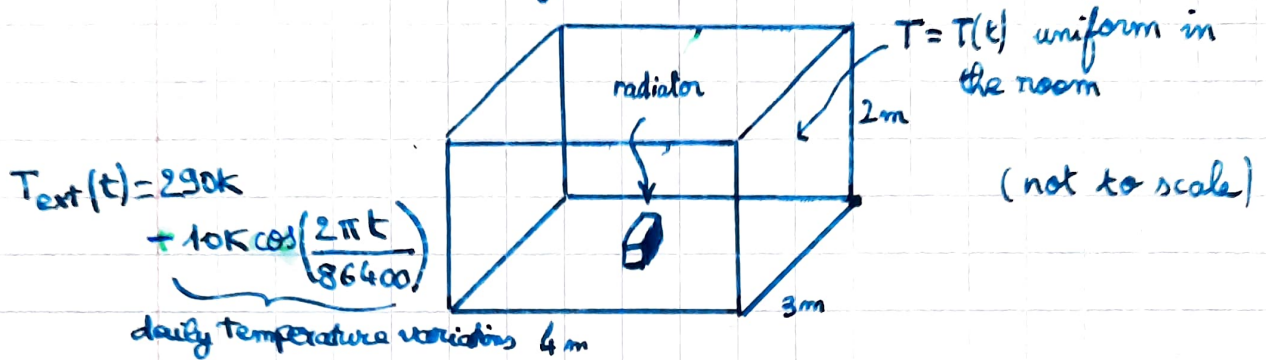


## - Optimal control of radiator power -

Wooden hut heated by a radiator:



4 walls and ceiling are made of 5 cm-thick wood with a thermal conductivity of  $0,15 \text{ W}\cdot\text{m}^{-1}\cdot\text{K}^{-1}$

Assumption: the floor is a thermal insulator

Total thermal resistance:  $R = \frac{0,05 \text{ m}}{0,15 \text{ W}\cdot\text{m}^{-1}\cdot\text{K}^{-1}} = 0,3 \text{ m}^2\cdot\text{K}\cdot\text{W}^{-1}$

Thermal capacity of the hut:  $C = V C_v = 3,1 \cdot 10^4 \text{ JK}^{-1}$

$\uparrow$  volume of the hut  
 $\uparrow$  air volumetric heat capacity  
 $1,3 \cdot 10^3 \text{ JK}^{-1}\cdot\text{m}^{-3}$

Power of the radiator:  $P(t)$

System = {hut}

First principle applied to this System:

$$\underbrace{H(t+dt) - H(t)}_{C(T(t+dt) - T(t))} = \delta Q = \left( P(t) - \underbrace{\phi_{\text{outgoing}}}_{\text{walls, ceiling}}(t) \right) dt$$

$$\rightarrow \frac{S}{R} (T(t) - T_{ext}(t))$$

wall + ceiling area  $S = 40 \text{ m}^2$   
 Fourier's phenomenological law

Therefore:

$$\dot{T}(t) = \frac{1}{C} P(t) - \frac{S}{CR} (T(t) - T_{ext}(t))$$

$$\dot{T}(t) = \alpha P(t) - \beta (T(t) - T_{ext}(t))$$

with:  $\alpha = \frac{1}{C} = 3,2 \cdot 10^{-5} \text{ K} \cdot \text{J}^{-1}$

$$\beta = \frac{S}{CR} = 4 \cdot 10^{-3} \text{ s}^{-1}$$

Goals: ① indoor temperature  $T_g(t) \approx 293 \text{ K}$   
 $\nearrow$  goal

② minimize  $P(t)$

Data/constraint: ① when we arrive (at  $t=0$ )  
 $T(0) = T_{ext}(0)$

② we stay 3 days:  
 $t \in [0; \underbrace{259200}_{t_f}]$

Therefore we want to solve:

$$\text{Min}_{T \in W^{1,\infty}(0, t_f, \mathbb{R})} \int_0^{t_f} \mathcal{L}(t, P(t), T(t)) dt + \varphi(T(t_f))$$

Sobolev:

$$W^{1,\infty}(0, t_f, \mathbb{R})$$

$$= \{y \in L^\infty(0, t_f, \mathbb{R})$$

$$\text{such that } y \in L^\infty(0, t_f, \mathbb{R})\}$$

s.t.



