

Let \vec{x} represent a point on a unit sphere.

a) The grand canonical partition function is given by

$$Z = \frac{1}{(h^3)^N N!} \int d\vec{p}_1 \dots d\vec{p}_N d\vec{x}_1 \dots d\vec{x}_N e^{-\hat{H}\beta + \mu N(\vec{p}_1, \dots, \vec{x}_N)\beta}$$

$$= \frac{1}{N! h^{3N}} \int d\vec{p}_1 \dots d\vec{p}_N d\vec{x}_1 \dots d\vec{x}_N \left[e^{-\sum_{i=1}^N \left(\frac{p_i^2}{2m} \beta - M_0 \vec{H} \cdot \hat{n}_i \beta \right) + \mu \beta N(\vec{p}_1, \dots, \vec{x}_N)} \right]$$

Since each i represents one particle

we can move $N(\vec{p}_1, \dots, \vec{x}_N)$ inside the summation.

$$= \frac{1}{N! h^{3N}} \int d\vec{p}_1 \dots d\vec{p}_N d\vec{x}_1 \dots d\vec{x}_N e^{\sum_{i=1}^N \left(\mu \beta + M_0 \beta \vec{H} \cdot \hat{n}_i - \frac{p_i^2}{2m} \beta \right)}$$

$$= \frac{1}{N! h^{3N}} \left(\int d\vec{p} d\vec{x} e^{\mu \beta + M_0 \beta \vec{H} \cdot \hat{n} - \frac{p^2}{2m} \beta} \right)^N$$

Pulling out the factor of $e^{\mu \beta}$ and performing the integral with respect to $d\vec{p}$ we get:

$$= \frac{1}{N!} \left(\sqrt{\frac{2\pi m}{\beta h^2}} \right)^{3N} e^{\mu \beta N} \left(\int d\vec{x} e^{M_0 \beta \vec{H} \cdot \hat{n}} \right)$$

Orienting our axes so that \vec{H} lies in the \hat{z} direction, and using spherical coordinates (unit sphere \Rightarrow radius = 1)

$$= \frac{1}{N!} \left(\sqrt{\frac{2\pi m}{\beta h^2}} \right)^{3N} e^{\mu \beta N} \left(\int_0^\pi d\theta \int_0^{2\pi} d\phi e^{M_0 \beta H \cos \theta} \sin \theta \right)^N$$

$$= \frac{(2\pi)^N}{N!} \left(\sqrt{\frac{2\pi M}{\beta h^2}} \right)^{3N} e^{M\beta N} \left(\frac{e^{\frac{H\beta M_0}{H\beta M_0}} - e^{-\frac{H\beta M_0}{H\beta M_0}}}{H\beta M_0} \right)^N$$

$$= \frac{(4\pi)^N}{N!} \left(\sqrt{\frac{2\pi M}{\beta h^2}} \right)^{3N} e^{M\beta N} \left(\frac{\sinh(H\beta M_0)}{H\beta M_0} \right)^N$$

Using Mathematica for the rest.

b)

We have found the Grand Partition Function to be:

$$\text{In}[15]:= Z[\beta_-] = 4 \frac{\pi}{n!} \left(\frac{2 \pi m}{\beta h^2} \right)^{3 n/2} \text{Exp}[\mu \beta n] \left(\frac{\text{Sinh}[H \beta \mu_0]}{H \beta \mu_0} \right)^n ;$$

The Grand Potential is given by:

$$\Omega = \frac{-1}{\beta} \text{Ln}(Z)$$

Expanding out all of the terms we get:

$$\begin{aligned} \text{In}[27]:= \Omega[\beta_-] &= \text{Expand} \left[\text{PowerExpand} \left[\frac{-1}{\beta} \text{Log}[Z[\beta]] \right] \right] \\ \text{Out}[27]= & -n \mu - \frac{2 \text{Log}[2]}{\beta} - \frac{3 n \text{Log}[2]}{2 \beta} + \frac{3 n \text{Log}[h]}{\beta} + \frac{n \text{Log}[H]}{\beta} - \frac{3 n \text{Log}[m]}{2 \beta} - \\ & \frac{\text{Log}[\pi]}{\beta} - \frac{3 n \text{Log}[\pi]}{2 \beta} + \frac{5 n \text{Log}[\beta]}{2 \beta} + \frac{n \text{Log}[\mu_0]}{\beta} + \frac{\text{Log}[n!]}{\beta} - \frac{n \text{Log}[\text{Sinh}[H \beta \mu_0]]}{\beta} \end{aligned}$$

c)

Finding the magnetism:

$$\begin{aligned} M &= -\frac{\partial \Omega}{\partial H} = -\frac{\partial}{\partial H} \left(\frac{n \text{Log}[H]}{\beta} - \frac{n \text{Log}[\text{Sinh}[H \beta \mu_0]]}{\beta} \right) \\ M &= -\frac{n}{H \beta} + n \mu_0 \text{Coth}[H \beta \mu_0] \end{aligned}$$