

# TKU211121

## Classical Mechanics

### Chapter #1 : Rotation of Rigid Body

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I. Rigid Body and Center of Mass

II. Rotational Energy and Moment of Inertia

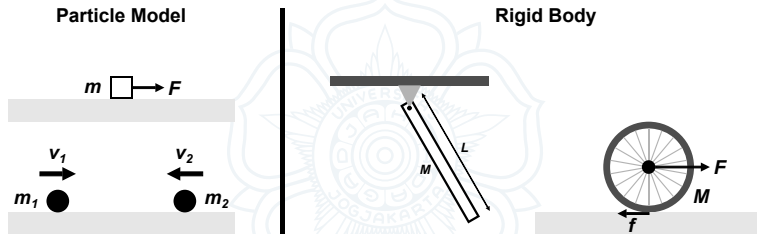
III. Torque and Static Equilibrium

IV. Rotational Dynamics and Angular Momentum

# Rigid Body and Center of Mass

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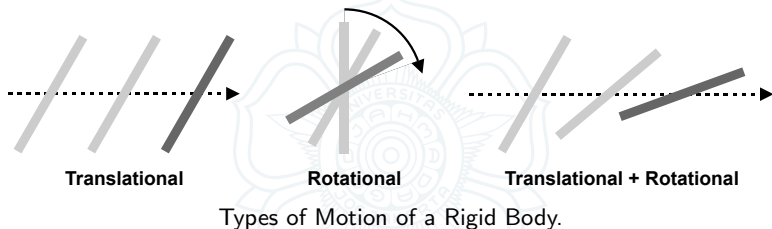
# Particle Model vs Rigid Body



Example of A System with Particle Model (left) and Rigid Body (right).

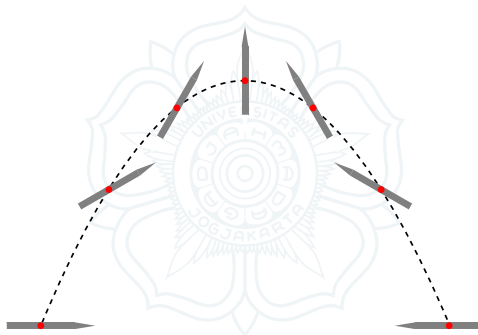
- **Particle Model** : An object is represented as a mass in a single point in space.
- **Rigid Body** : An extended object whose size and shape do not change as it moves.
- Example of Rigid Body : Wheel, Rod, Ball, Axle, etc.

# Particle Model vs Rigid Body



- In a particle model, the motion is very simple.
- Because it is modelled as a single point in space, a rotational motion cannot be defined.
- Only translational motion can be considered in a particle model.
- The types of motion in a rigid body is more complex.
- A rigid body can have a translational motion, a rotational motion, or the combination of both translational and rotational motion.

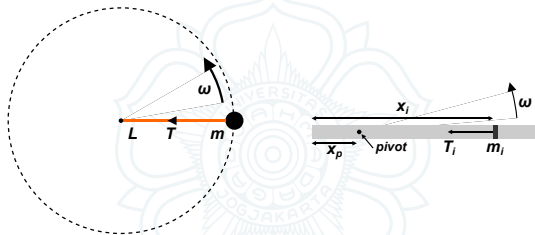
# Center of Mass



Motion of a Pencil thrown to the Air.

- A student throws a pencil to the air. He notices that the pencil is moving while rotating about a certain point. **About what point does it rotate?**
- An unconstrained object (i.e., one not on an axle or a pivot) tends to rotate about a point called the **Center of Mass**.

# Center of Mass

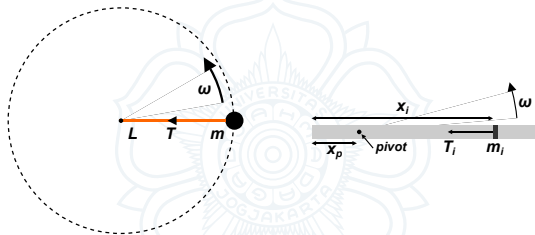


Calculation of Center of Mass Position.

- Where is the center of mass position of an object?
- First, let us consider a ball with mass  $m$  tied to a string with length  $L$ . The ball is rotating with constant angular velocity  $\omega$ .
- Because the ball is rotating, the ball is experiencing a centripetal acceleration

$$a_c = \omega^2 L$$

# Center of Mass



Calculation of Center of Mass Position.

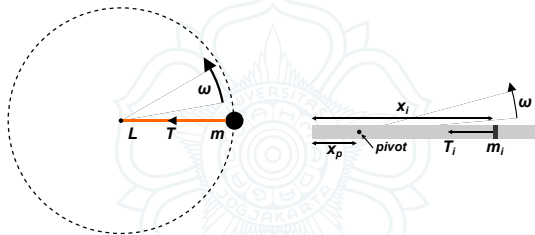
- Acceleration requires a force, which in this case is provided by the tension of the string  $T$  and can be calculated as follow

$$T = ma_c = m\omega^2 L$$

- Now let us consider a rigid body (a rod) is rotating with angular velocity  $\omega$  about a pivot point located at distance  $x_p$  from left end of the rod.



