

Euler's Method, Shortcomings, and Higher Order Methods

MATH 333

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Slides and Code

Why Study Numerical Methods?

Khan: There's this lovely part in the movie in which your character turns to “old” math — **Euler's method** — to figure out how to get John Glenn back down from orbit. Did that really happen?

Johnson: It seemed logical to me. I could see in my mind what I needed and sort of worked backwards.

– 2017 LA Times interview of Katherine Johnson by Amina Khan.



Figure: Katherine Johnson in 1983. Image in Public Domain.

Outline

Review of Euler's method

Error in Euler's Method

Higher Order Methods

Summary

Euler's Method

Last time, we learned a technique to compute an approximate numerical solution to an **initial-value problem (IVP)** of the form

$$\begin{cases} y'(x) = f(x, y), & x \text{ in } [a, b] \\ y(a) = y_0 \end{cases} \quad (1)$$

We also learned that this technique can be applied to *systems* of ODEs of the form

$$\begin{cases} \mathbf{y}'(x) = \mathbf{f}(x, \mathbf{y}), & x \text{ in } [a, b] \\ \mathbf{y}(a) = \mathbf{y}_0 \end{cases} \quad (2)$$

Euler's Method

Consider a single IVP $y'(x) = f(x, y)$ with x in $[a, b]$ subject to the initial condition $y(a) = y_0$.

The idea behind the numerical methods we will cover is that instead of computing a solution $y(x)$ at *every* x in $[a, b]$, we select a *discrete* set of points

$$a = x_0 < x_1 < \cdots < x_{N-1} < x_N = b,$$

called **nodes**, and compute *approximate* values u_n of the solution $y(x_n)$ for $n = 0, \dots, N$. That is,

$$u_n \approx y(x_n)$$

for $n = 0, \dots, N$. We will often write $y_n = y(x_n)$. The spacing between nodes is called the **step-size**, and we will assume a uniform step-size of

$$h = \frac{b - a}{N}.$$

Euler's Method

There is one node at which we know the exact value of the solution: the left endpoint $a = x_0$. So $u_0 = y_0$. We may then compute $u_1 \approx y_1$ by Taylor expanding y_1 about x_0 :

$$y_1 = y_0 + y'_0 h + R(h)$$

Throwing away the remainder term, $R(h)$, yields an approximation $u_1 \approx y_1$:

$$u_1 = y_0 + y'_0 h \approx y_1.$$

Lastly, writing $y'_0 = f(x_0, y_0)$ and $u_0 = y_0$, we obtain the first step of Euler's method:

$$u_1 = u_0 + hf(x_0, u_0).$$

Euler's Method

Repeating this process at each node yields Euler's method:

$$u_{n+1} = u_n + hf_n,$$

where $f_n = f(x_n, u_n)$. This is also known as the **forward Euler method**.

Geometric View of Euler's Method

Geometrically, Euler's method at step n can be described as follows:

1. Construct the line, $L(t)$, parallel to the line tangent to y at x_n that passes through (x_n, u_n) .
Thus¹

$$L(t) = u_n + (t - x_n)f_n$$

2. Walk h units along the tangent line starting from x_n , i.e., set $t = x_n + h$, and take u_{n+1} to be the y -coordinate of the point you land on. In other words,

$$u_{n+1} = L(x_n + h) = u_n + hf_n.$$

¹Remember that, given a coordinate point (x, y) , the function $f(x, y)$ gives us the derivative of y at x .

Code Demonstration

Let's look at [exp_model.m](#) to see Euler's method in action. In this program, we aim to solve the IVP

$$\begin{cases} y'(x) = ky(x), & x \text{ in } [0, 6] \\ y(0) = 1. \end{cases} \quad (3)$$

Task 1: Use the Forward Euler Method to solve this IVP with $k = 1$. First use 10 nodes, and then use 30 nodes.

What do you observe?

Code Demonstration

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$$\begin{cases} y'(x) = ky(x), & x \text{ in } [0, 6] \\ y(0) = 1. \end{cases} \quad (4)$$

Task 2: Now use the Forward Euler Method to solve this IVP with $k = -3$. First use 9 nodes, then 10, and then 11.

What do you observe?

What About Systems?

Consider the following setup:

1. We attach a unit mass ($m = 1$) to a spring on a surface with no friction.
2. We allow the mass to reach equilibrium, we then pull it back 1 unit from this position and release.

If $x(t)$ denotes the position of the mass relative to its equilibrium at time $t \geq 0$ since we set the mass in motion, Newton's second law tells us that $x(t)$ satisfies

$$x''(t) + x(t) = 0 \tag{5}$$

with $x(0) = 1$ and $x'(0) = 0$.

