Lesson 9

Euler-Bernoulli Beam

Minh-Chien Trinh, PhD

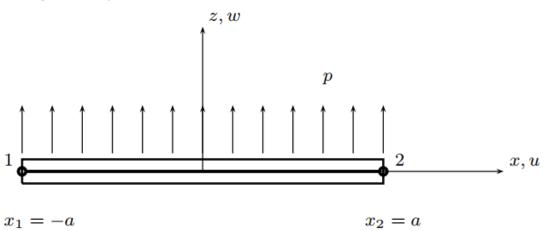
Division of Mechanical System Engineering Jeonbuk National University

Content

- 1. Bernoulli Beam
 - Formulation
 - Bernoulli Beam Problem
 - Bernoulli Beam Free Vibrations
- 2. Bernoulli 2D Frames
 - Formulation
 - 2D Frame Problem Python
 - 2D Frame in Free Vibrations Python
- 3. Bernoulli 3D Frames
 - Formulation
 - 3D Frame Problem Python
 - 3D Frame in Free Vibrations Python



- ❖ A classical beam theory where the transverse shear deformation is neglected.
- \bullet The beam is defined in the x-z plane, with constant cross-section area A.
- \diamond The transverse deflection w is the only degree of freedom of the model.
- The plane sections remain plane and normal to the axis of the beam after deformation.
- ❖ The in-plane rotation given by the derivative of the transverse deflection w.r.t the beam axis.



 \diamond The displacement u, at a distance z of the beam middle axis is given by

$$u = -z \frac{\partial w}{\partial x}$$

Strains are defined as $\varepsilon_x = \frac{\partial u}{\partial x} = -z \frac{\partial^2 w}{\partial x^2}; \quad \gamma_{xz} = \frac{\partial u}{\partial z} + \frac{\partial w}{\partial x} = 0$



The strain energy is

$$U = \frac{1}{2} \int_{V} \sigma_{x} \varepsilon_{x} dV = \frac{1}{2} \int_{V} E \varepsilon_{x}^{2} dV = \frac{1}{2} \int_{A-a}^{a} E \left(\frac{\partial^{2} w}{\partial x^{2}} \right)^{2} z^{2} dA dx = \frac{1}{2} \int_{-a}^{a} E I_{y} \left(\frac{\partial^{2} w}{\partial x^{2}} \right)^{2} dx$$

with I_v is the second moment of area of the beam cross-section.

The kinetic energy is

$$K = \frac{1}{2} \int_{V} (\rho \dot{u}^{2} + \rho \dot{w}^{2}) dV = \frac{1}{2} \int_{-a}^{a} \left(\rho I_{y} \left(\frac{\partial \dot{w}}{\partial x} \right)^{2} + \rho A \dot{w}^{2} \right) dx$$

The first term is the rotary inertia, can be neglected for thin beams. The second term is the vertical bulk inertia.

The external work due to the transverse pressure p and the axial load N_0 (that accounts for nonlinear Von Karman strains) is

$$\delta W = \int_{-a}^{a} p \, \delta w \, dx - \int_{-a}^{a} N_0 \, \frac{\partial w}{\partial x} \, \frac{\partial \delta w}{\partial x} \, dx$$

At each node we consider 2 DOFs: w and $\frac{\partial w}{\partial x}$, respectively are the transverse displacement and rotation of the cross section.

$$\mathbf{w}^{eT} = \begin{bmatrix} w_1 & \frac{\partial w_1}{\partial x} & w_2 & \frac{\partial w_2}{\partial x} \end{bmatrix}$$



