

CA Q6 Part 2

Q6 In this question we need to show are told that Part(2) for conductors in the ME's we can set the

(a) $\vec{j}_\perp = \vec{j}_\parallel = \sigma \vec{E}$ and therefore refer to the transverse component as the total component field. we then need to show

$$\vec{k} \cdot \vec{k} = 0, \quad \vec{k} \cdot \vec{H} = 0$$

~~$$\vec{k} \cdot \vec{E} = 0$$~~

$$\vec{k} \times \vec{E} = \omega \mu \vec{H}, \quad \vec{k} \times \vec{H} = \omega \epsilon(\omega) \vec{E}$$

from the longitudinal & transverse field components, we know

$$\nabla \cdot \vec{J} = -\frac{\partial \rho}{\partial t}$$

and from the continuity equation we have that

$$\vec{J} = \sigma \vec{E} \quad (\text{for conductor})$$

$$\nabla \cdot (\sigma \vec{E}) = -\frac{\partial \rho}{\partial t}$$

$$\nabla \cdot (\nabla \vec{E}) = -\frac{\partial \rho}{\partial t}$$

from the first Maxwell equation $\nabla \cdot \vec{E} = \rho / \epsilon_0$

$$\Rightarrow \frac{\partial \rho}{\partial t} = \frac{\sigma \rho}{\epsilon_0} = 0$$

By solving this we get

$$\rho = \rho_0 e^{-(\sigma/\epsilon) t}$$

6(a) now from the algebraic form of the MG's
6(a) we know that $\nabla \cdot E = \rho_f / \epsilon$ likewise, for
conductors $\rho_f = 0$, $\Rightarrow \nabla \cdot E = 0$.

if we now substitute $\nabla \cdot (E_0 e^{j(kx - \omega t)}) = 0$

$$\Rightarrow k \cdot E = 0.$$

Similarly $\nabla \cdot H = 0$
 $\Rightarrow H = H_0 e^{j(kx - \omega t)}$
 $\Rightarrow \boxed{k \cdot H = 0}$

for the 3rd maxwell equation

$$\nabla \times E = -\mu \partial H / \partial t$$

$$\Rightarrow \nabla \times (E_0 e^{j(kx - \omega t)}) = \frac{-\partial}{\partial t} (H_0 e^{j(kx - \omega t)})$$

$$\Rightarrow \boxed{k \times E = \omega \mu H.}$$

lastly from the fourth maxwell equation

$$\nabla \times H = \sigma E$$

$$k \times H = \omega \epsilon (\omega E)$$

$$\Rightarrow \boxed{E(\omega) = \epsilon (1 + i \sigma / \epsilon \omega)}$$

thus we have effectively proven what was required
from the question.

(b) in the second part of this question we know
that the complex and frequency dependant

EM constants. Also propagation constants in conductors are complex because of the conduction field.

$$\nabla^2 \phi - \sigma \mu \frac{\partial \phi}{\partial t} - \mu \epsilon \frac{\partial^2 \phi}{\partial t^2} = 0$$

If we now substitute $\phi = \phi_0 e^{i(kr - \omega t)}$
Then we get

$$k^2 = \epsilon \mu \omega^2 + i \sigma \mu \omega$$

other constants also depend on (k) . Hence it is written as $k^2 = (k'^2 - k''^2) + i(k' + ik'')$
and $k^2 = (k'^2 - k''^2) + i2(k'k'')$

If we now combine

$$(1) k'^2 - k''^2 = \mu \epsilon \omega^2 \quad (2) 2k'k'' = \sigma \mu \omega$$

Then we get

$$k' = \frac{\sqrt{\mu \epsilon} \cdot \omega \left(\sqrt{1 + (\sigma / \epsilon \omega)^2} + 1 \right)^{1/2}}{2}$$

$$k'' = \frac{\sqrt{\mu \epsilon} \cdot \omega \left(\sqrt{1 + (\sigma / \epsilon \omega)^2} - 1 \right)^{1/2}}{2}$$

Similarly we can calculate the others, given the quantities.