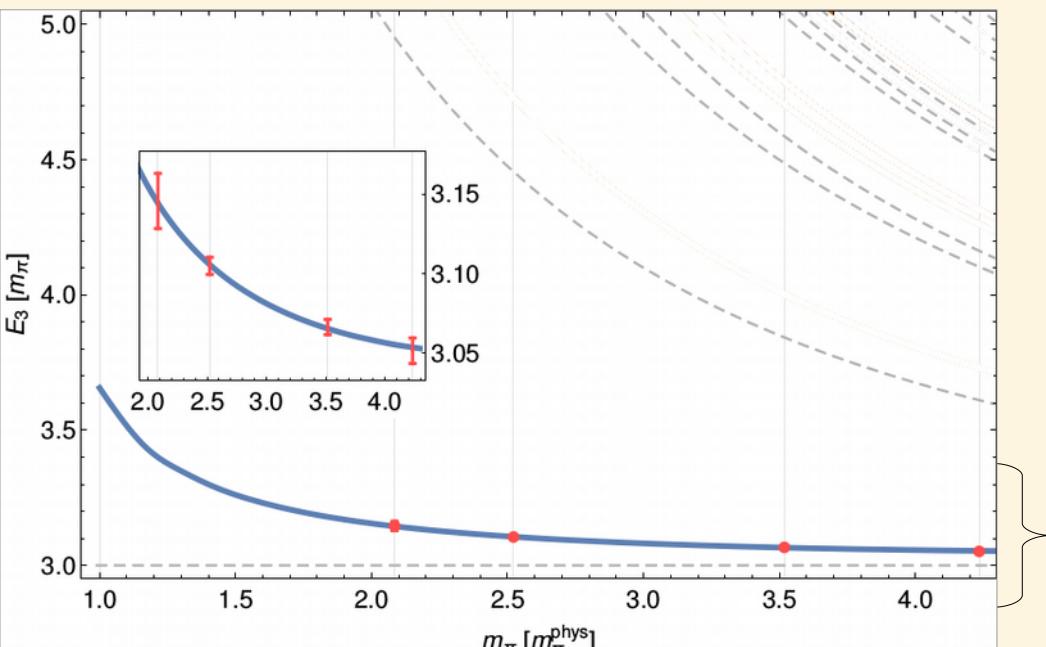
II. 3-body spectrum

Remaining unknown: C

3-BODY QUANTIZATION CONDITION

$$\text{Det} \left(B_{uu'}^{\Gamma ss'}(\mathbf{W}^2) + \frac{2E_s L^3}{\vartheta(s)} \tau_s(\mathbf{W}^2)^{-1} \delta_{ss'} \delta_{uu'} \right) = 0$$

- genuine (momenta-dependent) 3-body “force”
- simplest case: $C_{qp} = c \delta^{(3)}(\mathbf{p}-\mathbf{q})$



Fit c to NPLQCD ground state level
 $\rightarrow c=0.2 \pm 1.5 \cdot 10^{-10}$

PHYSICAL SYSTEM

$\pi^+\pi^+$ & $\pi^+\pi^+\pi^+$ ground level @ $L=2.5\text{ fm}$ & $m_\pi=291/352/491/591\text{ MeV}$

Detmold et al. (2008)

- ☞ *Q: does the “isobar” picture hold for repulsive channel?*
- ☞ *Bonus Q: chiral extrapolation in 3body system?*