The transverse displacement is interpolated using Hermite shape functions as

$$w = \mathbf{N}(\xi)\mathbf{w}^e$$

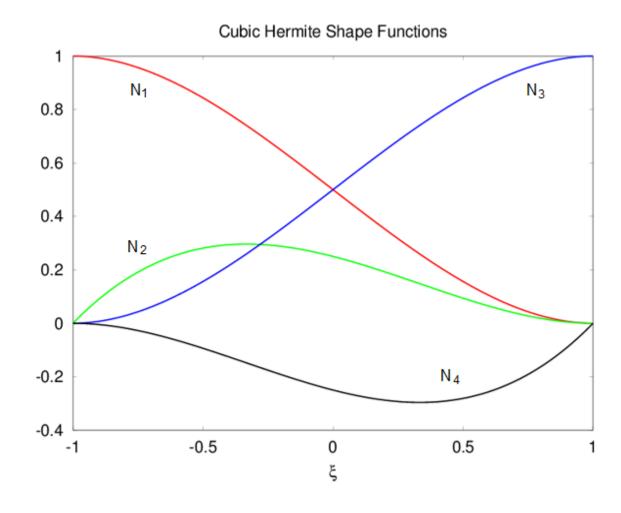
The shape functions are

$$N_{1}(\xi) = \frac{1}{4}(2 - 3\xi + \xi^{3})$$

$$N_{2}(\xi) = \frac{a}{4}(1 - \xi - \xi^{2} + \xi^{3})$$

$$N_{3}(\xi) = \frac{1}{4}(2 + 3\xi - \xi^{3})$$

$$N_{4}(\xi) = \frac{a}{4}(-1 - \xi + \xi^{2} + \xi^{3})$$



❖ The transverse displacement is interpolated using Hermite shape functions as

$$w = \mathbf{N}(\xi)\mathbf{w}^e$$

The shape functions are defined as

$$N_{1}(\xi) = \frac{1}{4}(2 - 3\xi + \xi^{3}); \quad N_{2}(\xi) = \frac{a}{4}(1 - \xi - \xi^{2} + \xi^{3})$$

$$N_{3}(\xi) = \frac{1}{4}(2 + 3\xi - \xi^{3}); \quad N_{4}(\xi) = \frac{a}{4}(-1 - \xi + \xi^{2} + \xi^{3})$$

with $\xi = x/a$ identifies the dimensionless axial coordinate, $\xi \in [-1, 1]$.

thus
$$\frac{\partial}{\partial x} = \frac{\partial}{\partial \xi} \frac{\partial \xi}{\partial x} = \frac{1}{a} \frac{\partial}{\partial \xi};$$
 $\frac{\partial}{\partial x} \left(\frac{\partial}{\partial x} \right) = \frac{1}{a} \frac{\partial}{\partial \xi} \left(\frac{1}{a} \frac{\partial}{\partial \xi} \right) = \frac{1}{a^2} \frac{\partial^2}{\partial \xi^2}$

The strain energy is obtained as

$$U = \frac{1}{2} \int_{-a}^{a} EI_{y} \left(\frac{\partial^{2} w}{\partial x^{2}} \right)^{2} dx = \frac{1}{2} \int_{-1}^{1} \frac{EI_{y}}{a^{4}} \left(\frac{\partial^{2} w}{\partial \xi^{2}} \right)^{2} a d\xi = \frac{1}{2} \mathbf{w}^{eT} \frac{EI_{y}}{a^{3}} \int_{-1}^{1} \mathbf{N}'' \mathbf{N}'' d\xi \mathbf{w}^{e} \quad \text{with} \quad \mathbf{N}'' = \frac{\partial^{2} \mathbf{N}}{\partial \xi^{2}}$$

The element stiffness matrix is obtained as

$$\mathbf{K}^{e} = \frac{EI_{y}}{a^{3}} \int_{-1}^{1} \mathbf{N}''^{T} \mathbf{N}'' d\xi = \frac{EI_{y}}{2a^{3}} \begin{vmatrix} 3 & 3a & -3 & 3a \\ 3a & 4a^{2} & -3a & 2a^{2} \\ -3 & -3a & 3 & -3a \\ 3a & 2a^{2} & -3a & 4a^{2} \end{vmatrix}$$



The kinetic energy takes the form

$$K = \frac{1}{2} \int_{-a}^{a} \rho A \dot{w}^2 dx = \frac{1}{2} \int_{-1}^{1} \rho A a \dot{w}^2 d\xi = \frac{1}{2} \dot{\mathbf{w}}^{eT} \int_{-1}^{1} \rho A a \mathbf{N}^T \mathbf{N} d\xi \dot{\mathbf{w}}^e$$