# Center of Mass

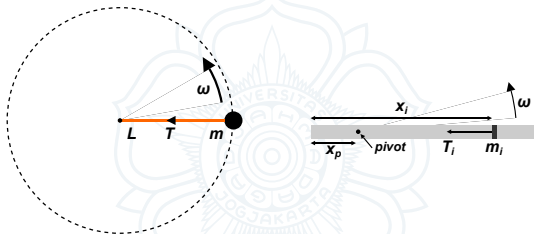


Calculation of Center of Mass Position.

- Consider a small portion of the rod with mass  $m_i$  located at distance  $\vec{x}_i$  from left end of the rod. It is experiencing a centripetal acceleration given by

$$\vec{a}_i = \omega^2(\vec{x}_p - \vec{x}_i)$$

# Center of Mass



Calculation of Center of Mass Position.

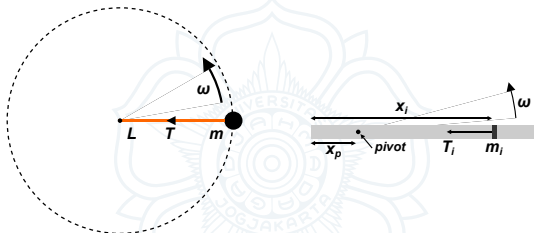
- Therefore the force  $\vec{T}_i$  for rotating this mass portion  $m_i$  is given by

$$\vec{T}_i = m_i \omega^2 (\vec{x}_p - \vec{x}_i)$$

- The sum of all tension forces is given by.

$$\sum_i \vec{T}_i = \sum_i m_i \omega^2 (\vec{x}_p - \vec{x}_i) = \omega^2 \sum_i m_i (\vec{x}_p - \vec{x}_i)$$

# Center of Mass



Calculation of Center of Mass Position.

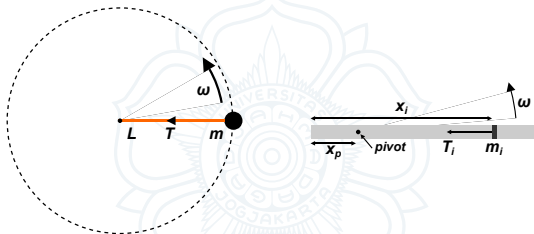
- When the rod is rotating about its center of mass ( $\vec{x}_p = \vec{x}_{cm}$ ), the sum of all tension forces are zero, therefore

$$\sum_i m_i (\vec{x}_{cm} - \vec{x}_i) = \vec{x}_{cm} \sum_i m_i - \sum_i m_i \vec{x}_i = 0$$

- As the result, the center of mass position is given by

$$\vec{x}_{cm} = \frac{\sum_i m_i \vec{x}_i}{\sum_i m_i} = \frac{m_1 \vec{x}_1 + m_2 \vec{x}_2 + m_3 \vec{x}_3 + \dots}{m_1 + m_2 + m_3 + \dots} \quad (1)$$

# Center of Mass



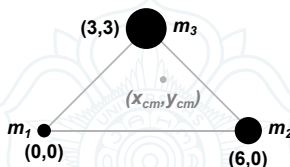
Calculation of Center of Mass Position.

- Using the same method, it can be proved also that the center of mass position in  $y$  and  $z$  directions are given by

$$y_{cm} = \frac{\sum_i m_i \vec{y}_i}{\sum_i m_i} = \frac{m_1 \vec{y}_1 + m_2 \vec{y}_2 + m_3 \vec{y}_3 + \dots}{m_1 + m_2 + m_3 + \dots} \quad (2)$$

$$z_{cm} = \frac{\sum_i m_i \vec{z}_i}{\sum_i m_i} = \frac{m_1 \vec{z}_1 + m_2 \vec{z}_2 + m_3 \vec{z}_3 + \dots}{m_1 + m_2 + m_3 + \dots} \quad (3)$$

# Center of Mass (Example)



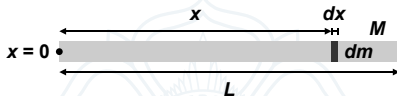
A Three-Bodies System.

- A system consists of three bodies with mass  $m_1 = 1\text{kg}$ ,  $m_2 = 2\text{kg}$  and  $m_3 = 3\text{kg}$  located at  $(0,0)$ ,  $(6,0)$  and  $(3,3)$ , respectively. Determine the  $x_{cm}$  and  $y_{cm}$  of this system.
- This problem can be solved easily by applying the  $x_{cm}$  and  $y_{cm}$  equations

$$x_{cm} = \frac{m_1 \vec{x}_1 + m_2 \vec{x}_2 + m_3 \vec{x}_3}{m_1 + m_2 + m_3} = \frac{(1)(0) + (2)(6) + (3)(3)}{1 + 2 + 3} = 3.5$$

$$y_{cm} = \frac{m_1 \vec{y}_1 + m_2 \vec{y}_2 + m_3 \vec{y}_3}{m_1 + m_2 + m_3} = \frac{(1)(0) + (2)(0) + (3)(3)}{1 + 2 + 3} = 1.5$$

# Center of Mass



The Center of Mass of a Rigid Object.

