

3. Prove that

$$\xi^1 = \sin \phi \partial_\theta + \cot \theta \cos \phi \partial_\phi \quad \xi^2 = -\cos \phi \partial_\theta + \cot \theta \sin \phi \partial_\phi \quad (1)$$

are the Killing vectors of the spherically symmetric spacetime

$$ds^2 = -A(r)dt^2 + B(r)dr^2 + r^2(d\theta^2 + \sin^2 \theta d\phi^2) \quad (2)$$

Answer: Killing vectors satisfy the Killing equation

$$\nabla_\alpha \xi_\beta + \nabla_\beta \xi_\alpha = 0. \quad (3)$$

Thus we need to show that $\nabla_\alpha \xi_\beta$ is an anti-symmetric tensor. We do it for ξ^1 . The contravariant components are

$$\xi_t^1 = \xi_r^1 = 0, \quad \xi_\theta^1 = g_{\theta\theta} \xi^{1\theta} = r^2 \sin \phi, \quad \xi_\phi^1 = g_{\phi\phi} \xi^{1\phi} = r^2 \sin^2 \theta \cot \theta \cos \phi = r^2 \sin \theta \cos \theta \cos \phi. \quad (4)$$

Since only ξ_θ and ξ_ϕ are non-vanishing, the relevant Christoffel symbols are $\Gamma_{\alpha\beta}^\theta$ and $\Gamma_{\alpha\beta}^\phi$. These we get from the geodesic equation for θ and ϕ using $\mathcal{L}^2 = (1/2)g_{\alpha\beta}\dot{x}^\alpha\dot{x}^\beta$ as the Lagrangian. For θ , we have

$$\ddot{\theta} + \frac{2}{r}\dot{r}\dot{\theta} - \sin \theta \cos \theta \ddot{\phi} = 0, \quad (5)$$

which provides

$$\Gamma_{r\theta}^\theta = \frac{1}{r}, \quad \Gamma_{\phi\phi}^\theta = -\sin \theta \cos \theta. \quad (6)$$

Similarly for ϕ ,

$$\ddot{\phi} + \frac{2}{r}\dot{r}\dot{\phi} + 2 \cot \theta \dot{\theta} \dot{\phi} = 0, \quad (7)$$

which provides

$$\Gamma_{r\phi}^\phi = \frac{1}{r}, \quad \Gamma_{\theta\phi}^\phi = \cot \theta. \quad (8)$$

Now that we have the required Christoffel symbols, we check that $\nabla_\alpha \xi_\beta^1$ is an anti-symmetric tensor.

$$\boxed{\nabla_t \xi_t^1 = \nabla_r \xi_r^1 = 0}$$

trivially. and

$$\boxed{\nabla_\theta \xi_\theta^1 = 0, \quad \nabla_\phi \xi_\phi^1 = \partial_\phi \xi_\phi^1 - \Gamma_{\phi\phi}^\phi = 0}$$

$$\nabla_t \xi_r^1 = -\nabla_r \xi_t^1 = 0. \quad \nabla_t \xi_\theta^1 = 0 = -\nabla_\theta \xi_t^1. \quad \nabla_t \xi_\phi^1 = 0 = -\nabla_\phi \xi_t^1.$$

$$\begin{aligned} \nabla_r \xi_\theta^1 &= \partial_r \xi_\theta^1 - \Gamma_{r\theta}^\theta \xi_\theta^1 \\ &= 2r \sin \phi - \frac{1}{r} r^2 \sin \phi = r \sin \phi. \end{aligned}$$

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$$\boxed{\nabla_\theta \xi_r^1 = -\Gamma_{r\theta}^\theta \xi_\theta^1 = -r \sin \phi = -\nabla_r \xi_\theta^1}.$$

$$\begin{aligned} \nabla_r \xi_\phi^1 &= \partial_r \xi_\phi^1 - \Gamma_{r\phi}^\phi \xi_\phi^1 \\ &= 2r \sin \theta \cos \theta \cos \phi - \frac{1}{r} r^2 \sin \theta \cos \theta \cos \phi \\ &= r \sin \theta \cos \theta \cos \phi. \end{aligned}$$

$$\boxed{\nabla_\phi \xi_r^1 = -\Gamma_{r\phi}^\phi \xi_\phi^1 = -r \sin \theta \cos \theta \cos \phi = -\nabla_r \xi_\phi^1}.$$

$$\begin{aligned} \nabla_\theta \xi_\phi^1 &= \partial_\theta \xi_\phi^1 - \Gamma_{\theta\phi}^\phi \xi_\phi^1 \\ &= r^2 \cos 2\theta \cos \phi - \cot \theta r^2 \sin \theta \cos \theta \cos \phi \\ &= r^2 \cos \phi (\cos 2\theta - \cos^2 \theta) \\ &= r^2 \cos \phi (-\sin^2 \theta) \end{aligned}$$

$$\begin{aligned} \nabla_\phi \xi_\theta^1 &= \partial_\phi \xi_\theta^1 - \Gamma_{\theta\phi}^\phi \xi_\phi^1 \\ &= r^2 \cos \phi - \cot \theta r^2 \sin \theta \cos \theta \cos \phi \\ &= r^2 \cos \phi (1 - \cos^2 \theta) \\ &= r^2 \cos \phi \sin^2 \theta \\ &\Rightarrow \boxed{\nabla_\phi \xi_\theta^1 = -\nabla_\theta \xi_\phi^1} \end{aligned}$$

Similarly, we can check for ξ^2 .