with  $\begin{cases} \mathcal{L} \\ \varphi \end{cases}$  uniformly quasi  $C^\infty$

derivative in the sense of distributions (cf L. Schwartz)

general form of cost function  $J(P(t), T(t))$



and satisfying  $\dot{T}(t) = f(t, \overbrace{P(t)}^{\text{control}}, \overbrace{T(t)}^{\text{state}})$  for a.a.  $t \in [0, t_f]$   
(equality in  $L^\infty(0, t_f, \mathbb{R})$ )

with  $\begin{cases} t_f > 0 \\ \text{initial condition } T(0) = T_{\text{ext}}(0) = 280 \text{ K} \\ \text{dynamics: } f: \mathbb{R} \times \mathbb{R} \times \mathbb{R} \rightarrow \mathbb{R} \end{cases} \begin{cases} \text{Lipschitz w.r.t. } T \\ \text{uniformly quasi-concave} \\ (n \geq 1) \end{cases}$

In optimal control theory, deriving the following Lagrangian

$$J(P, T) + \int_0^{t_f} \underbrace{\left\langle \overbrace{p(t)}^{\text{lower-case; lagrange multiplier}}, f(t, P(t), T(t)) - \dot{T}(t) \right\rangle_{\mathbb{R}', \mathbb{R}}}_{\substack{\text{dual notation (not very} \\ \text{useful here because the state} \\ \text{space is } \mathbb{R})}} dt + \left\langle q, T^0 - T(0) \right\rangle_{\mathbb{R}', \mathbb{R}}$$

gives us the costate dynamics equations

$$\begin{cases} -\dot{\bar{p}}(t) = \nabla_{\bar{p}} H(t, \bar{P}(t), \bar{T}(t), \bar{p}(t)) \text{ for a.a. } t \in [0, t_f] \\ \bar{p}(t_f) = \nabla_{\bar{p}} \varphi(\bar{T}(t_f)) \end{cases}$$

where  $H(t, P, T, p) = l(t, P, T) + \langle p, f(t, P, T) \rangle_{\mathbb{R}', \mathbb{R}}$

is called « pre-hamiltonian »

From the derivative of optimality conditions:  $J(P, t)$  we can obtain necessary  
reduced cost mapping called  $\nabla_p \bar{H}(t)$

$$\begin{aligned} \bar{P} \text{ local minimum of } J(\underbrace{P}_{\text{fixed}}, \bar{T}[P, T^0]) &\Rightarrow \nabla_p H(t, \bar{P}(t), \bar{T}(t), \bar{p}(t)) = 0 \text{ for a.a. } t \in [0, t_f] \\ \bar{T}^0 \text{ local minimum of } J(\underbrace{P}_{\text{fixed}}, \bar{T}[P, T^0]) &\Rightarrow \bar{p}(0) = 0 \end{aligned}$$

## PROOF

We remind  $J(P, T)$ :

$$J(P, T) = \int_0^{t_f} \ell(t, P(t), T(t)) dt + \varphi(T(t_f))$$

Its derivative is

$$J'(P, T)(\tilde{P}, \tilde{T}) = \int_0^{t_f} \ell'(t, P(t), T(t))(\tilde{P}(t), \tilde{T}(t)) dt + \langle \nabla \varphi(T(t_f)), \tilde{T}(t_f) \rangle_{\mathbb{R}', \mathbb{R}}$$

Why? Because we have:

$$J(P + \tilde{P}, T + \tilde{T}) := J(P, T) + J'(P, T)(\tilde{P}, \tilde{T}) + o(|(\tilde{P}, \tilde{T})|)$$

$$\int_0^{t_f} \ell(t, P(t) + \tilde{P}(t), T(t) + \tilde{T}(t)) dt + \varphi(T(t_f) + \tilde{T}(t_f))$$

$$J(P, T) + \int_0^{t_f} \ell'(t, P(t), T(t))(\tilde{P}(t), \tilde{T}(t)) dt + \underbrace{\varphi'(T(t_f))(\tilde{T}(t_f))}_{\text{Jacobian notation}} + o(|(\tilde{P}, \tilde{T})|)$$

$$\underbrace{\langle \nabla \varphi(T(t_f)), \tilde{T}(t_f) \rangle_{\mathbb{R}', \mathbb{R}}}_{\text{dual notation with gradient}}$$



because it's a costate equation

But

$$\underbrace{= p(t_f)}_{\substack{\text{But} \\ \uparrow}} \quad \langle \nabla \psi(t_f), \tilde{T}(t_f) \rangle_{R', R} = \langle \nabla p(0), \tilde{T}(0) \rangle_{R', R} + \int_0^{t_f} \left\langle \frac{d}{dt} \nabla p(t), \tilde{T}(t) \right\rangle_{R', R} dt$$