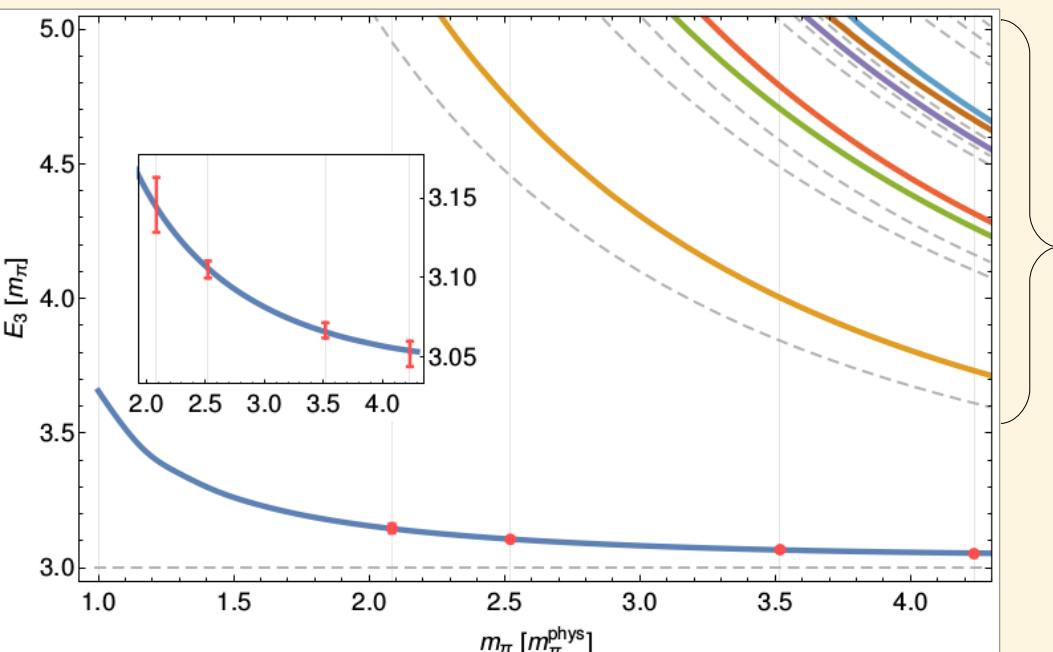
II. 3-body spectrum

Remaining unknown: $\textcolor{blue}{C}$

- genuine (momenta-dependent) 3-body “force”
- simplest case: $C_{qp} = c \delta^{(3)}(\mathbf{p}-\mathbf{q})$

3-BODY QUANTIZATION CONDITION

$$\text{Det} \left(B_{uu'}^{\Gamma ss'} (\mathbf{W}^2) + \frac{2E_s L^3}{\vartheta(s)} \tau_s (\mathbf{W}^2)^{-1} \delta_{ss'} \delta_{uu'} \right) = 0$$



Predict excited spectrum:

- novel pattern
- 1/1 of interacting/non-interacting lvs
- shifted block wise
- all genuine poles are simple
- chiral extrapolation to phys. point

SUMMARY

3-body amplitude in infinite volume

- Parametrization via 2-body sub-channel amplitudes \leftrightarrow “isobars”
- Relativistic integral equation
- Equivalent to Khuri-Treiman equations*
- Phenomenological applications in progress...

Analysis of finite-volume spectra

- discretization & projection to irreps of O_h :
a relativistic 3-body quantization condition
- numerical (ρ -like) toy-examples explored
First num. appl. to physical system: $\pi^+\pi^+\pi^+$
- (excited spectrum) of $\pi^+\pi^+$ & $\pi^+\pi^+\pi^+$ predicted
- ground level compared with NPLQCD results
- new 3-body fin-vol. spectrum features & chiral extrapolation...

OUTLOOK

- include: spin isobars, multiple isobars, unequal masses
- further practical studies: $a_1(1260)$, *Roper* ...

* Ian Aitchison, private communication

SUMMARY

3-body amplitude in infinite volume

- Parametrization via 2-body sub-channel amplitudes \leftrightarrow “isobars”
- Relativistic integral equation
- Equivalent to Khuri-Treiman equations*
- Phenomenological applications in progress...

THANK YOU

Analysis of finite-volume spectra

- discretization & projection to irreps of O_h :
a relativistic 3-body quantization condition
- numerical (ρ -like) toy-examples explored
First num. appl. to physical system: $\pi^+\pi^+\pi^+$
- (excited spectrum) of $\pi^+\pi^+$ & $\pi^+\pi^+\pi^+$ predicted
- ground level compared with NPLQCD results
- new 3-body fin-vol. spectrum features & chiral extrapolation...

OUTLOOK

- include: spin isobars, multiple isobars, unequal masses
- further practical studies: $a_1(1260)$, *Roper* ...

