THE a1(1260)-RESONANCE FROM LATTICE QCD 2107.03973 [hep-lat]

-50Im \sqrt{s} [MeV] -100224 -150-200P D G -2501100 1000 1200 1300 Re \sqrt{s} [MeV]

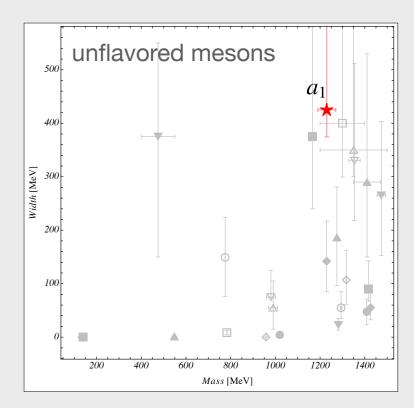
Maxim Mai, A. Alexandru, R. Brett, C. Culver M. Döring, F. Lee, D. Sadasivan [GWQCD]

QCD SPECTRUM

Many states of QCD have large coupling to 3-body channels

- $\omega(782)$, $a_1(1260)$...
- exotic mesons: $\pi_1(1600)$, ... exp. searches @ COMPASS, GlueX
- Roper resonance *N*(1440)*

This work: a₁(1260) from lattice QCD



QCD SPECTRUM

unflavored mesons

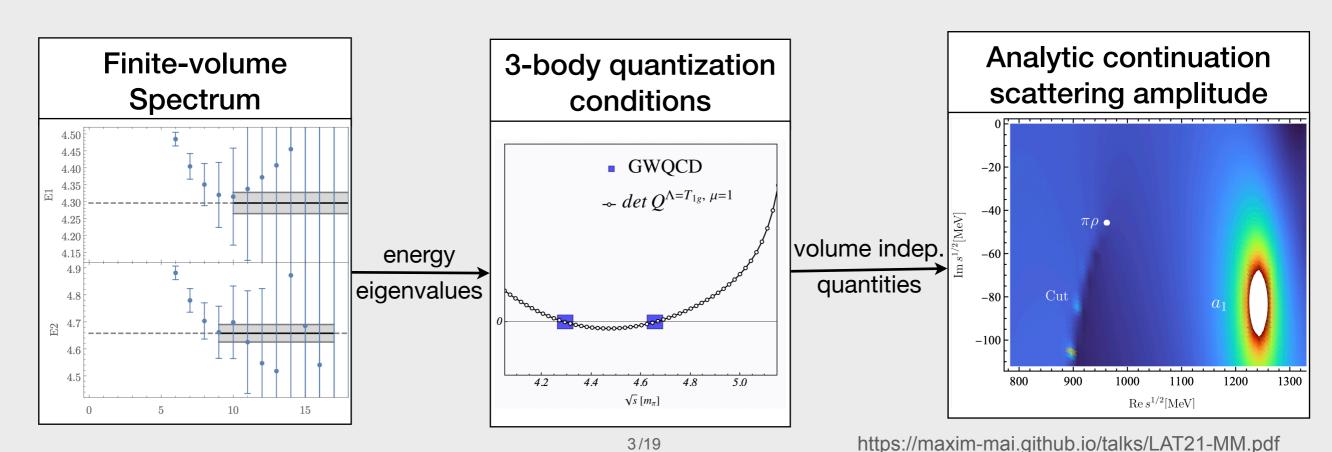
Mass [MeV]

Many states of QCD have large coupling to 3-body channels

- $\omega(782)$, $a_1(1260)$...
- exotic mesons: $\pi_1(1600)$, ... exp. searches @ COMPASS, GlueX
- Roper resonance *N*(1440)*

This work: a₁(1260) from lattice QCD

- Universal parameters from poles on the Riemann surface
- 3 step procedure:



FINITE-VOLUME SPECTRUM

GWQCD ensemble used for 2/3 pion calculations Alexandru, Brett, Culver, Guo, Lee, Pelissier (2013-2020) PRD87,PRD94,PRD98,PRD96,PRL117,PRD100

Some key details: (more in the next talk -- Ruairí Brett)

