

LECTURE 9:

- RATES OF DEFORMATION
- COVARIANT CONTRAVARIANT VECTORS
- PUSH FORWARD PULL BACK
- INTRO TO STRESS

- N -

RATES OF DEFORMATION

$$\underline{L} = \underline{\nabla}_x \underline{v}(x, t) \quad \leftarrow \text{SPATIAL VELOCITY GRADIENT}$$

$$\dot{\underline{F}} = \frac{d}{dt}(\underline{F}) \Big|_{\underline{x}} = \underline{\nabla}_x \left(\frac{d\phi}{dt} \right) = \underline{\nabla}_x \underline{v}(x, t) \quad \leftarrow \text{MAT VEL GRAD}$$

$$\underline{L} = \underline{\nabla}_x \underline{v} = \underline{\nabla}_x \left(\underline{v}(\underline{\phi}^{-1}(x, t), t) \right) = \underline{\nabla}_x \underline{v} \underline{\nabla}_x \underline{\phi}^{-1} = \dot{\underline{F}} \underline{F}^{-1}$$

\underline{L} CAN BE ADDITIVELY DECOMPOSED

$$\underline{L} = \underline{d} + \underline{w}$$

$$\underline{d} = \text{SYM}(\underline{L}) \quad \leftarrow \text{RATE OF DEFORMATION TENSOR}$$

$$\underline{w} = \text{SKEW}(\underline{L}) \quad \leftarrow \text{SPIN (VORTICITY) TENSOR}$$

BOTH \underline{d} & \underline{w} ARE SPATIAL QUANTITIES

IMPORTANT:

$$\dot{J} = \frac{dJ}{dt} \Big|_{\underline{F}} = J \underline{F}^T : \dot{\underline{F}} = J \underline{1} : \dot{\underline{F}} \underline{F}^{-1} = J \underline{1} : \underline{L} = J \underline{\nabla}_x \underline{v}$$

THIS IMPLIES THAT FOR VOLUME PRESERVING DEF

$$J=1 \Rightarrow \dot{J}=0 \Rightarrow \underline{\nabla}_x \underline{v} = 0$$

COVARIANT & CONTRAVARIANT VECTORS

RECALL THAT WE SAID TANGENT VECTORS \underline{A} IN THE REFERENCE CONFIGURATION CAN BE MAPPED INTO THE DEFORMED AS

$$\underline{a} = \underline{F} \underline{A}$$

NOTE: $\underline{F} = \underline{F}(\underline{x}, t) \Rightarrow \underline{a} = \underline{a}(\underline{x}, t)$

ALSO WE SAW (NANSON'S FORMULA) THAT NORMAL VECTORS MAP AS

$$\underline{n} = \underline{F}^{-T} \underline{N}$$

$$(\underline{T} \cdot \underline{N} = \underline{t} \cdot \underline{n} = 0 \Rightarrow (\underline{T} \cdot \underline{N} - \underline{t} \cdot \underline{n}) = 0 \Rightarrow (\underline{T} \cdot \underline{N} - \underline{t} \cdot \underline{n}) = 0$$

$$\Rightarrow (\underline{F}^{-1} \underline{t} \cdot \underline{N} - \underline{t} \cdot \underline{n}) = 0 \Rightarrow \underline{t} \cdot (\underline{F}^{-T} \underline{N} - \underline{n}) = 0 \Rightarrow \underline{F}^{-T} \underline{N} = \underline{n}$$

WHERE $\underline{n}(\underline{x}, t)$

A SPATIAL VECTOR FIELD \underline{a} IS SAID TO CONVECT WITH THE BODY IF THERE EXIST A REFERENCE VECTOR FIELD \underline{A} SUCH THAT

$$\underline{a}(\underline{x}, t) = \underline{F}(\underline{x}, t) \underline{A}(\underline{x})$$

← COVARIANT

OR

$$\underline{a}(\underline{x}, t) = \underline{F}^{-T}(\underline{x}, t) \underline{A}(\underline{x})$$

← CONTRAVARIANT

THUS, IF A VECTOR FIELD IS COVARIANT

$$\frac{d}{dt} \underline{F}^{-1} \underline{a} = \frac{d}{dt} \underline{A} = 0$$