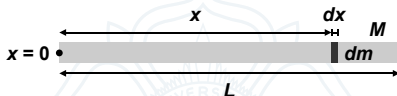
- Now let us consider a rigid object (rod) with mass  $M$  and length  $L$ . How to determine the center of mass position of this kind of object?
- First, we need to divide this object into many small cell with mass  $dm$ , where its position is  $x$ . The center of mass position of this object can be calculated as follow

$$x_{cm} = \frac{\int \vec{x} \cdot dm}{\int dm} \quad (4)$$

- Similarly, it can be proved also that the center of mass position in  $y$  and  $z$  directions are given by

$$y_{cm} = \frac{\int \vec{y} \cdot dm}{\int dm} \quad z_{cm} = \frac{\int \vec{z} \cdot dm}{\int dm} \quad (5)$$

# Center of Mass (Example)



The Center of Mass of a Rigid Object.

- As an example, we will try to determine the center of mass of a rod using the aforementioned equation
- First, we need to determine the  $dm$  of this rod

$$dm = \left( \frac{M}{L} \right) dx$$

- Therefore the center of mass position can be calculated as follow

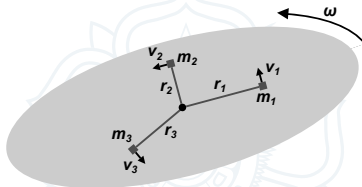
$$\vec{x}_{cm} = \frac{\int \vec{x} \cdot dm}{\int dm} = \frac{(M/L)(\int_0^L x dx)}{(M/L)(\int_0^L dx)} = \frac{L^2/2}{L} = \frac{L}{2}$$

# Rotational Energy and Moment of Inertia

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# Rotational Kinetic Energy

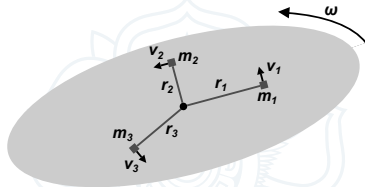


Rotational Kinetic Energy of a Rigid Body.

- A rotating rigid body has kinetic energy because all of the particles are in motion. The kinetic energy due to rotation is called a **Rotational Kinetic Energy**.
- All of the particles in a rigid body are rotating with the same angular velocity  $\omega$ . However, because the distance of each particle from the rotation axis are different, each has different velocity  $v_i$  which is given by

$$v_i = \omega r_i$$

# Rotational Kinetic Energy



Rotational Kinetic Energy of a Rigid Body.

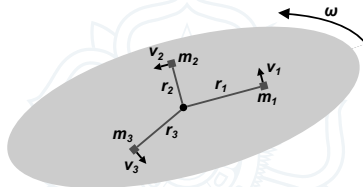
- The rotational kinetic energy of an object is the sum of all the particles' kinetic energy

$$\begin{aligned} K_{rot} &= \frac{m_1 v_1^2}{2} + \frac{m_2 v_2^2}{2} + \frac{m_3 v_3^2}{2} + \dots \\ &= \frac{m_1 r_1^2 \omega^2}{2} + \frac{m_2 r_2^2 \omega^2}{2} + \frac{m_3 r_3^2 \omega^2}{2} + \dots \end{aligned}$$

Therefore

$$K_{rot} = \left( \sum_i m_i r_i^2 \right) \frac{\omega^2}{2} = \frac{I \omega^2}{2} \quad (6)$$

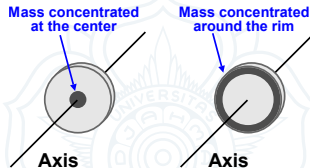
# Rotational Kinetic Energy



Rotational Kinetic Energy of a Rigid Body.

- Parameter  $I = \sum_i m_i r_i^2$  in the previous equation is defined as the object's **moment of inertia**.
- Notice that  $K_{rot} = I\omega^2/2$ , which indicates that for the same angular velocity  $\omega$ , the object with higher moment of inertia  $I$  carries higher rotational kinetic energy. Therefore, it requires more work to accelerate an object with higher moment of inertia.
- Therefore, this fact indicates that it is more difficult to accelerate an object with larger moment of inertia.

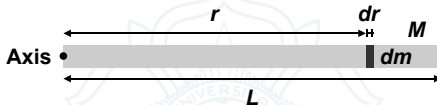
# Moment of Inertia



Effect of Mass Distribution on the Moment of Inertia.

- The *moment of inertia* in rotational motion indicates how difficult to *rotate* an object. It is analogous with *mass* (or *inertial mass*) in translational motion, which indicates how difficult to *move* an object.
- The main different between *moment of inertia* and *inertial mass* is that the *moment of inertia* of an object also depends on how the mass is distributed.
- Two wheels with the same mass  $M$  and radius  $R$ , it is easier to rotate the wheel whose mass is concentrated at the center rather than concentrated around the rim.