- $N_f = 2$ dynamical fermions, LapH smearing
- $P=(0,0,0), m_{\pi}=224 \text{ MeV}, m_{\pi}L=3.3$
- GEVP with one-, two-, three-meson operators

Geometry	P	Λ	$J^P (I^G = 1^-)$
Cubic	$\mathbf{P} = (0, 0, 0)$	T_{1g}	$1^+, 3^+, \dots$
		A_{1u}	$0^-, 4^-, \dots$

• Relevant irrep(O_h) for **a1(1260)** IG (JPC) = 1-(1++): T_{1g}

FINITE-VOLUME SPECTRUM

GWQCD ensemble used for 2/3 pion calculations Alexandru, Brett, Culver, Guo, Lee, Pelissier (2013-2020) PRD87,PRD94,PRD98,PRD96,PRL117,PRD100

Some key details: (more in the next talk -- Ruairí Brett)

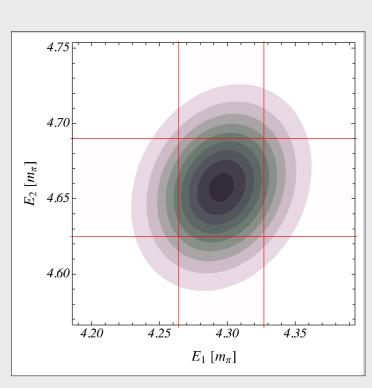
- $N_f = 2$ dynamical fermions, LapH smearing
- $P=(0,0,0), m_{\pi}=224 \text{ MeV}, m_{\pi}L=3.3$
- GEVP with one-, two-, three-meson operators

Geometry	P	Λ	$J^P \ (I^G = 1^-)$
Cubic	$\mathbf{P} = (0, 0, 0)$	T_{1g}	$1^+, 3^+, \dots$
		A_{1u}	$0^-, 4^-, \dots$

• Relevant irrep(O_h) for **a1(1260)** IG (JPC) = 1- (1++): T_{1g}

Key insights:

- 3-meson operators stabilize the excited state extraction c.f. need for $\rho\pi$ operators in pioneering 2-meson a_1 -calculation Lang et al. JHEP 04, 162 (2014)
- high-momentum states are required: $\pi(0,0,0)\pi(1,1,0)\pi(-1,-1,0)$ etc..
- two interacting levels exists below 5π threshold



Discrete, real finite-volume (lattice) spectrum → continuous complex-valued amplitudes

• established in 2-body: Lüscher's method, extensions...

Lüscher, Gottlieb, Rummukainen, Feng, Li, Liu, Döring, Briceño, Bernard, Meißner, Rusetsky...

3-body methods matured (this session)

Bedaque, Blanton, Briceño, Davoudi, Döring, Grießhammer, Guo, Hammer, Hansen, MM, Meißner, Müller, Pang, Polejaeva, Romero-López, Rusetsky, Sharpe, Wu

Reviews: Hansen/Sharpe(2019) MM/Döring/Rusetsky(2021)

Discrete, real finite-volume (lattice) spectrum → continuous complex-valued amplitudes

established in 2-body: Lüscher's method, extensions...

Lüscher, Gottlieb, Rummukainen, Feng, Li, Liu, Döring, Briceño, Bernard, Meißner, Rusetsky...

3-body methods matured (this session)

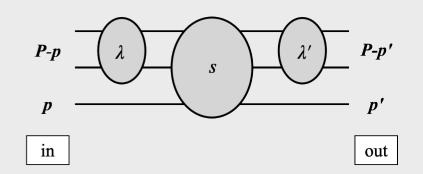
Bedaque, Blanton, Briceño, Davoudi, Döring, Grießhammer, Guo, Hammer, Hansen, MM, Meißner, Müller, Pang, Polejaeva, Romero-López, Rusetsky, Sharpe, Wu

Reviews: Hansen/Sharpe(2019) MM/Döring/Rusetsky(2021)

Finite Volume Unitarity MM, Döring EPJA (2017) PRL (2019)

· basic idea:

$$0 = \det \left[B(s) + C(s) - 2L^3 E_{\mathbf{p}} \left(\tilde{K}_2^{-1}(s) - \Sigma_2^L(s) \right) \right]_{(\lambda'\lambda)(\mathbf{p}'\mathbf{p})}^{\Lambda}$$