DOES NOT DEPEND ON TIME

$$\Rightarrow \frac{d}{dt} \underline{F}^{-1} \underline{a} = \frac{d}{dt} \underline{F}^{-1} \underline{a} + \underline{F}^{-1} \dot{\underline{a}} \Rightarrow \dot{\underline{a}} = -\underline{F} \frac{d}{dt} \underline{F}^{-1} \underline{a} = \underline{L} \underline{a}$$

$$(\underline{L} = \dot{\underline{F}} \underline{F}^{-1} = \frac{d}{dt} \underline{F} \underline{F}^{-1} - \underline{F} \frac{d}{dt} \underline{F}^{-1} = -\underline{F} \dot{\underline{F}}^{-1})$$

& IF A SPATIAL VECTOR FIELD IS CONTRAVARIANT

$$\frac{d}{dt} \underline{\underline{F}}^T \underline{\underline{a}} = 0 \Rightarrow \underline{\underline{\dot{F}}}^T \underline{\underline{a}} + \underline{\underline{F}}^T \underline{\underline{\dot{a}}} = 0 \Rightarrow \underline{\underline{\dot{a}}} = - \underline{\underline{\dot{F}}}^T \underline{\underline{F}}^T \underline{\underline{a}} = - \underline{\underline{L}}^T \underline{\underline{a}}$$

COVARIANT & CONTRAVARIANT BASIS

LET $\underline{\underline{q}}_i$ BE A SET OF LINEARLY INDEPENDENT (BASIS)

SPATIAL VECTORS THAT CONVECT WITH THE BODY TANGENTIALLY

THEN WE KNOW $\exists \underline{\underline{G}}_i$ SUCH THAT

$$\underline{\underline{q}}_i = \underline{\underline{F}} \underline{\underline{G}}_i \quad \& \quad \underline{\underline{\dot{G}}}_i = 0$$

SINCE $\underline{\underline{F}}$ IS INVERTIBLE $\underline{\underline{G}}_i$ IS ALSO A SET OF BASIS

$$\underline{\underline{\dot{q}}}_i = \underline{\underline{L}} \underline{\underline{q}}_i$$

LET $\underline{\underline{q}}^j$ BE THE DUAL BASIS TO $\underline{\underline{q}}_i$ SUCH THAT $\underline{\underline{q}}^j \cdot \underline{\underline{q}}_i = \delta^j_i$

IT FOLLOWS

$$\frac{d}{dt} \underline{\underline{q}}^j \cdot \underline{\underline{q}}_i = 0 \Rightarrow \underline{\underline{\dot{q}}}^j \cdot \underline{\underline{q}}_i + \underline{\underline{q}}^j \cdot \underline{\underline{\dot{q}}}_i = 0 \Rightarrow \underline{\underline{L}} \underline{\underline{q}}_i \cdot \underline{\underline{q}}^j + \underline{\underline{q}}^j \cdot \underline{\underline{\dot{q}}}_i = 0$$

$$\Rightarrow \underline{\underline{q}}^j \cdot (\underline{\underline{L}}^T \underline{\underline{q}}^i + \underline{\underline{\dot{q}}}^i) = 0 \Rightarrow \underline{\underline{\dot{q}}}^j = - \underline{\underline{L}}^T \underline{\underline{q}}^j$$

THUS, IF $\underline{\underline{q}}^j$ IS A DUAL BASIS TO A TANGENTIALLY CONVECTIVE (COVARIANT) BASIS $\underline{\underline{q}}_i$, THEN $\underline{\underline{q}}^j$ IS CONTRAVARIANT (CONVECS NORMALLY) $\underline{\underline{G}}^j$ & $\exists \underline{\underline{G}}^j$

$$\underline{\underline{G}}^j = \underline{\underline{F}}^T \underline{\underline{q}}^j, \quad \underline{\underline{\dot{G}}}^j = 0$$

WE CAN WRITE ANY TENSOR $\underline{\underline{A}} = A_{ij} \underline{\underline{q}}^i \otimes \underline{\underline{q}}^j = A^{ij} \underline{\underline{q}}_i \otimes \underline{\underline{q}}_j$

WHERE A_{ij} \leftarrow COVARIANT COMPONENTS

A^{ij} \leftarrow CONTRAVARIANT COMPONENTS

COVARIANT & CONTRAVARIANT TENSORS

CONSIDER TWO COVARIANT VECTOR FIELDS $\underline{a} = \underline{F}\underline{A}$, $\underline{b} = \underline{F}\underline{B}$
LET \underline{G} BE A SPATIAL TENSOR

$$\underline{a} \cdot \underline{G} \underline{b} = \underline{F}\underline{A} \cdot \underline{G}\underline{F}\underline{B} = \underline{A} \cdot \underline{F}^T \underline{G} \underline{F} \underline{B}$$

