

# Taylor Expansions of UCC

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# Unitary Coupled Cluster Theory

- Coupled cluster theory is (kind of) accurate and size-consistent, but non-variational
- Unitary coupled cluster (UCC) is variational, but classically intractable
- Its Trotterized form(s) can be implemented on quantum hardware

$$|\psi_{CC}\rangle = e^{\hat{T}} |0\rangle$$

Make ansatz unitary  $\downarrow$

$$|\psi_{UCC}\rangle = e^{\hat{T} - \hat{T}^\dagger} |0\rangle$$

Trotterization  $\downarrow$

$$|\psi_{tUCC}\rangle = \left( \prod_{ij\dots}^{ab\dots} e^{(t_{ab\dots}^{ij\dots} \hat{a}_{ij\dots}^{ab\dots} - \text{H.C.})} \right) |0\rangle$$

## 2nd-Order UCCSD

- BCH Expansion of UCC:

$$\begin{aligned} E_{UCCSD} &= \langle H \rangle + \langle [\hat{H}, \hat{T} - \hat{T}^\dagger] \rangle + \langle [[\hat{H}, \hat{T} - \hat{T}^\dagger], \hat{T} - \hat{T}^\dagger] \rangle, \dots \\ &= E_{ref} + 2 \langle 0 | \hat{H}_N \hat{T} | 0 \rangle + \langle 0 | \hat{T}^\dagger \hat{H}_N \hat{T} | 0 \rangle \\ &\quad + \langle 0 | \hat{H}_N \hat{T}^2 | 0 \rangle - \langle 0 | \hat{H}_N \hat{T}^\dagger \hat{T} | 0 \rangle + \dots \end{aligned}$$

- Truncating at second order in  $\hat{T}$  and only including singles and doubles yields the Simons-Hoffmann “SHUCC” equations for minimizing the energy. For canonical orbitals:

Not in LCCSD

$$\begin{aligned} 0 &= [\hat{H}_N \mathbf{t}]_a^i + \sum_{ijab} \langle 0 | \hat{H}_N | \phi_{ij}^{ab} \rangle t_b^j \\ - \langle 0 | \hat{H}_N | \phi_{ij}^{ab} \rangle &= [\hat{H}_N \mathbf{t}]_{ab}^{ij} \end{aligned}$$

# Why Does This Look Like LCCSD?

- Linearized coupled cluster (LCC) usually refers to variational minimization of a linearized version of the CC functional:

$$\epsilon[\mathbf{t}, \boldsymbol{\lambda}] = E_{ref} + \langle 0 | (1 + \hat{\Lambda}) [H_N (1 + \hat{T})]_C | 0 \rangle$$

- For an HF reference with only singles and doubles,  $\hat{\Lambda} = \hat{T}^\dagger$  and we get a similar linear stationary condition

$$\begin{aligned} 0 &= [\hat{H}_N \mathbf{t}]_a^i \\ -\langle 0 | \hat{H}_N | \phi_{ij}^{ab} \rangle &= [\hat{H}_N \mathbf{t}]_{ab}^{ij} \end{aligned}$$

# Commonalities (For Canonical Orbitals)

- Both LCCSD and SHUCCSD are size-consistent
- For STO-3G  $H_2$ , singles are irrelevant, so the methods are equivalent and amount to solving a two-level problem:

$$E_c = -\frac{|\langle ij||ab\rangle|^2}{E_1 - E_0}$$

- $E_c$  diverges as the energies of  $|0\rangle$  ( $E_0$ ) and  $|\phi_{ij}^{ab}\rangle$  ( $E_1$ ) become degenerate

- We can include “diagonal” terms up to third order in  $\hat{T}$  to treat degenerate systems. For canonical orbitals:

