The Coarse Scattering Method

Pablo A. Vaquer

22 October 2018

../graphic/TAMUlogo.png

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In the coarse scattering (CS) method, we make the following substitution to the MG transport equation to reduce the size of the scattering matrices

$$\sum_{g'} \Sigma_{s,\ell,g' \to g} \phi_{\ell,g'} \quad \to \quad S_{\ell,e \to g} \sum_{e'} \Sigma_{s,\ell,e' \to e} \phi_{\ell,e'}$$

where each fine-group g is a subset of a coarse-element e and

$$\begin{split} \mathcal{S}_{\ell,e \to g} &= \frac{\sum_{g'} \sum_{s,\ell,g' \to g} \phi_{\ell,g'}}{\sum_{e'} \sum_{s,\ell,e' \to e} \phi_{\ell,e'}} \\ \sum_{s,\ell,e' \to e} &= \sum_{g \in e} \sum_{g' \in e'} \sum_{s,\ell,g' \to g} \\ \phi_{\ell,e} &= \sum_{g \in e} \phi_{\ell,g} \end{split}$$

The CS method can also be applied to the fission matrix

For fission, we make the following substitution

$$\sum_{g'} \Sigma_{f,g' \to g} \phi_{g'} \quad \to \quad F_{e \to g} \sum_{e'} \Sigma_{f,e' \to e} \phi_{e'}$$

where

$$F_{e \to g} = \frac{\sum_{g'} \sum_{f,g' \to g} \phi_{g'}}{\sum_{e'} \sum_{f,e' \to e} \phi_{e'}}$$

$$\sum_{f,e' \to e} = \sum_{g \in e} \sum_{g' \in e'} \sum_{f,g' \to g} \phi_{f}$$

$$\phi_{e} = \sum_{g \in e} \phi_{g}$$

Example of fine transfer matrix being decomposed into a coarse transfer matrix and mapping operator

 $../{\tt graphic/matrix_decomposition.png}$

Recomputing the scattering and fission spectra

ullet Not necessary to recompute both $S_{\ell,e o g}$ and $F_{e o g}$ every iteration

$$\begin{split} \left[\frac{1}{v_g}\frac{\partial}{\partial t} + \mu \frac{\partial}{\partial x} + \Sigma_{t,g}(\vec{r},t)\right] \psi_g^{(i+1)}(\vec{r},\mu,t) &= q_g(\vec{r},\hat{\Omega},t) + \\ S_{\ell,e\to g}^{(i-s)} \sum_{\ell=0}^{L} \frac{2\ell+1}{2} \sum_{e'} \Sigma_{s,\ell,e'\to e}(\vec{r},t) \phi_{\ell,e'}^{(i)}(\vec{r},t) + \\ \frac{F_{e\to g}^{(i-f)}}{2} \sum_{e'} \Sigma_{f,e'\to e}(\vec{r},t) \phi_{e'}^{(i)m}(\vec{r},t) \end{split}$$

- Asymptotic computational cost of solving the MG transport equation is order G^2
- Asymptotic computational cost of the CS method is order G in iterations when $S_{\ell,e\to g}$ and $F_{e\to g}$ are not recomputed and order G^2 only in iterations where $S_{\ell,e\to g}$ and $F_{e\to g}$ are recomputed

Particle balance can be maintained

Particle balance is maintained if following residual ρ is less than machine epsilon in the last iteration,

$$\rho = \left[\sum_{g'} \left(\Sigma_{s,0,g' \to g} + \Sigma_{f,g' \to g} \right) \phi_{g'} \right] - \left[S_{0,e \to g} \sum_{e'} \Sigma_{s,0,e' \to e} \phi_{e'} + F_{e \to g} \sum_{e'} \Sigma_{f,e' \to e} \phi_{e'} \right]$$

Futhermore, it may be possible to perserve particle balance every iteration if the CS method is only applied to Legendre moments greater than $\boldsymbol{0}$

Preliminary CS Results for a 1D slab of uranium (20% enriched)

| Simulation | MG | CS | CS | CS |
|--|----------------------|----------------------|----------------------|----------------------|
| Number of Groups | 200 | 200 / 25 | 200 / 5 | 200 / 1 |
| Recomputations $(S_{e	o g}\ /\ F_{e	o g})$ | | 2 / 4 | 2 / 4 | 2 / 4 |
| k _{eff} | 1.15420 | 1.15420 | 1.15420 | 1.15420 |
| RHS Cost | 3.26×10^{8} | 1.89×10^{8} | 2.03×10^{8} | $2.34{	imes}10^{8}$ |
| Total Cost | 3.32×10^{8} | 1.95×10^{8} | 2.11×10^{8} | 2.44×10^{8} |

Properties of Spherical Harmonics

$$\begin{split} &\int_{4\pi} d\Omega \, \frac{2\ell+1}{4\pi} Y_\ell^m(\hat{\Omega}) Y_{\ell'}^{m'}(\hat{\Omega}) = \delta_{\ell\ell'} \delta_{mm'} \\ &\int_{4\pi} d\Omega \, \frac{2\ell+1}{4\pi} Y_\ell^m(\hat{\Omega}) Y_0^0(\hat{\Omega}) = \delta_{\ell0} \delta_{m0} \\ &\int_{4\pi} d\Omega \, \sum_{\ell=0}^L \sum_{m=-\ell}^\ell \frac{2\ell+1}{4\pi} Y_\ell^m(\hat{\Omega}) Y_0^0(\hat{\Omega}) = 1 \end{split}$$

Particle balanced is maintain if coarse scattering is only used for higher moments

$$\int_{4\pi} d\Omega \, Y_0^0(\hat{\Omega}) \sum_{\ell=0}^L \sum_{m=-\ell}^\ell \frac{2\ell+1}{4\pi} Y_\ell^m(\hat{\Omega}) \sum_{g'}^G \Sigma_{s,\ell,g'\to g} \phi_{\ell,g'}^m = \Sigma_{s,0,g'\to g} \phi_{0,g'}^m$$

$$\begin{split} \int_{4\pi} d\Omega \, Y_0^0(\hat{\Omega}) \Big\{ \frac{1}{4\pi} Y_0^0(\hat{\Omega}) \sum_{g'}^G \Sigma_{s,0,g' \to g} \phi_{0,g'} \, + \\ \sum_{\ell=1}^L \sum_{m=-\ell}^\ell \frac{2\ell+1}{4\pi} Y_\ell^m(\hat{\Omega}) S_{e \to g} \sum_{e'}^E \Sigma_{s,\ell,e' \to e} \phi_{\ell,e'}^m \Big\} &= \Sigma_{s,0,g' \to g} \phi_{0,g'} \end{split}$$

Thank you

Questions?