MSc Computational Science joint programme UvA/VU







MODELLING DIFFUSIVE SIGNALLING IN ASPERGILLUS SPP. GERMINATION INHIBITION

INTERMEDIATE PRESENTATION - APRIL

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2 Single-spore experiments: revision

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Diffusion and permeation consta
Inhibitor concentrations
Single-spore release scenarios

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Homogeneous vs. cell-wall-bound inhibitor
Extreme cases
Functional relationship

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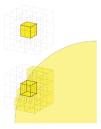






- » The topic of permeation and diffusion can be approached on different scales and with different formulations. The following models have been developed so far.
- » Numerical models (produce simulations)
 - Macroscopic scale the spore is 4 times smaller than the lattice subdivision (volume considered negligible),
 releases/absorbs concentration at a source site
 - Spore scale the spore is a special node on the lattice.
 - Cell wall scale the spore is a voxelised sphere.
- » Analytical models (produce mathematical formulas)
 - Simple permeation the spore is a spherical volume that "leaks" based on the trans-barrier concentration drop.
 - Permeation + saturation under high spore densities, the outside medium saturates







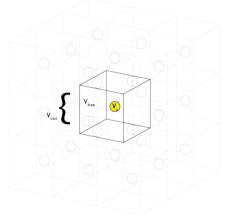








- » The latest analytical model yields results closest to all other models.
- » At a given spore density ρ_s , a spore has a designated volume $V_{\rm cell}$ around it, consisting of its own volume and the volume of the medium $(V_s + V_{\rm free})$.
- » The volume fraction $\phi = \rho_s V_s$ represents the part of the space occupied by the spore.
- It is assumed that the inhibitor diffuses fast in the medium, so its dynamics in V_{free} can be simplified.









Volume-based saturation formula

Consider the flux through a barrier, driven by the concentration drop:

$$J(t) = P_s A \left[c_{\rm in}(t) - c_{\rm out}(t) \right]. \tag{1}$$

>>

The inside and outside concentrations change over time as

$$\tau = \frac{V_s}{P_s A} \tag{5}$$

$$\dot{c}_{\rm in} = \tau^{-1} \left(c_{\rm in} - c_{\rm out} \right) \tag{2}$$

»

 $c_{\rm in}(0)=c_0 \tag{6}$

$$\dot{c}_{
m out} = au^{-1} rac{V_s}{V_{
m free}} \left(c_{
m in} - c_{
m out}
ight).$$
 (3)

 $c_{\text{out}}(0) = c_{\text{ex}} \tag{7}$

» They reach the long-time equilibrium limit

$$c_{
m eq} =
ho_s \left[V_s \, c_0 + \left(rac{1}{
ho_s} - V_s
ight) \, c_{
m ex}
ight] = \phi c_0 + (1 - \phi) c_{
m ex}$$
 (4)









The system of ODEs has an eigenvalue

$$\lambda_1 = - au^{-1}\left(1 + rac{V_s}{V_{ ext{free}}}
ight) = - au^{-1}\left(rac{1}{1-\phi}
ight).$$
 (8)

This means that the characteristic time for the concentration release (modulated by density) is

$$\tau_{\rm eff} = -\frac{1}{\lambda_1} = \tau(1 - \phi) \tag{9}$$

From Equation 13 it follows that

$$c_{\rm in}(t) = c_{\rm eq} + (c_0 - c_{\rm eq})e^{-t/\tau_{\rm eff}},$$
 (10)

$$c_{
m in}(t) = \phi c_0 + (1-\phi) \left[c_{
m ex} + (c_0-c_{
m ex}) \, e^{-rac{t}{ au(1-\phi)}}
ight]$$
 (11)

The deviation of any concentration from equilibrium.

$$\delta(t) = c_{
m in}(t) - c_{
m eq},$$
 (12)

relaxes exponentially:

$$\delta(t) = \delta(0)e^{-t/\tau_{\rm eff}},$$
 (13)

$$\delta(0) = c_0 - c_{\rm eq} \tag{14}$$







Comparison to spore-scale numerical scheme

- » This new analytical formula is the closest mathematical representation so far of the saturation phenomena from the numerical simulations.
- » The deviation could be coming from
 - a numerical error:
 - the representation of the spore as a cube

in the numerical simulation.

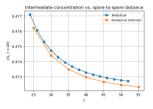


Figure: Centre-to-centre distance vs. concentration at t = 1 h

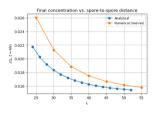


Figure: Centre-to-centre distance vs. concentration at t = 4 h





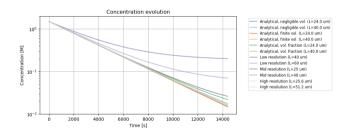




Overall model comparison

» It can be concluded that:

- at low ρ_s, all models align with the analytical simple permeation model;
- the **super-spore scale** model becomes too coarse at high ρ_s , cannot represent saturation well;
- the cell-wall scale model may suffer from rounding errors due to its fine granularity;
- the spore-scale numerical model and the volume-based analytical model exhibit the best alignment, despite coming from different formulations.