# Moment of Inertia

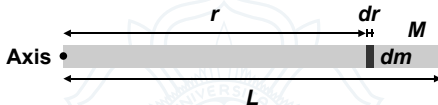


Moment of Inertia Calculation of a Rigid Object.

- Let us once again consider a rigid object (rod) with mass  $M$  and length  $L$ . How to determine the moment of inertia of this kind of object?
- First, we need to divide this object into many small cell with mass  $dm$ , where its position is  $r$  from the rotation axis.
- The moment of inertia of this object can be calculated as follow

$$I = \int r^2 \cdot dm \quad (7)$$

# Moment of Inertia (Example)



Moment of Inertia Calculation of a Rigid Object.

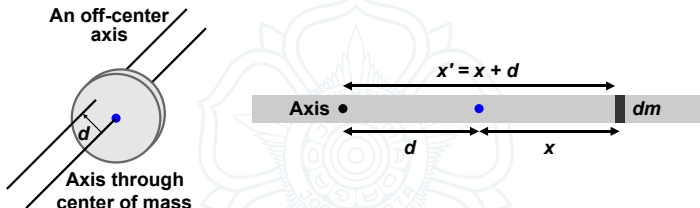
- As an example, we will try to determine the moment of inertia of a rod using the aforementioned equation
- First, we need to determine the  $dm$  of this rod

$$dm = \left(\frac{M}{L}\right) dr$$

- Therefore the moment of inertia can be calculated as follow

$$I = \int r^2 \cdot dm = \left(\frac{M}{L}\right) \left(\int_0^L r^2 dr\right) = \left(\frac{M}{L}\right) \left(\frac{L^3}{3}\right) = \frac{ML^2}{3}$$

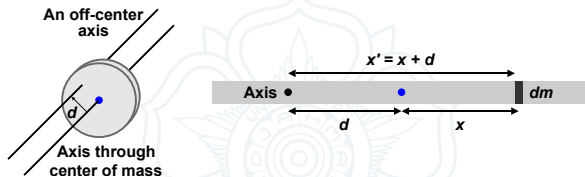
# Parallel Axis Theorem



Rotation about an Off-center Axis.

- A rigid body with mass  $M$  has moment of inertia  $I_{cm}$  when rotating about its center of mass. When the rotation axis is moved toward a new axis with distance  $d$  from its center of mass, what is the new moment of inertia of this object?
- Let us take an example of a rod, whose rotation axis is moved toward a new axis with distance  $d$  from its center of mass.

# Parallel Axis Theorem



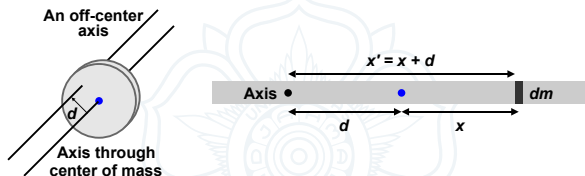
Rotation about an Off-center Axis.

- The moment of inertia of this case can be calculated as follow

$$\begin{aligned} I &= \int (x')^2 \cdot dm = \int (x^2 + 2xd + d^2) \cdot dm \\ &= \int (x^2 + 2d(x' - d) + d^2) \cdot dm \\ &= \left( \int x^2 \cdot dm \right) + 2d \left( \int x' \cdot dm - \int d \cdot dm \right) + d^2 \left( \int dm \right) \\ &= I_{cm} + 2d(x'_{cm} - d)M + d^2(M) \end{aligned}$$



# Parallel Axis Theorem



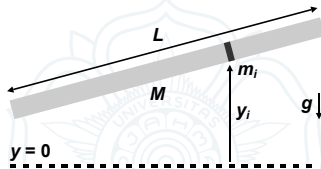
Rotation about an Off-center Axis.

- Note that  $x'_{cm}$  is the object's center of mass position assuming that the rotation axis is the center of coordinate (0,0,0). Therefore  $x'_{cm} = d$  and the moment of inertia about the new axis is given by

$$I = I_{cm} + Md^2 \quad (8)$$

- This equation is known as the **Parallel Axis Theorem**.

# Gravitational Potential Energy of a Rigid Body



Gravitational Potential Energy of a Rigid Body.

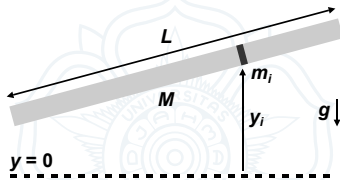
- How to calculate the gravitational potential energy of a rigid body?
- Let us once again consider a rod with mass  $M$  and length  $L$ . Assume that a particle with mass  $m_i$  is located at height  $y_i$ , therefore its gravitational potential energy is

$$U_i = m_i g y_i$$

- The total gravitational potential energy of the rod is the sum of the potential energy of all particles, therefore

$$U = \sum_i U_i = g \sum_i m_i y_i$$

# Gravitational Potential Energy of a Rigid Body



Gravitational Potential Energy of a Rigid Body.