BACKUP

Cancellations:

→ fin. vol. normalization of δ -distribution!

$$\begin{aligned}\bar{T}_{nm}^{A_1^+}(s) &= \tau_n(s) T_{nm}^{A_1^+}(s) \tau_m(s) - 2E_n \tau_n(s) \frac{L^3}{\vartheta(n)} \delta_{nm} \\ T_{nm}^{A_1^+}(s) &= B_{nm}^{A_1^+}(s) - \frac{1}{L^3} \sum_{x \in set_8} \vartheta(x) B_{nx}^{A_1^+}(s) \frac{\tau_x(s)}{2E_x} T_{xm}^{A_1^+}(s)\end{aligned}$$

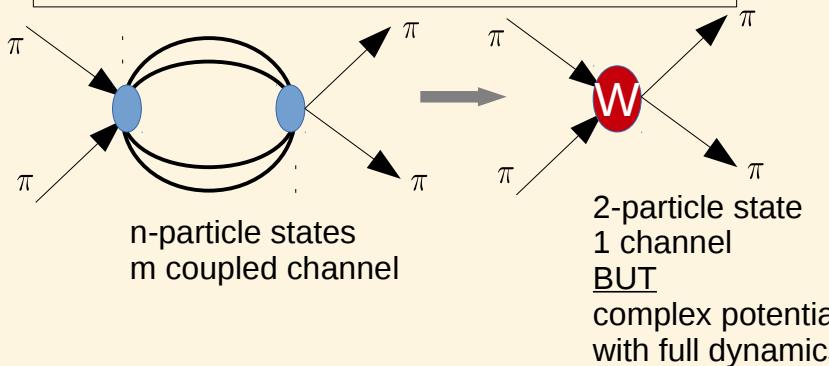
$B^{A_1^+}$ singular at $W^+ = E_m + E_n + E(\mathbf{q}_{nj} + \mathbf{p}_{mi})$

τ_m^{-1} singular at $W^{\pm\pm} = E_m \pm E((2\pi/L)\mathbf{y}) \pm E((2\pi/L)\mathbf{y} + \mathbf{p}_{mi})$ for $\mathbf{y} \in \mathbb{Z}^3$

– when isobar-momenta are discretized in the 3-body cms momenta

$$\tau = \sigma(k) - M_0^2 - \frac{1}{(2\pi)^3} \int d^3\ell \frac{\lambda^2}{2E_\ell(\sigma(k)-4E_\ell^2+i\epsilon)}$$