The mass matrix is

$$\mathbf{M}^{e} = \int_{-1}^{1} \rho A a \mathbf{N}^{T} \mathbf{N} d\xi = \frac{\rho A a}{210} \begin{bmatrix} 156 & 44a & 54 & -26a \\ 44a & 16a^{2} & 26a & -12a^{2} \\ 54 & 26a & 156 & -44a \\ -26a & -12a^{2} & -44a & 16a^{2} \end{bmatrix}$$

The work done by the distributed forces p is

$$\delta W_1^e = \int_{-a}^a p \, \delta w \, dx = \int_{-1}^1 p \, \delta w \, ad \, \xi = \delta \mathbf{w}^{eT} a \int_{-1}^1 p \mathbf{N}^T d \, \xi$$

The vector of nodal forces equivalent to distributed uniform forces p is obtained as

$$\mathbf{f}^e = ap \int_{-1}^{1} \mathbf{N}^T d\xi = \frac{ap}{3} \begin{vmatrix} 3 \\ a \\ 3 \\ -a \end{vmatrix}$$

 \bullet The work done by the axial force N^0 is defined as

$$\delta W_2^e = \int_{-a}^a N_0 \frac{\partial w}{\partial x} \frac{\partial \delta w}{\partial x} dx = \int_{-1}^1 \frac{N_0}{a^2} \frac{\partial w}{\partial \xi} \frac{\partial \delta w}{\partial \xi} a d\xi = \delta \mathbf{w}^{eT} \frac{N_0}{a} \int_{-1}^1 \mathbf{N}'^T \mathbf{N}' d\xi \mathbf{w}^e$$

The stability matrix is defined as

$$\mathbf{G}^{e} = \frac{1}{a} \int_{-1}^{1} \mathbf{N}'^{T} \mathbf{N}' d\xi = \frac{1}{60a} \begin{bmatrix} 36 & 6a & -36 & 6a \\ 6a & 16a^{2} & -6a & -4a^{2} \\ -36 & -6a & 36 & -6a \\ 6a & -4a^{2} & -6a & 16a^{2} \end{bmatrix}$$

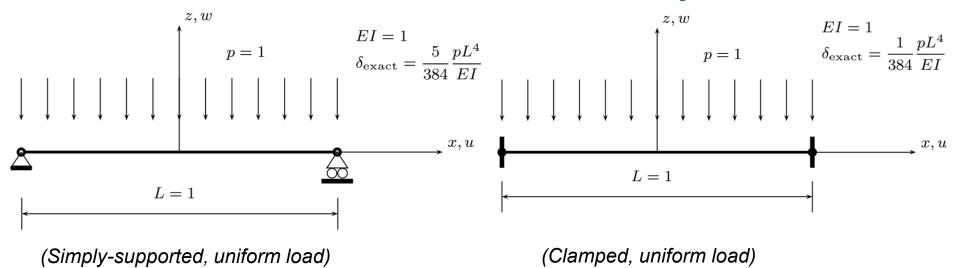
After assembly, the algebraic solving system is

$$\mathbf{M}\ddot{\mathbf{u}} + \mathbf{K}\mathbf{u} - N_0 \mathbf{G}^e \mathbf{u} = \mathbf{f}$$

From this, the static, buckling, and free vibration problems of Bernoulli beam can be solved. The buckling and free vibration problems are both solved in the form of an eigenvalue problem.

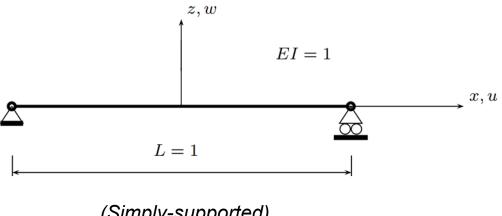


1. Bernoulli Beam – Bernoulli Beam Problem - Python



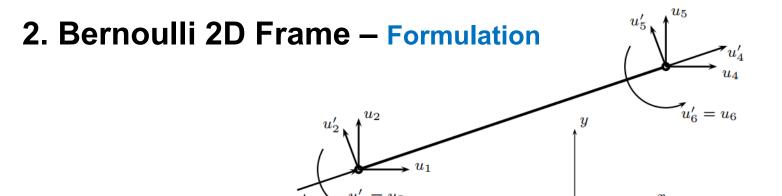
In the Python codes, students can define the number of elements.