IF

$\underline{F}^T \underline{G} \underline{F} = 0$ THEN \underline{G} IS TERMED A COVARIANT TENSOR

AND

$\underline{F}^T \underline{G} \underline{F}$ IS THE COVARIANT PULL BACK OF \underline{G}

NOTE THE ABOVE IS EQUIVALENT TO $\dot{\underline{G}}_{ij} = 0$

EQUIVALENTLY LET $\underline{a} = \underline{F}^{-T} \underline{A}$, $\underline{b} = \underline{F}^{-T} \underline{B}$ &

$$\underline{a} \cdot \underline{G} \underline{b} = \underline{A} \cdot \underline{F}^{-1} \underline{G} \underline{F}^{-T} \underline{B}$$

IF

$\underline{F}^{-1} \underline{G} \underline{F}^{-T} = 0$ THEN \underline{G} IS TERMED A CONTRAVARIANT TENSOR

AND

$\underline{F}^{-1} \underline{G} \underline{F}^{-T}$ IS THE CONTRAVARIANT PULL BACK OF \underline{G}

NOTE THE ABOVE IS EQUIVALENT TO $\dot{\underline{G}}^{ij} = 0$

NOTE THAT THE COVARIANT PULL BACK (AND PUSH FORWARD) ARE OFTEN DENOTED BY $\phi^*(\cdot)^b$ ($\phi_*(\cdot)^b$)

NOTE THAT THE CONTRAVARIANT PULL BACK (AND PUSH FORWARD) ARE OFTEN DENOTED BY $\phi^*(\cdot)^\#$ ($\phi_*(\cdot)^\#$)

COROTATIONAL VECTOR FIELD

IF A SPATIAL VECTOR FIELD SPINS WITH THE BODY IS CALLED COROTATIONAL. NAMELY

$$\dot{\underline{a}} = \underline{\underline{w}} \underline{a}$$

IF \underline{a} & \underline{b} ARE COROTATIONAL VECTOR FIELDS

$$\frac{d}{dt}(\underline{a} \cdot \underline{b}) = \underline{\dot{a}} \cdot \underline{b} + \underline{a} \cdot \underline{\dot{b}} = \underline{\underline{w}} \underline{a} \cdot \underline{b} + \underline{a} \cdot \underline{\underline{w}} \underline{b} = \underline{b} \cdot \underline{\underline{w}} \underline{a} + \underline{b} \cdot \underline{\underline{w}}^T \underline{a} = \underline{b} \cdot \underline{\underline{w}} \underline{a} - \underline{b} \cdot \underline{\underline{w}} \underline{a} = 0$$

HENCE THE ANGLE BETWEEN THE VECTORS REMAIN UNCHANGED

NOW LET \underline{e}_i BE A SET OF COROTATIONAL ORTHONORMAL BASIS

$$G_{ij} = \underline{e}_i \cdot \underline{\underline{G}} \underline{e}_j$$

THEN

$$\begin{aligned} \dot{G}_{ij} &= \frac{d}{dt}(\underline{e}_i \cdot \underline{\underline{G}} \underline{e}_j) = \underline{\dot{e}_i} \cdot \underline{\underline{G}} \underline{e}_j + \underline{e}_i \cdot \underline{\underline{w}}^T \underline{\underline{G}} \underline{e}_j + \underline{e}_i \cdot \underline{\underline{G}} \underline{\dot{e}_j} = \\ &= \underline{e}_i \cdot (\underline{\dot{\underline{G}}} - \underline{\underline{w}} \underline{\underline{G}} + \underline{\underline{G}} \underline{\underline{w}}) \underline{e}_j \end{aligned}$$

IF $\dot{G}_{ij} = 0$ OR $\underline{\dot{\underline{G}}} - \underline{\underline{w}} \underline{\underline{G}} + \underline{\underline{G}} \underline{\underline{w}} = 0$ THEN $\underline{\underline{G}}$ IS A COROTATIONAL TENSOR

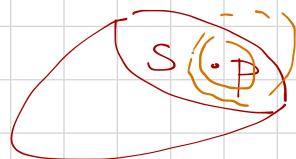
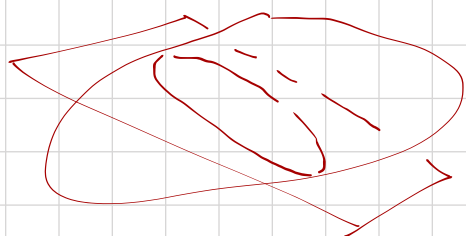
EFFECTIVELY THE TENSOR ROTATES WITH THE BODY WHICH IS IMPORTANT NOTION FOR INVARIANCE

STRESS

SHORT RANGE FORCES \rightarrow ATOMIC INTERACTIONS

LONG RANGE FORCES \rightarrow GRAVITATIONAL FORCES ETC.

