

THIS BRINGS US TO THE S.H.O.

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$$\hat{H} = \frac{\hat{p}^2}{2m} + \frac{1}{2} m \omega^2 \hat{x}^2$$

WANT TO WRITE AS A PRODUCT OF 2 OPERATORS

IF WE HAVE SOMETHING OF THE FORM

$$\begin{aligned} a^2 + b^2 &= (a - ib)(a + ib) \\ &= a^2 - \underbrace{iba + iab}_{\text{SCALARS}} + b^2 \end{aligned}$$

SO TRY IT

$$\hat{H} = \frac{\hat{p}^2}{2m} + \frac{1}{2} m \omega^2 \hat{x}^2 = \frac{m \omega^2}{2} \left( \frac{\hat{p}^2}{m^2 \omega^2} + \hat{x}^2 \right) \quad (1)$$

LET'S CALCULATE THE PRODUCT

$$\begin{aligned} \frac{m \omega^2}{2} \left( \hat{x} - \frac{i \hat{p}}{m \omega} \right) \left( \hat{x} + \frac{i \hat{p}}{m \omega} \right) &= \frac{m \omega^2}{2} \left( \hat{x}^2 + \frac{i}{m \omega} (\hat{x} \hat{p} - \hat{p} \hat{x}) + \frac{\hat{p}^2}{m^2 \omega^2} \right) \\ &= \frac{m \omega^2}{2} \left( \frac{\hat{p}^2}{m^2 \omega^2} + \hat{x}^2 \right) + \frac{m \omega^2 i}{2 m \omega} [\hat{x}, \hat{p}] \end{aligned}$$

AN EXTRA TERM

$$= \hat{H} + i \omega \left( \frac{i \hbar}{2} \right)$$

$$\hat{H} = \frac{m \omega^2}{2} \left( \hat{x} - \frac{i \hat{p}}{m \omega} \right) \left( \hat{x} + \frac{i \hat{p}}{m \omega} \right) + \frac{\hbar \omega}{2}$$

WANTED

$$\hat{H} = \hat{O}^\dagger \hat{O}$$

HAVE

$$\hat{H} = \frac{m\omega^2}{2} \left( X - \frac{iP}{m\omega} \right) \left( X + \frac{iP}{m\omega} \right) + \frac{\hbar\omega}{2}$$

X IS HERMITIAN

P IS HERMITIAN ... NOT AS OBVIOUS

TO SHOW COMPARE  $\langle \phi | P | \psi \rangle^\dagger$  TO  $\langle \phi | P | \psi \rangle$   
 HINT INTEGRATE BY PARTS **PROBLEM #1**

IF WE SAY

$$\hat{O} = X + \frac{iP}{m\omega}$$

$$\hat{O}^\dagger = X^\dagger + \left( \frac{iP}{m\omega} \right)^\dagger = X - \frac{iP}{m\omega}$$

$$\hat{H} = \frac{m\omega^2}{2} (\hat{O}^\dagger \hat{O}) + \frac{\hbar\omega}{2}$$

THIS EXTRA CONSTANT IS OK

$$\langle \psi | \frac{\hbar\omega}{2} | \psi \rangle = \frac{\hbar\omega}{2}$$

RECOGNIZE AS THE GROUND STATE ENERGY

AS LAST TIME WANT TO MOVE TO A  
DIMENSIONLESS  $H$  (ENERGY)

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TO DO THIS LET'S CALCULATE A COMMUTATOR

COMMUTATORS ARE USEFUL AS WE SHALL SEE

$$[O, O^\dagger] = \left(X + \frac{iP}{m\omega}\right) \left(X - \frac{iP}{m\omega}\right) - \left(X - \frac{iP}{m\omega}\right) \left(X + \frac{iP}{m\omega}\right)$$

$X^2, P^2$  TERMS VANISH

$$= \frac{i}{m\omega} [X, P] - \frac{i}{m\omega} [X, P]$$

$$= -\frac{2i}{m\omega} [X, P] = -2i \frac{i\hbar}{m\omega}$$

$$[O, O^\dagger] = \frac{2\hbar}{m\omega}$$

SO IF WE DEFINE

$$\hat{a} \triangleq \sqrt{\frac{m\omega}{2\hbar}} \hat{O}$$

$$\Rightarrow [\hat{a}, \hat{a}^\dagger] = [a, a^\dagger] = 1$$

$$\begin{aligned}\hat{a} &= \sqrt{\frac{m\omega}{2\hbar}} \left( \hat{X} + \frac{i\hat{P}}{m\omega} \right) \\ \hat{a}^\dagger &= \sqrt{\frac{m\omega}{2\hbar}} \left( \hat{X} - \frac{i\hat{P}}{m\omega} \right)\end{aligned}$$

