# Cracking RSA with Quantum Computing

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#### Outline

- The Setting
- 2 Classical Computers
- Quantum Computers
- 4 Shor's Algorithm

RSA is a commonly used set of algorithms that provides security when sending encrypted messages.

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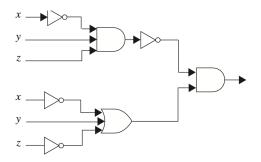
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Quantum computers *are* able to factor in polynomial time. This talk will focus on explaining how quantum algorithms work, building up to Shor's famous algorithm for factoring.

## What Computers Look Like

Equivalent to a circuit whose representation can be quickly computed:



# What Computers Look Like

We can think of a circuit as having n registers, each of which contain 0 or 1. A possible *state* of these n registers is an element of  $\{0,1\}^n$  (n-length bitstring). Then the action of a gate can be described as a matrix of 0s and 1s.

NOT gate:

$$NOT(0) = 1$$

$$NOT(1) = 0$$

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Let's relabel 
$$0 = \begin{bmatrix} 1 \\ 0 \end{bmatrix}$$
,  $1 = \begin{bmatrix} 0 \\ 1 \end{bmatrix}$ .

Elements of a 2-element vector space, as there are 2 possibilities for bits.

NOT gate:

$$NOT \begin{pmatrix} 1 \\ 0 \end{pmatrix} = \begin{bmatrix} 0 \\ 1 \end{bmatrix} \\
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This strongly suggests we can consider it as a matrix:

$$NOT = \begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix}$$

AND gate:

AND 
$$(0,0) = 0$$

$$AND(0,1)=0$$

$$AND(1,0)=0$$

$$AND(1,1) = 1$$

For two input bits, there are  $4 = 2^2$  possible states.

$$(0,0) = |00\rangle = |0\rangle \otimes |0\rangle$$

$$\textbf{(0,1)}=|01\rangle=|0\rangle\otimes|1\rangle$$

$$(1,0)=|10\rangle=|1\rangle\otimes|0\rangle$$

$$(0,1)=|11\rangle=|1\rangle\otimes|1\rangle$$

$$\mathrm{AND}\left(|00\rangle\right)=|0\rangle$$

$$AND(|01\rangle) = |0\rangle$$

$$AND(|10\rangle) = |0\rangle$$

$$\mathrm{AND}\left(|11\rangle\right)=|1\rangle$$

AND 
$$(1 |00\rangle + 0 |01\rangle + 0 |10\rangle + 0 |11\rangle) = 1 |0\rangle + 0 |1\rangle$$
  
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Okay, we have a vector space now. What about linear combinations?

$$\begin{bmatrix} 1 & 1 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{bmatrix} \left( \frac{1}{2} \left| 00 \right\rangle + 0 \left| 01 \right\rangle + \frac{1}{4} \left| 10 \right\rangle + \frac{1}{4} \left| 11 \right\rangle \right) = ???$$

The state of an *n*-bit classical computer is a vector in  $\mathbb{R}^{2^n}$  with only one coefficient nonzero. (we already saw this in explaining classical computers)

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The state of an n-qubit quantum computer is a vector in  $\mathbb{C}^{2^n}$  that is normalized (this reflects the underlying quantum property of superposition):

$$a_0\left|00\right\rangle+a_1\left|01\right\rangle+a_2\left|10\right\rangle+a_3\left|11\right\rangle$$

with  $a_i \in \mathbb{C}$  and  $\sum_{i=0}^{n-1} |a_i|^2 = 1$  (unit vectors)



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This produces  $|00\rangle$  with probability  $|a_0|^2$ ,  $|01\rangle$  with probability  $|a_1|^2$ , etc.

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$$\begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 1 \\ 0 & 0 & 1 & 0 \end{bmatrix}$$

Phase Rotation

$$\begin{bmatrix} 1 & 0 \\ 0 & e^{i\frac{\pi}{4}} \end{bmatrix}$$

# Comparison to Classical

$$\mathrm{NOT}\left(1\left|0\right\rangle+0\left|1\right\rangle\right)=0\left|0\right\rangle+1\left|1\right\rangle$$

$$\mathrm{NOT}\left(0\left|0\right\rangle+1\left|1\right\rangle\right)=1\left|0\right\rangle+0\left|1\right\rangle$$

# Comparison to Classical

$$\begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix} (1 |0\rangle + 0 |1\rangle) = 1 |0\rangle + 0 |1\rangle$$
$$\begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix} (0 |0\rangle + 1 |1\rangle) = 0 |0\rangle + 1 |1\rangle$$

## Comparison to Classical

Superposition!

$$egin{bmatrix} 0 & 1 \ 1 & 0 \end{bmatrix} \left( rac{1}{\sqrt{2}} \ket{0} + rac{1}{\sqrt{2}} \ket{1} 
ight) = rac{1}{\sqrt{2}} \left( \ket{0} + \ket{1} 
ight)$$

Like applying the gate twice 'at the same time'!

