

UNIVERSITY OF CALIFORNIA SAN DIEGO

**A Search for New Physics Using the Stransverse Mass Variable in
All-Hadronic Final States Produced in Proton-Proton Collisions With a
Center of Mass Energy of 13 TeV**

A dissertation submitted in partial satisfaction of the
requirements for the degree
Doctor of Philosophy

in

Physics

by

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2018

DEDICATION

To my family, who has always supported me in all my pursuits.

EPIGRAPH

For knowledge comes slowly, and when it comes, it is often at great personal expense.

— Paul Auster, *Ghosts*

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In Chapter 3, figures from the CMS Technical Design Report and Particle Flow

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PUBLICATIONS

Search for new physics with the M_{T2} variable in all-jets final states produced in pp collisions at $\sqrt{s} = 13$ TeV, CMS Collaboration, J. High Energ. Phys. **10 (2016) 006, arXiv:1603.04053 [hep-ex]**

Search for new physics in the one soft lepton final state using 2015 data at $\sqrt{s} = 13$ TeV, CMS Collaboration, Physics Analysis Summary (2016), CMS-PAS-SUS-16-011, cds.cern.ch/record/2161097

Search for new phenomena with the M_{T2} variable in the all-hadronic final state produced in proton-proton collisions at $\sqrt{s} = 13$ TeV, CMS Collaboration, Eur. Phys. J. C **77 (2017) no. 10, 710, arXiv:1705.04650 [hep-ex]**

ABSTRACT OF THE DISSERTATION

A Search for New Physics Using the Stransverse Mass Variable in All-Hadronic Final States Produced in Proton-Proton Collisions With a Center of Mass Energy of 13 TeV

by

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Professor Frank Würthwein, Chair
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A search for physics beyond the Standard Model (SM) is performed in events with final states including hadronic activity, missing energy, and significant momentum imbalance as measured with the M_{T2} variable. The results are based on data collected by the Compact Muon Solenoid detector at the Large Hadron Collider, and correspond to a total integrated luminosity of 35.9 fb^{-1} of proton-proton collisions at a center-of-mass energy of 13 TeV. No significant excess above the predicted SM background is observed. The results

are interpreted as 95% confidence-level exclusion limits on the masses of hypothesized particles in a variety of simplified models of R -parity conserving supersymmetry (SUSY). Additional techniques for extending the search to target final states with low-momentum leptons is discussed, interpreted in the context of SUSY models with compressed mass spectra, and compared to the exclusion strength of a typical all-hadronic search targeting such models.

Chapter 1

Introduction

This dissertation documents a search for new physics in proton-proton collisions at the Large Hadron Collider, with data collected by the Compact Muon Solenoid at a center-of-mass energy of 13 TeV. The state of particle physics and motivations for such a search are summarized in Chapter 2, including the current understanding of the Standard Model of particle physics and theorized extensions of the Standard Model which motivate these physics searches.

Chapter 3 describes the experimental apparatus used to collect the data analyzed in this analysis. Section 3.1 describes the physics of the Large Hadron Collider which delivers protons to the collision point, and section 3.2 details the different subsystems of the Compact Muon Solenoid detector which collect particle data from collisions. The data are reconstructed into abstract physics objects as described in section 3.3, which are subsequently analyzed in this search for new physics.

Chapters 4, 5, and 6 describe the analysis design and execution. In chapter 4, the

overall analysis strategy and signal definitions are discussed, and chapter 5 describes the background predictions in detail. Chapter 6 presents the results of the search, including interpretations constraining parameters of new physics models.

Finally, chapter 7 presents an extension of the all-hadronic search that targets specific signatures motivated by new physics models and limitations of the typical all-hadronic analysis. A pilot analysis of this type was presented in [22].

The all-hadronic analyses detailed in this dissertation were published in [23] and [30].

Chapter 2

Theory and Motivation

Particle physics is concerned with the study of the most fundamental constituents of nature and the rules that govern their interactions. Over the course of the twentieth century, physicists developed models using quantum field theory to describe the fundamental forces binding elementary particles together, and experimentalists discovered an abundance of particles predicted by such theories. Today, the Standard Model of particle physics represents the best experimentally verified theoretical framework for describing the elementary components of the universe, and experimental particle physicists work in tandem with theorists to identify possible extensions to the Standard Model which may explain some theoretical and experimental issues with the current framework.

2.1 The Standard Model of Particle Physics

The Standard Model (SM) is a quantum field theory which describes the fundamental particles which make up the universe and the various interactions between those particles [36, 52, 47]. It is widely regarded as one of the most successful theories ever constructed to describe nature, and has been rigorously verified through experiment. The announcement in 2012 by the ATLAS [6] and CMS [19] detectors of a new particle consistent with the Higgs Boson marked the discovery of the last fundamental particle predicted by the SM over 50 years before its discovery, and some precision measurements of Standard model parameters have been verified to better than one part in a billion.

The SM classifies all the elementary particles into two categories, depending on their intrinsic angular momentum, *spin*. Particles with a half-integer spin are known as *fermions* while those with integer spin are referred to as *bosons*. Fermions obey the Pauli exclusion principle (where identical fermions cannot occupy the same quantum state), and are often thought of as the constituents of matter. Bosons are not subject to the Pauli exclusion principle, and are often thought of as the “force carriers” which mediate different fundamental forces between different particles. The various fundamental particles of the SM and their properties are detailed in figure 2.1.