Optical potential: The formal rewriting of a complicated scattering problem



Lattice: measure eigenvalues,
map to the optical potential

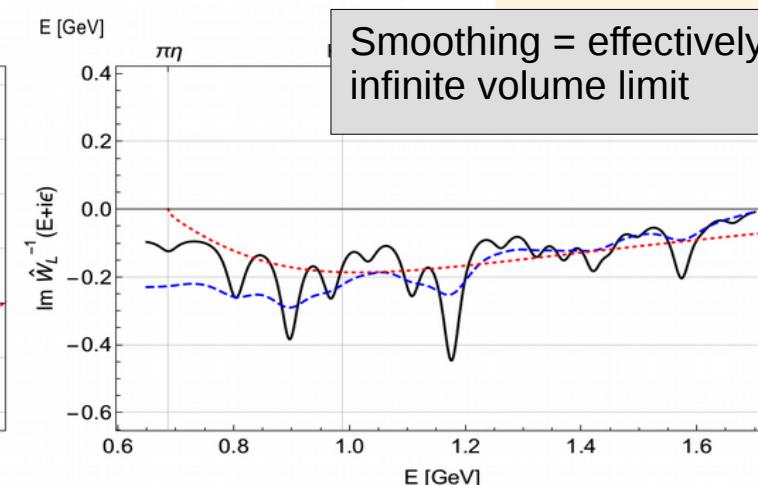
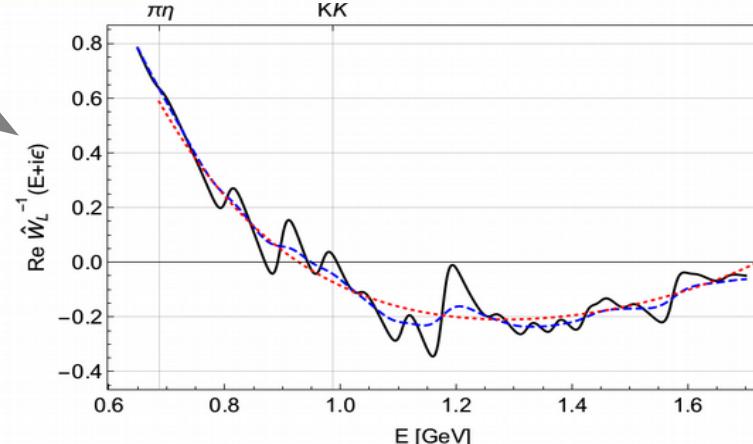
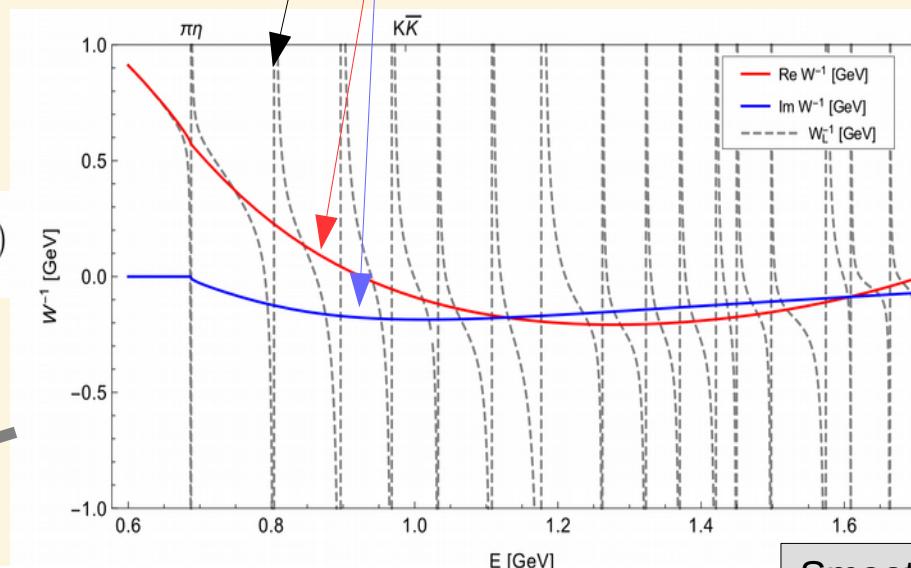
$$E \uparrow$$