1. Bernoulli Beam - Bernoulli Beam Free Vibrations



(Simply-supported)

In the free vibration, there is no external load.



- The previous Bernoulli beam formulations in local coordinates will be generalized.
- ❖ The local stiffness and mass matrices are transformed by a matrix of rotation which is a function of the beam slope w.r.t the horizontal axis.
- ❖ Each node has 3 DOFs (i.e., 2 displacements and 1 rotation)
- The vector of local displacements is given $\mathbf{u'}^T = \begin{bmatrix} u_1' & u_4' & u_2' & u_3' = u_3 & u_6' = u_6 \end{bmatrix}$
- The vector of global displacements is given $\mathbf{u}^T = \begin{bmatrix} u_1 & u_4 & u_2 & u_5 & u_3 = u_3' & u_6 = u_6' \end{bmatrix}$
- ightharpoonup The relation between local and global displacement $\mathbf{u}' = \mathbf{L}\mathbf{u}$

with
$$\mathbf{L} = \begin{bmatrix} l & 0 & m & 0 & 0 & 0 \\ 0 & l & 0 & m & 0 & 0 \\ -m & 0 & l & 0 & 0 & 0 \\ 0 & -m & 0 & l & 0 & 0 \\ 0 & 0 & 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 0 & 0 & 1 \end{bmatrix}$$
 and

$$l = \cos \theta = \frac{x_2 - x_1}{L_e}$$

$$m = \sin \theta = \frac{y_2 - y_1}{L_e}$$

$$L_e = \sqrt{(x_2 - x_1)^2 + (y_2 - y_1)^2}$$



In local coordinates, the DOFs of the frame element are the combination of the DOFs of a bar element u'_1 and the DOFs of a Bernoulli beam element u'_2 and u'_3 . Thus, the local stiffness matrix of the frame element is the combination of the stiffness of the bar

stiffness matrix of the frame element is the combination of the stiffness of the bar element and the Bernoulli beam element.
$$\mathbf{K}'^e = \frac{E}{L^3}\begin{bmatrix} AL^2 & -AL^2 & 0 & 0 & 0 & 0 \\ AL^2 & 0 & 0 & 0 & 0 \\ & & 12I & -12I & 6IL & 6IL \\ & & & & 12I & -6IL & -6IL \\ & & & & & 4IL^2 & 2IL^2 \\ sym & & & & & 4IL^2 \end{bmatrix} \qquad \mathbf{N}$$

$$\mathbf{N}$$

$$\mathbf$$

$$\mathbf{w}^{eT} = \begin{bmatrix} u_2' = w_1 & u_3' = \frac{\partial w_1}{\partial x} & u_5' = w_2 & u_6' = \frac{\partial w_2}{\partial x} \end{bmatrix}$$
$$\mathbf{u}^T = \begin{bmatrix} u_1 & u_4 & u_2 & u_5 & u_3 = u_3' & u_6 = u_6' \end{bmatrix}$$

The strain energy is transformed to global coordinates as

$$U^e = \frac{1}{2} \mathbf{u'}^T \mathbf{K'}^e \mathbf{u'} = \frac{1}{2} \mathbf{u}^T \mathbf{L}^T \mathbf{K'}^e \mathbf{L} \mathbf{u} = \frac{1}{2} \mathbf{u}^T \mathbf{K} \mathbf{u} \quad \text{with the global stiffness of the frame element} \quad \mathbf{K} = \mathbf{L}^T \mathbf{K'}^e \mathbf{L}$$

The local mass matrix of the frame element is the combination of the mass of the bar element and the Bernoulli beam element.