Fermions can be divided into two groups, *leptons* and *quarks*. Both leptons and quarks can be further divided into two subgroups based on their electromagnetic charge. Charged leptons carry either positive or negative charge e (where e is the fundamental charge constant), whereas the neutral leptons known as *neutrinos* carry no electromag-

netic charge. Different species of quarks can be either positively charged ($+2/3e$) or negatively charged ($-1/3e$). Both leptons and quarks can also be divided into three “flavors” of particles grouped into different families with similar electromagnetic properties but different masses. The leptons in order of increasing mass are *electron*, *muon*, or *tau* flavored (with both charged lepton and neutrino species), and the quarks are grouped into *up/down*, *charmed/strange*, and *top/bottom* pairs (that are positively and negatively charged, respectively).

The bosons of the standard model are associated with the different fundamental forces by which particles (including the bosons themselves) interact. The Standard Model Lagrangian is invariant under local transformations of an $SU(3) \times SU(2) \times U(1)$ gauge symmetry, and the various gauge symmetries correspond to different forces through which particles interact. $SU(3)$ is associated with the Quantum Chromodynamics (QCD) color charge (or strong force), and mediated by the *gluon* which binds quarks together (leptons are represented by an $SU(3)$ singlet and do not interact via the strong force). The $SU(2) \times U(1)$ groups are associated with weak isospin and hypercharge respectively, and responsible for electroweak interactions between particles, mediated by the W^\pm and Z bosons and the *photon*, γ .

While the gauge symmetries of the standard model are preserved, explicit mass terms are forbidden in the Lagrangian and the associated gauge bosons might otherwise remain massless. However, the SM Lagrangian also includes a complex scalar doublet which allows for the Higgs mechanism; the neutral component of the scalar doublet acquires a non-zero vacuum expectation value which manifests as the *Higgs boson*, and the other

degrees of freedom become Goldstone bosons which grant mass to the W^\pm and Z bosons. In this manner the electroweak $SU(2) \times U(1)$ is spontaneously broken into the familiar weak and electromagnetic forces.

2.2 Issues with the Standard Model

Despite the great success of the Standard Model in describing many of the fundamental interactions of particles, it is insufficient to explain all physical phenomena. There are several outstanding theoretical questions left unanswered by the Standard Model, and experimental evidence to suggest it is incomplete.

One of the most perplexing theoretical features of the SM is the fact that we have been able to detect all the elementary particles at energies accessible in our experiments. Because the bare masses of particles in the SM Lagrangian receive corrections from quantum loop diagrams, it is not clear why other particles (including undiscovered particles, which could be very massive) do not contribute significantly to these corrections. In particular, the quantum loop corrections to the Higgs boson mass include all massive particles, and contributions from massive particles could drive the Higgs mass far beyond what is currently accessible, yet this is not observed. This is often referred to as the *hierarchy problem*.

Experimental evidence also suggests there are physical phenomena not explained by the SM. Measurements of the rotational velocity of galaxies compared to the visible matter indicate there is an abundance of *dark matter* in the universe which cannot be