SHORT RANGE FORCES ARE TRANSMITTED THROUGH THE BODY



THE TRACTION AT A POINT P IN OUR CONTINUUM

$$\underline{t}_s(P) = \lim_{t \rightarrow 0} \frac{\Delta \underline{F}}{|S \cap B_t(P)|} = \lim_{t \rightarrow 0} \frac{\Delta \underline{F}}{\Delta S}$$

EFFECTIVELY A TRACTION IS AN AVERAGE OF FORCES ACTING ON AN INFINITESIMAL AREA ELEMENT

$$d\underline{f} = \underline{t} ds = \underline{T} ds$$

FOLLOWING CAUCHY'S PRINCIPLE $\underline{t} = \underline{t}(\underline{x}, \underline{n})$

IE. THE VALUE OF THE TRACTION VECTOR IS INHERENTLY LOCAL & DEPEND ON POSITION & THE SURFACE NORMAL

LASTLY NOTE THAT ON OPPOSITE FACES OF THE CUT

$$\Delta \underline{F}(\underline{n}) = -\Delta \underline{F}(-\underline{n}) \Rightarrow \underline{t}(\underline{x}, \underline{n}) = -\underline{t}(\underline{x}, -\underline{n})$$

(NOTE WE OMIT DEPENDENCE ON TIME FOR SIMPLICITY)

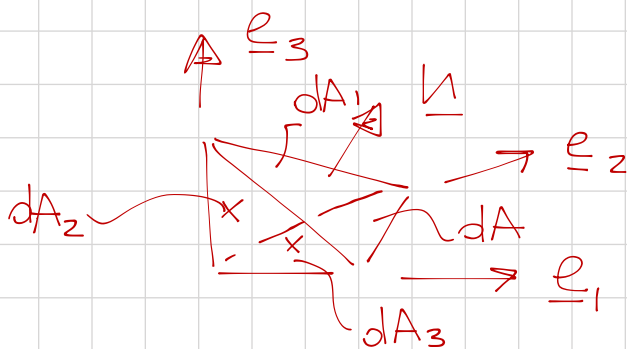
CAUCHY'S TETRAHEDRON THEOREM

$\underline{t}(\underline{x}, \underline{n})$ IS LINEAR IN \underline{n}

BY DEFINITION $\underline{t}_x(\)$ IS A TENSOR

THEREFORE CAUCHY'S TETRAHEDRON THEOREM STATES THAT $\exists \underline{\nabla} \in \mathbb{R}^{d \times d}$ SUCH THAT

$$\underline{t} = \underline{\nabla} \underline{n}$$



$$\underline{t}(-\underline{e}_3) dA_3 + \underline{t}(-\underline{e}_1) dA_1 + \underline{t}(-\underline{e}_2) dA_2 + \underline{t}(\underline{n}) dA = 0$$

$$dA_i = \underline{n} \cdot \underline{e}_i dA$$

$$\underline{t}(-\underline{e}_3) \underline{n} \cdot \underline{e}_3 + \underline{t}(-\underline{e}_2) \underline{n} \cdot \underline{e}_2 + \underline{t}(-\underline{e}_1) \underline{n} \cdot \underline{e}_1 + \underline{t}(\underline{n}) = 0$$

$$\text{WITH } \underline{t}(-\underline{n}) = -\underline{t}(\underline{n})$$

$$\underline{t}(\underline{n}) = \underline{t}(\underline{e}_1) \overbrace{\underline{e}_1 \cdot \underline{n}}^{n_1} + \underline{t}(\underline{e}_2) \overbrace{\underline{e}_2 \cdot \underline{n}}^{n_2} + \underline{t}(\underline{e}_3) \overbrace{\underline{e}_3 \cdot \underline{n}}^{n_3} = \underline{t}(n_1 \underline{e}_1 + n_2 \underline{e}_2 + n_3 \underline{e}_3)$$

\underline{t} IS LINEAR IN \underline{n} !

SYMMETRY OF CAUCHY'S STRESS TENSOR