$$\mathbf{M'}^e = \frac{\rho AL}{420} \begin{bmatrix} 140 & 70 & 0 & 0 & 0 & 0 \\ 140 & 0 & 0 & 0 & 0 \\ & 156 & 54 & 22L & -13L \\ & & 156 & 13L & -22L \\ sym & & & 4L^2 \end{bmatrix}$$
Note: the new numbering of global DOFs in the frame element and the Bernoulli beam element when assembling the stiffness matrix.
$$\mathbf{w}^{eT} = \begin{bmatrix} u_2' = w_1 & u_3' = \frac{\partial w_1}{\partial x} & u_5' = w_2 & u_6' = \frac{\partial w_2}{\partial x} \end{bmatrix}$$

$$\mathbf{u}^T = \begin{bmatrix} u_1 & u_4 & u_2 & u_5 & u_3 = u_3' & u_6 = u_6' \end{bmatrix}$$
In global coordinates, the kinetic energy is given by

Note: the new numbering of global DOFs in the frame element and the Bernoulli beam element

$$\mathbf{w}^{eT} = \begin{bmatrix} u_2' = w_1 & u_3' = \frac{\partial w_1}{\partial x} & u_5' = w_2 & u_6' = \frac{\partial w_2}{\partial x} \end{bmatrix}$$
$$\mathbf{u}^T = \begin{bmatrix} u_1 & u_4 & u_2 & u_5 & u_3 = u_3' & u_6 = u_6' \end{bmatrix}$$

In global coordinates, the kinetic energy is given by

$$M^e = \frac{1}{2} \mathbf{u'}^T \mathbf{M'}^e \mathbf{u'} = \frac{1}{2} \mathbf{u}^T \mathbf{L}^T \mathbf{M'}^e \mathbf{L} \mathbf{u} = \frac{1}{2} \mathbf{u}^T \mathbf{M} \mathbf{u}$$
 with the global mass of the frame element $\mathbf{M} = \mathbf{L}^T \mathbf{M'}^e \mathbf{L}$

The global load vector is also defined according to the order of the global DOFs

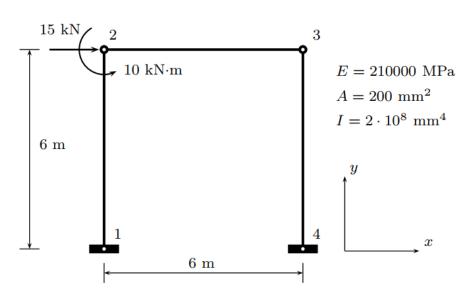
$$\mathbf{F} = \begin{bmatrix} F_{x,1} & \dots & F_{x,n} & F_{y,1} & \dots & F_{y,n} & M_1 & \dots & M_n \end{bmatrix}$$

 F_x, F_y, M are the horizontal and vertical concentrated forces and moments applied at the nodes. is the number of nodes in the mesh

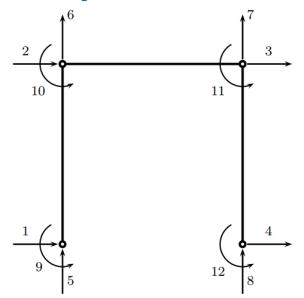
Static and free vibration problems when F = 0.



2. Bernoulli 2D Frame – 2D Frame Problem - Python

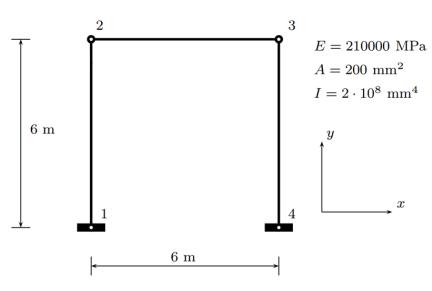


(A 2D frame: geometry, materials, and loads)

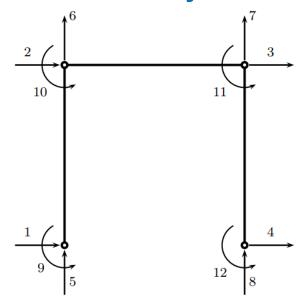


(DOF ordering)