Standard Model of Elementary Particles

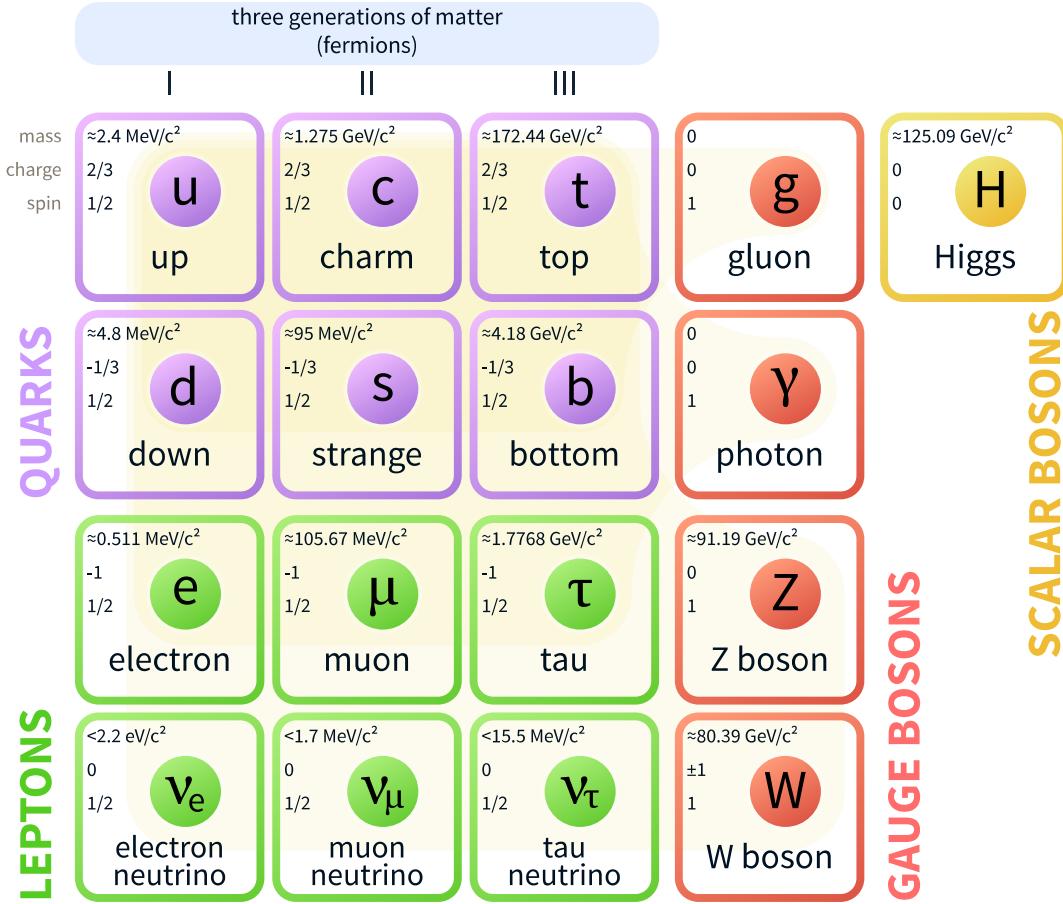


Figure 2.1: An illustration of the elementary particles in the Standard Model and their properties. The three left columns represent the different generations of fermions, with quarks in purple and leptons in green. The gauge bosons are represented in the fourth column in red, and the only scalar boson, the Higgs, in the fifth column in yellow. [33]

accounted for by the particle content of the SM [31]. Furthermore, while the SM predicts neutrinos to be massless, measurements of neutrino flavor oscillations imply that neutrino masses are very small, but nonzero [1]. Such measured phenomena indicate there may be additional particle content beyond what is posited by the SM, and additional interaction terms between SM particles and a “dark sector” of weakly-interacting particles.

2.3 Beyond the Standard Model: Supersymmetry

In order to solve many of the theoretical and experimental issues with the SM, theorists had considered extending SM particle content and interactions to include additional species which might satisfy some of the theoretical constraints of the SM in a consistent manner with observation. However, in 1975 the Haag-Łopuszański-Sohnius theorem [38] demonstrated that the only non-trivial extensions of quantum theories do not only include internal symmetries and the Poincaré symmetry, but also a non-trivial extension of the Poincaré algebra known as *supersymmetry* [45, 42, 37, 43, 35, 54, 53, 51].

A quantum field theory of supersymmetry (SUSY) can be visualized as a dual version of a typical field theory. For each particle, there exists a superpartner with different spin; fermions have boson-like superpartners and bosons have fermion-like superpartners. This is a particularly appealing and elegant solution to many of the issues with the SM; loop diagram corrections to particle mass are partially cancelled out by contributions from superpartner loop diagrams, providing a natural solution to the hierarchy problem [41]. In typical SUSY theories which conserve R-parity, particle interactions also conserve “SUSY-

ness” in decays, and thus when a superparticle decays into SM particles, the cascade will always end in the lightest supersymmetric particle (LSP). The LSP must be stable, and because it interacts very weakly with the SM sector provides a suitable dark matter candidate.

While there are many beyond-SM (BSM) theories that seek to explain phenomena beyond the scope of the SM, SUSY provides both elegant solutions to the theoretical concerns of the SM as well as a robust framework to unexplained physical phenomena. The analysis described here is designed to be model independent in a search for new physics, but also sets constraints on SUSY models as it is one of the more promising BSM theories that might be realized by nature.

Chapter 3

The Large Hadron Collider and the CMS Detector

In order to probe the most fundamental constituents of nature, accelerators ramp particles up to great energies and then collide them to probe particle interactions. Detectors situated at the collision point can measure the outgoing products from particle collisions, and physicists can study the interactions to measure SM parameters, determine particle properties, quantify interaction rates, or search for new particles and processes.

The Large Hadron Collider is the most powerful particle accelerator ever constructed, and collides protons at a rate of nearly one billion times per second at a record energy of 13 Teraelectronvolts. The Compact Muon Solenoid is one of the all-purpose physics detectors that monitors the particles produced in proton-proton collisions delivered by the LHC. In addition to the ATLAS detector (A Toroidal LHC ApparatuS), the CMS detector is responsible for collecting data to measure a wide range of physics phe-

nomena.

3.1 The Large Hadron Collider

The Large Hadron Collider (LHC) is the world's largest and most powerful particle accelerator ever constructed [34]. The accelerator was constructed and is operated by the European Organization for Nuclear Research (CERN), and straddles the border between Switzerland and France near Geneva. The accelerator tunnels of the LHC are 27 kilometers in circumferences and 50-175 meters underground, as well as additional smaller rings used to accelerate protons before they are injected into the LHC. An aerial view of the LHC layout and location of the different experiments is shown in figure 3.1.

The protons delivered by the LHC are accelerated in stages before reaching final collision energy. First, protons are accelerated to 50 MeV by a linear accelerator before being injected into the Proton Synchrotron Booster (PSB). The PSB accelerates protons to 1.4 GeV before they are injected into the Proton Synchrotron (PS) and accelerated to 26 GeV, then injected into the Super Proton Synchrotron (PSP) and accelerated to 450 GeV. Finally the protons are injected into the LHC and directed in two parallel beamlines in opposite directions to a beam energy of 6.5 TeV. A schematic representation of the different boost rings is illustrated in figure 3.2.