- From the definition of center of mass in  $y$  direction, we know that

$$y_{cm} = \frac{\sum_i m_i y_i}{\sum_i m_i} = \frac{\sum_i m_i y_i}{M} \rightarrow \sum_i m_i y_i = M y_{cm}$$

- Therefore, the gravitational potential energy of a rigid body is given by

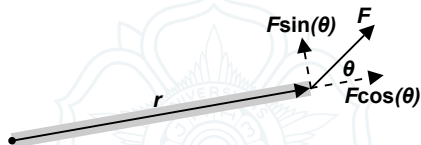
$$U = M g y_{cm} \quad (9)$$

which only depends on its total mass and the height of its center of mass.

# Torque and Static Equilibrium

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# Torque

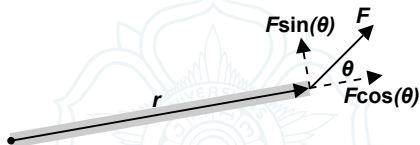


A Force  $\vec{F}$  acting on a Rigid Body.

- **Torque** is a quantity that measures the "effectiveness" of a force at causing an object to rotate about a pivot.
- Torque is the rotational equivalent of force.
- Now let us consider a force  $\vec{F}$  is working on an object at a distance  $\vec{r}$  from its axis/pivot. The *torque* working on this object is defined as

$$\vec{\tau} = \vec{r} \times \vec{F} \quad (10)$$

# Torque



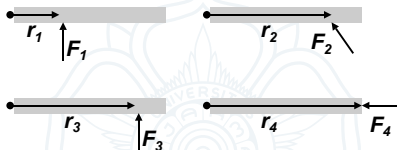
A Force  $\vec{F}$  acting on a Rigid Body.

- Remember that what really matter in a cross product between two vectors is the vector's component which is perpendicular to each other.
- Therefore, if there exists an angle  $\theta$  between  $\vec{r}$  and  $\vec{F}$ , the  $|\vec{\tau}|$  is defined as follow

$$|\vec{\tau}| = |\vec{r}| |\vec{F}| \sin \theta \quad (11)$$

which indicates that the produced torque is a multiplication between  $|\vec{r}|$  and the component of  $\vec{F}$  which is perpendicular to  $\vec{r}$ .

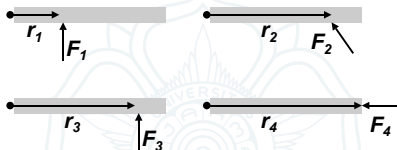
# Torque (Example)



Four Different Forces act on a Door.

- Let us consider the case when someone is pushing a door. There are 4 different forces  $\vec{F}_1$ ,  $\vec{F}_2$ ,  $\vec{F}_3$  and  $\vec{F}_4$  working on a door at 4 different positions and angles. If all forces have the same magnitude, which force produces the strongest torque? Which one produces the weakest torque?
- For  $\vec{F}_1$  and  $\vec{F}_3$ , the angle between  $\vec{r}$  and  $\vec{F}$  are  $\theta = 90^\circ$  (perpendicular). However, because  $|\vec{r}_1| < |\vec{r}_3|$ , therefore  $0 < |\vec{\tau}_1| < |\vec{\tau}_3|$ .

# Torque (Example)



Four Different Forces act on a Door.

- $\vec{F}_2$  is working at the same location as  $\vec{F}_3$  ( $\vec{r}_2 = \vec{r}_3$ ). However, because the angle  $\theta$  between  $\vec{r}_2$  and  $\vec{F}_2$  is between  $90^\circ$  and  $180^\circ$ , therefore  $0 < |\vec{\tau}_2| < |\vec{\tau}_3|$ .
- Finally, for  $\vec{F}_4$ , because the angle between  $\vec{r}_4$  and  $\vec{F}_4$  is  $180^\circ$ , therefore  $|\vec{\tau}_4| = 0$ .
- Therefore we can conclude that the strongest torque is produced by  $\vec{F}_3$ , and  $\vec{F}_4$  produces the weakest torque.



# Gravitational Torque



Gravitational Torque acting on a Rigid Body.

- Gravity exerts torque on many objects. Let us consider an object shown in the figure above. A torque due to gravity will cause it to rotate around the axle.
- Every particle inside an object contributes to the gravitational torque. Consider a particle with mass  $m_i$  which is located at the position  $(x_i, y_i)$  (assuming that the rotation axis is at coordinate  $(0,0)$ ). The torque produces by this particle is

$$\tau_i = -m_i g x_i$$

Note that the *negative sign* appears because the torque causes the object to rotate in *clockwise direction*.

# Gravitational Torque



Gravitational Torque acting on a Rigid Body.