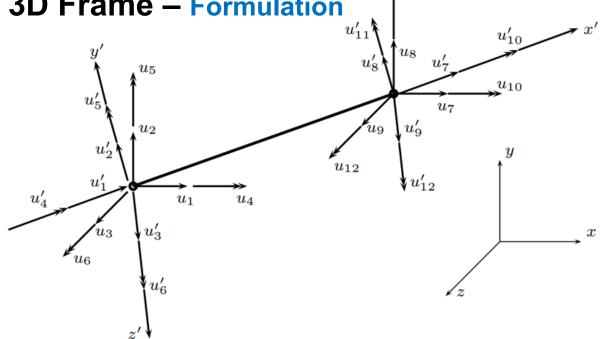
2. Bernoulli 2D Frame - 2D Frame in Free Vibration - Python



Density: $\rho = 8.05 \cdot 10^{-9} \ ton/mm^3$



(DOF ordering)



- ❖ In global coordinates, each node has 3 displacement DOFs and 3 rotation DOFs.
- More complex than 2D beam element due to the orientation of the beam in 3D space.
- \diamond The local axis x' is aligned with the beam major axis.
- \diamond The direction cosines according to axis x' is straightforward and follows that of 2D frame.

$$C_{xx'} = \frac{x_2 - x_1}{L_e}; \quad C_{yx'} = \frac{y_2 - y_1}{L_e}; \quad C_{zx'} = \frac{z_2 - z_1}{L_e}; \quad L_e = \sqrt{\left(x_2 - x_1\right)^2 + \left(y_2 - y_1\right)^2 + \left(z_2 - z_1\right)^2}$$

The vector rotation matrix in 3D space

$$\mathbf{r} = \begin{bmatrix} C_{xx'} & C_{yx'} & C_{zx'} \\ C_{xy'} & C_{yy'} & C_{zy'} \\ C_{xz'} & C_{yz'} & C_{zz'} \end{bmatrix} = R_{\alpha}R_{\beta}R_{\gamma} \quad \text{with} \quad R_{\alpha}, R_{\beta}, R_{\gamma} \quad \text{are three rotation matrices.}$$

$$\alpha, \beta, \gamma \quad \text{are the rotation angles about} \quad x', y', z' \quad \text{axes.}$$

• Rotation about z axis, R_{ν} is given by

$$R_{\gamma} = \begin{bmatrix} \cos \gamma & \sin \gamma & 0 \\ -\sin \gamma & \cos \gamma & 0 \\ 0 & 0 & 1 \end{bmatrix} \quad \text{with} \quad \cos \gamma = C_{xx'}/C_{xy}, \quad \sin \gamma = C_{yx'}/C_{xy}, \quad C_{xy} = \sqrt{C_{xx'}^2 + C_{yx'}^2}$$

• Rotation about y axis, R_{β} is given by

$$R_{\beta} = \begin{bmatrix} \cos \beta & 0 & \sin \beta \\ 0 & 1 & 0 \\ -\sin \beta & 0 & \cos \beta \end{bmatrix} \quad \text{with} \quad \cos \beta = C_{xy}, \quad \sin \beta = C_{yx'}$$

• Combining the rotation about y and z axis, the vector rotation matrix is $\mathbf{r} = R_{\beta}R_{\gamma} = \begin{bmatrix} C_{xx'} & C_{yx'} & C_{zx'} \\ -C_{yx'}/C_{xy} & C_{xx'}/C_{xy} & 0 \\ -C_{xx'}C_{zx'}/C_{xy} & -C_{yx'}C_{zx'}/C_{xy} & C_{xy} \end{bmatrix}$

When $\beta = 90^{\circ}$ or $\beta = 270^{\circ}$, the 2-node beam element change only along z axis, thus the vector rotation matrix takes a special form.