The total amount of proton-proton collisions delivered by the LHC is often measured by the *integrated luminosity*. The *instantaneous luminosity* (L) is a measure of the number of interactions that can be produced per unit area per second, and the integrated luminosity



Figure 3.1: An aerial view of CERN and the LHC complex. The colored lines indicate the location of the underground tunnels of the LHC and other booster rings used to accelerate protons. [33]

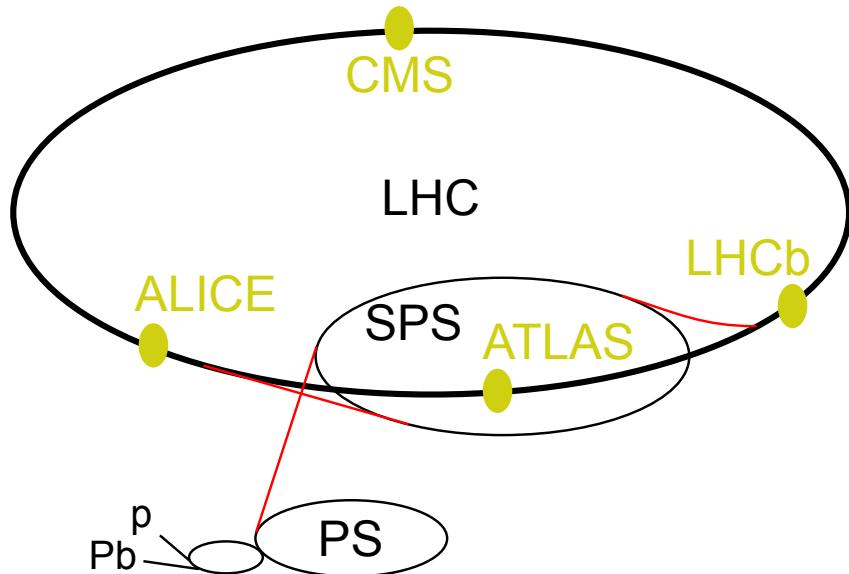


Figure 3.2: A schematic of the different accelerator rings used by the LHC. Protons are accelerated in a series of booster rings – the Proton Synchrotron Booster (PSB, not labeled), Proton Synchrotron (PS), and Super Proton Synchrotron (SPS) – before they are injected into the LHC and ramped up to their collision energy. Particles can be directed to collide at any of the interaction points along the ring where different detectors are situated. [33]

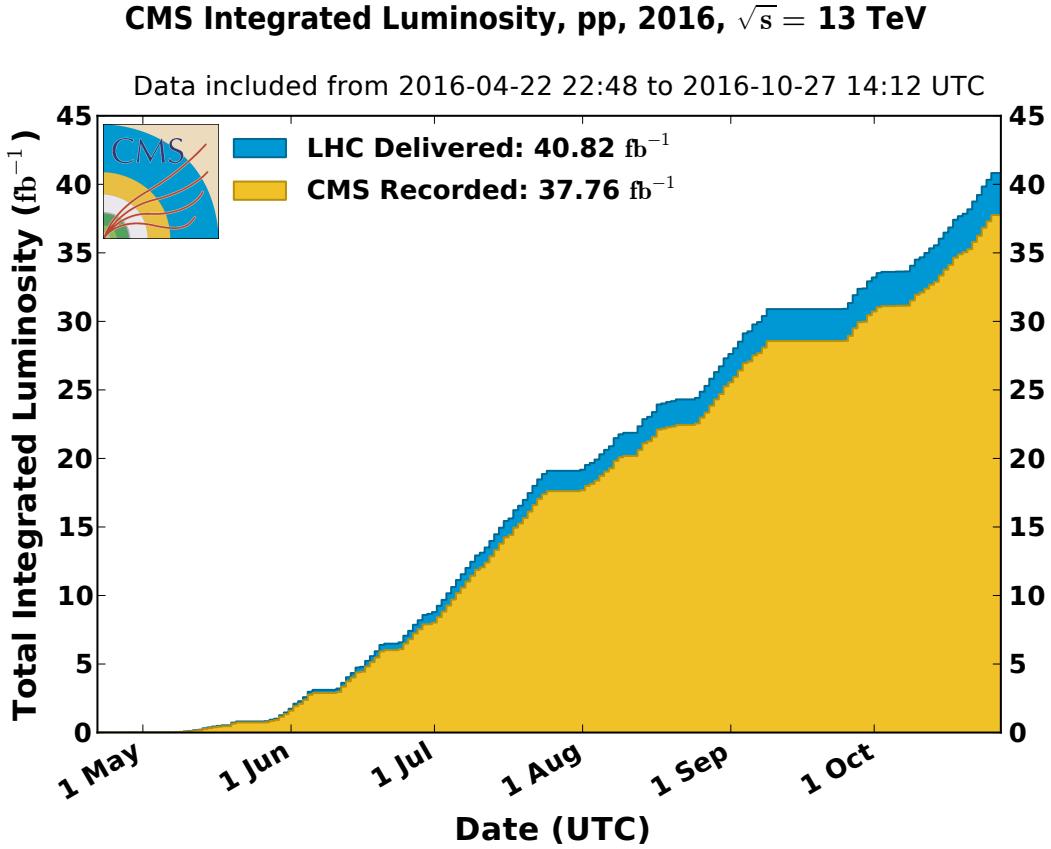


Figure 3.3: The total integrated luminosity delivered by the LHC at a center-of-mass energy of 13 TeV compared to that recorded by CMS through 2016.