- extended to higher spin and coupled-channels: new degree of freedom (λ)
- ∞ -dim. determinant equation in $\mathbf{p} \in \frac{2\pi}{L} \mathbf{Z}^3 \to \text{practical applications require truncation}$
 - → common to all quantization conditions

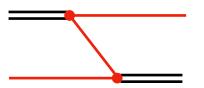
see discussion in e.g. MM/Döring/Rusetsky(2021)

$$0 = \det \left[\frac{B(s) + C(s) - 2L^3 E_{\mathbf{p}} \left(\tilde{K}_2^{-1}(s) - \Sigma_2^L(s) \right) \right]_{(\lambda'\lambda)(\mathbf{p}'\mathbf{p})}$$

$$0 = \det \left[\frac{B(s) + C(s) - 2L^3 E_{\mathbf{p}} \left(\tilde{K}_2^{-1}(s) - \Sigma_2^L(s) \right) \right]_{(\lambda'\lambda)(\mathbf{p}'\mathbf{p})}$$

one-particle exchange

• fixed by 3b-unitarity

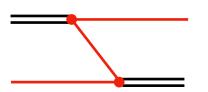


• no free parameters

$$0 = \det \left[\frac{B(s) + C(s) - 2L^3 E_{\mathbf{p}} \left(\tilde{K}_2^{-1}(s) - \Sigma_2^L(s) \right) \right]_{(\lambda'\lambda)(\mathbf{p}'\mathbf{p})}$$

one-particle exchange

• fixed by 3b-unitarity



• no free parameters

two-body self-energy

• fixed by 2b-unitarity

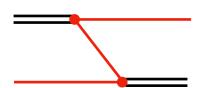


no free parameters

$$0 = \det \left[\frac{B(s) + C(s) - 2L^3 E_{\mathbf{p}} \left(\tilde{K}_2^{-1}(s) - \Sigma_2^L(s) \right) \right]_{(\lambda'\lambda)(\mathbf{p}'\mathbf{p})}$$

one-particle exchange

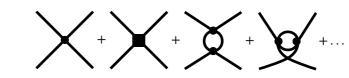
• fixed by 3b-unitarity



• no free parameters

two-body kernel

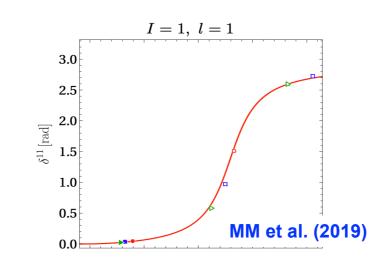
• dynamics of I=1 $\pi\pi$ system



regular function ⇒ polynomial

$$\tilde{K}_n^{-1}(s) = \sum_{i=0}^{n-1} a_i \cdot \sigma_p^i$$

• parameters (a_0,a_1) from crosschannel fit to $\pi\pi$ GWQCD levels



two-body self-energy

• fixed by 2b-unitarity

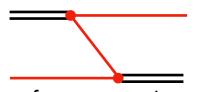


no free parameters

$$0 = \det \left[\frac{B(s) + C(s) - 2L^3 E_{\mathbf{p}} \left(\tilde{K}_2^{-1}(s) - \Sigma_2^L(s) \right) \right]_{(\lambda'\lambda)(\mathbf{p}'\mathbf{p})}$$

one-particle exchange

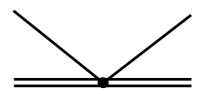
• fixed by 3b-unitarity



no free parameters

three-body force

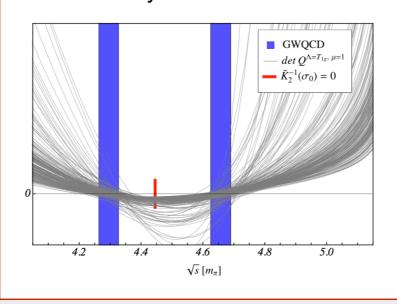
• dynamics of $\rho\pi$ system



regular function ⇒ Laurent series

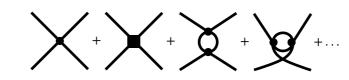
$$C_{\ell'\ell}(s, \mathbf{p}', \mathbf{p}) = \sum_{i=-1}^{\infty} \frac{c_{\ell'\ell}^{(i)}}{(\mathbf{p}', \mathbf{p})(s - \mathbf{m}_{a_1}^2)^i}$$

