

## Delta-Gap Model Justification

We assume the antenna is excited at a single angular frequency  $\omega$ , so all fields are time-harmonic. This means the electric field can be written as:

$$\vec{E}(\vec{r}, t) = \text{Re} \left\{ \vec{E}(\vec{r}) e^{i\omega t} \right\}$$

In the frequency domain, Maxwell's equations give:

$$\vec{E}(\vec{r}) = -\nabla V(\vec{r}) - i\omega \vec{A}(\vec{r})$$

First of all, no tangential electric field can exist on the conductor surface. But a voltage perturbation (i.e. point-like voltage source) is introduced at the center of the antenna  $z = 0$ ...

We now integrate the  $z$ -component of this expression along the wire axis (at the surface) over a small interval  $[-\epsilon, \epsilon]$  centered at the feed point:

$$\int_{-\epsilon}^{+\epsilon} E_z(z) dz = - \int_{-\epsilon}^{+\epsilon} \frac{\partial V}{\partial z} dz - i\omega \int_{-\epsilon}^{+\epsilon} A_z(z) dz$$

We make the following assumptions:

- The vector potential  $\vec{A}$  is generated by smooth, finite currents on the antenna (recall plots of current in standing wave antenna section). Thus,  $A_z(z)$  is continuous near  $z = 0$ , and

$$\int_{-\epsilon}^{+\epsilon} A_z(z) dz \rightarrow 0 \quad \text{as } \epsilon \rightarrow 0$$

- The scalar potential  $V(z)$  has a discontinuity across the small feed gap, representing an imposed voltage  $V_1$ . Thus,

$$\int_{-\epsilon}^{+\epsilon} \frac{\partial V}{\partial z} dz = V(-\epsilon) - V(+\epsilon) = -V_1$$

Therefore, we find:

$$\int_{-\epsilon}^{+\epsilon} E_z(z) dz = V_1$$

This motivates modeling the (idealized) electric field near the feed point as a delta function:

$$E_z(z) = -V_1 \delta(z)$$

This is the essence of the **delta-gap source model**: a localized voltage drop across a small region of the antenna is idealized as a delta-function electric field, consistent with Maxwell's equations in the limit as  $\epsilon \rightarrow 0$ .