

# Max Proft U5190335

## Assignment 4

### Q1

The critical point occurs when  $\frac{\partial p}{\partial v} = \frac{\partial^2 p}{\partial v^2} = 0$

Finding  $\frac{\partial p}{\partial v}$

$$p[v] = \frac{k T}{v - b} \text{Exp}\left[\frac{-a}{k T v}\right];$$

`FullSimplify[p'[v]]`

$$-\frac{e^{-\frac{a}{k T v}} (a (b - v) + k T v^2)}{(b - v)^2 v^2}$$

Setting  $\frac{\partial p}{\partial v} = 0$  we get:

$$a (b - v) + k T v^2 = 0$$

$$\Rightarrow T = \frac{-a(b-v)}{v^2 k}$$

Finding and setting  $\frac{\partial^2 p}{\partial v^2} = 0$

`FullSimplify[p''[v]]`

$$\frac{e^{-\frac{a}{k T v}} (a^2 (b - v)^2 - 2 a k T (b - 2 v) (b - v) v + 2 k^2 T^2 v^4)}{k T v^4 (-b + v)^3}$$

$$\Rightarrow a^2 (b - v)^2 - 2 a k T (b - 2 v) (b - v) v + 2 k^2 T^2 v^4 = 0$$

Substituting in the expression for T in order to find  $v_c$ :

$$\Rightarrow 3 a^2 (b - v_c)^2 + \frac{2 a^2 (b - 2 v_c) (b - v_c)^2}{v_c} = 0$$

$$\Rightarrow 3 v_c + 2 (b - 2 v_c) = 0$$

$$\Rightarrow v_c = 2 b$$

Substituting this back into T to find  $T_c$ :

$$T_c = \frac{a}{4 b k}$$

Substituting this into p to find  $p_c$ :

$$p_c = \frac{k}{2 b - b} \frac{a}{4 b k} \text{Exp}\left[\frac{-a}{k 2 b} \frac{4 b k}{a}\right]$$

$$p_c = \frac{a}{4 b^2 e^2}$$

$$\frac{p_c v_c}{k T_c} = \frac{\frac{a}{4 b^2 e^2} 2 b}{k \frac{a}{4 b k}} = \frac{2}{e^2}$$

$$\bar{p} = \frac{p}{p_c} = \frac{2 \bar{T} v_c}{2 v - v_c} \frac{e^2}{2} \text{Exp}\left[\frac{-a}{k \bar{T} T_c v v_c}\right] = \frac{2 \bar{T}}{2 \bar{V} - 1} \frac{e^2}{2} \text{Exp}\left[\frac{-2}{\bar{T} \bar{V}}\right]$$

Finding the expansion coefficients:

Along the coexistance curve:  $\bar{p}_g = \bar{p}_L$ , and  $\bar{T}_g = \bar{T}_L$

$$\frac{2\bar{T}}{2V_g-1} \frac{\epsilon^2}{2} \text{Exp}\left[\frac{-2}{\bar{T}V_g}\right] = \frac{2\bar{T}}{2V_L-1} \frac{\epsilon^2}{2} \text{Exp}\left[\frac{-2}{\bar{T}V_L}\right]$$

$$\text{Exp}\left[\frac{-2}{\bar{T}}\left(\frac{1}{V_g} - \frac{1}{V_L}\right)\right] = \frac{2\bar{V}_g-1}{2\bar{V}_L-1}$$

$$\bar{T} = \frac{-2\left(\frac{1}{V_g} - \frac{1}{V_L}\right)}{\ln\left(\frac{2\bar{V}_g-1}{2\bar{V}_L-1}\right)}$$

Setting  $\bar{V}_g \approx 1 + \epsilon/2$  and  $\bar{V}_L \approx 1 - \epsilon/2$

$$\bar{T} = \frac{-2\left(\frac{\bar{V}_L-\bar{V}_g}{\bar{V}_g\bar{V}_L}\right)}{\ln(1+\epsilon)-\ln(1-\epsilon)} = \frac{-2\left(\frac{1-\epsilon/2-(1+\epsilon/2)}{(1-\epsilon/2)(1+\epsilon/2)}\right)}{\epsilon-(-\epsilon)} = \frac{-2\left(\frac{-\epsilon}{1-\epsilon^2/4}\right)}{2\epsilon} = \frac{1}{1-\epsilon^2/4}$$