• fit to 3-body levels



two-body kernel

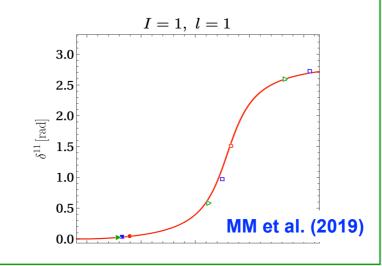
• dynamics of I=1 $\pi\pi$ system



regular function ⇒ polynomial

$$\tilde{K}_n^{-1}(s) = \sum_{i=0}^{n-1} a_i \cdot \sigma_p^i$$

• parameters (a_0,a_1) from crosschannel fit to $\pi\pi$ GWQCD levels



two-body self-energy

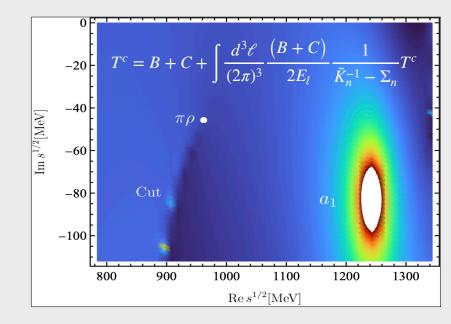
• fixed by 2b-unitarity



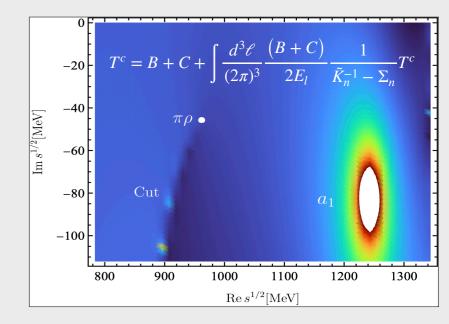
· no free parameters

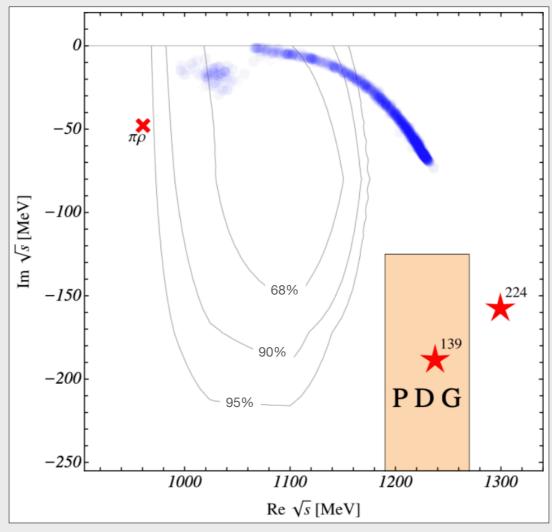
- ∞-vol. scattering equation via contour deformation of spectator momenta

 Döring et al.(2009) Sadasivan et al. (2020)
- various forms of the 3-body term C tested:
 - pole is generated with or without explicit pole-term
 - best description via $C_{\ell'\ell} = g_{\ell'} |\mathbf{p}'|^{\ell'} \frac{1}{\mathbf{s} \mathbf{m}_{\mathbf{a}_1}^2} \mathbf{g}_{\ell} |\mathbf{p}|^{\ell} + \mathbf{c} \, \delta_{\ell' \mathbf{0}} \delta_{\ell \mathbf{0}}$...with large correlations