- The total gravitational torque is the sum of all the contribution of each particles

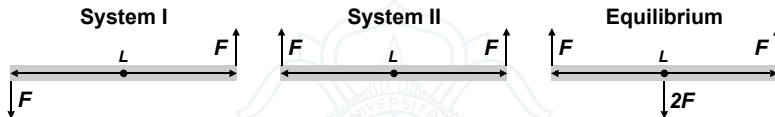
$$\tau_{grav} = - \sum_i m_i g x_i = -g \left( \sum_i m_i x_i \right)$$

- From the definition of center of mass,  $\sum_i m_i x_i = M x_{cm}$ , therefore

$$\tau_{grav} = -M g x_{cm} \quad (12)$$

- This equation indicates that the gravitational torque of an object is calculated by **treating the object as if all the mass is concentrated at its center of mass.**

# Static Equilibrium



Static Equilibrium of a Rigid Body.

- From Newton's 1<sup>st</sup> Law we know that a rest object will remain at rest when

$$\sum \vec{F} = 0 \rightarrow \sum \vec{F}_x = 0 \quad \sum \vec{F}_y = 0 \quad \sum \vec{F}_z = 0$$

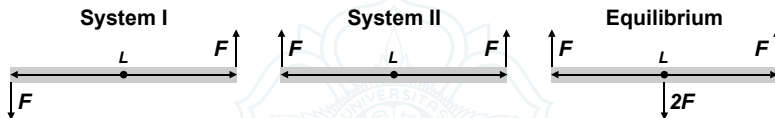
- As torque is a rotational equivalent of force, therefore a rest object (no rotation) will remain at rest (no rotation) when

$$\sum \vec{\tau} = 0 \rightarrow \sum \vec{\tau}_x = 0 \quad \sum \vec{\tau}_y = 0 \quad \sum \vec{\tau}_z = 0$$

- Therefore we can conclude that for an object to be in *static equilibrium* (no movement and no rotation) when

$$\sum \vec{F} = 0 \quad \text{and} \quad \sum \vec{\tau} = 0 \quad (13)$$

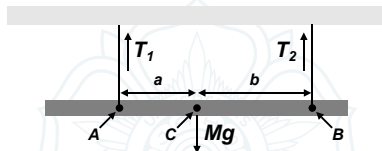
# Static Equilibrium



Static Equilibrium of a Rigid Body.

- Now let us consider the systems shown in the figure above.
- For the first system, we can see that  $\sum \vec{F} = 0$  and  $\sum \vec{\tau} = FL \neq 0$ . As the result, the rod is still able to rotate about its center of mass.
- On the other hand, for the second system, we can see that  $\sum \vec{F} = 2F \neq 0$  and  $\sum \vec{\tau} = 0$ . As the result, the rod has no rotation, but it is still able to have a translational motion.
- This example emphasizes that for achieving a *static equilibrium*, both  $\sum \vec{F} = 0$  and  $\sum \vec{\tau} = 0$  conditions have to be satisfied.

# Static Equilibrium (Example)



System under a Static Equilibrium Condition.

- Consider a system shown in the figure above. How much the tension feels by string 1 and string 2?
- From the equilibrium conditions  $\sum \vec{F} = 0$  and  $\sum \vec{\tau} = 0$

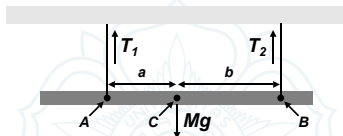
$$\sum F = T_1 + T_2 - Mg = 0 \quad \sum \tau = T_2 b - T_1 a = 0$$

where the torque is calculated assuming that the pivot point is at point C (center of mass).

- Solving those two equations we will get

$$T_1 = Mg \left( \frac{b}{a+b} \right) \quad T_2 = Mg \left( \frac{a}{a+b} \right)$$

# Static Equilibrium (Example)



System under a Static Equilibrium Condition.

- We can also solve this problem by analyzing the torque at point A and B.
- Let us analyze the  $\sum \vec{\tau} = 0$  condition by assuming that the pivot point is at point B, therefore

$$\sum \tau = Mgb - T_1(a + b) = 0 \rightarrow T_1 = Mg \left( \frac{b}{a + b} \right)$$

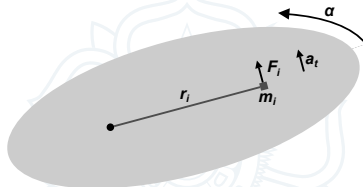
- Next, we analyze the  $\sum \vec{\tau} = 0$  condition by assuming that the pivot point is at point A, therefore

$$\sum \tau = Mga - T_2(a + b) = 0 \rightarrow T_2 = Mg \left( \frac{a}{a + b} \right)$$

# Rotational Dynamics and Angular Momentum

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# Rotational Dynamics



Effect of a Force on a Rigid Body.

- Consider an object is rotating freely about a pivot point as shown in the figure above. The forces  $F_i$  acting on a particle with mass  $m_i$  in the object. This force causes the particle to have an acceleration  $a_t$ , therefore

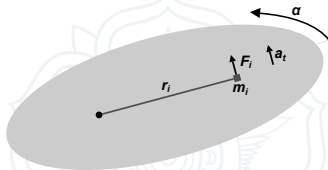
$$F_i = m_i a_t$$

- The relation between tangential acceleration  $a_t$  and angular acceleration  $\alpha$  is

$$a_t = \alpha r_i$$



# Rotational Dynamics



Effect of a Force on a Rigid Body.

