## white dwarfs

# Course on Compact Objects: Assignment 4

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#### Problem 1

Show that the pressure exerted by a gas of particle with isotropic momentum distribution n(p) is given by

$$P=rac{1}{3}\int_0^\infty p v_p n(p) dp$$

where  $v_p$  is the velocity associated with momentum p.

Let us consider the pressure due to a particle moving at an angle  $\theta$  with the z direction. The component of velocity along z would be given by  $v\cos\theta$ 

$$v_z = v \cos \theta$$

and therefore,

$$p_z = p\cos\theta.$$

The rate of change of momentum is given by

$$rac{p_z}{\Delta z/v_z} = rac{p_z v_z}{\Delta z} = rac{p v \cos^2 heta}{\Delta z}$$

So pressure per particle moving at an angle  $\theta$  is

$$\frac{pv\cos^2\theta}{\Delta z\Delta x\Delta y} = \frac{pv\cos^2\theta}{V}.$$

The total pressure for all particles moving at an angle  $\theta$  would be given by

$$P = \int_0^\infty rac{pv\cos^2 heta}{V} N(p, heta) dp$$

where  $N(p,\theta)$  is the number of particles with momentum between p and p+dp and with angle  $\theta$  to  $\theta+d\theta$ . This number would be

$$N(p, heta)=rac{d^3p}{h^3}g_sf(p)=g_sf(p)rac{2\pi p^2\sin heta d heta dp}{h^3},$$

where  $g_s$  is the degeracy of the state and f(p) is the probabilty that the state would be occupied.

Thus,

$$P=\int_0^\infty rac{pv\cos^2 heta}{V}g_sf(p)rac{2\pi p^2\sin heta d heta dp}{h^3}dp.$$

To get the average total pressure we integrate over  $\theta$ 

$$P=\int_0^\pi d heta\cos^2 heta\sin heta\int_0^\infty rac{pv}{V}rac{g_sf(p)2\pi p^2}{h^3}dp=rac{1}{3}\int_0^\infty rac{pv}{V}rac{g_sf(p)4\pi p^2}{h^3}dp=rac{1}{3}\int_0^\infty pvn(p)dp$$

where  $n(p)=g_sf(p)4\pi p^2/V$  .

### Problem 2

Argue why we are justified in using a 'cold' degenerate equation of state to describe a white dwarf with a temperature  $T\sim 10^4$ K (Hint: Show that the degeneracy parameter  $\mu/kT\gg 0$ , where  $\mu$  is the chemical potential and k the Boltzmann constant. The density of the white dwarf is  $\sim 10^6 g/cm^3$  and the chemical potential  $\sim$  the Fermi energy. Assume that  $\mu_e=2$ ).

We have to basically show that the thermal energy of the electron gas is negligible compared to Fermi energy

```
· using Unitful, UnitfulAtomic
```

```
kB = k
• kB = u"k_au"
```

```
Temp = 10000.0 K
- Temp = 1e4 * u"K"
```

```
ThermalEnergy = 2.0709735e-12 erg

ThermalEnergy = uconvert(u"erg", (3/2) * kB * Temp)
```

```
FermiEnergy (generic function with 1 method)
```

```
function FermiEnergy(ne, me, h)
pF = h * (3 * ne / (8 * pi))^(1/3)
sqrt(pF^2 * c^2 + me^2 * c^4)
end
```

```
h = 6.283185307179586 h

h = 2 * pi * u"h_au"

µe = 2

· µe = 2

amu = 1.66054e-24 g

· amu = 1.66054e-24 * u"g"

ne = 3.011068688498922e29 cm^-3

· ne = 1e6 * u"g/cm^3" / (µe * amu)

me = 9.1090000000000001e-31 kg

· me = 9.109 * 1e-31 * u"kg"

c = c

· c = u"c"

Fenergy = 1.0488022231371987e-6 erg

· Fenergy = uconvert(u"erg", FermiEnergy(ne, me, h))
```

#### 506429.5719559901

Fenergy/ThermalEnergy

### Problem 3

Derive Lane-Emden equation

## Lane-Emden equation

$$rac{dm(r)}{dr}=4\pi r^2
ho(r), rac{dp(r)}{dr}=-rac{Gm(r)
ho(r)}{r^2}$$

. These two equations can be combined to give

$$rac{1}{r^2}rac{d}{dr}igg(rac{r^2}{
ho(r)}rac{dp(r)}{dr}igg) = -4\pi G
ho(r)$$

This can be further written in terms of a dimensionless form using

$$ho = 
ho_c heta^n, r = a \xi, \Gamma = 1 + rac{1}{n}$$

Then we get the following equation

$$\frac{1}{\xi^2} \frac{d}{d\xi} \left( \xi^2 \frac{d\theta}{d\xi} \right) = -\theta^n$$

This the Lane-Emden equation.