What About Systems?

To use Euler's method, we first recast this **second order** differential equation as a **first order** system. Let's say $v = x'$. Then $x''(t) + x(t) = 0$ is equivalent to

$$\begin{cases} x'(t) = v(t), & x(0) = 1 \\ v'(t) = -x(t), & v(0) = 0. \end{cases} \quad (6)$$

Written in matrix-vector form:

$$\begin{bmatrix} x(t) \\ v(t) \end{bmatrix}' = \begin{bmatrix} 0 & 1 \\ -1 & 0 \end{bmatrix} \begin{bmatrix} x(t) \\ v(t) \end{bmatrix}, \quad \begin{bmatrix} x(0) \\ v(0) \end{bmatrix} = \begin{bmatrix} 1 \\ 0 \end{bmatrix} \quad (7)$$

We thus obtain an equation of the form (2) with $\mathbf{y}(t) = [x(t) \ v(t)]^T$ and $\mathbf{f}(\mathbf{x}, \mathbf{y}) = \begin{bmatrix} 0 & 1 \\ -1 & 0 \end{bmatrix} \mathbf{y}$.

What About Systems?

Let's jump into [harmonic_oscillator.m](#) to see Euler's method applied to the system (7).

Task 1: Use the Forward Euler Method to solve equation (7) on the interval $[0, 5]$ using 50, 100, and 200 nodes.

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What About Systems?

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Task 1: Use the Forward Euler Method to solve equation (7) on the interval $[0, 5]$ using 50, 100, and 200 nodes.

What do you observe?

Task 2: Use the Forward Euler Method to solve equation (7) on the interval $[0, 20]$ using 50, 100, and 200 nodes.

What do you observe?

Code demonstration observations

1. As we increase the number of nodes, or, equivalently, decrease the step-size h , Forward Euler appears to yield more accurate results.
2. When $[a, b]$ is "large," we seem to need a smaller step-size to obtain a reasonable solution over the entire interval.

Let's see if we can explain these observations theoretically.

Error in Euler's Method

To understand the error in Euler's method, let's sketch a picture. On your own paper, please do the following.

1. Sketch the function $y = e^x$.
2. Starting from the coordinate point $(0, 1)$, draw the first two steps of Euler's method. Using a larger step-size will be helpful for what comes next.

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1. Sketch the function $y = e^x$.
2. Starting from the coordinate point $(0, 1)$, draw the first two steps of Euler's method. Using a larger step-size will be helpful for what comes next.
3. Suppose that we know $y(x_1) = y_1$ exactly. Draw an Euler step from (x_1, y_1) .

Error in Euler's Method

To make our picture more formal, if we let $u_n^* = y_{n-1} + hf(x_{n-1}, y_{n-1})$, then we can write the **error** $E_n = y_n - u_n$ at step n as

$$E_n = (y_n - u_n^*) + (u_n^* - u_n) \quad (8)$$

1. $y_n - u_n^* = y_n - y_{n-1} - hf(x_{n-1}, y_{n-1})$ is the error due to **truncation**, i.e., cutting off the Taylor Expansion after the first derivative term.
2. $u_n^* - u_n = y_{n-1} + hf(x_{n-1}, y_{n-1}) - (u_{n-1} + hf(x_{n-1}, u_{n-1}))$ is the error accrued by using our approximate value u_{n-1} for the solution at x_{n-1} in place of the exact value y_{n-1} .

Error in Euler's Method

If we assume that $df(x, y)/dx = y''(x)$ is continuous, then the remainder term in the Taylor expansion

$$y_n = y_{n-1} + hf(x_{n-1}, y_{n-1}) + R(h)$$

may be bounded by

$$|R(h)| \leq Mh^2,$$

where M is a constant that is independent of h determined by the max value of $df(x, y(x))/dx$ in $[a, b]$.

So one part of the overall error, $y_n - u_n^*$, the **local truncation error**, can be bounded by

$$|y_n - u_n^*| \leq Mh^2.$$

The other part of the error, $u_n^* - u_n$, on the other hand, has a significant effect on the maximum error over $[a, b]$.

Error in Euler's Method

Indeed, one can show that the maximum error across all nodes x_0, x_1, \dots, x_N , which we denote by E_{\max} , is bounded by

$$|E_{\max}| \leq \left[\frac{M}{L} (e^{L(b-a)} - 1) \right] h \quad (9)$$

where

- ▶ M is the same constant from earlier determined by $df(x, y(x))/dx$, and
- ▶ $L = \max |\partial f(x, y)/\partial y|$ over a region containing our solution curve.