- Therefore the  $F_i$  equation can be re-written as follow

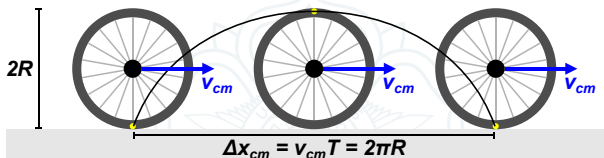
$$F_i = m_i \alpha r_i \quad \rightarrow \quad F_i r_i = m_i r_i^2 \alpha \quad \rightarrow \quad \tau_i = (m_i r_i^2) \alpha$$

- By summing all of the contribution of each individual particles, we get

$$\sum_i \tau_i = \left( \sum_i m_i r_i^2 \right) \alpha \quad \rightarrow \quad \sum \tau = I \alpha \quad (14)$$

which is known as *Newton's second law for rotational motion*

# Rolling Motion



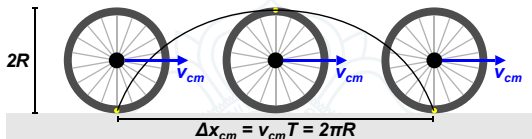
Rolling Motion without Slipping.

- **Rolling Motion** is a combination motion on which an object rotates about an axis that is moving along a straight-line trajectory.
- Consider a ball shown in the figure above is rolling without slipping. After the ball underwent a full rotation, the center of mass position had moved as far as

$$\Delta x_{cm} = v_{cm} T = 2\pi R \quad \rightarrow \quad v_{cm} = \left( \frac{2\pi}{T} \right) R$$

where  $T$  is the rotation period.

# Rolling Motion



Rolling Motion without Slipping.

- When the ball is rotating with constant angular velocity  $\omega$ , then  $\omega T = 2\pi$ . Therefore

$$v_{cm} = \omega R \quad (15)$$

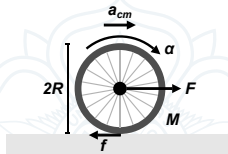
which is known as rolling constraint, the basic link between translation and rotation for an objects that roll without slipping.

- By differentiating both sides of the previous equation, we can derive the relation between translational acceleration  $a_{cm}$  and angular acceleration  $\alpha$  as follow

$$a_{cm} = \alpha R \quad (16)$$

- These 2 equations are the most important equations that are usually used when analyzing an object that rolls without slipping.

# Rolling Motion (Example)



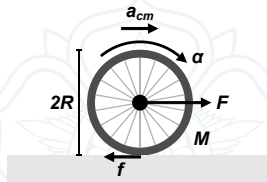
A Wheel is Rolling without Slipping.

- A force  $F$  is working at the center of mass of a wheel with mass  $M$ , radius  $R$  and moment of inertia  $I$ . The wheel is rolling without slipping on a rough surface with coefficient of friction  $\mu$ . Determine the wheel's translational acceleration ( $a_{cm}$ ).
- To solve this problem, first, we need to draw all the forces acting on the wheel. Then we need apply the Newton's second law for both translational and rotational motion

$$\sum F = Ma_{cm} \rightarrow F - f = Ma_{cm}$$

$$\sum \tau = I\alpha \rightarrow fR = I\alpha$$

# Rolling Motion (Example)



A Wheel is Rolling without Slipping.

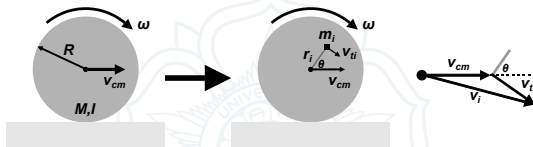
- It is important to note that because the wheel is rotating about its center of mass, therefore the torque is calculated by assuming the pivot point is at its center of mass.
- Finally, because the wheel is rolling without slipping, therefore

$$a_{cm} = \alpha R$$

- By solving those equations, we can get the  $a_{cm}$  equation as follow

$$a_{cm} = \frac{F}{M + (I/R^2)}$$

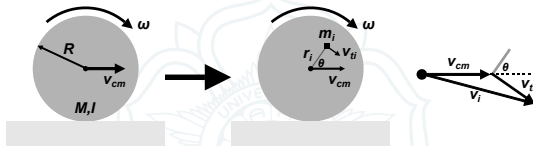
# Kinetic energy of a Rolling Object



Kinetic energy of a Rolling Object.

- Consider an object with mass  $M$  and moment of inertia  $I$ . Its center of mass is moving with speed  $v_{cm}$  while rotating about its center of mass with angular velocity  $\omega$ . How much is the kinetic energy of this object?
- Let us consider a particle in the object with mass  $m_i$ . In order to calculate the kinetic energy contribution of this particle, we need to calculate the speed of this particle relative to the ground ( $v_i$ ).

# Kinetic energy of a Rolling Object



Kinetic energy of a Rolling Object.