4. A particle with electric charge e moves in a spacetime with metric $g_{\alpha\beta}$ in the presence of a vector potential A_α . The equations of motion are $u^\beta \nabla_\beta u_\alpha = e F_{\alpha\beta} u^\beta$, where u^α is the four-velocity and $F_{\alpha\beta} = \nabla_\alpha A_\beta - \nabla_\beta A_\alpha$. It is assumed that the spacetime possesses a Killing vector ξ^α , so that $\mathcal{L}_\xi g_{\alpha\beta} = \mathcal{L}_\xi A_\alpha = 0$. Prove that

$$(u_\alpha + e A_\alpha) \xi^\alpha \tag{9}$$

is constant on the world line of the of the charged particle.

Answer:

$$\begin{aligned} \frac{d}{d\lambda} [(u_\alpha + e A_\alpha) \xi^\alpha] &= u^\beta \nabla_\beta [(u_\alpha + e A_\alpha) \xi^\alpha] \\ &= u^\beta (\nabla_\beta u^\alpha) \xi^\alpha + \underbrace{u^\beta u^\alpha}_{\text{symmetric}} \times \underbrace{\nabla_\beta \xi^\alpha}_{\text{anti-symmetric}} + e \xi^\alpha u^\beta \nabla_\beta A_\alpha + e A_\alpha u^\beta \nabla_\beta \xi^\alpha \\ &= e u^\beta (\nabla_\alpha A_\beta - \nabla_\beta A_\alpha) \xi^\alpha + \cancel{e \xi^\alpha u^\beta \nabla_\beta A_\alpha} + e A_\alpha u^\beta \nabla_\beta \xi^\alpha \\ &= e u^\beta (\xi^\alpha \nabla_\alpha A_\beta + A_\alpha \nabla_\beta \xi^\alpha) \\ &= e u^\beta \mathcal{L}_\xi A_\alpha \\ &= 0 \end{aligned}$$

5. A particle moving on a circular orbit in a stationary, axially symmetric spacetime is subjected to a dissipative force which drives it to another, slightly smaller, circular orbit. During the transition, the particle loses an amount $\delta \tilde{E}$ of orbital energy (per unit rest mass) and an amount $\delta \tilde{L}$ of orbital angular momentum (per unit rest mass). Show that these quantities are related by $\delta \tilde{E} = \Omega \delta \tilde{L}$, where Ω is the particle's original angular velocity.

Hints: Express the four-velocity u^α of the particle in terms of the Killing vectors, energy angular momentum and orbital velocity. Find the variation δu^α . Use the normalization condition $u_\alpha u^\alpha = -1$.

Answer: The spacetime possesses the Killing vectors $t^\alpha = \delta_t^\alpha$ and $\phi^\alpha = \delta_\phi^\alpha$. In terms of these we can write down the four-velocity u^α as

$$u^\alpha = u^t t^\alpha + u^\phi \phi^\alpha = u^t (t^\alpha + \Omega \phi^\alpha), \quad \Omega = \frac{u^\phi}{u^t} \quad (10)$$

as the particle moves in a circular orbit. Using the normalization condition $u_\alpha u^\alpha = -1$ we find

$$u_\alpha u^\alpha = u_t u^t + u_\phi u^\phi \Omega = -1 \quad (11)$$

or,

$$u^t (u_t + u_\phi \Omega) = -1 \quad (12)$$

Now the two constants of motion from the two Killing vectors are $E = -t^\alpha u_\alpha = -u_t$ and $L = \phi^\alpha u_\alpha = u_\phi$. Thus in terms of E and L , u^t becomes

$$u^t = \frac{1}{E - \Omega L} \quad (13)$$

and hence the four-velocity could be written as

$$u^\alpha = \frac{(t^\alpha + \Omega \phi^\alpha)}{(E - \Omega L)} \quad (14)$$

The variation of the four-velocity would be given by

$$\delta u^\alpha = -\frac{(t^\alpha + \Omega \phi^\alpha)}{(E - \Omega L)^2} (\delta E - \Omega \delta L) = -u^\alpha u^t (\delta E - \Omega \delta L) \quad (15)$$

contracting the above with u_α gives

$$u_\alpha \delta u^\alpha = u^t (\delta E - \Omega \delta L) \quad (16)$$

Now, since $u_\alpha u^\alpha = -1$, variation of it should vanish which implies $u_\alpha \delta u^\alpha = 0$. This provides the relation

$$\boxed{\delta E = \Omega \delta L}. \quad (17)$$