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- The goal is to test if interface permeability theory (spore as a "leaky balloon") can be used to estimate meaningful properties of a non-replenishable germination inhibitor released by the spore.
- » The release from a single spore is simulated in isolation using numerical and analytical methods.









Glucose as a reference molecule

- » Glucose has a size very similar to 1-octen-3-ol and has been widely studied.
- » Known diffusivities:
 - the diffusion constant of glucose in water is $D = 6 \times 10^{-6} \text{ cm}^2/\text{s} = 600 \,\mu\text{m}^2/\text{s}$ [4];
 - the diffusion constant of glucose in an agarose film is $D = 6.38 \times 10^{-6} \text{ cm}^2/\text{s} = 638 \,\mu\text{m}^2/\text{s}$ [9];
 - the diffusion constant of glucose in cellulose nanofibrils is $D=1.7\times 10^{-7}~{\rm cm^2/s}=17~\mu{\rm m^2/s}$ [3].



Figure: Glucose









» The diffusion constant of 1-Octen-3-ol in water can be obtained from the Stokes-Finstein relation

$$D = \frac{k_B T}{6\pi \eta a},\tag{15}$$

considering a temperature of T=303.15 K, a thermodynamic constant of $k_BT=4.18\times 10^{-21}$ J and a water viscosity $\eta=0.797$ mPa s.



» The Stokes radius computed from its molecular volume is

$$a(1\text{-octen-3-ol}) = \left(\frac{3}{4\pi}V_M\right)^{1/3} = 0.393 \text{ nm}.$$
 (16)

Figure: 1-octen-3-ol

» Therefore,

$$D(1\text{-octen-3-ol in water}) \approx 6.9016 \times 10^{-6} \, \text{cm}^2/\text{s} = 690.16 \, \mu\text{m}^2/\text{s}.$$
 (17)









» Alternatively, one can use the ratio of the Stokes radii of 1-octen-3-ol (0.393 nm) and glucose (0.36 nm) to scale the *empirically measured* diffusivity of glucose to that of 1-octen-3-ol:

$$D(\text{1-octen-3-ol in water}) = \frac{a(\text{glucose})}{a(\text{1-octen-3-ol})} \times D(\text{glucose in water})$$

$$\approx 6.55 \times 10^{-6} \text{ cm}^2/\text{s} = 655 \ \mu\text{m}^2/\text{s}. \tag{18}$$

-phydydon

Figure: 1-octen-3-ol

» This is slightly lower than the other estimate, but it has to be noted that the viscosity is somewhat sensitive to the temperature assumption.









Diffusion constants

- » This means other diffusion constants can be estimated like this as well!
 - *D*(1-octen-3-ol in agarose) ≈ *D*(1-octen-3-ol in water),
 - $D(1 ext{-octen-3-ol in cellulose}) \approx \frac{a(\text{glucose})}{a(1 ext{-octen-3-ol})} \times D(\text{glucose in cellulose}) = 1.56 \times 10^{-7} \text{ cm}^2/\text{s} = 15.6 \ \mu\text{m}^2/\text{s}.$

Mahahaha.

Other volatile organic compounds (VOCs) like 3-octanone and 3-octanol have a very similar Stokes radius and can be considered to have the same diffusivities. Figure: 3-octanone

- » A potential heat-labile peptide (a = 6 nm as to permeate the cell wall [2]) has
 - $D({
 m heat}$ -labile peptide in water $)pprox 4.54855 imes 10^{-7}~{
 m cm}^2/{
 m s}=45.4855~\mu{
 m m}^2/{
 m s},$
 - $D(\text{heat-labile peptide in cellulose}) \approx \frac{a(\text{glucose})}{a(\text{heat-labile peptide})} \times D(\text{glucose in cellulose}) = 1.02 \times 10^{-8} \, \text{cm}^2/\text{s} = 1.02 \, \mu \text{m}^2/\text{s}.$



Figure: 3-octanol











- » The permeation constant of glucose in synthetic lipid bylayer found in literature is $P_s = 1.4 \times 10^{-10} \text{ cm/s}$ [1] or $P_s = 2.4 \times 10^{-10} \text{ cm/s}$ [5]. A rough average of $P_s = 2.0 \times 10^{-10} \text{ cm/s}$ is taken.
- » In lipid membranes, the permeation constant is linked to the diffusion constant via the Meyer-Overton rule:

$$P_s = \frac{KD}{d},\tag{19}$$

where d = 3.8 nm in liquid-phase lipid bilayers [8].