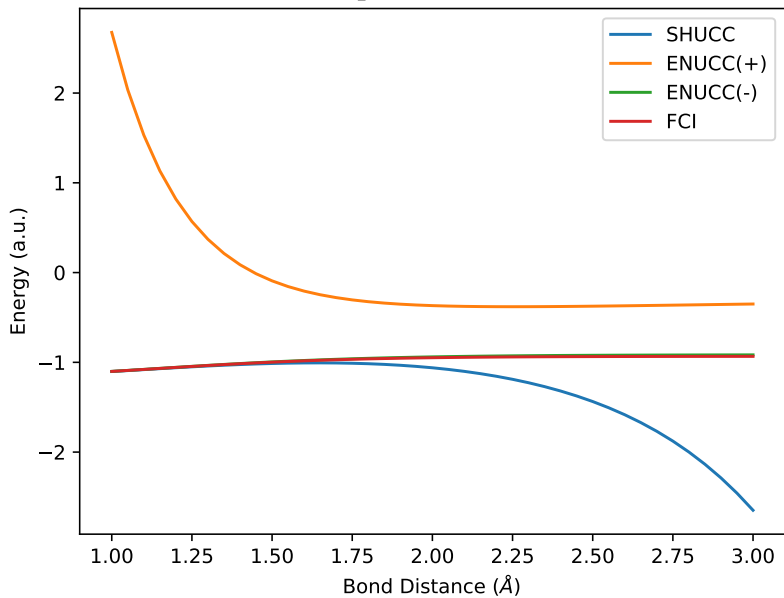
$$\epsilon_{ENUCC} = \epsilon_{SHUCC} - \frac{4}{3} \sum_{i < j}^{a < b} \langle ij || ab \rangle (t_{ab}^{ij})^3$$

- When  $|\phi_{ij}^{ab}\rangle$  and  $|0\rangle$  approach degeneracy, this gives finite stationary  $H_2$  energies:

$$E = E_0 \pm \frac{2\sqrt{2}}{3} \langle ij || ab \rangle$$

- The high-energy ENUCC solution appears to diverge in the weakly coupled regime, but the low-energy solution does not

## H<sub>2</sub>, STO-3G PES



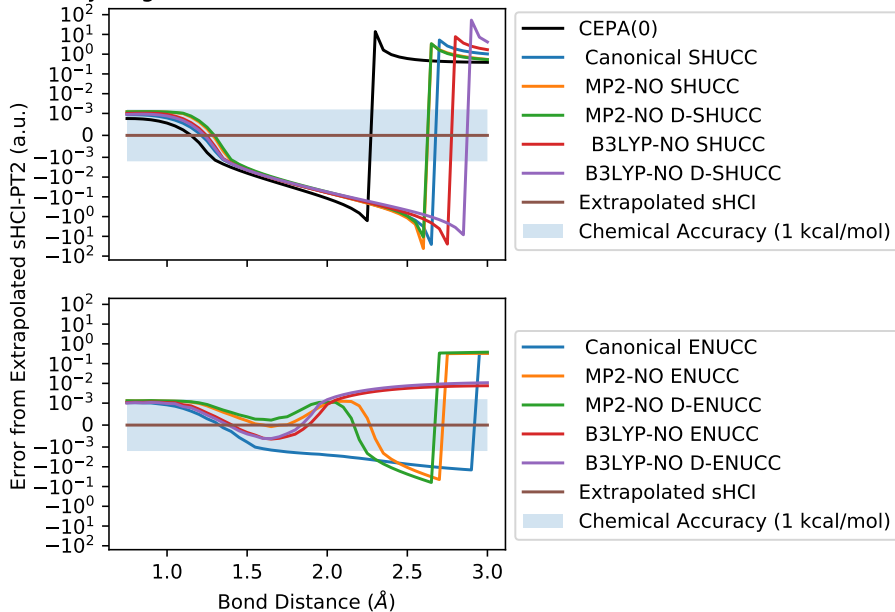
- Both SHUCCSD and ENUCCSD become size-inconsistent if we mix the occupied and virtual canonical orbitals
- We can define disentangled (D-) SHUCCSD and ENUCCSD to avoid this, redefining the energy functional as:

$$\epsilon[\mathbf{t}] = \langle 0 | e^{-\hat{T}_2 + \hat{T}_2^\dagger} e^{-\hat{T}_1 + \hat{T}_1^\dagger} \hat{H} e^{\hat{T}_1 - \hat{T}_1^\dagger} e^{\hat{T}_2 - \hat{T}_2^\dagger} | 0 \rangle$$