### Problem 4

# Solve Lane-Emden Equation

To solve the lane-emden equation we write it as two first order odes by defining  $\psi=rac{d heta}{d ilde{\xi}}$ 

$$rac{d heta}{d\xi}=\psi, rac{d\psi}{d\xi}=-rac{2\psi}{\xi}- heta^n$$

With the following initial values

$$\theta(0) = 1, \psi(0) = 0$$

```
    using DifferentialEquations
```

using StaticArrays

```
laneEmden (generic function with 1 method)
```

```
    function laneEmden(u, n, ξ)
    θ, ψ = u
    dθ = ψ
    dψ = - (2 * ψ / ξ) - θ^n
    @SVector [dθ, dψ]
    end
```

```
u0 = ▶StaticArrays.SArray{Tuple{2},Float64,1,2}: [1.0, 0.0]
• u0 = @SVector [1., 0.]
```

```
ξspan1 = ▶ (1.0e-10, 10)

• ξspan1 = (1e-10, 10)
```

```
problem1 =
2[36mODEProblem2[0m with uType 2[36mStaticArrays.SArray{Tuple{2},Float64,1,2}2[0m and tType timespan: (1.0e-10, 10.0)
u0: [1.0, 0.0]
```

```
• problem1 = ODEProblem(laneEmden, u0, ξspan1, 1.)
```

sol1 = timestamp value1 value2

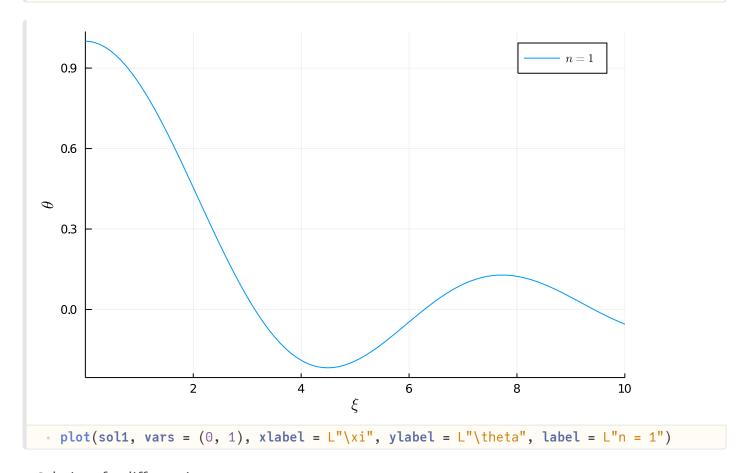
	timestamp	value1	value2
1	1.0e-10	1.0	0.0
2	3.05653e-5	1.0	-9.0812e-6
3	4.38094e-5	1.0	-1.40642e-5
4	0.000107578	1.0	-3.57691e-5
5	0.000187092	1.0	-6.2334e-5
6	0.000467431	1.0	-0.000155805
7	0.000989918	1.0	-0.000329971
8	0.00256049	0.999999	-0.000853496
9	0.00736913	0.999991	-0.00245636

sol1 = solve(problem1)

```
▶ Plots.GRBackend()
```

• using Plots; gr()

#### using LaTeXStrings



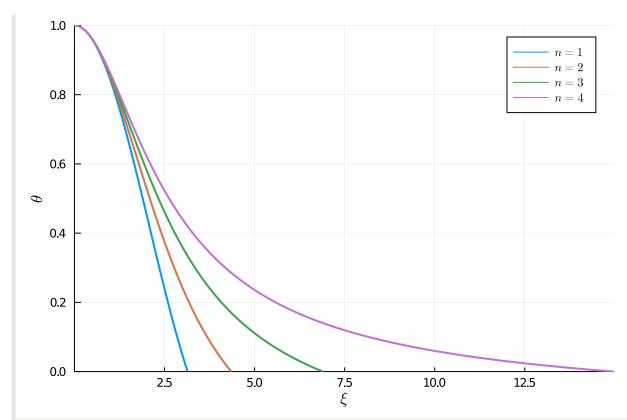
Solutions for different integer n

```
    condition(u, ξ, integrator) = u[1];
    affect!(integrator) = terminate!(integrator);
```

```
ξspan = ▶ (1.0e-10, 20)

• ξspan = (1e-10, 20)
```

