

THE SPACE OF STATES AKA THE HILBERT SPACE (1)

$|4\rangle$ LIVES THERE IN THE "SPACE OF STATES"

↑
KET

WHATEVER THIS SPACE OF STATES IS

THERE'S THE IDEA OF CALCULATING
AN "OVERLAP" BETWEEN STATES

$(|\phi\rangle, |4\rangle)$

↑
(IDEA OF)
AN INNER PRODUCT

FAMILIAR INNER PRODUCT; IN A VECTOR SPACE

VECTORS \vec{a}, \vec{b}

INNER PRODUCT = $\vec{a} \cdot \vec{b}$ "THE DOT PRODUCT"

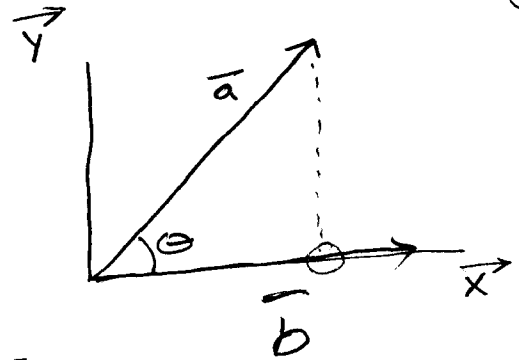
WHAT IS IT?

WHAT IS THIS INNER PRODUCT

(2)

$\vec{a} \cdot \vec{b} =$ A REAL NUMBER
SCALAR
 ≥ 0

$$\vec{a} \cdot \vec{b} = |\vec{a}| |\vec{b}| \cos \theta$$



= LENGTH OF \vec{a} * LENGTH OF \vec{b} * PORTION PROJECTED

$(|\phi\rangle, |\psi\rangle)$ CORRESPONDS TO

$$\langle \phi | \psi \rangle$$

"THE AMOUNT TO WHICH
 $|\phi\rangle$ AND $|\psi\rangle$ OVERLAP"

$$(|\phi\rangle, |\psi\rangle) = \langle \phi | \psi \rangle$$

$|\psi\rangle \in \mathcal{E}$ CORRESPONDS TO $\psi(x) \in \mathcal{H}$
 \uparrow HILBERT SPACE \uparrow SPACE OF FUNCTIONS

$\psi(x)$ IN THE CONTINUOUS x BASIS

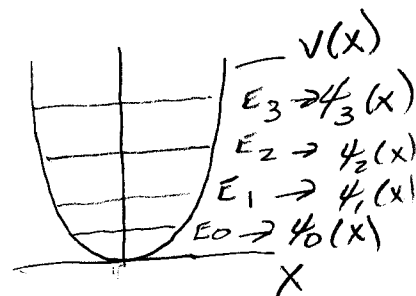
$$\langle \phi | \psi \rangle \triangleq \int_{-\infty}^{\infty} \phi^*(x) \psi(x) dx$$

$$\begin{aligned}\langle \phi | \psi \rangle &= \int_{-\infty}^{\infty} \phi^*(x) \psi(x) dx \\ &= \int_{-\infty}^{\infty} \phi^*(p) \psi(p) dp\end{aligned}$$

SUPPOSE

$|\psi\rangle$ LIVED IN A DISCRETE SET OF STATES

$$|\psi\rangle = \sum_{\text{ALL } n} a_n |E_n\rangle$$



$$|\phi\rangle = \sum_{\text{ALL } m} b_m |E_m\rangle$$

$$\langle \phi | \psi \rangle = \sum_m b_m^* \langle E_m | \sum_n a_n |E_n\rangle$$

BUT THE ENERGY STATES ARE EIGENSTATES
AND ORTHOGONAL SOM

$$\begin{aligned}\langle E_m | E_n \rangle &= 0 \quad \text{IF } m \neq n \\ &= 1 \quad \text{IF } m = n\end{aligned} = \delta_{mn}$$

SO

$$\langle \phi | \psi \rangle = \sum_m b_m^* a_m = \text{A COMPLEX NUMBER}$$

WITHOUT HAVING TO CALCULATE

$$\langle \phi | \psi \rangle \rightarrow \int \left(\left(\frac{\pi \omega}{\pi \hbar} \right)^{1/4} e^{-\frac{m\omega^2 x^2}{2\hbar}} \right) \left(x \left(\frac{m\omega}{\pi \hbar} \right)^{1/4} \frac{1}{\sqrt{2}} e^{-\frac{m\omega}{2\hbar} x^2} \right) dx$$

$$|4\rangle$$

↑
KET

$$\langle \phi |$$

↑
BRA

$$\langle \phi | 4 \rangle$$

└──────────┘
BRA (C) KET

WHAT IS THIS

$$\langle \phi | \quad ?$$

DOES IT LIVE IN THE HILBERT SPACE? \mathcal{P}

NO, IT HANGS OUT IN SOMETHING
CALLED THE DUAL SPACE \mathcal{E}^*
("A SLIGHTLY LARGER NEIGHBORHOOD")

BUT WHAT IS $\langle \phi |$?