- » Then,
 - $P_s(1 ext{-octen-3-ol through lipid bilayer}) = \frac{K(1 ext{-octen-3-ol})\cdot D(glucose\ in\ lipid\ bilayer)}{d} imes \frac{a(glucose)}{a(1 ext{-octen-3-ol})} \approx 9.04 imes 10^{-5}\ cm/s = 0.904\ \mu m/s.$









- » In porous barriers, filled with water, that weakly interact with the solute, the partition coefficient can be approximated to $K \approx 1$, indicating a lack of preference for the molecule to reside inside or outside the barrier.
- » Other factors like the porosity of the barrier can be absorbed in an effective permeation constant $P_{\text{eff}} \approx \frac{D}{d}$.
- » Considering a 400 nm thick barrier (upper estimate [7]), we can compute
 - $^ P_{\rm eff}(1\text{-octen-3-ol through agarose}) = \frac{D(1\text{-octen-3-ol in agarose})}{d} \approx \frac{D(1\text{-octen-3-ol in water})}{d} = \frac{690~\mu\text{m}^2/\text{s}}{0.4~\mu\text{m}} = 0.1725~\text{cm/s} = 1725~\mu\text{m/s},$
 - $P_{\rm eff}(\text{1-octen-3-ol through cellulose}) = \frac{\textit{D(1-octen-3-ol in cellulose)}}{\textit{d}} = \frac{15.6 \ \mu\text{m}^2/\text{s}}{0.4 \ \mu\text{m}} = 0.0039 \ \text{cm/s} = 39 \ \mu\text{m/s}.$









- » Proteins usually permeate cell membranes via channels or transporters, so a cell-membrane permeability is hard to define.
- » In polysaccharides,
 - $\begin{array}{l} \ P_{\rm eff}({\rm heat\mbox{-}labile\ peptide\ through\ agarose}) = \frac{D({\rm heat\mbox{-}labile\ peptide\ in\ agarose})}{d} \approx \frac{D({\rm heat\mbox{-}labile\ peptide\ in\ water})}{d} = \\ \frac{45.4855}{0.4\ \mu{\rm m}} = 0.011371375\ {\rm cm/s} = 113.71375\ \mu{\rm m/s}, \end{array}$
 - $P_{\text{eff}}(\text{heat-labile peptide through cellulose}) = \frac{D(\text{heat-labile peptide in cellulose})}{d} = \frac{1.02 \ \mu \text{m}^2/\text{s}}{0.4 \ \mu \text{m}} = 2.55 \times 10^{-4} \text{ cm/s} = 2.55 \ \mu \text{m/s}.$









	1-octen-3-ol	heat-labile peptide
water	$D = 6.9016 \times 10^{-6} \mathrm{cm}^2/\mathrm{s}$	$D = 4.54855 \times 10^{-7} \mathrm{cm}^2/\mathrm{s}$
lipid bilayer membrane	$P_s=2.0 imes10^{-10}\ \mathrm{cm/s}$	N/A
agarose-like cell wall (400 nm)	$P_{ m eff}=0.1725~{ m cm/s}$	$P_{ m eff} = 0.011371375 \ { m cm/s}$
cellulose-like cell wall (400 nm)	$P_{\rm eff}=0.0039~{ m cm/s}$	$P_{\rm eff} = 2.55 \times 10^{-4} { m cm/s}$

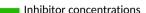
Table: Diffusion coefficients and permeabilities for different substances

» Assumption: there is a layer in the cell wall, with the respective properties, that is diffusion limiting (determines the permeation rate of the entire barrier).









- » It is important to narrow down
 - what is the **initial inhibitor concentration** in the spore,
 - what is the concentration threshold that allows for germination.
- » Studies have shown that, after rupturing the cell wall with a freeze-dry cycle, 10^9 A. flavus conidia in a $300~\mu$ L suspension release not more than $1~\mu$ M 1-octen-3-ol [6]. That equals $c_0 \approx 1.5 \times 10^{-5}$ M per conidium.
- » A protein in moderate abundance is usually found in the micromolar order ($c_0 \approx 10^{-6}$ M)
- » It can be assumed that the inhibitor stops being effective when there are about a hundred molecules left in the spore, at the least. That equals $c_T \approx 2.54 \times 10^{-9}$ M.

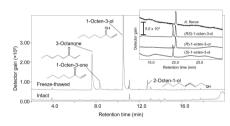


Figure: SPME-GC/MS analysis of volatiles formed from A. flavus conidia [6].