From Binomial Expansion:

$$\bar{T} \approx 1 + \epsilon^2/4$$

Since  $\epsilon = V_g - V_L$ , if we rearrange we get:

$$(V_g - V_L) \sim (T - T_c)^{-0.5}$$

And so we get  $\beta = -0.5$

Along an isotherm:

$$p = \frac{kT}{v-b} \text{Exp}\left[\frac{-a}{kT(v-b)}\right]$$

If we taylor expand this about  $v = v_c$  we get (By definition of the critical point, the first and second derivatives are zero):

$$\begin{aligned} p &= p_c + \frac{d^3\left(\frac{kT}{v-b} \text{Exp}\left[\frac{-a}{kT(v-b)}\right]\right)}{dv^3} \Big|_{v=v_c} (v - v_c)^3 \\ p &= p_c - \frac{\text{Exp}\left[-\frac{a}{kT v_c}\right] (a(b-v_c) + k v_c^2)}{(b-v_c)^2 v_c^2} (v - v_c)^3 \\ p - p_c &= \frac{4k(-1+T_c)}{\epsilon^2 v_c^2} (v_c - v)^3 \end{aligned}$$

And so we get  $\delta = 3$

$$\kappa = \frac{-1}{v} \frac{\partial v}{\partial p} \Big|_T$$

Finding  $\frac{\partial p}{\partial v}$ , and evaluating this at  $v = v_c$

$$\text{FullSimplify}\left[D\left[\frac{kT}{v-v_c/2} \text{Exp}\left[\frac{-Tc*2vc k}{kT v}\right], v\right]\right]$$

$$-\frac{4 e^{-\frac{2 T c v c}{T v}} k \left(T v^2 + T c v c (-2 v + v c)\right)}{v^2 (-2 v + v c)^2}$$

$$\text{FullSimplify}\left[-\frac{4 e^{-\frac{2 T c v c}{T v}} k \left(T v^2 + T c v c (-2 v + v c)\right)}{v^2 (-2 v + v c)^2} / . v \rightarrow v c\right]$$

$$\frac{4 e^{-\frac{2 T c}{T}} k (-T + T c)}{v c^2}$$

Hence we get:

$$K = \frac{-1}{V_c} \frac{V_c}{4 e^{-\frac{2T_c}{T}} k (-T+T_c)} = \frac{e^{2/T}}{4 k} \frac{1}{T-T_c}$$

And hence  $\gamma=1$

## Q2

Throughout this question, assume that  $a, b, c$  are functions of  $T$ .

$$b < 0$$

$$c > 0$$

Stable solutions occur when  $\frac{\partial F}{\partial m} = 0$  and  $\frac{\partial^2 F}{\partial m^2} > 0$

Solving  $\frac{\partial F}{\partial m} = 0$  we get either  $m=0$  or:

$$a + 2 b m^2 + 3 c m^4 = 0$$

$$m^2 = \frac{-b \pm \sqrt{b^2 - 3ac}}{3c}$$

If  $b^2 - 3ac < 0$ , we will only have 1 real solution,  $m=0$ .

$$a > \frac{b^2}{3c} > 0$$

Suppose that  $b^2 - 3ac > 0$  and that  $-b - \sqrt{b^2 - 3ac} < 0$ .

$$\Rightarrow b^2 - 3ac > b^2$$

$$\Rightarrow ac < 0$$

$$\Rightarrow a < 0$$

This means we have 3 solutions:

$$m=0$$

$$m = \pm \sqrt{\frac{-b \pm \sqrt{b^2 - 3ac}}{3c}}$$

Suppose that  $b^2 - 3ac > 0$  and that  $-b - \sqrt{b^2 - 4ac} > 0$ .

$$\Rightarrow b^2 - 3ac < b^2$$

$$\Rightarrow ac > 0$$

$$\Rightarrow a > 0$$

This means we have 5 solutions:

$$m=0 \text{ or:}$$

$$m = \pm \sqrt{\frac{-b \pm \sqrt{b^2 - 3ac}}{3c}}$$

I assume that  $a \neq 0$ , as this case can be considered as the limit of the above cases.