$$W_L^{-1}(E) \doteq \frac{2}{\sqrt{\pi}L} Z_{00}(1; q_{K\bar{K}}^2)$$

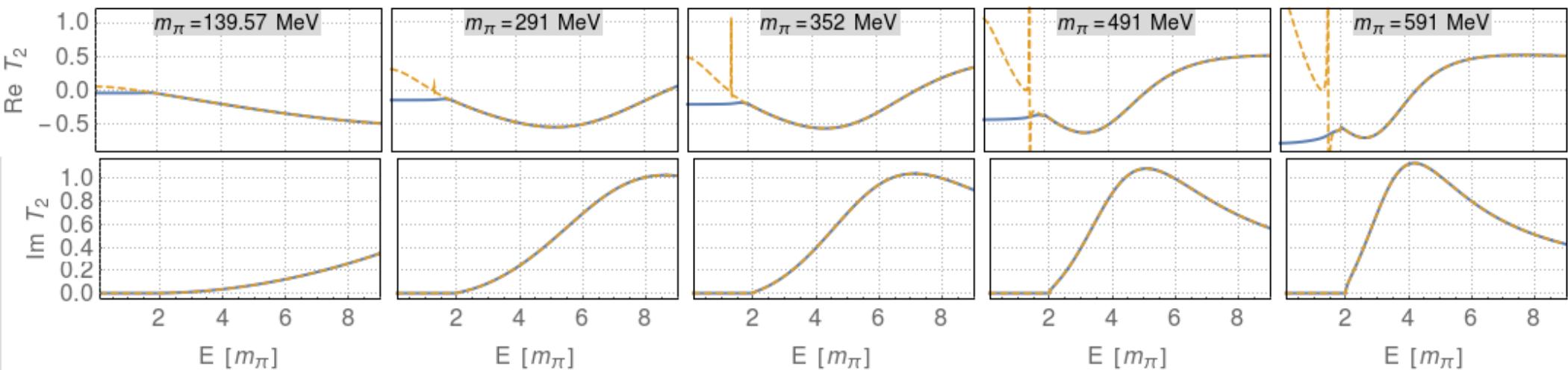
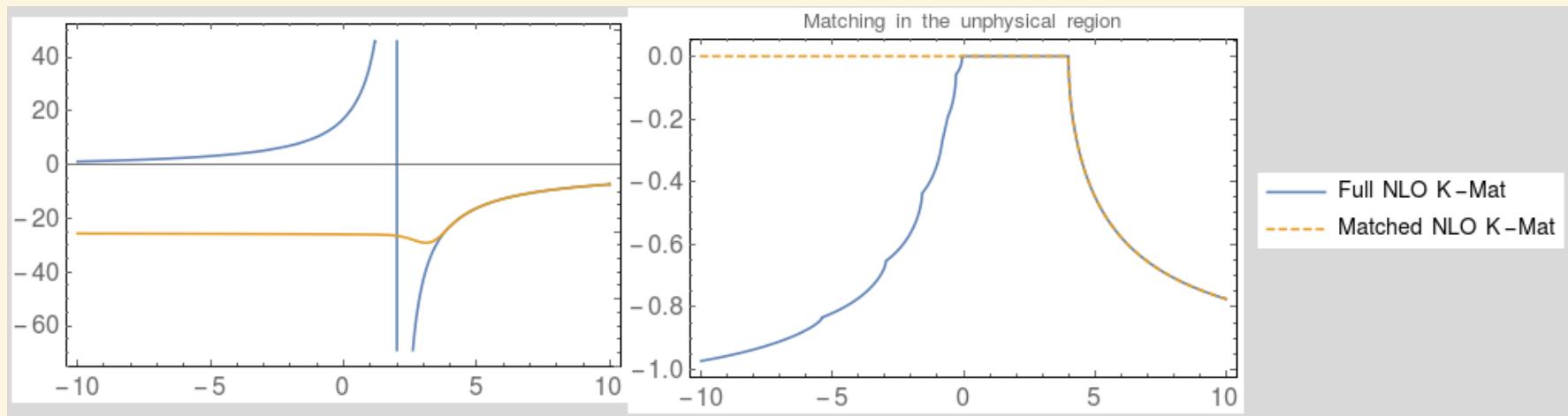
$$W^{-1} = \lim_{\epsilon \rightarrow 0} \lim_{L \rightarrow \infty} W_L^{-1}$$