- Equivalently, we could define the ansatz as any Trotterized UCCSD ansatz where the singles are all closer to  $\hat{H}$  than the doubles
- It might also be possible to understand D-SHUCCSD as a Newton step toward orbital-optimized UCCD



# Hydrogen Fluoride, 6-31G\*\* Error in PES



# Understanding ENUCC

- Ignoring the singles for simplicity, ENUCC bears a similarity to the non-averaged CEPA variants, e.g. Malrieu's (SC)<sup>2</sup>CI:

$$\frac{1}{4} \sum_{klcd} \langle \phi_{ij}^{ab} | \hat{H}_N | \phi_{kl}^{cd} \rangle t_{cd}^{kl} - \frac{1}{4} \sum_{klcd}^{EPV} \langle kl || cd \rangle t_{cd}^{kl} t_{ab}^{ij} = - \langle ij || ab \rangle$$

- Only summing over terms where the occupied orbitals violate the exclusion principle gives CEPA(3)
- Only including the term in the summand where **every** orbital is EPV-violating gives the ENUCC stationary condition (down to a factor of 2):

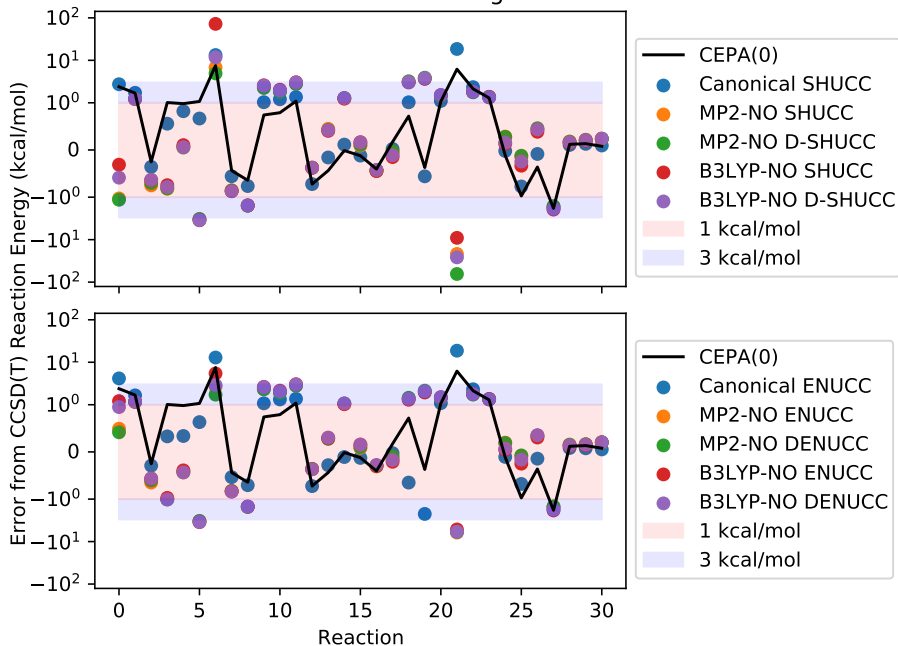
$$\frac{1}{4} \sum_{klcd} \langle \phi_{ij}^{ab} | \hat{H}_N | \phi_{kl}^{cd} \rangle t_{cd}^{kl} - \overset{\substack{\text{Only} \\ \text{in} \\ \text{ENUCC}}}{\underset{\text{2}}{\square}} \langle ij || ab \rangle (t_{ab}^{ij})^2 = - \langle ij || ab \rangle$$