$$\mathbf{r} = \begin{bmatrix} 0 & 0 & C_{zx'} \\ 0 & 1 & 0 \\ -C_{zx'} & 0 & 0 \end{bmatrix} \qquad \text{for } z_2 > z_1 \text{ or } \beta = 90^0 \text{, } C_{zx'} = 1$$

$$\text{for } z_1 > z_2 \text{ or } \beta = 270^0 \text{, } C_{zx'} = -1$$

- $\mathbf{r} = \begin{bmatrix} 0 & 0 & C_{zx'} \\ 0 & 1 & 0 \\ -C_{zx'} & 0 & 0 \end{bmatrix} \quad \text{for } z_2 > z_1 \text{ or } \beta = 90^{\circ}, \ C_{zx'} = 1$ $\text{for } z_1 > z_2 \text{ or } \beta = 270^{\circ}, \ C_{zx'} = -1$ If the extra rotation α is included, the rotation matrix is $R_{\alpha} = \begin{bmatrix} 1 & 0 & 0 \\ 0 & \cos \alpha & -\sin \alpha \\ 0 & \sin \alpha & \cos \alpha \end{bmatrix}$
- Finally, the vector rotation matrix becomes

$$\mathbf{r} = R_{\alpha}R_{\beta}R_{\gamma} = \begin{bmatrix} C_{xx'} & C_{yx'} & C_{zx'} \\ (C_{xx'}C_{zx'}\sin\alpha - C_{yx'}\cos\alpha)/C_{xy} & (C_{yx'}C_{zx'}\sin\alpha + C_{xx'}\cos\alpha)/C_{xy} & -C_{xy}\sin\alpha \\ (-C_{xx'}C_{zx'}\cos\alpha + C_{yx'}\sin\alpha)/C_{xy} & (-C_{yx'}C_{zx'}\cos\alpha + C_{xx'}\sin\alpha)/C_{xy} & C_{xy}\cos\alpha \end{bmatrix}$$

Special case of vertical members $\beta = 90^{\circ}$ and $\beta = 270^{\circ}$ can be derived as

$$\mathbf{r} = \begin{bmatrix} 0 & 0 & C_{zx'} \\ C_{zx'} \sin \alpha & \cos \alpha & 0 \\ -C_{zx'} \cos \alpha & \sin \alpha & 0 \end{bmatrix}$$



In the local coordinate system, the stiffness matrix is given by **

$$\mathbf{K'} = \begin{bmatrix} EA/L & 0 & 0 & 0 & 0 & 0 & -EA/L & 0 & 0 & 0 & 0 & 0 \\ 12EI_z/L^3 & 0 & 0 & 0 & 6EI_z/L^2 & 0 & -12EI_z/L^3 & 0 & 0 & 0 & 6EI_z/L^2 \\ 12EI_y/L^3 & 0 & -6EI_y/L^2 & 0 & 0 & 0 & -12EI_y/L^3 & 0 & -6EI_y/L^2 & 0 \\ GJ/L & 0 & 0 & 0 & 0 & 0 & -GJ/L & 0 & 0 & 0 \\ 4EI_y/L & 0 & 0 & 0 & 6EI_z/L^2 & 0 & 2EI_y/L & 0 \\ EA/L & 0 & -6EI_z/L^2 & 0 & 0 & 0 & 2EI_z/L \\ EA/L & 0 & 0 & 0 & 0 & 0 & -6EI_z/L^2 \\ & & & & & & & & & & & & & & & \\ & & & & & & & & & & & & \\ & & & & & & & & & & & & \\ & & & & & & & & & & & & \\ & & & & & & & & & & & \\ & & & & & & & & & & & \\ & & & & & & & & & & & \\ & & & & & & & & & & & \\ & & & & & & & & & & \\ & & & & & & & & & & \\ & & & & & & & & & & \\ & & & & & & & & & & \\ & & & & & & & & & & \\ & & & & & & & & & \\ & & & & & & & & & \\ & & & & & & & & & \\ & & & & & & & & & \\ & & & & & & & & & \\ & & & & & & & & \\ & & & & & & & & \\ & & & & & & & & \\ & & & & & & & & \\ & & & & & & & & \\ & & & & & & & & \\ & & & & & & & & \\ & & & & & & & & \\ & & & & & & & & \\ & & & & & & & & \\ & & & & & & & & \\ & & & & & & & & \\ & & & & & & & & \\ & & & & & & & & \\ & & & & & & & & \\ & & & & & & & & \\ & & & & & & & & \\ & & & & & & & & \\ & & & & & & & \\ & & & & & & & \\ & & & & & & & \\ & & & & & & & \\ & & & & & & \\ & & & & & & & \\ & & & & & & & \\ & & & & & & & \\ & & & & & & & \\ & & & & & & & \\ & & & & & & & \\ & & & & & & & \\ & & & & & & & \\ & & & & & & & \\ & & & & & & & \\ & & & & & & & \\ & & & & & & \\ & & & & & & & \\ & & & & & & \\ & & & & & & & \\ & & & & & \\ & & & & & & \\ & & & & & & \\ & & & & & & \\ & & & & & &$$