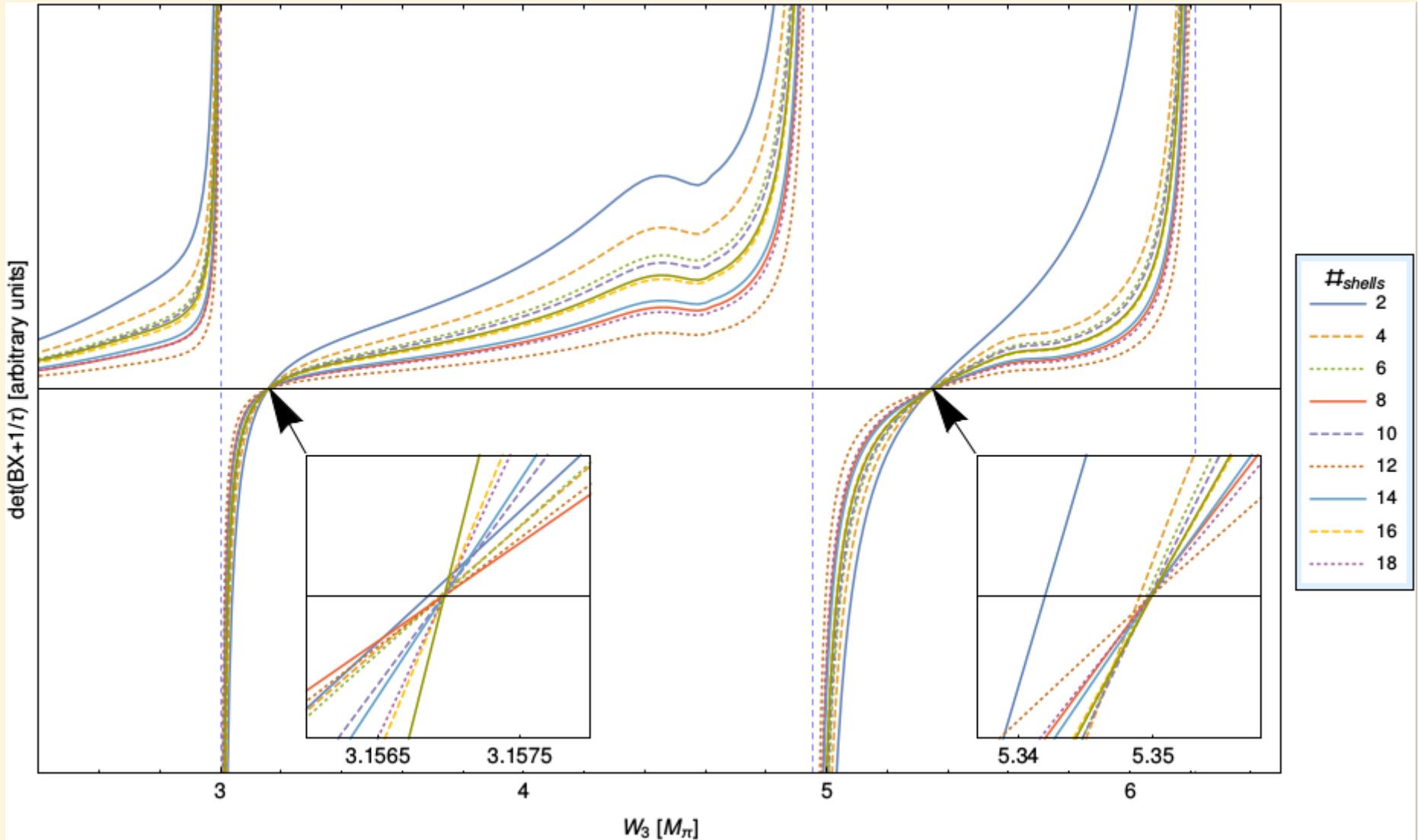
- Measured finite-volume optical potential
- Poles/functional form contain full multi-channel/multi-particle dynamics
- How to efficiently measure this function → later