(\mathcal{L}) is the total luminosity integrated over some time interval. Integrated luminosity is thus a unit of inverse area and typically measured in barns (1 barn = 100 fm²). The total number of a particular type of event measured by a detector N_{events} is the product of the integrated luminosity, the cross section for that particular process (σ), and the detector-determined efficiency (ϵ) as described in equation 3.1. The total integrated luminosity delivered by the LHC compared to that recorded by CMS is illustrated in figure 3.3.

$$N_{\text{events}} = \mathcal{L} \cdot \sigma \cdot \epsilon \quad (3.1)$$

3.1.1 Proton Bunches and Pileup

The protons delivered by the LHC are not in a continuous stream, but separated into “bunches” so that interactions occur at discrete intervals. Each proton bunch is centimeters long and contains over 100 billion protons at the beginning of the machine fill. Proton bunches are injected into the LHC and stored with a bunch spacing of 25ns (or 40 MHz), where they can be accelerated and circulated for hours while collisions occur at the interaction points.

Because protons are delivered in discrete bunches at such a high rate, any bunch crossing results in multiple interactions and outgoing particles. This is known as *pileup*, and presents a significant challenge for physics detectors which must be able to associate outgoing particles to a specific interaction. *In-time pileup* is caused by multiple interactions per bunch crossing (the majority of interaction in a given bunch crossing are soft proton-proton inelastic scattering), and *out-of-time pileup* is due to interactions from a different bunch crossing. The latter effect is due to the readout time of some of the detector electronics being longer than the 25ns bunch spacing. A distribution illustrating the mean number of interactions per bunch crossing is shown in figure 3.4. An effective detector needs to correctly identify which particles are associated with the interaction of interest in any given event, and mitigate the effects of pileup in design and construction.

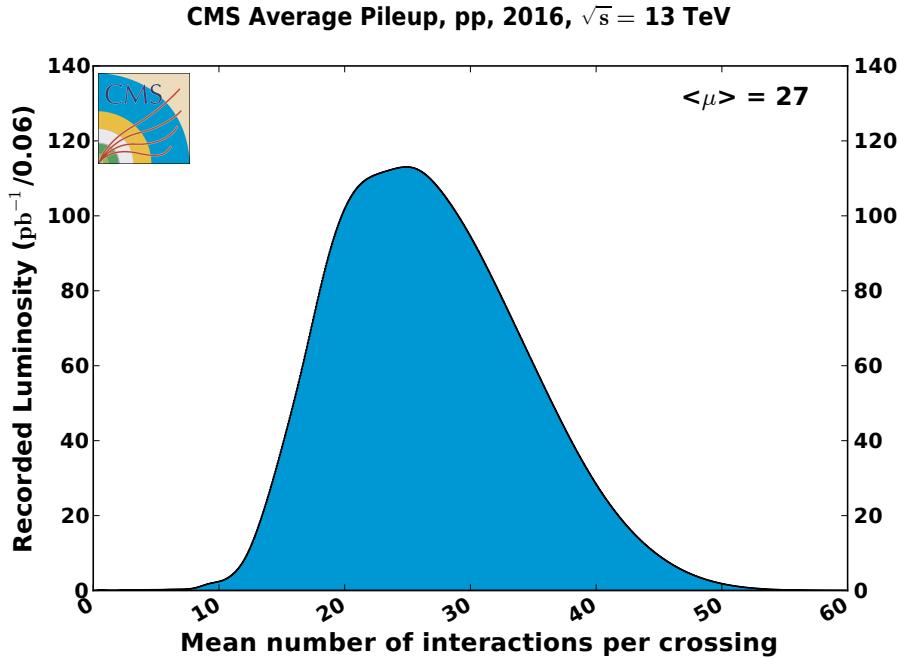


Figure 3.4: The mean number of interactions per bunch crossing for the 2016 run at 13 TeV. The cross section is taken to be 80 mb.

3.2 The Compact Muon Solenoid

The Compact Muon Solenoid (CMS) is a general-purpose physics detector at the LHC, situated at one of the five collision points along the main ring [15, 12]. The detector encapsulates the collision point with layers of various subsystems designed to interact with the outgoing particles of the proton-proton collisions, and measure the position and energies of the collision products. Because of the extremely high rate of interactions at the collision point (on the order of one billion interactions per second), saving data from every bunch crossing would be unsustainable, and so the detector is also equipped with a system of hardware and software implemented *triggers* which identify events of interest for physics analyses to be saved to disk for further analysis.

The physical construction of the detector is motivated by the different interaction of particles with different types of materials, and consists of several subsystems layered as coaxial cylinders around the interaction point. Each subsystem consists of different components covering the fiducial area coaxial with the beamline (the *barrel*) and also the ends of the cylinder (the *endcap*). The innermost subsystem of CMS is the silicon tracker, which consists of many layered silicon pixels and strips designed to pinpoint the locations of charged particles while minimally interacting with the particle's trajectory. The layer beyond the tracker is the electromagnetic calorimeter (ECAL), a grid of lead-tungstate crystals which scintillate to measure the energies of electromagnetic particles. Beyond the ECAL is the hadronic calorimeter (HCAL), a sampling calorimeter designed to measure the energies of hadronic particles (which deposit minimal energy in the ECAL). The final, outer layer is the CMS muon detector, where the muon detection stations are interweaved with the magnetic return yoke that generates the toroidal 3.8T magnetic field inside the detector volume. The total dimensions of the detector are 21.6m long and 14.6m in diameter, weighing over 12,500 tons.