- The global stiffness matrix is $\mathbf{K} = \mathbf{R}^T \mathbf{K}' \mathbf{R}$ with $\mathbf{R} = \begin{vmatrix} \mathbf{r} & \mathbf{0} & \mathbf{0} & \mathbf{0} \\ \mathbf{0} & \mathbf{r} & \mathbf{0} & \mathbf{0} \\ \mathbf{0} & \mathbf{0} & \mathbf{r} & \mathbf{0} \end{vmatrix}$ and \mathbf{r} is defined previously.
- After the static problem is solved, the reactions can be computed as $\mathbf{F} = \mathbf{K}\mathbf{U}$.
- The element nodal forces can be evaluated by axes transformation as $\mathbf{f}_e = \mathbf{K}' \mathbf{R} \mathbf{U}_e$ with \mathbf{f}_e is the element nodal forces vector, \mathbf{U}_e is the global vector of displacement of element e.



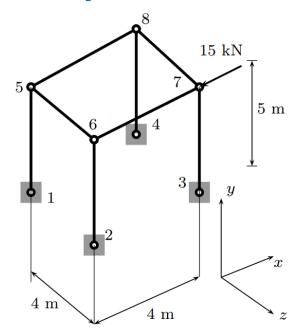
In the local coordinate system, the consistent and lumped mass matrices are defined as

with $r_x^2 = (I_y' + I_z')/A$ and I_y' , I_z' are the second moment of area of the cross-section about the principal y' and z' axes.

The mass matrix in the global coordinate system takes the form $\mathbf{M} = \mathbf{R}^T \mathbf{M}' \mathbf{R}$ with \mathbf{R} is the rotation matrix defined previously.

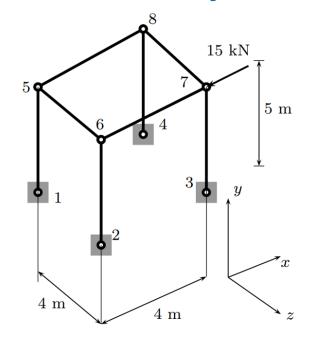
3. Bernoulli 3D Frame – 3D Frame Problem - Python

$$E = 210 \, GPa$$
, $G = 84 \, GPa$, $A = 2 \cdot 10^{-2} \, m^2$
 $I_y = 1 \cdot 10^{-4} \, m^4$, $I_z = 2 \cdot 10^{-4} \, m^4$, $J = 0.5 \cdot 10^{-4} \, m^4$



3. Bernoulli 3D Frame – 3D Frame in Free Vibrations - Python

$$E = 210 \, GPa$$
, $G = 84 \, GPa$, $A = 2 \cdot 10^{-2} \, m^2$
 $I_y = 1 \cdot 10^{-4} \, m^4$, $I_z = 2 \cdot 10^{-4} \, m^4$, $J = 0.5 \cdot 10^{-4} \, m^4$
 $\rho = 7850 \, kg/m^3$



Content

- 1. Bernoulli Beam
 - Formulation
 - Bernoulli Beam Problem
 - Bernoulli Beam Free Vibrations
- 2. Bernoulli 2D Frames
 - Formulation
 - 2D Frame Problem Python
 - 2D Frame in Free Vibrations Python
- 3. Bernoulli 3D Frames
 - Formulation
 - 3D Frame Problem Python
 - 3D Frame in Free Vibrations Python



