

4. (a) Consider a spacetime with a Killing vector ξ . Show that if the energy-momentum tensor $T^{\mu\nu}$ satisfies the conservation equation

$$\nabla_\mu T^{\mu\nu} = 0, \quad (1)$$

then the four-vector $Q^\mu := T^{\mu\nu}\xi_\nu$ also satisfies a conservation law $\nabla_\mu Q^\mu = 0$.

- (b) Now let us consider a flow consisting of perfect fluid and governed by adiabatic equation of state. Using first law of thermodynamics and specific enthalpy $h := (p + \varepsilon)/\rho$, where p is pressure density, ε is the internal energy density and ρ is the fluid density, show that the conservation law in Eq. (1) provides

$$u^\mu \nabla_\mu (h u_\nu) = \frac{\nabla_\nu p}{\rho} \quad (2)$$

- (c) Using Eq. (2), find the expression for $\mathcal{L}_u(h u_\mu)$. Contracting it with ξ , find the expression for $\mathcal{L}_u(h u_\mu \xi^\mu)$ in terms of $\mathcal{L}_\xi p$ and $\mathcal{L}_\xi h$.
- (d) Let's also assume that the flow is invariant under the same symmetry group of the spacetime which provides the Killing vector ξ . Mathematically this gives the condition

$$\mathcal{L}_\xi(B) = 0, \quad (3)$$

where B is any tensor field associated with the flow, e.g., pressure, density etc. Show that this implies that $h u_\mu \xi^\mu$ is conserved along the flow lines.

- (e) Find the conserved quantity along the flow lines when the spacetime is stationary and the flow is stationary. This is the *general relativistic Bernoulli constant*.
- (f) Find the Newtonian limit and check that the quantity reduces to the familiar expression of Bernoulli equation.

Answer:

$$(a) \quad \nabla_\mu (T^{\mu\nu} \xi_\nu) = \underbrace{\nabla_\mu T^{\mu\nu}}_{=0 \text{ due to conservation law}} \times \xi_\nu + \underbrace{T^{\mu\nu}}_{\text{symmetric}} \times \underbrace{\nabla_\mu \xi_\nu}_{\text{anti-symmetric}} = 0.$$

- (b) The energy-momentum tensor for perfect fluid is given by $T^{\mu\nu} = (p + \varepsilon)v^\mu v^\nu + p g^{\mu\nu}$.

$$\begin{aligned} \nabla_\mu T^{\mu\nu} &= \nabla_\mu [(p + \varepsilon)v^\mu v^\nu] + g^{\mu\nu} \nabla_\mu p \quad \text{metric is divergenceless} \\ &= \nabla_\mu (h v^\nu \times \rho v^\mu) + g^{\mu\nu} \nabla_\mu p \\ &= \nabla_\mu (h v^\nu) \rho v^\mu + g^{\mu\nu} \nabla_\mu p \quad \text{using } \nabla_\mu (\rho v^\mu) = 0 \\ &= 0. \end{aligned}$$

This gives

$$\boxed{v^\mu \nabla_\mu (h v^\nu) = -\frac{g^{\mu\nu} \nabla_\mu p}{\rho} = -\frac{\nabla^\nu p}{\rho}.} \quad (4)$$

- (c)

$$\begin{aligned} \mathcal{L}_v(h v_\mu) &= v^\nu \nabla_\nu (h v_\mu) + h v_\nu \nabla_\mu v^\nu \\ &= -\frac{\nabla_\mu p}{\rho} \end{aligned}$$

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The second term in first line vanishes due to the $v^\mu v_\mu = -1$. Contracting with ξ gives

$$\xi^\mu \mathcal{L}_v(hv_\mu) = -\frac{1}{\rho} \xi^\mu \nabla_\mu p$$

Therefore,

$$\begin{aligned} \mathcal{L}_v(\xi^\mu hv_\mu) &= \xi^\mu \mathcal{L}_v(hv_\mu) + hv_\mu \mathcal{L}_v \xi^\mu \\ &= -\frac{\xi^\mu}{\rho} \nabla_\mu p + hv_\mu v^\nu \nabla_\nu \xi^\mu - hv_\mu \xi^\nu \nabla_\nu v^\mu \\ &= -\frac{\xi^\mu}{\rho} \nabla_\mu p \\ &= -\frac{\mathcal{L}_\xi p}{\rho}. \end{aligned}$$

In the second line, the second term vanishes because $\nabla_\nu \xi^\mu$ is anti-symmetric but $v_\mu v^\nu$ is symmetric. The third term vanishes using $v^\mu v_\mu = -1$.

- (d) Using $\mathcal{L}_\xi p = 0$ gives $\mathcal{L}_v(\xi^\mu hv_\mu) = 0$.

$$\mathcal{L}_v(\xi^\mu hv_\mu) = v^\nu \nabla_\nu (\xi^\mu hv_\mu) = 0$$

which implies that $\xi^\mu hv_\mu$ is constant along the flow line.

- (e) In stationary spacetime we have the stationary Killing vector $\xi^\mu = \delta_t^\mu$. This provides the constant hv_t along the flow lines.

- (f)

$$\begin{aligned} hv_t &= \frac{\varepsilon + p}{\rho} v_t \\ &= \frac{\rho(c^2 + \epsilon) + p}{\rho} v_t \\ &= (c^2 + \epsilon + \frac{p}{\rho}) v_t \quad \epsilon \ll c^2, \text{ in non-relativistic limit} \\ &= (c^2 + \frac{p}{\rho}) (-1 - \frac{\phi}{c^2} - \frac{1}{2} \frac{v_i v^i}{c^2}) \\ &= -c^2 - \phi - \frac{1}{2} v_i v^i - \frac{p}{\rho} + \mathcal{O}(1/c^2) \end{aligned}$$

ϵ is the thermal energy and ϕ is the potential energy. This gives the newtonian limit of Bernoulli constant $\boxed{\phi + (1/2)v_i v^i + p/\rho}$.