Finding out which of these solutions are stable:

i.e.  $\frac{\partial^2 F}{\partial m^2} > 0$

$$\frac{\partial^2 F}{\partial m^2} = 2 a + 12 b m^2 + 30 c m^4$$

1 solution:

$$m=0 \Rightarrow \frac{\partial^2 F}{\partial m^2} = 2 a$$

Since  $a > \frac{b^2}{3c} > 0$ , this solution is stable.

3 Solutions:

$$m=0 \Rightarrow \frac{\partial^2 F}{\partial m^2} = 2 a$$

Since  $a < 0$ , this is unstable.

$$\begin{aligned} m &= \pm \sqrt{\frac{-b + \sqrt{b^2 - 3ac}}{3c}} \\ \frac{\partial^2 F}{\partial m^2} &= 2a + 12b \frac{-b + \sqrt{b^2 - 3ac}}{3c} + 30c \left( \frac{-b + \sqrt{b^2 - 3ac}}{3c} \right)^2 \\ &= 2a + 4b \frac{-b + \sqrt{b^2 - 3ac}}{c} + 10 \left( \frac{b^2 - 2b\sqrt{b^2 - 3ac} + b^2 - 3ac}{3c} \right) \\ &= 2a - \frac{4b^2}{c} + \frac{4b}{c} \sqrt{b^2 - 3ac} + \frac{20b^2}{3c} - \frac{20b}{3c} \sqrt{b^2 - 3ac} - 10a \\ &= -8a + \frac{8b^2}{3c} - \frac{8b\sqrt{b^2 - 3ac}}{3c} \end{aligned}$$

Let  $\Delta = \sqrt{b^2 - 3ac}$

$$\frac{\partial^2 F}{\partial m^2} = \frac{8\Delta^2}{3c} - 8 \frac{b\Delta}{3c}$$

Since  $\Delta > 0$ ,  $c > 0$ ,  $b < 0$ , both of these solutions have  $\frac{\partial^2 F}{\partial m^2} > 0$ , and so they are stable.

5 Solutions:

$$m=0 \Rightarrow \frac{\partial^2 F}{\partial m^2} = 2 a$$

Since  $a > 0$ , this is stable.

$$m^2 = \frac{-b \pm \sqrt{b^2 - 3ac}}{3c}$$

The + solution is the same as in the 3 solution case, and so we get these solutions to be stable.

Now consider only the - solution

$$\begin{aligned} \frac{\partial^2 F}{\partial m^2} &= 2a + 12b \frac{-b - \sqrt{b^2 - 3ac}}{3c} + 30c \left( \frac{-b - \sqrt{b^2 - 3ac}}{3c} \right)^2 \\ &= 2a + 4b \frac{-b - \sqrt{b^2 - 3ac}}{c} + 10 \left( \frac{b^2 + 2b\sqrt{b^2 - 3ac} + b^2 - 3ac}{3c} \right) \\ &= 2a - \frac{4b^2}{c} - \frac{4b}{c} \sqrt{b^2 - 3ac} + \frac{20b^2}{3c} + \frac{20b}{3c} \sqrt{b^2 - 3ac} - 10a \\ &= -8a + \frac{8b^2}{3c} + \frac{8b\sqrt{b^2 - 3ac}}{3c} \end{aligned}$$

Let  $\Delta = \sqrt{b^2 - 3ac}$

$$\frac{\partial^2 F}{\partial m^2} = \frac{8\Delta^2}{3c} + 8 \frac{b\Delta}{3c} = \frac{8\Delta}{3c} (\Delta + b)$$