#### c-NOT

Want a gate to 'conditionally' apply its effect. Control bit controls whether gate act, and the gate acts on the second bit.

$$\begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 1 \\ 0 & 0 & 1 & 0 \end{bmatrix} (|0\rangle \otimes |0\rangle) = |0\rangle \otimes |0\rangle$$

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#### Phase estimation

Phase-estimate

**Input**:  $\omega \in [0,1]$  unknown. We get a state of the form

$$|\psi\rangle = \frac{1}{\sqrt{2}^n} \sum_{i=0}^{2^n-1} e^{2\pi i \omega y} |y\rangle$$

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What does the output even look like? Need to approximate  $\omega$  with n bits when measured.

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**Output**: An estimate of  $\omega = 0.x_1x_2...x_n = \frac{x}{2^n}$  for  $x \in \{0, 1, 2, ..., 2^n - 1\}$ 

It turns out

$$\frac{1}{\sqrt{2}^{n}} \sum_{i=0}^{2^{n}-1} e^{2\pi i \omega y} |y\rangle = \frac{1}{\sqrt{2}^{n}} \left( |0\rangle + e^{2\pi i 0.x_{n}} |1\rangle \right) \otimes \left( |0\rangle + e^{2\pi i 0.x_{n-1}x_{n}} |1\rangle \right)$$

$$\otimes \dots$$

$$\otimes \left( \left| 0 \right\rangle + e^{2\pi i 0.x_2 x_3...x_{n-1}} \left| 1 \right\rangle \right) \otimes \left( \left| 0 \right\rangle + e^{2\pi i 0.x_1 x_2...x_n} \left| 1 \right\rangle \right)$$

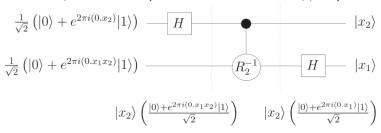
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Idea: 'pull'  $x_i$  out one by one, then remove them from the remaining bits.

#### Implementation (we will examine what happens):



$$H = \frac{1}{\sqrt{2}} \begin{bmatrix} 1 & 1 \\ 1 & -1 \end{bmatrix}$$

Looking at

$$Hrac{1}{\sqrt{2}}(\ket{0}+e^{2\pi i(0.x_2)}\ket{1})=Hrac{1}{\sqrt{2}}(\ket{0}+(-1)^{x_2}\ket{1})=\ket{x_2}$$

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So H 'pulls'  $x_2$  out for us, into a qbit!

We have  $|x_2\rangle$  after the first H, and we want to 'eliminate' it from the 2nd qubit:

$$\frac{1}{\sqrt{2}} \left( |0\rangle + e^{2\pi i (0.x_2)} |1\rangle \right) \qquad H \qquad |x_2\rangle$$

$$\frac{1}{\sqrt{2}} \left( |0\rangle + e^{2\pi i (0.x_1x_2)} |1\rangle \right) \qquad R_2^{-1} \qquad H \qquad |x_1\rangle$$

$$|x_2\rangle \left( \frac{|0\rangle + e^{2\pi i (0.x_1x_2)} |1\rangle}{\sqrt{2}} \right) \qquad |x_2\rangle \left( \frac{|0\rangle + e^{2\pi i (0.x_1)} |1\rangle}{\sqrt{2}} \right)$$

$$R_2^{-1} = \begin{bmatrix} 1 & 0 \\ 0 & e^{-2\pi i(0.01)} \end{bmatrix}$$

 $R_2$  only rotates the phase on  $|1\rangle$ .

Action when  $x_2 = 1$ :

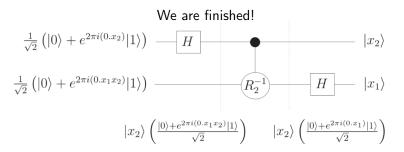
$$R_2^{-1} \frac{1}{\sqrt{2}} (|0\rangle + e^{2\pi i (0.x_1 1)} |1\rangle) = \frac{1}{\sqrt{2}} (|0\rangle + e^{2\pi i (0.x_1 0)} |1\rangle)$$

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So we want to perform  $R_2^{-1}$  only when  $x_2 = 1$ ! So we use c- $R_2^{-1}$  instead.



### Phase estimation technicalities

This analysis is really incomplete:

• We only saw the action  $\omega = 0.x_1x_2$ . What if  $\omega$  is actually like 2/3, and cannot be written as  $x/2^n$ ?

### Phase estimation technicalities

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- We only saw the action  $\omega = 0.x_1x_2$ . What if  $\omega$  is actually like 2/3, and cannot be written as  $x/2^n$ ?
- Circuit gets more and more complicated for more qubits, but follows the same idea of 'pulling' then repeatedly eliminating.