$\hat{X}$   
+  
 $\hat{P}$  IN TERMS OF  $\hat{a}, \hat{a}^\dagger$

SHOW BOTH ARE HERMITIAN

**PROBLEM #2**

$$\hat{H} = \frac{m\omega^2}{2} (\hat{O}^\dagger \hat{O}) + \frac{\hbar\omega}{2}$$

$$\begin{aligned}\hat{H} &= \frac{m\omega^2}{2} \left( \sqrt{\frac{2\hbar}{m\omega}} a^\dagger \sqrt{\frac{2\hbar}{m\omega}} a \right) + \frac{\hbar\omega}{2} \\ &= \frac{m\omega^2 (2\hbar)}{2m\omega} (a^\dagger a) + \frac{\hbar\omega}{2}\end{aligned}$$

$$\boxed{\hat{H} = \hbar\omega \left( \hat{a}^\dagger \hat{a} + \frac{1}{2} \right)}$$

LET'S LOOK AT THE ENERGY OF A STATE

$$\begin{aligned}\langle \hat{H} \rangle_\psi &= \langle \psi | \hat{H} | \psi \rangle = \langle \psi | \hbar\omega \left( \hat{a}^\dagger \hat{a} + \frac{1}{2} \right) | \psi \rangle \\ &= \hbar\omega \left( \langle \psi | \hat{a}^\dagger \hat{a} | \psi \rangle + \frac{1}{2} \langle \psi | \psi \rangle \right) \\ &= \hbar\omega \left( \langle \psi | \hat{a}^\dagger \hat{a} | \psi \rangle + \frac{\hbar\omega}{2} \right) \\ &= \hbar\omega \left( \langle \psi | \hat{a}^\dagger | \hat{a} | \psi \rangle + \frac{\hbar\omega}{2} \right) \\ &= \hbar\omega \left( \langle \hat{a} \psi | \hat{a} \psi \rangle + \frac{\hbar\omega}{2} \right) \\ &= \hbar\omega \left( \langle \phi | \phi \rangle + \frac{\hbar\omega}{2} \right)\end{aligned}$$

$$\boxed{\langle \hat{H} \rangle_\psi \geq \frac{\hbar\omega}{2}}$$

$$\begin{aligned}&\int_{-\infty}^{\infty} \phi^*(x) \phi(x) dx \\ &\geq 0\end{aligned}$$

SO THINKING ABOUT ENERGY EIGENSTATES

ALL

$$\langle H \rangle_{\psi} \geq \frac{\hbar \omega}{2}$$

WHAT'S THE GROUND STATE

$$\begin{aligned} \langle H \rangle_{\psi_0} &= \hbar \omega \langle \psi_0 | a^\dagger a | \psi_0 \rangle + \frac{\hbar \omega}{2} \\ &= \hbar \omega \langle a \psi_0 | a \psi_0 \rangle + \frac{\hbar \omega}{2} \end{aligned}$$

$$|\phi\rangle = a|\psi_0\rangle$$

$$\langle \phi | \phi \rangle = \int \phi^*(x) \phi(x) dx \geq 0$$

$$\text{ONLY } = 0 \quad \text{IF } \phi(x) = 0$$

$$\Rightarrow a|\psi_0\rangle = 0$$

$$\left( \hat{x} + i \frac{\hat{p}}{m\omega} \right) \psi_0(x) = 0$$

$$\left( \hat{x} + \frac{i}{m\omega} \left( -i\hbar \frac{\partial}{\partial x} \right) \right) \psi_0(x) = 0$$

$$x\psi_0(x) + \frac{\hbar}{m\omega} \frac{d\psi_0(x)}{dx} = 0$$

$$\frac{d\psi_0(x)}{dx} = -\frac{m\omega x}{\hbar} \psi_0(x)$$

1ST ORDER O.D.E.

$$\Rightarrow \boxed{\psi_0(x) = C_0 e^{-\frac{m\omega}{2\hbar} x^2}}$$

$$C_0 = \left( \frac{m\omega}{\pi\hbar} \right)^{1/4}$$

$$a |\psi_0\rangle = 0$$

↑

ANNIHILATION  
OR  
DESTRUCTION  
OR  
LOWERING  
OPERATOR

$a^+$  = CREATION OPERATOR  
OR  
RAISING

WHY ?