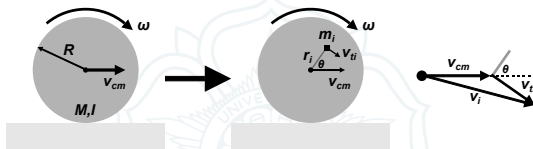
- The particle is moving with speed  $v_{ti} = \omega r_i$  relative to the center of mass, where the center of mass is moving with speed  $v_{cm}$  relative to the ground. Therefore

$$\vec{v}_i = \vec{v}_{cm} + \vec{v}_{ti}$$

- It can be shown that

$$\begin{aligned} v_i^2 &= (v_{cm} + \omega r_i \sin \theta)^2 + (\omega r_i \cos \theta)^2 \\ &= v_{cm}^2 + 2v_{cm}\omega(r_i \sin \theta) + \omega^2 r_i^2 \\ &= v_{cm}^2 + 2v_{cm}\omega y_i + \omega^2 r_i^2 \end{aligned}$$

# Kinetic energy of a Rolling Object



Kinetic energy of a Rolling Object.

- The total kinetic energy of this object can be calculated as follow

$$\begin{aligned}
 K &= \sum_i \frac{m_i v_i^2}{2} = \frac{v_{cm}^2}{2} \left( \sum_i m_i \right) + v_{cm} \omega \left( \sum_i m_i y_i \right) + \frac{\omega^2}{2} \left( \sum_i m_i r_i^2 \right) \\
 &= \frac{M v_{cm}^2}{2} + M v_{cm} \omega y_{cm} + \frac{I \omega^2}{2}
 \end{aligned}$$

- Because  $y_{cm} = 0$  (why?), therefore the total kinetic energy is given by

$$K = \frac{M v_{cm}^2}{2} + \frac{I \omega^2}{2} = K_{translation} + K_{rotation} \quad (17)$$



# Angular Momentum

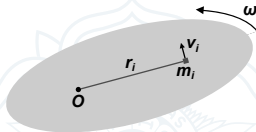


Angular Momentum of a Moving Particle.

- We have studied that a particle with mass  $m$  moving with velocity  $\vec{v}$  carries a momentum  $\vec{p} = m\vec{v}$ .
- Similar also for a particle rotating about a certain axis, it also carries an **Angular Momentum**.
- Consider a particle with linear momentum  $\vec{p}$  is moving at distance  $\vec{r}$  relative to a point O. The *angular momentum* of this particle relative to point O is defined as

$$\vec{L} = \vec{r} \times \vec{p} \quad (18)$$

# Angular Momentum



Angular Momentum of a Rigid Body.

- How about the angular momentum of a rigid body? To calculate the angular momentum of a rigid body, let us consider a rigid body shown in the figure above.
- A particle with mass  $m_i$  is located at a distance  $r_i$  from the rotation axis. Its angular momentum can be calculated as follow

$$L_i = m_i v_i r_i = m_i r_i^2 \omega$$

- The total angular momentum of this object can be calculated as follow

$$L = \sum L_i = \left( \sum_i m_i r_i^2 \right) \omega \rightarrow L = I \omega \quad (19)$$

# Angular Momentum and Torque

- From the definition of angular momentum  $\vec{L} = \vec{r} \times \vec{p}$ , let us differentiate both sides of this equation

$$\frac{d\vec{L}}{dt} = \vec{r} \times \frac{d\vec{p}}{dt}$$

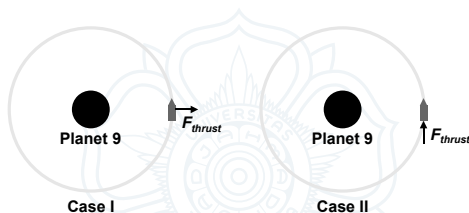
assuming that  $\vec{r}$  is constant. It is known that  $\frac{d\vec{p}}{dt} = \vec{F}$ . Therefore

$$\frac{d\vec{L}}{dt} = \vec{r} \times \vec{F} \rightarrow \frac{d\vec{L}}{dt} = \vec{\tau} \quad (20)$$

- This equation shows that when a torque  $\vec{\tau}$  is acting on an object, it will change its angular momentum  $\vec{L}$ .
- On the other hand, when there is no torque acting on an object, its angular momentum is constant (conserved), i.e.

$$\vec{L}_{initial} = \vec{L}_{final} \quad \text{when} \quad \sum \tau = 0 \quad (21)$$

# Conservation of Angular Momentum (Example)



A Space Shuttle Orbits a Planet.

- Consider the case of a space shuttle orbiting a planet. In case I, a thrust  $F_{thrust}$  is working on the space shuttle in a radially outward direction. In case II, the thrust is working in the tangential direction. In which case is the angular momentum conserved?
- For case A, because  $F_{thrust}$  is working in radially outward direction, therefore the net torque is zero. Therefore its angular momentum is conserved.
- For case B, the force  $F_{thrust}$  causes a torque working on the space shuttle. Therefore, its angular momentum is *not* conserved.