3.2.1 Silicon Tracker

The silicon vertex tracker (SVT) is a series of silicon pixels and strips designed to measure the position of charged particles in the detector, while disturbing their path as little as possible. The position of particles in the interior is of particular importance in event reconstruction; charged particles traveling in a magnetic field will deflect in a curved path with radius proportional to the particles momentum as described in equation 3.2,

and so the track reconstructed can be used to not only determine a particle's momentum with high precision, but also the sign of its charge based on the direction of curvature.

$$p = qrB \quad (3.2)$$

As the innermost detector subsystem, the SVT experiences the highest flux of particle radiation. In the barrel region, the tracker layers are oriented in 3 coaxial layers. Closest to the interaction point where particle flux is the greatest, very precise silicon pixels are used, measuring $100 \times 150\mu\text{m}^2$, whereas in other layers the flux is low enough to use microstrip detectors, measuring $10\text{cm} \times 80\mu\text{m}$ and $25\text{cm} \times 180\mu\text{m}$ in the middle and outer layers respectively. In the endcaps, the pixel strips are arranged in a turbine-like pattern in two separate layers on each end. This configuration allows for the precise measurement of particle position for track reconstruction, while minimizing the amount of material which might deflect particles from their original trajectories. A partial geometry of the pixel layers in the SVT is shown in figures 3.5 and 3.6.

3.2.2 Electromagnetic Calorimeter

The ECAL is used to measure the energies of particles which interact electromagnetically, both absorbing the incident particles and scintillating to provide an energy-readout to photodiodes attached to each crystal. Constructed of lead-tungstate (PbWO_4), electromagnetically interacting particles (such as electrons or photons) will interact with the crystal material, losing energy through a cascade of electromagnetic interactions including electron-positron pair production and bremsstrahlung as shown in figure 3.7. This phe-

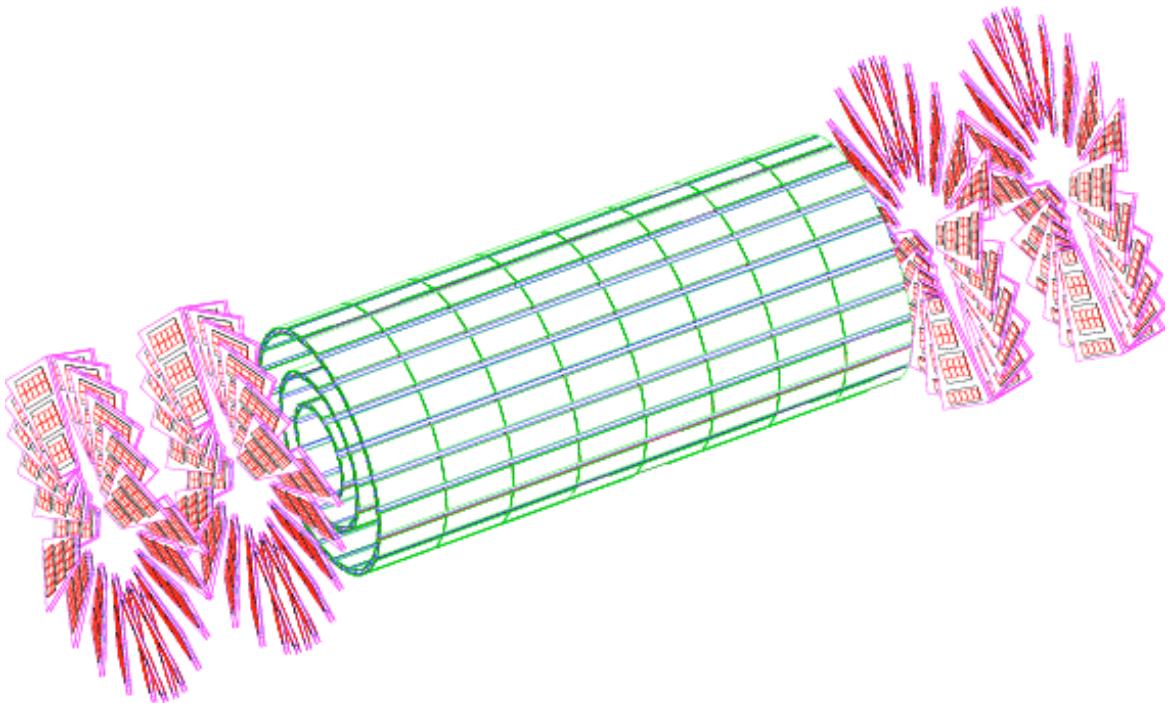


Figure 3.5: Geometry of pixel tracker inner layers in CMS.