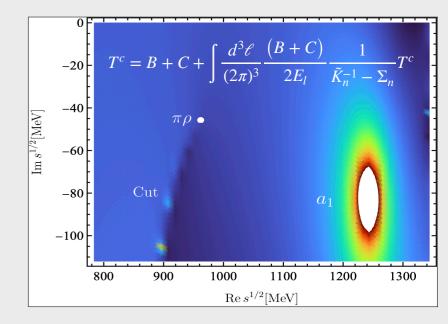
- ∞-vol. scattering equation via contour deformation of spectator momenta
 Döring et al.(2009) Sadasivan et al. (2020)
- various forms of the 3-body term C tested:
 - pole is generated with or without explicit pole-term
 - best description via $C_{\ell'\ell} = g_{\ell'} |\mathbf{p}'|^{\ell'} \frac{1}{\mathbf{s} \mathbf{m}_{\mathbf{a}_1}^2} \mathbf{g}_{\ell} |\mathbf{p}|^{\ell} + \mathbf{c} \, \delta_{\ell' \mathbf{0}} \delta_{\ell \mathbf{0}}$...with large correlations

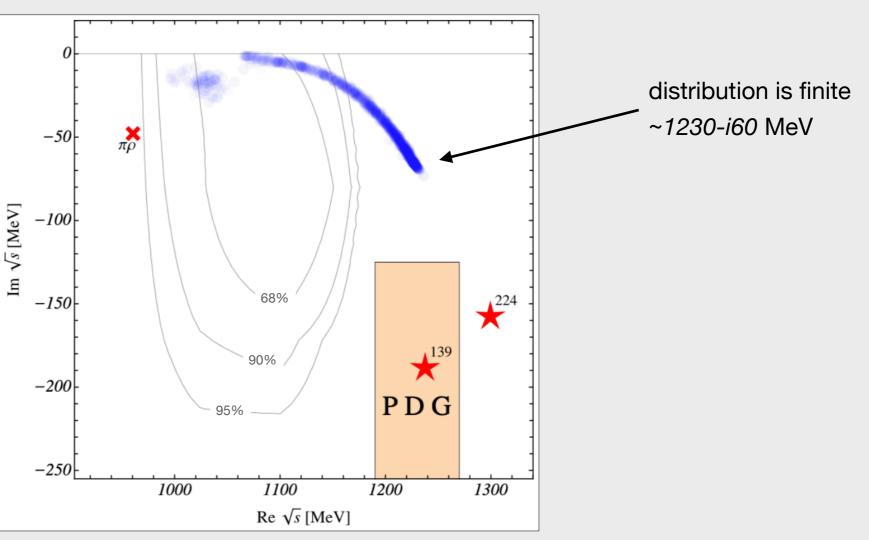




- ∞-vol. scattering equation via contour deformation of spectator momenta

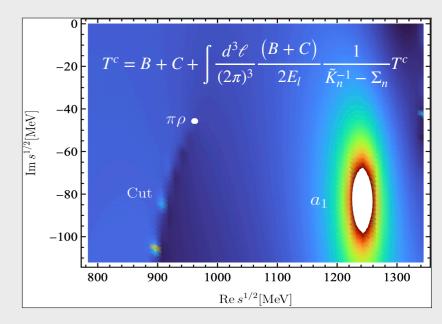
 Döring et al.(2009) Sadasivan et al. (2020)
- various forms of the 3-body term C tested:
 - pole is generated with or without explicit pole-term
 - best description via $C_{\ell'\ell} = g_{\ell'} |\mathbf{p}'|^{\ell'} \frac{1}{\mathbf{s} \mathbf{m}_{\mathbf{a}_1}^2} \mathbf{g}_{\ell} |\mathbf{p}|^{\ell} + \mathbf{c} \, \delta_{\ell'\mathbf{0}} \delta_{\ell\mathbf{0}}$...with large correlations