Both constants are independent of the step-size h .

Error in Euler's Method

Indeed, one can show that the maximum error across all nodes x_0, x_1, \dots, x_N , which we denote by E_{\max} , is bounded by

$$|E_{\max}| \leq \left[\frac{M}{L} (e^{L(b-a)} - 1) \right] h \quad (10)$$

This bound on the maximum error explains both observations we made earlier:

1. Since the term $M(e^{L(b-a)} - 1)/L$ is independent of h , $|E_{\max}|$ decreases as h decreases.
2. The term $e^{L(b-a)}$ significantly magnifies the error, with greater magnification the larger $b - a$ is.

Note: we may perform the same analysis and obtain an identical bound for the error in the case of systems. The only difference is we replace the absolute value $|\cdot|$ with a vector **norm** $\|\cdot\|$.

The Order of a Numerical Method

When analyzing numerical methods, we often obtain an expression of the form

$$|E_{\max}| \leq Ch^p$$

where C is some constant independent of h , and h is parameter we can adjust. In the the case of Euler's method, and the other numerical schemes we will see today, h is the step-size.

The number p is a positive integer, and if the error for a given method satisfies, $|E_{\max}| \leq Ch^p$, we say that the method is **order p** .

Based on our earlier analysis, we can therefore say that Euler's method is of order 1, or it's a **first order** method.

Numerically Verifying the Order

How can we “see” the order of convergence?

If we consolidate $M(e^{L(b-a)} - 1)/L$ into a single constant, C , then we have $|E_{\max}| \leq Ch$ for Euler's method. It follows that, if $\log(x)$ is the base-10 logarithm,

$$\log(|E_{\max}|) \leq \log(C) + \log(h).$$

This means that if we plot our error using the base-10 logarithm scale for each axis, we should see a line of slope 1.

Let's check out [error_rate.m](#) to see what we see.

Higher Order Methods

Our goal now is to study methods that are of a higher order than Euler's method.

Higher Order Methods

How does one develop a more accurate method?

Higher Order Methods

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Let's briefly revisit Euler's method from a different perspective. When solving a differential equation, we are ultimately undoing a derivative, i.e., we are integrating.

Consider the IVP $y'(x) = f(x, y)$ over $[a, b]$ subject to $y(a) = y_0$. We again select equally spaced nodes $a = x_0, \dots, x_N = b$ with step-size $h = (b - a)/N$. Starting from $u_0 = y_0$, we can obtain $u_1 \approx y_1$ by integrating our ODE and using the Fundamental Theorem of Calculus:

$$\begin{aligned}\int_{x_0}^{x_1} y'(x) dx &= \int_{x_0}^{x_1} f(x, y(x)) dx \\ y_1 - u_0 &= \int_{x_0}^{x_1} f(x, y(x)) dx.\end{aligned}$$

Higher Order Methods

If we approximate the integral of f using a **left-endpoint approximation**, we obtain the first step of Euler's method

$$y_1 \approx u_1 = u_0 + f(x_0, y(x_0))h.$$

This suggests one way to obtain a more accurate approximation: *use a more accurate integral approximation!*

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- ▶ The error E from using left and right-endpoint approximations satisfies

$$E \leq Ch^2,$$

where C depends on the function and h is the step-size.

- ▶ The **trapezoid rule**, on the other hand, has error $E_T \leq Ch^3$. This is better!
- ▶ In the homework, you will explore what happens if we use another integrator of similar accuracy to the trapezoid rule: the **midpoint rule**.

Higher Order Methods

Going back to

$$y_1 = u_0 + \int_{x_0}^{x_1} f(x, y(x)) dx,$$

approximating the integral with the Trapezoid rule yields

$$y_1 \approx u_1 = u_0 + \frac{h}{2} [f(x_0, y_0) + f(x_1, u_1)],$$

and for $n \geq 0$:

$$u_{n+1} = u_n + \frac{h}{2} [f(x_n, u_n) + f(x_{n+1}, u_{n+1})],$$

Higher Order Methods

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Higher Order Methods

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$$u_{n+1} = u_n + \frac{h}{2} [f(x_n, u_n) + f(x_{n+1}, u_{n+1})],$$

we may notice that:

1. We are using derivative information at *two* points, x_n and x_{n+1} , rather than at one point like we did with Euler's method.
2. The value we want to compute at step n , u_{n+1} , appears on *both* sides of the equation. That is, we have expressed u_{n+1} as an **implicit** function of u_n and x_n .