This is stable if  $\Delta > -b$

$$\Rightarrow -3ac > 0$$

$$\Rightarrow a < 0$$

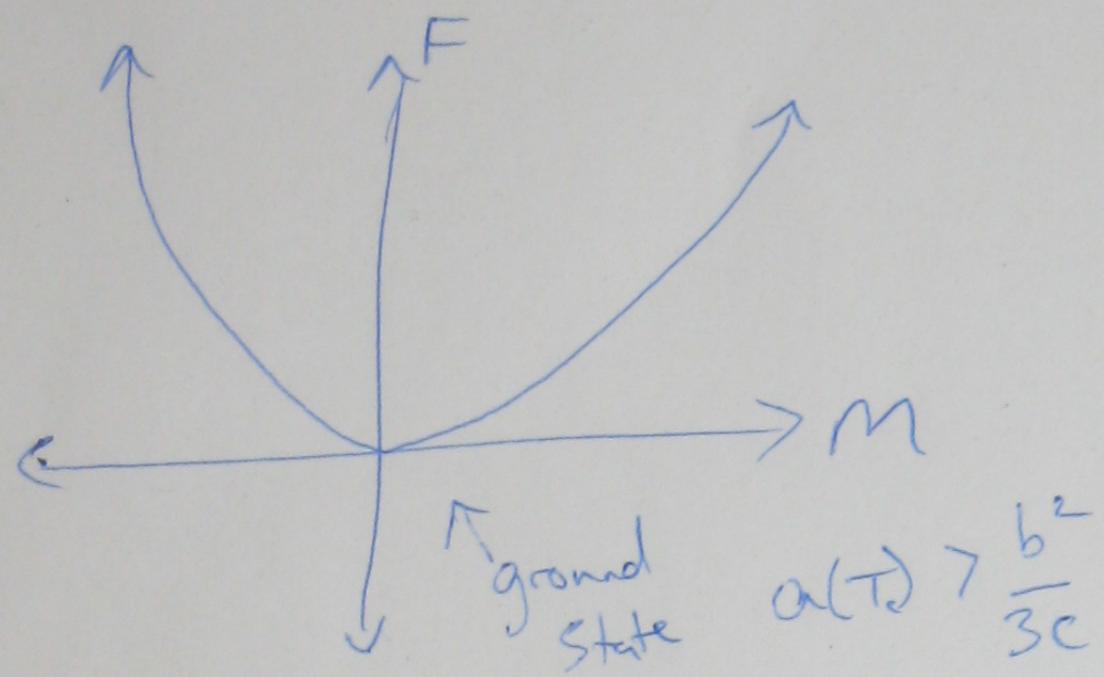
But for this solution, we have  $a > 0$ .

Hence it cannot be stable.

Sketches attached at the end.