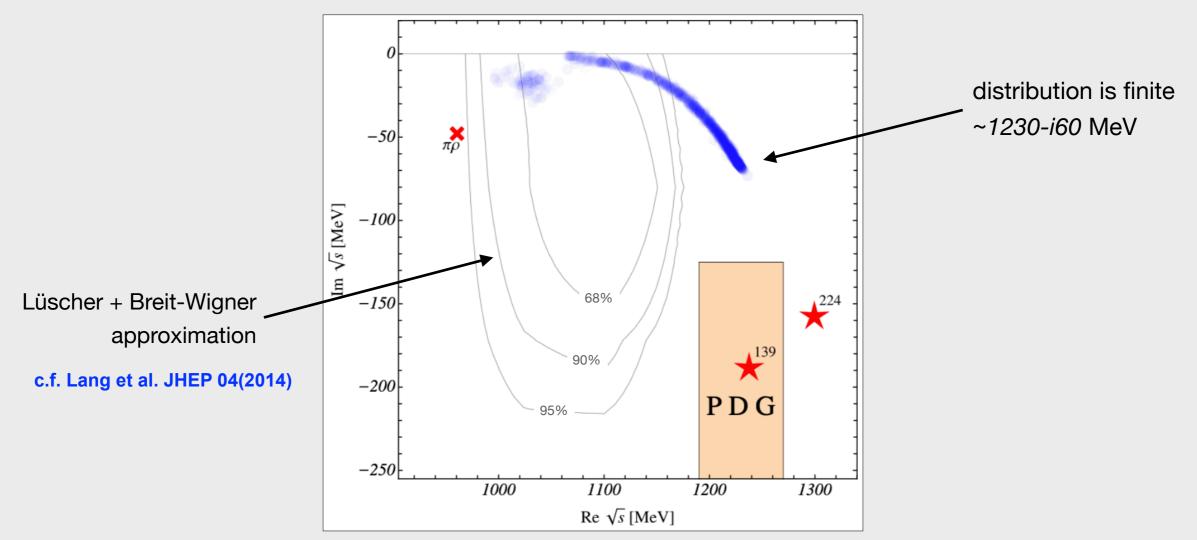




- ∞-vol. scattering equation via contour deformation of spectator momenta

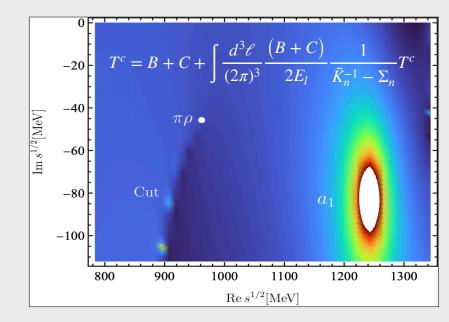
 Döring et al.(2009) Sadasivan et al. (2020)
- various forms of the 3-body term C tested:
 - pole is generated with or without explicit pole-term
 - best description via $C_{\ell'\ell} = \mathbf{g}_{\ell'} |\mathbf{p}'|^{\ell'} \frac{1}{\mathbf{s} \mathbf{m}_{\mathbf{a}_1}^2} \mathbf{g}_{\ell} |\mathbf{p}|^{\ell} + \mathbf{c} \, \delta_{\ell'\mathbf{0}} \delta_{\ell\mathbf{0}}$...with large correlations