$\langle \phi |$ IS A LINEAR FUNCTIONAL ON $|\psi\rangle$

(4)

↑
SECTION 1.3 IN Q.V.T. FOR GIFTED AMATEUR

A FUNCTION $f(x) = x^2$ TAKES IN
A NUMBER (x) AND RETURNS A NUMBER

A FUNCTIONAL TAKES IN A FUNCTION
 $\psi(x) = \left(\frac{m\omega}{\pi\hbar} \right)^{1/4} e^{-\frac{m\omega}{2\hbar} x^2}$ FOR EXAMPLE

AND RETURNS A NUMBER (MAY BE COMPLEX)

2nd POSTULATE WANT TO MEASURE λ (OBSERVABLE QUANTITY) (5)
 IS DESCRIBED BY $A|\psi\rangle$ ← KET IN \mathcal{E}
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 OPERATOR = OBSERVABLE
 BUT TO BE AN OBSERVABLE A MUST BE HERMITIAN

REVIEW

CONJUGATE TRANSPOSE = $^+$ (DAGGER)

$$7^+ = 7$$

$$(a+ib)^+ = a-ib$$

$$\hat{M} = \begin{pmatrix} \frac{\hbar\omega}{2} & i \\ -i & -\frac{\hbar\omega}{2} \end{pmatrix}$$

$$\hat{M}^+ = \left(\begin{array}{c} \nearrow \end{array} \right)^* = \hat{M}^T^*$$

$$= \begin{pmatrix} \frac{\hbar\omega}{2} & -i \\ i & -\frac{\hbar\omega}{2} \end{pmatrix}^* = \begin{pmatrix} \frac{\hbar\omega}{2} & i \\ -i & -\frac{\hbar\omega}{2} \end{pmatrix}$$

$$(|\psi\rangle)^+ \rightarrow \langle\psi|$$

$^+$ = CONJUGATE TRANSPOSE = ADJOINT

PROPERTIES

6

$$(\hat{M}^\dagger)^\dagger = \hat{M}$$

$$(\hat{M}\hat{N})^\dagger = \hat{N}^\dagger\hat{M}^\dagger$$

OPERATOR \hat{M} IS SAID TO BE HERMITIAN

$$\hat{M}^\dagger = \hat{M}$$

IN PARTICULAR
WE'LL LOOK FOR

$$\hat{H}^\dagger = \hat{H}$$

THE OPERATOR = IT'S ADJOINT
(MATRIX)

$\hat{O}|4\rangle$ CHANGES STATE
COULD EQUIVALENTLY WRITE

$$\underset{\substack{\nwarrow \\ \text{OLD} \\ \text{STATE}}}{\hat{O}|4\rangle} = |\underset{\substack{\nwarrow \\ \text{NEW STATE}}}{\hat{O}}\rangle$$

$$(|\hat{O}4\rangle)^\dagger = \boxed{\langle\hat{O}4| \triangleq \langle 4|\hat{O}^\dagger}$$

TO MOVE THE \hat{O} , MUST TAKE THE ADJOINT

$$\langle \hat{O}\psi | \hat{O}\psi \rangle = \text{A NUMBER} > 0$$

$$\langle \hat{O}\psi | \hat{O}\psi \rangle$$

$$\langle \psi | \underbrace{\hat{O}^+ \hat{O}}_{\hat{H}} | \psi \rangle$$

SO IF WE CAN WRITE

$$\boxed{\hat{H} = \hat{O}^+ \hat{O}} = \text{A HERMITIAN OPERATOR}$$

BECAUSE

$$(\hat{M}\hat{N})^+ = \hat{N}^+ \hat{M}^+$$

$$\hat{M} = \hat{O}^+ \quad \hat{N} = \hat{O}$$

$$(\hat{O}^+ \hat{O})^+ = \hat{O}^+ (\hat{O}^+)^+$$

$$\hat{H}^+$$

$$\hat{H}$$

$$\boxed{\hat{O}^+ \hat{O}} = \hat{H}$$