$$= \langle \nabla p(0), \tilde{T}(0) \rangle_{R', R}$$

$$+ \int_0^{t_f} \left\langle \dot{p}(t), \tilde{T}(t) \right\rangle_{R', R} + \left\langle p(t), \dot{\tilde{T}}(t) \right\rangle_{R', R} dt$$

$$= \langle p(0), \tilde{T}(0) \rangle_{R', R}$$

$$+ \int_0^{t_f} \left\langle -\nabla_T \bar{H}(\bar{E}), \tilde{T}(t) \right\rangle_{R', R} + \left\langle p(t), f_t(\bar{E}, p, T)(\tilde{P}, \tilde{T}) \right\rangle_{R', R} dt$$

$$= \langle p(0), \tilde{T}(0) \rangle_{R', R}$$

$$+ \int_0^{t_f} \left\langle -\nabla_T \bar{L}(\bar{E}) - \bar{f}'_T(\bar{E}) p(t), \tilde{T}(t) \right\rangle_{R', R} dt$$

$$+ \int_0^{t_f} \left\langle p(t), \bar{f}'_p(\bar{E}, p, T) \tilde{P} + \bar{f}'_T(\bar{E}, p, T) \tilde{T} \right\rangle_{R', R} dt$$

Therefore:

$$J'_R(p, T^0)(\tilde{P}, \tilde{T}^0) = \langle p(0), \tilde{T}^0 \rangle_{R', R} + \int_0^{t_f} \underbrace{\left( \bar{L}'_p(\bar{E}) \tilde{P}(t) + p(t)^T \bar{f}'_p(\bar{E}) \tilde{P} \right)}_{\left\langle \nabla_P \bar{H}(\bar{E}), \tilde{P}(t) \right\rangle_{R', R}} dt$$

which concludes.

Let's apply this to

$$J(P, T) = \int_0^t l(t, P(t), T(t)) dt + \varphi(T(t_f))$$

with  $\varphi = 0 \in \mathbb{R}$

and  $l(t, P, T) = \frac{1}{2} (|P|^2 + |T - T_g|^2)$

to minimize  
the consumption  
of power

uniformly  
quasi  $C^0$

Costate dynamics:

lower  
case

$$-\dot{\bar{P}}(t) = \bar{T}(t) - T_g(t) + \beta \bar{P}(t)$$

lower-case

lower  
case  $\rightarrow \bar{P}(T) = 0$

Necessary optimality condition:

$$\bar{P}(t) + \alpha \bar{P}(t) = 0$$

uppercase

lower  
case

And of course,  $\bar{T}$  must satisfy the differential equation:

$$\dot{\bar{T}}(t) = \alpha \bar{P}(t) - \beta (T(t) - T_{ext}(t))$$

Therefore, with  $Y(t) := \begin{pmatrix} T(t) \\ P(t) \end{pmatrix}$  we obtain:

$$\dot{Y}(t) = \underbrace{\begin{pmatrix} -\beta & -\alpha^2 \\ -1 & \beta \end{pmatrix}}_A Y(t) + \underbrace{\begin{pmatrix} \beta T_{ext}(t) \\ T_g \end{pmatrix}}_{B(t)}$$



This can be solved analytically, but it's very time-consuming!

Here is the beginning of the analytical resolution:

there exists a vector  $\lambda = \begin{pmatrix} \lambda_1 \\ \lambda_2 \end{pmatrix} \in \mathbb{R}^2$  s.t.

$$Y(t) = e^{At} \left( c + \int_0^t e^{-As} B(s) ds \right)$$

From Cayley-Hamilton theorem we know that

$$e^{At} = c_0(t) I_2 + c_1(t) A$$

with

$$\begin{cases} c_0(t) = \frac{e^{\xi_1 t} - e^{\xi_2 t}}{\xi_1 - \xi_2} \\ c_1(t) = \frac{e^{\xi_1 t} \xi_2 - e^{\xi_2 t} \xi_1}{\xi_2 - \xi_1} \end{cases}$$

where  $\xi_1, \xi_2$  are the eigenvalues of  $A = \begin{pmatrix} -\beta & -\alpha^2 \\ -1 & \beta \end{pmatrix}$

We stop here the calculations, because we use a numerical solution with Runge Kutta. See the file .ipynb

On that file we discover a very strange phenomenon: high frequency oscillations of  $p(t)$  and  $T(t)$ .  
Another «strange» phenomenon:  $P(t)$  is negative for some values of  $t$ . This shows we should add constraints:

③  $0 \leq P(t) \leq P_{\max} \approx 1$

④  $\mathcal{L}(t, P, T) = \frac{1}{2} (|P|^2 + |\dot{P}|^2 + |T - T_g|^2)$

to avoid high frequency oscillations