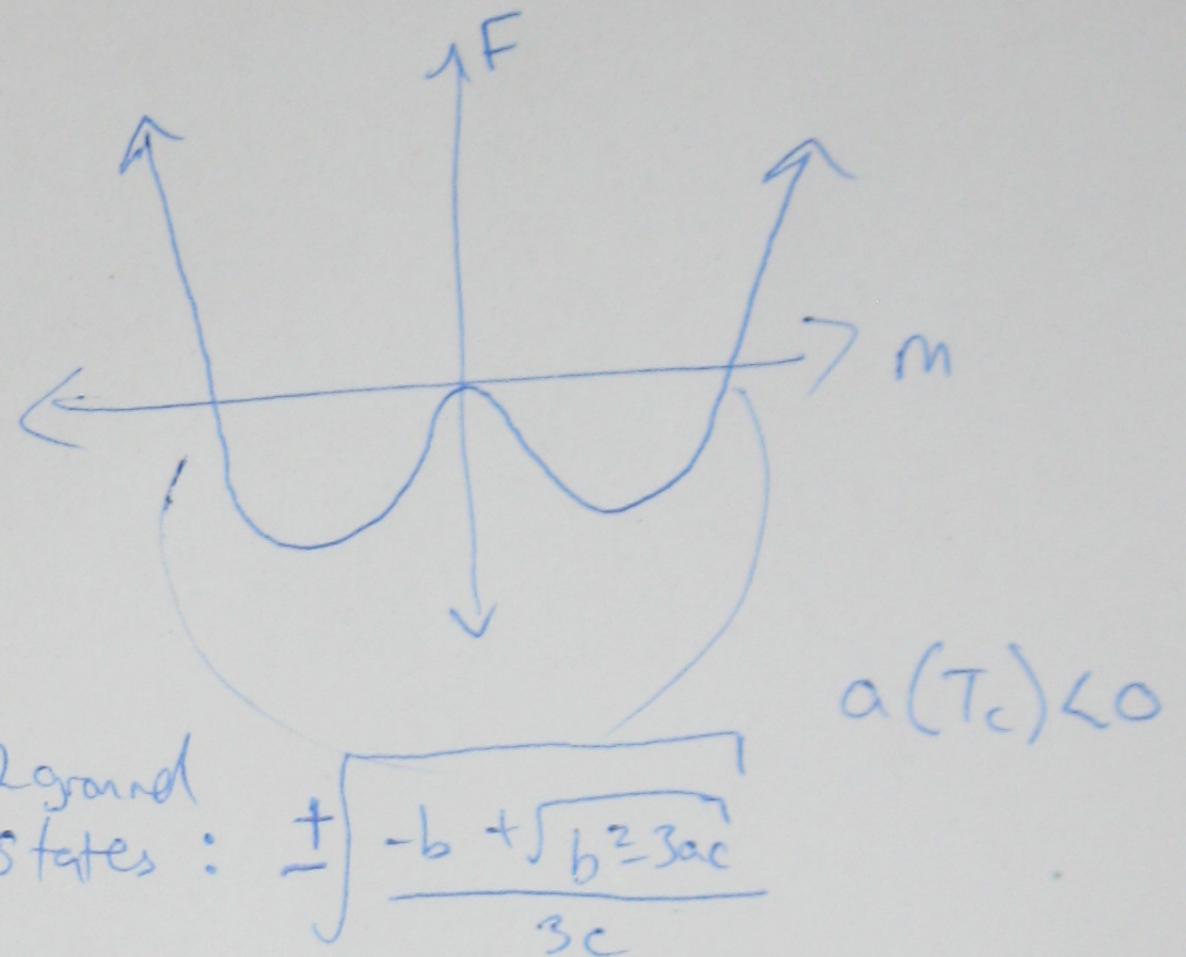
Q2

1 Solution



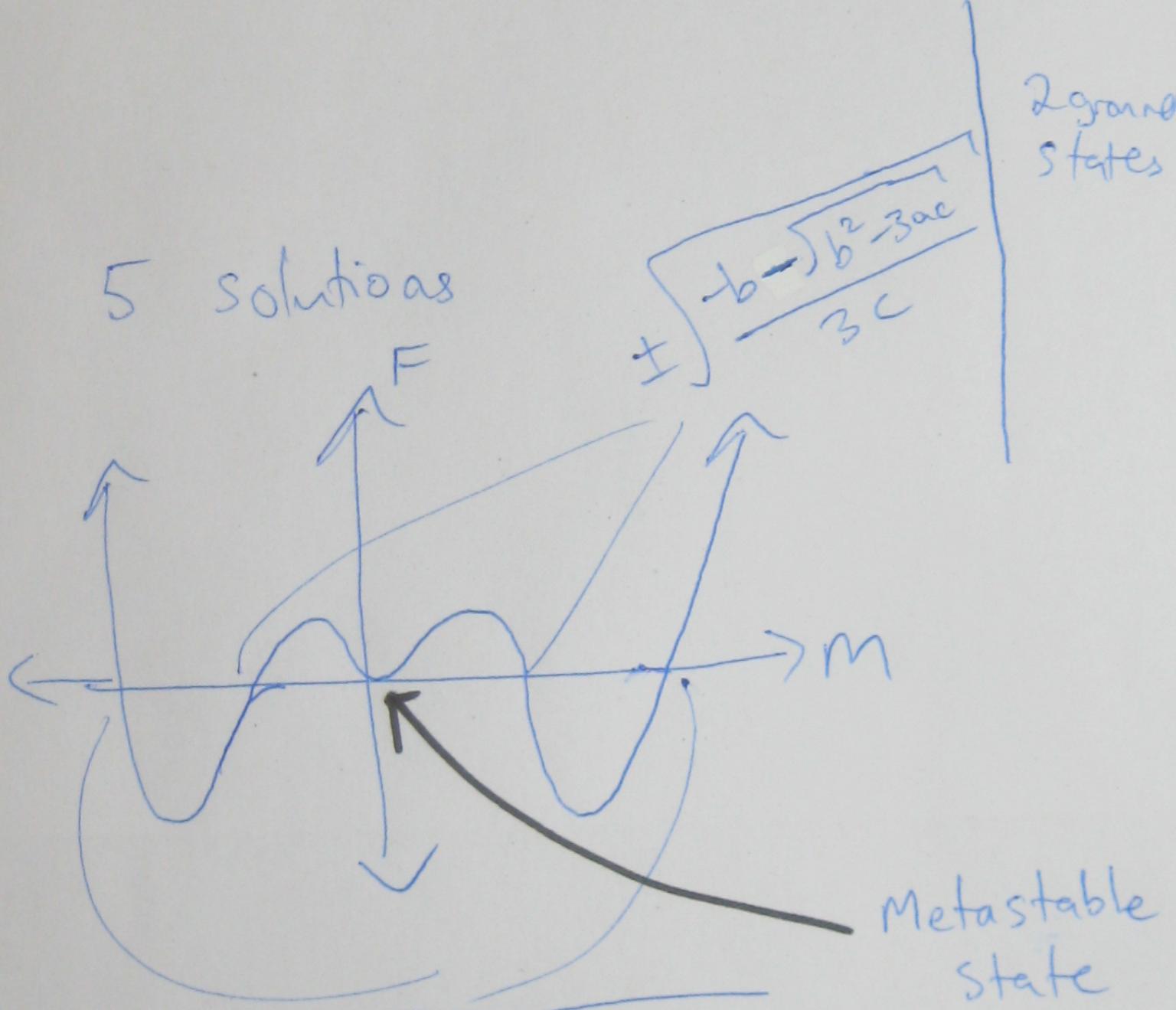
$$a(T) > \frac{b^2}{3c}$$

3 solutions



$$a(T_c) < 0$$

5 solutions



Metastable state

$$2 \text{ ground states: } \pm \sqrt{\frac{-b + \sqrt{b^2 - 3ac}}{3c}}$$

$$0 < a(T_c) < \frac{b^2}{3c}$$

Q3. Since Isotropic, we can treat this as a 1D problem.  
 let  $f = f_0(v) + f_1(x, v, t)$

$$E = E_0 + E_1, \quad E_0 = 0, \quad B = 0.$$

~~$$\frac{\partial f}{\partial t} + v \cdot \nabla f + \frac{q}{m} (E + v \times B) \cdot \frac{\partial f}{\partial v} = 0$$~~

~~Please~~ keeping 1st order terms only:

$$\Rightarrow \partial_t f_1 + v \partial_x f_1 + \frac{q}{m} E_1 \cancel{\partial_v f_0} = 0$$

$$\Rightarrow (\partial_t + v \partial_x) f_1 = -\frac{q}{m} E_1 \cancel{\partial_v f_0}$$

$$= -\frac{q}{m} E_1 \frac{-2va}{\pi(a^2+v^2)^2}$$

$$= \frac{2va\pi q E_1}{am} f_0(v)^2$$

$$\Rightarrow \frac{am}{2\pi va f_0(v)^2} (\partial_t + v \partial_x) f_1 = E_1$$

If there is a perturbation of angular frequency  $\omega$ ,