• cb = ContinuousCallback(condition, affect!);



```
begin

p = plot();
for n in 1:4

problem = ODEProblem(laneEmden, u0, \xispan, n)

sol = solve(problem, callback = cb)

plot!(p, sol, vars = (0, 1), lw = 2, label = L"n = %\xi")

end

plot!(xlabel = L"\xi", ylabel = L"\theta", ylim = (0, 1))

p

end
end
```

for fractional n the solver would throw error as negative value of  $\theta$  would give complex value for  $\theta^n$ . To avoid this problem we can just change  $\theta$  in our set of odes to  $abs(\theta)$ .

```
laneEmden2 (generic function with 1 method)

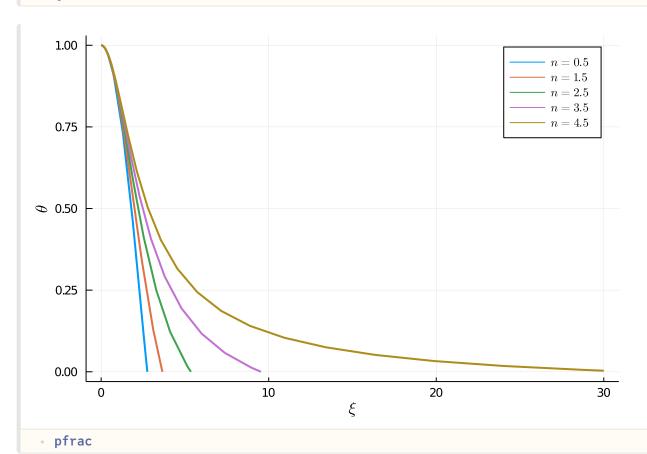
• function laneEmden2(u, n, ξ)
```

```
function laneEmden2(u, n, ξ)
    θ, ψ = u
    dθ = ψ
    dψ = - (2 * ψ / ξ) - abs(θ)^n
    @SVector [dθ, dψ]
    end
```

```
begin
    pfrac = plot();
    ns = 0.5:1.0:4.5
    ξspan2 = [1.0e-10, 30.0]
    ξ1s = []
    ws = []
    for n in ns
        problem = ODEProblem(laneEmden2, u0, ξspan2, n)
        sol = solve(problem, callback = cb)
        push!(ξ1s, sol.t[end])
        push!(ψs, sol[2, :][end])
        plot!(pfrac, sol.t, sol[1, :], lw = 2, xlabel = L"\xi", ylabel = L"\theta", label = L"n = %$n")
    end
end
```

▶ Any[2.75439, 3.65392, 5.3553, 9.53613, 30.0]

• ξ1s



#### using DataFrames

data =

	n	ξ1	ψ1
1	0.5	2.75439	-0.493092
2	1.5	3.65392	-0.203681
3	2.5	5.3553	-0.0762582

	n	ξ1	ψ1
4	3.5	9.53613	-0.0207888
5	4.5	30.0	-0.00193084

data = DataFrame(n = ns, 
$$\xi$$
1 =  $\xi$ 1s,  $\psi$ 1 =  $\psi$ s)

### Problem 5

Radius and Mass of white dwarf

Mass of the star is given by

$$M=\int_0^\infty 4\pi 
ho r^2 dr$$

now,

$$r=\xiigg(rac{4\pi G{
ho_c}^{1-rac{1}{n}}}{K(n+1)}igg)^{-rac{1}{2}}$$

$$\rho = \rho_c \theta^n$$

$$heta^n = -rac{1}{\xi^2}rac{d}{d\xi}igg(\xi^2rac{d heta}{d\xi}igg)$$

Thus,

$$M = 4\pi 
ho_c \Biggl(rac{4\pi G 
ho_c^{-1-rac{1}{n}}}{K(n+1)}\Biggr)^{-3/2} \int_0^{\xi_1} digg(\xi^2 rac{d heta}{d\xi}igg) = 4\pi 
ho_c \Biggl(rac{4\pi G 
ho_c^{-1-rac{1}{n}}}{K(n+1)}\Biggr)^{-3/2} \Bigl(\xi^2 rac{d heta}{d\xi}\Bigr)_{\xi_1}$$

or

$$M=4\pi
ho_c^{(3-n)/2n}igg(rac{4\pi G}{K(n+1)}igg)^{-3/2}igg(\xi^2rac{d heta}{d\xi}igg)_{\xi_1}$$