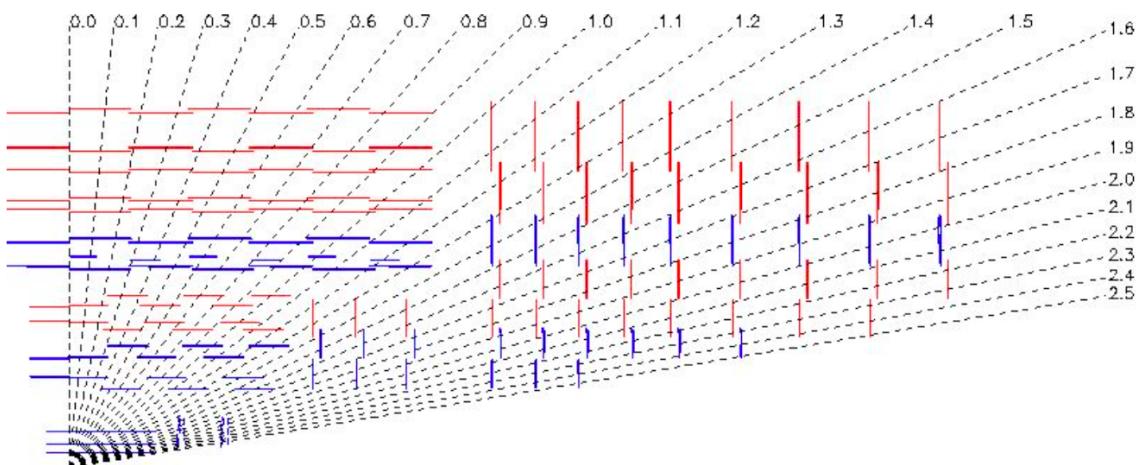


Figure 3.6: One quadrant of the tracker layout in the transverse plane. The dashed lines indicate the pseudorapidity coverage of different layers.

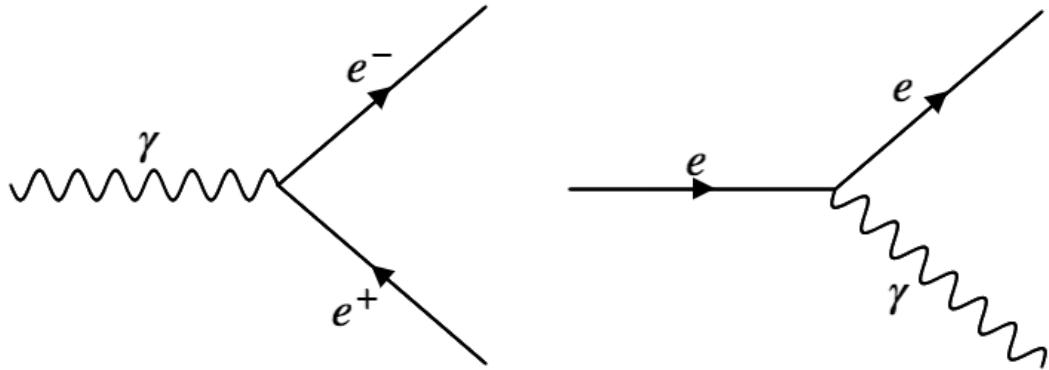


Figure 3.7: Feynman diagrams depicting some processes causing showers in the ECAL. The left diagram illustrates electron-positron pair production from a photon, and the right diagram depicts bremsstrahlung, where an electron radiates energy away through a photon.

nomenon — also referred to as "showering" — causes the crystals to scintillate proportional to the energy deposited in the crystal, which is then measured by various photodiodes to extract an accurate measurement of the particle energy, now fully absorbed by the calorimeter.

The fundamental principle of the calorimeter measurement relies on the energy loss of particles interacting with matter. In general, the energy of a particle traveling a distance X through some material is given by equation 3.3, where E_0 is the initial energy of the particle and X_0 is the material-dependent radiation length.

$$E(x) = E_0 e^{\frac{-x}{X_0}} \quad (3.3)$$

The design of the calorimeter is motivated by the choice of a scintillating, radiation-hard material with short X_0 such that incident electromagnetic particles deposit all their energy and are stopped by the ECAL. The resolution of the energy measurement is also dependent on the "stochastic term", which parametrizes the uncertainty due to statistical

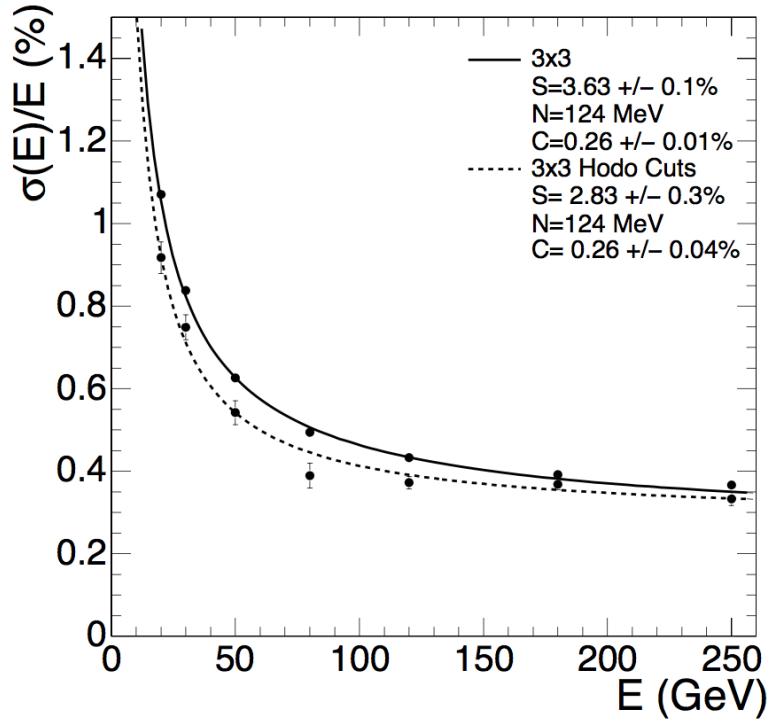


Figure 3.8: Energy resolution σ/E of the ECAL as a function of electron energy measured using a test beam. The energy was measured in a 3×3 crystal array with electrons incident on the center crystal, with electrons falling in a $4 \times 4\text{mm}^2$ region (lower points) and $20 \times 20\text{mm}^2$ region (upper points).