### Phase estimation backwards: QFT

When  $\omega \approx 0.x_1x_2, \dots x_n$ , phase estimation now gives us

$$\frac{1}{\sqrt{2}^n}\sum_{i=0}^{2^n-1}e^{2\pi i\omega y}|y\rangle\mapsto|x_1x_2\ldots x_n\rangle$$

With high probability.

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This replicates the behavior of the discrete fourier transform in a quantum implementation, hence QFT. Phase estimation is called "QFT $^{-1}$ ".

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FACTOR

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**Output**: A factor of N, which is either p or q.

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The idea of reductions is very common in computer science, we write  $FACTOR \rightarrow ORDER$ -FIND



# Order finding

Order-find

**Input**: a, N with gcd(a, N) = 1.

**Output**: Minimum r such that  $a^r = 1 \pmod{N}$ 

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$$a^r \equiv 1 \pmod{N}$$
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We've just factorized N!  $a^{r/2} + 1$  and  $a^{r/2} - 1$  are almost always nontrivial (meaning not 1 and N) factors.



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$$U_a \colon \ket{s} \mapsto \ket{sa} \pmod{N}$$

$$U_a^r\colon \ket{s}\mapsto \ket{sa^r}\pmod{\mathsf{N}}=\ket{s}$$

So  $U_a^r = I!$ 

This means its eigenvalues are all rth roots of unity, i.e. complex numbers of the form

$$e^{2\pi ik/r}, k \in \{0,\ldots,r-1\}$$



For an eigenvector of  $U_a$ ,  $|\psi\rangle$ 

$$U_{a}\left|\psi
ight
angle =e^{2\pi ik/r}\left|\psi
ight
angle$$

For an eigenvector of  $U_{a}$ ,  $|\psi\rangle$ 

$$U_{\mathsf{a}}\ket{\psi} = \mathsf{e}^{2\pi i k/r}\ket{\psi}$$

Modulo some details (k,  $\psi$ ), we already have a way to estimate  $\omega = k/r$  from phases!!

# Eigenvalue estimation

#### EIGENVALUE-ESTIMATE

**Input**: A unitary operator U implemented in quantum gates, and

an eigenvector  $|\psi\rangle$ 

**Output**:  $\omega$  such that  $U|\psi\rangle = e^{2\pi i\omega}|\psi\rangle$ 

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Idea: Apply  $\,U$  repeatedly, so we get a quantum state where we can estimate  $\,\omega$  from our phase estimation algorithm

$$U|\psi\rangle = e^{2\pi i 0.x_1...x_n}|\psi\rangle$$

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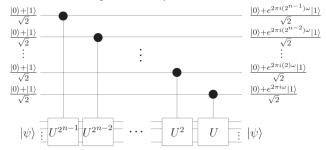
$$U^2 |\psi\rangle = (e^{2\pi i 0.x_1...x_n})^2 |\psi\rangle = e^{2\pi i x_1.x_2...x_{n-1}} |\psi\rangle = e^{2\pi i 0.x_2...x_{n-1}} |\psi\rangle$$

$$\begin{split} U |\psi\rangle &= e^{2\pi i 0.x_1...x_n} |\psi\rangle \\ U^2 |\psi\rangle &= \left(e^{2\pi i 0.x_1...x_n}\right)^2 |\psi\rangle = e^{2\pi i x_1.x_2...x_{n-1}} |\psi\rangle = e^{2\pi i 0.x_2...x_{n-1}} |\psi\rangle \\ U^{2^j} |\psi\rangle &= e^{2\pi i 0.x_j...x_{n-j}} |\psi\rangle \end{split}$$

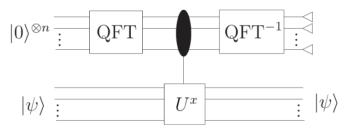
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So we can use  $c-U^{2^j}$  to get the individual qu-bits in the state we want for eigenvalue estimation.

#### So here is our way to 'set up' the states for estimation:



It turns out  $\mathrm{QFT}$  sets up 0 states to  $\frac{|0\rangle+|1\rangle}{2}$ , so here is our entire diagram (including actually doing phase estimation):



# Shor's 'entire' algorithm

We have talked about a way to implement Shor's algorithm (ignoring crucial details like time complexity and correctness):

 $FACTOR \rightarrow ORDER-FIND \rightarrow$ 

Eigenvalue-estimate  $\rightarrow$  Phase-estimate

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This isn't a huge deal yet. The largest number factored with Shor's algorithm is 21, RSA numbers are on the order of  $2^{1024}$ .

But be careful what you tell people about in 20 years.

### End

Thank you!

Diagrams and much of material from 'An Introduction to Quantum Computing', Kaye, Laflamme, Mosca