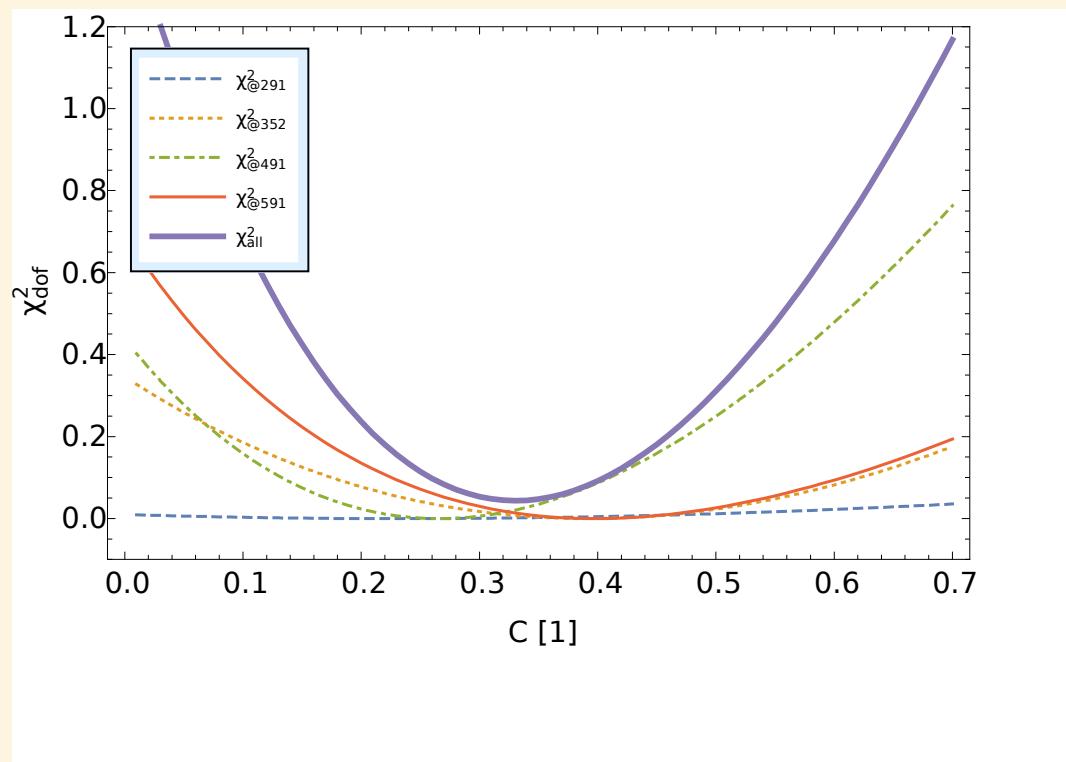
How to reconstruct true OP (complex)
from finite volume OP (real)?



Smoothing = effectively taking infinite volume limit







m_π [MeV]	139.57	291	352	491	591
E_2^1 [m_π]	$2.1228^{+0.0068}_{-0.0069}$	$2.0437^{+0.0071}_{-0.0086}$	$2.0334^{+0.0076}_{-0.0086}$	$2.0233^{+0.0105}_{-0.0098}$	$2.0204^{+0.0200}_{-0.0106}$
Refs. [24, 25]	—	2.0471(27)(65)	2.0336(22)(22)	2.0215(16)(13)	2.0171(16)(19)
E_2^2 [m_π]	—	—	$3.6245^{+0.0746}_{-0.0299}$	$2.9556^{+0.0728}_{-0.0263}$	$2.7045^{+0.0827}_{-0.0271}$
E_2^3 [m_π]	—	—	—	$3.7114^{+0.1482}_{-0.0737}$	$3.2911^{+0.1241}_{-0.0688}$
E_2^4 [m_π]	—	—	—	—	$3.6802^{+0.0707}_{-0.0902}$
E_2^5 [m_π]	—	—	—	—	$3.9829^{+0.0500}_{-0.0299}$
E_3^1 [m_π]	$3.6564^{+0.1014}_{-0.0847}$	$*3.1444^{+0.0171}_{-0.0192}$	$*3.1058^{+0.0091}_{-0.0147}$	$*3.0655^{+0.0029}_{-0.0095}$	$*3.0537^{+0.0048}_{-0.0119}$
Refs. [24, 25]	—	3.1458(49)(125)	3.1050(27)(27)	3.0665(26)(22)	3.0516(27)(53)
E_3^2 [m_π]	—	—	$4.7301^{+0.1577}_{-0.1027}$	$4.0031^{+0.0196}_{-0.1836}$	$3.7315^{+0.0309}_{-0.0742}$
E_3^3 [m_π]	—	—	—	$4.7043^{+0.0126}_{-0.5923}$	$4.2621^{+0.0001}_{-0.1739}$
E_3^4 [m_π]	—	—	—	$4.7890^{+0.0506}_{-0.1722}$	$4.3155^{+0.0837}_{-0.1341}$
E_3^5 [m_π]	—	—	—	—	$4.5913^{+0.0001}_{-0.1995}$
E_3^6 [m_π]	—	—	—	—	$4.6634^{+0.0001}_{-0.1070}$
E_3^7 [m_π]	—	—	—	—	$4.6995^{+0.0001}_{-0.0661}$