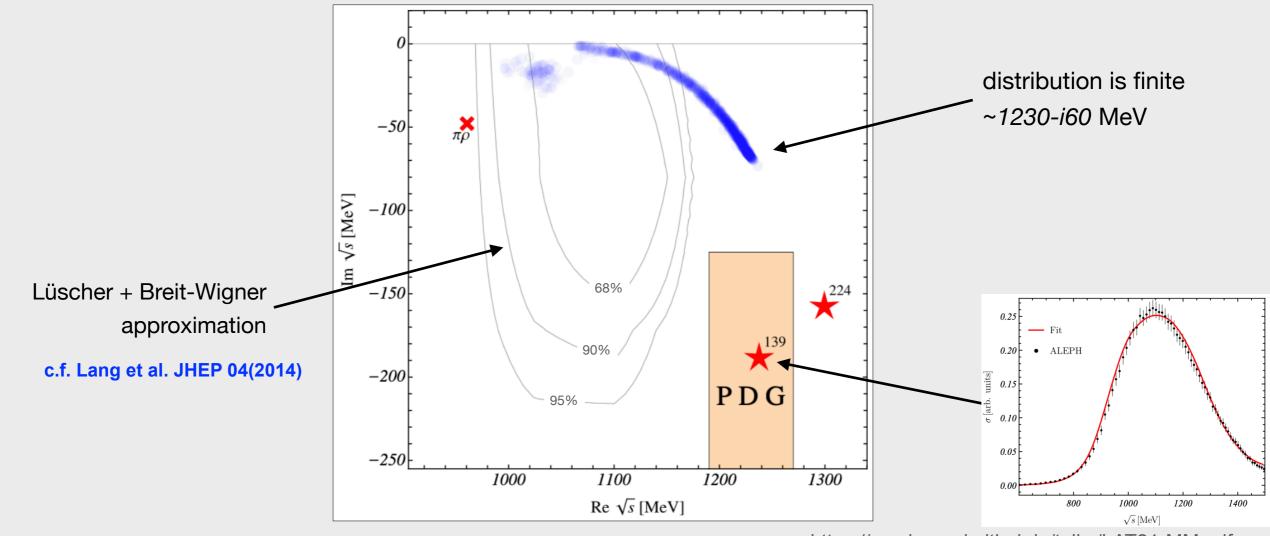




- ∞-vol. scattering equation via contour deformation of spectator momenta

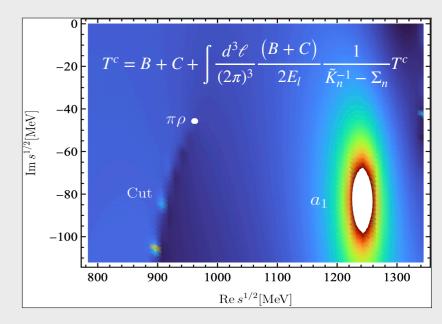
 Döring et al.(2009) Sadasivan et al. (2020)
- various forms of the 3-body term C tested:
 - pole is generated with or without explicit pole-term
 - best description via $C_{\ell'\ell} = g_{\ell'} |\mathbf{p}'|^{\ell'} \frac{1}{\mathbf{s} \mathbf{m}_{\mathbf{a}_1}^2} \mathbf{g}_{\ell} |\mathbf{p}|^{\ell} + \mathbf{c} \, \delta_{\ell' \mathbf{0}} \delta_{\ell \mathbf{0}}$...with large correlations