REMEMBER

$$[a, a^+] = 1 = aa^+ - a^+a$$

$$(aa^+ - a^+a) |\psi_0\rangle = 1 |\psi_0\rangle$$

$$aa^+ |\psi_0\rangle = |\psi_0\rangle \Rightarrow a^+ |\psi_0\rangle$$

MUST NOT  
BE THE GROUND  
STATE

$$\hat{H} = \hbar\omega \left( \hat{a} \hat{a}^+ + \frac{1}{2} \right)$$

$$= \hbar\omega \left( \hat{N} + \frac{1}{2} \right)$$

↑ THE NUMBER OPERATOR

RECALL

$$E_n = \hbar\omega \left( n + \frac{1}{2} \right) \quad \text{FROM SCHRÖDINGER'S EQUATION}$$

$$\boxed{\hat{N} = \hat{a} \hat{a}^+}$$

$$[\hat{N}, \hat{a}] = [a^\dagger a, a]$$

$$= a^\dagger a a - \underbrace{a a^\dagger a}$$

$$= (a^\dagger a - a a^\dagger) a$$

$$= [a^\dagger, a] a$$

$$\boxed{[\hat{N}, \hat{a}] = -\hat{a}}$$

$$[\hat{N}, \hat{a}^\dagger] = [a^\dagger a, a^\dagger]$$

$$= a^\dagger [a, a^\dagger]$$

$$\boxed{[\hat{N}, \hat{a}^\dagger] = a^\dagger}$$

$$[\hat{a}, a^{+K}] = a a^{+K} - a^{+K} a$$

MOVE THIS  $a$  ACROSS  $a^{+K}$

REMEMBER  $[a, a^\dagger] = 1$   $aa^\dagger - a^\dagger a = 1$

$$\underline{aa^\dagger} = 1 + a^\dagger a$$

$$a a^{+K} = (1 + a^\dagger a) a^{+K-1}$$

$$= (1 + a^\dagger a) a^{+K-1} - a^{+K} a$$

$$= a^{+K-1} + a^\dagger (1 + a^\dagger a) a^{+K-2} - a^{+K} a$$

$$= 2 a^{+K-1} + a^{\dagger 2} a a^{+K-2} - a^{+K} a$$

$\vdots$

$$\boxed{[\hat{a}, a^{+K}] = K a^{+K-1}}$$

OTHERS

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$$[\hat{a}^{\dagger}, \hat{a}^k] = ?$$

$$[\hat{N}, \hat{a}^k] = -k \hat{a}^k$$

$$[\hat{N}, \hat{a}^{\dagger k}] = k \hat{a}^{\dagger k}$$

$N$  GIVES THE NUMBER OF CREATION OR DESTRUCTION OPERATORS

$$|\psi_1'\rangle \triangleq \hat{a}^{\dagger} |\psi_0\rangle$$

IF  $|\psi\rangle$  IS AN EIGENSTATE OF  $\hat{N}$  AN EIGENSTATE OF  $\hat{H}$

$$\hat{N} |\psi_1'\rangle = \hat{N} \hat{a}^{\dagger} |\psi_0\rangle$$

$$= [\hat{a}^{\dagger} \hat{a} \hat{a}^{\dagger} - \hat{a}^{\dagger} \hat{a}^{\dagger} \hat{a}] |\psi_0\rangle$$

$$= [\hat{N}, \hat{a}^{\dagger}] |\psi_0\rangle$$

$$\hat{N} |\psi_1'\rangle = \hat{a}^{\dagger} |\psi_0\rangle = (1) |\psi_1'\rangle$$

NORMALIZED?

$$\langle \psi_1' | \psi_1' \rangle = \langle \hat{a}^{\dagger} \psi_0 | \hat{a}^{\dagger} \psi_0 \rangle$$

$$= \langle \psi_0 | \hat{a} \hat{a}^{\dagger} | \psi_0 \rangle$$

$$= \langle \psi_0 | [\hat{a}, \hat{a}^{\dagger}] | \psi_0 \rangle = \langle \psi_0 | \psi_0 \rangle$$

$$\Rightarrow |\psi_1\rangle = |\psi_1'\rangle$$



$$| \psi_2' \rangle = a^\dagger a^\dagger | \psi_0 \rangle$$

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$$N | \psi_2' \rangle = \hat{N} a^{\dagger 2} | \psi_0 \rangle$$

↓

$$N | \psi_2' \rangle = [\hat{N}, a^{\dagger 2}] | \psi_0 \rangle$$

$$= 2 a^{\dagger 2} | \psi_0 \rangle$$

$$= 2 | \psi_2' \rangle$$

IS

$| \psi_2' \rangle$  NORMALIZED

$$\langle \psi_2' | \psi_2' \rangle = \langle a^\dagger a^\dagger \psi_0 | a^\dagger a^\dagger \psi_0 \rangle$$

$$= \langle \psi_0 | \underbrace{a a a^\dagger a^\dagger}_{K a^{K-1}} | \psi_0 \rangle$$

$$= \langle \psi_0 | a 2 a^\dagger | \psi_0 \rangle$$

$$\langle \psi_2' | \psi_2' \rangle = 2 \langle \psi_0 | [a, a^\dagger] | \psi_0 \rangle = 2$$

$$\Rightarrow | \psi_2 \rangle = \frac{1}{\sqrt{2}} a^{\dagger 2} | \psi_0 \rangle$$

$$\Rightarrow \boxed{ | \psi_n \rangle = \frac{1}{\sqrt{n!}} a^{\dagger n} | \psi_0 \rangle } \quad 2.34$$