Order 2 Runge-Kutta

To obtain an expression for u_{n+1} as an **explicit** function of u_n and x_n , we may replace u_{n+1} with a Forward Euler step. This gives us the **Order 2 Runge-Kutta method**, or **RK2**:

$$u_{n+1} = u_n + \frac{h}{2} [f(x_n, u_n) + f(x_n + h, u_n + hf_n)] \quad (11)$$

We may also write step n of RK2 as

1. $s_1 = f(x_n, u_n)$
2. $s_2 = f(x_n + h, u_n + hs_1)$
3. $u_{n+1} = u_n + h \left(\frac{s_1 + s_2}{2} \right)$

Benefit of explicit methods: no need to solve a (possibly nonlinear) equation for u_{n+1} .

RK2 Error

The error in RK2 is very similar to that of Euler's method, except that we gain an extra h :

$$\max \text{RK2 error} \leq \left[\frac{M}{L} (e^{L(b-a)} - 1) \right] h^2. \quad (12)$$

- ✓ Order 2 in h , and therefore significantly more accurate than Euler.
- ✓ Simple to implement.
- ✓ Explicit—no equation to solve.
- ✗ Requires two evaluations of f . This could be expensive for some problems.
- ✗ Still have $e^{L(b-a)}$ in error bound, which means our solution will become unreliable for long intervals.

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Let's see how RK2 and Euler compare in practice with the harmonic oscillator!

Can we do even better?

The success of RK2 suggest that *more derivative information yields more accurate methods*.

This idea leads to a whole of family of methods. Each member of this family is characterized by the derivatives chosen, and how much weight we assign them.

In RK2, we use two derivatives: $s_1 = f(x_n, u_n)$ and $s_2 = f(x_{n+1}, u_n + hf_n)$, and gave them equal weight.

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What if we used *more* derivatives?

Order 4 Runge-Kutta

The **fourth order Runge-Kutta method**, **RK4**, or simply **the** Runge-Kutta method, is defined by the following process at each step n :

1. $s_1 = f(x_n, u_n)$
2. $s_2 = f(x_n + h/2, u_n + hs_1/2)$
3. $s_3 = f(x_n + h/2, u_n + hs_2/2)$
4. $s_4 = f(x_n + h, u_n + hs_3)$
5. $u_{n+1} = u_n + h \left(\frac{s_1 + 2s_2 + 2s_3 + s_4}{6} \right)$

As the name suggests, this is a *fourth order method*:

$$\max \text{RK4 error} \leq \left[\frac{M}{L} (e^{L(b-a)} - 1) \right] h^4$$

Order 4 Runge-Kutta

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- ✓ ✓ Order 4 in h , and therefore significantly more accurate than both Euler and RK2.
- ✓ Relatively simple to implement.
- ✓ Explicit—no equation to solve.
- ✗ Requires *four* evaluations of f . This could be expensive for some problems.
- ✗ Still have $e^{L(b-a)}$ in error bound, which means our solution will become unreliable for long intervals, but RK4 is accurate for much longer than RK2 or Euler.

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Order 4 Runge-Kutta

Task: Compare the solutions obtained from Forward Euler, RK2, and RK4 for the Harmonic Oscillator using $N = 50$ nodes.

Task: What happens if you increase the length of the interval, but keep $N = 50$? Try setting $b = 60$, and then $b = 100$.

Which method to use?

- ▶ In general, there is no single “best” numerical method for every problem you may encounter. Each method has its strengths and weaknesses.
- ▶ Of the methods discussed here, RK4 is by far the most accurate, but it requires additional function evaluations that may be too costly for some problems.
- ▶ Forward Euler and RK2 require fewer function evaluations, but are less accurate.
- ▶ RK2 remains accurate for longer than Euler with the same step-size.

What We Haven't Said

There are many, many numerical methods for ODEs not discussed here. For example,

- ▶ We have discussed the *forward* Euler's method, but there is also the **Backward** Euler method that is order one like forward Euler, but is much more *stable*.
- ▶ If in the derivation of RK2, we kept u_{k+1} in place of $u_n + hf_n$, we obtain what is known as the **Crank-Nicolson** method.
- ▶ ODE solver packages in MATLAB such as the popular ode23 and ode45 use combinations of Runge-Kutta methods to solve an ODE efficiently and guarantee that a user-specified tolerance is reached.

Note: Both Backward Euler and Crank-Nicolson are included in the [Github repository](#) if you would like to experiment with them.