Now, the radius would be given by value of r at  $\xi=\xi_1$  which is

$$R_\star = \xi_1 \Biggl(rac{4\pi G {
ho_c}^{1-rac{1}{n}}}{K(n+1)}\Biggr)^{-rac{1}{2}}$$

From this we can express  $ho_c$  in terms of  $R_{\star}$  and  $\xi_1$  as

$$ho_c = \left(rac{K(n+1)\xi_1^2}{4\pi G R_\star^2}
ight)^{rac{n}{n-1}}$$

So,

$$M = 4\pi R_{\star}^{(3-n)/(1-n)} igg(rac{K(n+1)}{4\pi G}igg)^{n/(n-1)} ig\xi_{1}^{-(3-n)/(1-n)} ig\xi_{1}^{2} rac{d heta}{d\xi}igg|_{ig\xi_{1}}$$

#### Problem 6

Compute mass and radius of white dwarfs

From problem 4, we get that for n=3/2,  $\xi_1 pprox 3.655$ . Now given that

$$ho_c = 10^6 - 10^9 g/cm^3, K pprox 10^{13} \mu_e^{-5/3}, \mu_e = 2.$$

MassWhiteDwarf (generic function with 1 method)

```
    function MassWhiteDwarf(R, K, G, ξ1, ψ1, n)
    4 * pi * R^((3. -n)/(1. - n))*((K*(n+1))/(4*pi*G))^(n/(n-1))*ξ1^(-(3. - n)/(1. - n))*ξ1^2 * abs(ψ1)
    end
```

Rstar (generic function with 1 method)

```
function Rstar(G, ξ1, ρc, K, n)
return ξ1 * (4 * pi * G * ρc^(1.0 - (1.0 / n))/(K * (n + 1)))^(-1.0 / 2.0)
end
```

```
\xi 1 = 3.6539245994932976
```

•  $\xi 1 = data.\xi 1[2]$ 

```
\psi 1 = -0.20368099291868041
```

•  $\psi$ 1 = data. $\psi$ 1[2]

using UnitfulAstro

```
G = GM⊙ M⊙^-1

• G = u"GMsun"/u"Msun"
```

```
\mu = 2.0
• \mu = 2.0
```

```
n = 1.5
n = 3/2
```

```
K = 3.149802624737183e12 dyn cm<sup>3</sup> g<sup>-3752999689475413/2251799813685248</sup>
```

```
• K = 1e13 * \mu^{(-5.0/3.0)} * u''dyn/cm^2''/(u''g/cm^3'')^{(1+(1/n))}
ρcs =
▶ Unitful.Quantity{Float64,M L^-3,Unitful.FreeUnits{(g, cm^-3),M L^-3,nothing}}[1.0e6 g cm^
 • \rho cs = [1e6, 1e7, 1e8, 1e9] * u"g/cm^3"
1.0e7 g cm^-3
 ρcs[2]
r1 = 11196.014751530298 \text{ km}

    r1 = uconvert(u"km", Rstar(G, ξ1, ρcs[1], K, n))

mass1 = 2.922862216500595e26 dyn^3 cm^9 Mo^3 GMo^-3 g^-5 km^-3
 mass1 = MassWhiteDwarf(r1, K, G, ξ1, ψ1, n)
0.49440774095900636 Mo
 uconvert(u"Msun", mass1)
 begin
       masses = []
       radii = []
        for pc in pcs
            radius = uconvert(u"km", Rstar(G, ξ1, ρc, K, n))
mass = uconvert(u"Msun", MassWhiteDwarf(radius, K, G, ξ1, ψ1, n))
            push!(masses, mass)
            push!(radii, radius)
        end
 end
massRadius =
                                    central density
                                                        radius
                                                                       mass
                                                                   0.494408 Mo
                                    1.0e6 g cm^-3
                                                      11196.0 km
                                   1.0e7 g cm^-3
                                                      7627.76 km
                                                                   1.56345 M☉
```

1.0e8 g cm^-3

1.0e9 g cm^-3

massRadius = DataFrame(central\_density = ρcs, radius = radii, mass = masses)

5196.73 km

3540.49 km

4.94408 Mo

15.6345 M⊙