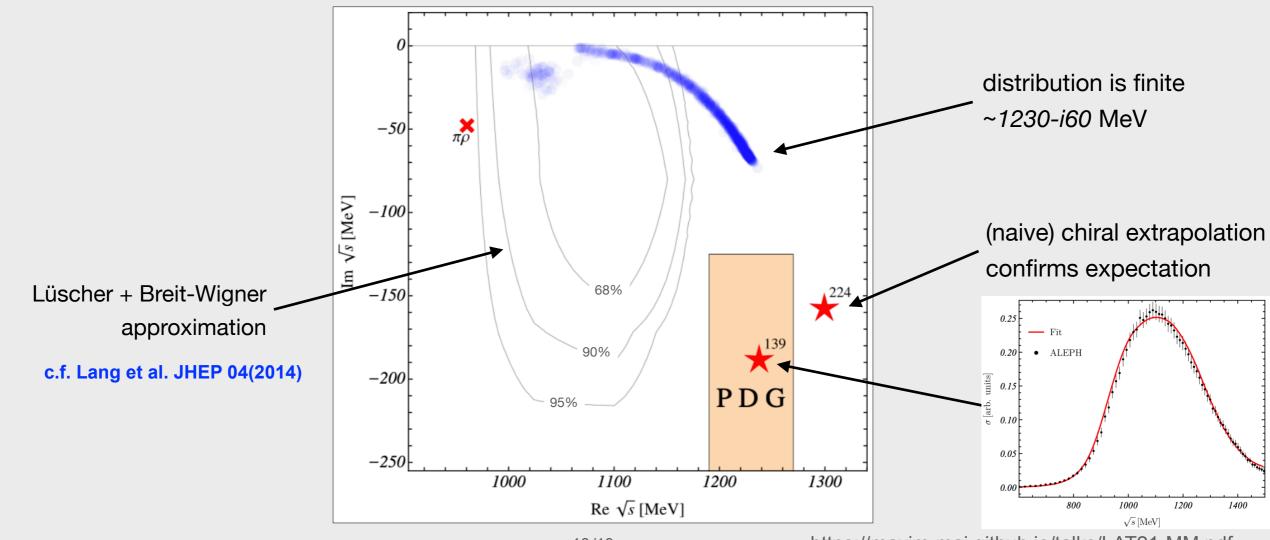




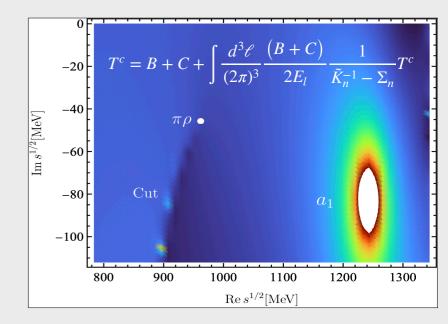
- ∞-vol. scattering equation via contour deformation of spectator momenta

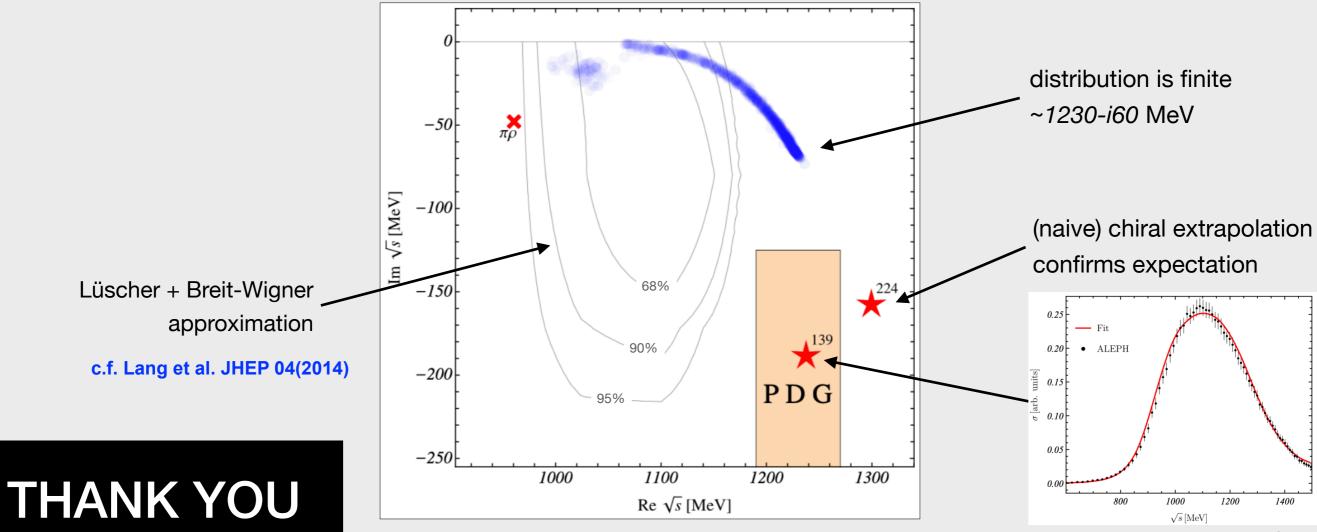
 Döring et al.(2009) Sadasivan et al. (2020)
- various forms of the 3-body term C tested:
 - pole is generated with or without explicit pole-term
 - best description via $C_{\ell'\ell} = g_{\ell'} |\mathbf{p}'|^{\ell'} \frac{1}{\mathbf{s} \mathbf{m}_{\mathbf{a}_1}^2} \mathbf{g}_{\ell} |\mathbf{p}|^{\ell} + \mathbf{c} \, \delta_{\ell' \mathbf{0}} \delta_{\ell \mathbf{0}}$...with large correlations





- ∞-vol. scattering equation via contour deformation of spectator momenta Döring et al.(2009) Sadasivan et al. (2020)
- various forms of the 3-body term C tested:
 - pole is generated with or without explicit pole-term
 - best description via $C_{\ell'\ell} = g_{\ell'} |\mathbf{p}'|^{\ell'} \frac{1}{\mathbf{s} - \mathbf{m}_{a_1}^2} \mathbf{g}_{\ell} |\mathbf{p}|^{\ell} + \mathbf{c} \, \delta_{\ell'0} \delta_{\ell 0}$...with large correlations





https://maxim-mai.github.io/talks/LAT21-MM.pdf