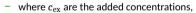




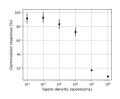
Intermezzo: inferring c_0 from experimental data

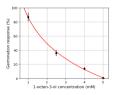
- » Studies with A. nidulans have documented the effect of spore crowding and added 1-octen-3-ol on germination.
- » Assuming that the analytical model is representative, one can equate the two effects on the concentration at the measured time ($t = 10 \, \text{h}$):

$$\phi c_0 + (1 - \phi) \left(c_0 e^{-\frac{t}{\tau(1 - \phi)}} \right) = \phi' c_0 + (1 - \phi') \left[c_{\text{ex}} + (c_0 - c_{\text{ex}}) e^{-\frac{t}{\tau(1 - \phi')}} \right], \tag{20}$$



- $-\phi$ are the volume fractions under varying spore densities,
- $-\phi'$ is the volume fraction when the inhibitor was exogenously added.
- » Solving for c_0 in all ϕ and $c_{\rm ex}$ combinations, we get an average $\langle c_0 \rangle \approx 2.58 \times 10^{-9}$ M (quite low but physically sensible).





Diffusive Signalling in Aspergillus Germination Inhibition

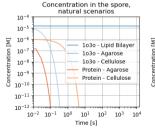


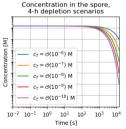


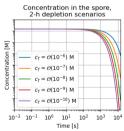




- » Testing
 - $c_0 = 1.5 \times 10^{-5}$ M for 1-octen-3-ol and $c_0 = 10^{-6}$ M for a heat-labile peptide,
 - candidate permeabilities / diffusivities and a range of valid germination thresholds,
- » it becomes evident that the appropriate permeability lies between cellulose and a lipid bilayer.







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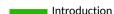
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SPORE CLUSTER EXPERIMENTS: REVISION









- » Previously, it was found that neighbour clusters affect diffusivity, but this effect could not be reproduced across all models.
- » An error was found in the high-resolution solver and was corrected.
- » Hence, the cluster simulations were repeated.











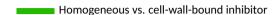


SPORE CLUSTER EXPERIMENTS: REVISION

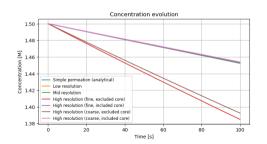


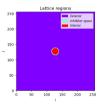


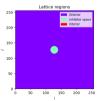




- It was found that, if the inhibitor is only in the cell wall, it is released faster compared to a **homogenous distribution** in the spore volume.
- The reason for this is that there is a higher concentration "pressure" for the same amount of molecules.







SPORE CLUSTER EXPERIMENTS: REVISION









- » The inhibitor release was compared between a single spore and a full cluster of 12+1 spores.
- » The results show that the neighbours in direct contact still obstruct inhibitor release, but to a much lesser degree (7% more residual concentration at $t=4\,\mathrm{h}$).
- » If the inhibitor is only in the cell wall and the spore interiors are inaccessible, the blocking effect is stronger (20% more residual concentration at $t=4\,\mathrm{h}$).

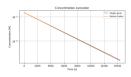


Figure: Empty spore interior.

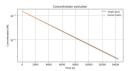


Figure: Full spore interior.



SPORE CLUSTERS VS. RELEASE EXPONENT







Functional relationship

- » The neighbour arrangements produce fluctuating results in the measured release exponent.
- » The results are sensitive to the arrangement of neighbours around the central spore (symmetric vs. asymmetric).
- » Finite lattice size with absorbing boundary possibly distorts the results.
- » A fitted power-law yields a relationship $c(t) \sim e^{\alpha t/\tau}$, where $\alpha \sim M^{-0.004}$, i.e. a very slow decrease with increasing number of neighbours M.
- » In general, ast diffusion through the cluster gaps still depletes the inhibitor locally. Thus, **clustering is not critical in slowing down inhibitor release**.
- » It is to be investigated whether the cluster effect is amplified by the combination of density-driven inhibitor saturation.

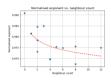


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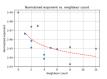


Figure: Full spore interior.

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NEXT STEPS



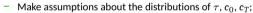




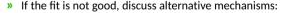
» Turn the **deterministic** model

$$c_{
m in}(t) = \phi c_0 + (1-\phi) \left[c_{
m ex} + (c_0 - c_{
m ex}) \, e^{-rac{t}{ au \, (1-\phi)}}
ight]$$
 (21)

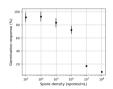
into a **statistical** one to explain experimental data from homogeneous spore cultures.

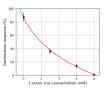


- Express the probability $P(c(t_{\max}) > c_T | \rho_s)$;
- Observe how the probability changes with varying ρ_s ;
- Fit to data.



- Cell wall adsorption of inhibitor;
- Inhibitor interactions.





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