

APMA 2610: Recent Applications of Probability and Statistics

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1.1 Maximum Entropy Principle

A strange though experiment of Gibbs: Imagine a physical system S (say a gas) in an “infinite bath”. Let x be the state of every particle (positions, velocities, ...) in S .

For simplicity, let S be 3 particles in \mathbb{Z}^2 with $x \in \mathbb{Z}^6$ being the positions. Let s be the number of states of particles in S .

What is $p(x)$, the probability that S has state x ?

In the simplest case (each particle is independent and the state distribution is uniform), we trivially have $P(x) = \frac{1}{s}$. But in general, these are incredibly strong assumptions.

We can create some constraints to do better.

1. Assume that the average kinetic energy \mathcal{E} of the infinite heat bath is some constant θ .

In this case, we expect the average kinetic energy of S is approximately θ :

$$\sum_x p(x) \mathcal{E}(x) = \theta$$

2. Trivially, p is a probability distribution, so

$$\sum_x p(x) = 1$$

But still this is far from enough: this gives us only 2 constraints for s many unknowns!

However, we can approximate with the LLN. Sample $n \gg s \gg 1$ iid copies of S , S_1, S_2, \dots, S_n with positions x_1, x_2, \dots, x_n .

Define the **empirical distribution**

$$\hat{p}_x = \frac{\#\{i : X_i = x\}}{n}$$

So with large n , $\hat{p} = p$, and

$$\sum_x \hat{p}(x) \mathcal{E}(x) \approx \theta$$

Claim: The vast majority of assignments of states to X_1, \dots, X_n yield a single empirical distribution \hat{p} .

Consider $C(\hat{p})$, the number of ways to assign a state to each of n systems that would yield \hat{p} . Then, with $\hat{n}_x = \hat{p}_x \cdot n = \#\{i : X_i = x\}$,

$$C(\hat{p}) = \binom{n}{\prod_{i=1}^s \hat{n}_i}$$