# Understanding ENUCC

- ENUCC's ability to handle HF dissociation is consistent with Malrieu's analysis- averaged (or 0) corrections to the EPVs fail for degenerate cases, even in single bond breaking
- Aside from the factor of 2, the comparison breaks down when one notes that ENUCCD and LCCD family of methods have different energy expressions for their stationary wavefunctions:

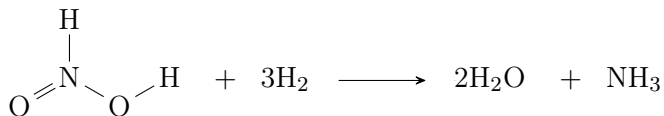
$$E_c = \frac{1}{4} \sum_{ijab} \left( \langle ij || ab \rangle t_{ab}^{ij} + \overset{\text{Only in ENUCC}}{\frac{2}{3} \langle ij || ab \rangle \left( t_{ab}^{ij} \right)^3} \right)$$

# Test Set of Reaction Energies

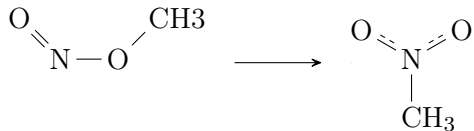


# Pathological Reactions

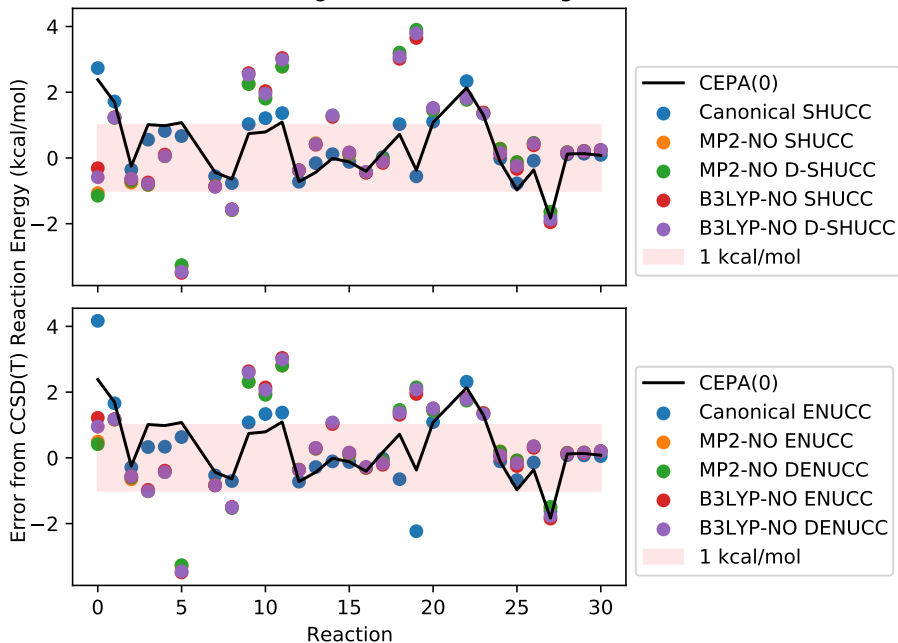
## ① Reductive Cleavage of Multiple Bonds in Nitrous Acid:



## ② Isomerization of Methyl Nitrite into Nitromethane:



# Test Set of Reaction Energies Without Pathological Cases



# Apparent Conclusions

- SHUCCSD and LCCSD have similar weaknesses
- ENUCCSD can help with degeneracies, but still struggles with breaking double bonds
- It might be possible to formulate it as a CEPA variant which only includes one EPV term per determinant
- The disentangled variant does not appear to generally improve accuracy, but does give a size-consistent, Hermitian functional for arbitrary orbitals