~~$$E_1 \propto e^{-(i\omega t - ikx)}$$~~

$$\Rightarrow f_1 \propto e^{-(i\omega t - ikx)}$$

$$\Rightarrow E_1 = \frac{am}{2\pi va f_0(v)^2} (-i\omega + ikv) f_1(x, v, t)$$

$$\Rightarrow f_1(x, v, t) = \frac{2\pi va E_1 f_0(v)^2}{iam(kv - \omega)}$$

From  $\nabla \cdot E = \frac{P}{\epsilon_0}$  we get:

To first order we get  $\nabla \cdot \vec{E}_0 = e n_{i0} - e n_{e0} \Rightarrow n_{i0} = n_{e0}$

~~$i k E_1 \epsilon = e n_i - e n_e$~~  (assuming ion has charge +e)

~~since  $f_{i0}$  and  $f_{e0}$  are normalized~~

~~(free parameter)~~

To second order we get:

$$i k E_1 \epsilon = e n_i - e n_e$$

$$= e(n_{i0} + n_{i1}) - e(n_{e0} + n_{e1}) = e(n_{i0} - n_{e1})$$

$$= e \int f_{i1} dv - e \int f_{e1} dv$$

Substituting in for  $f_{i1}$  and  $f_{e1}$ , and dividing both sides by  $E_1 \epsilon$  we get our dispersion relation:

$$\frac{i k \epsilon}{\epsilon} = \frac{2\pi e}{i a M_i} \int_{-\infty}^{\infty} \frac{v f_{i1}(v)^2}{kv - \omega} dv + \frac{2\pi e}{i a m_e} \int_{-\infty}^{\infty} \frac{v f_{e1}(v)}{kv - \omega} dv$$