and measurement fluctuations in the calorimeter, and is given by equation 3.4, where S is the stochastic term, N the noise, and C the constant term.

$$\left(\frac{\sigma}{E}\right)^2 = \left(\frac{S}{\sqrt{E}}\right)^2 + \left(\frac{N}{E}\right)^2 + C^2 \quad (3.4)$$

The energy resolution can be measured by a test beam of known energy, as shown in figure 3.8.

The construction of the calorimeter is also divided into two sections by the cylindrical geometry, the ECAL barrel section (EB) and ECAL endcap sections (EE). The EB consists of 61,200 crystals arranged into 36 "supermodules", each spanning half the barrel length, and uses silicon avalanche photodiodes (APDs) as photodetectors. The individual

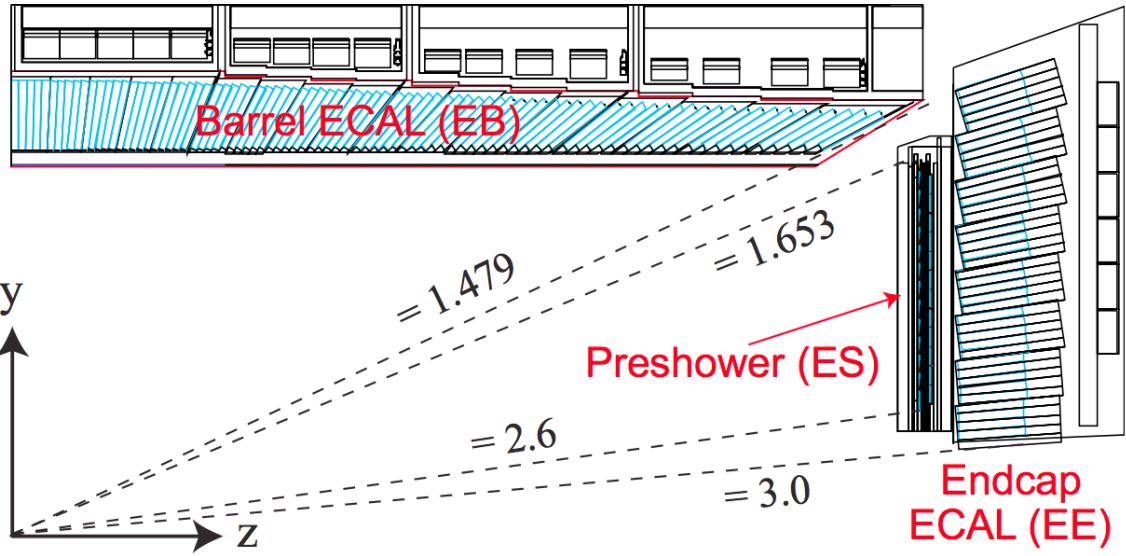


Figure 3.9: A cross section of the ECAL geometry, with the dashed lines marking the pseudorapidity values η covered by the various subsystems.

crystals are tilted slightly (3°) in an $\eta - \phi$ grid with respect to the nominal interaction point, with a front-facing area of $22 \times 22\text{mm}^2$ and a length of 230mm. The EE instead uses vacuum phototriodes (VPTs) as photodetectors, and consists of approximately 15,000 crystals clustered in 5×5 units, also offset from the interaction point but arranged in an $x - y$ grid, with a cross section of $28.6 \times 28.6\text{mm}^2$ and a length of 220mm. The EE is also equipped with a “preshower” device placed in front of the crystal calorimeter, consisting of two strips of silicon strip detectors to enhance π^0 rejection. The layout of the ECAL can be seen in figure 3.9. Because of the depth of the ECAL crystals (which are $\sim 25X_0$), and the confining properties of the crystals (which have a Moliere radius of 2.2cm, the radius of a cylinder containing 90% of a shower’s energy on average), electrons and photons are typically easily identified and well reconstructed in CMS, except in the transition region where EB and EE meet.

3.2.3 Hadronic Calorimeter

The CMS HCAL is a sampling calorimeter. Designed with alternating layers of scintillating and absorbing material, incident hadronic particles (such as charged pions, kaons, protons, etc.) interact with the absorber material and consequently shower into electromagnetic particles, whose energy can be read out by photodiodes connected to the scintillating material. Brass is used as the absorber material for both its interaction length and non-magnetic properties, and plastic scintillator tiles connected to embedded wavelength-shifting fibers carry the light to a readout system.

As with the ECAL, the energy loss of hadronic particles in the absorber is characterized by the (hadronic) interaction length and equation 3.3. However, unlike the ECAL, the HCAL contains both hadronic and electromagnetic showers. Electromagnetic particles generated in hadronic showers often fail to escape the absorber layers, and thus some electromagnetic energy is lost in the absorbers. The CMS HCAL is sometimes referred to as a *non-compensating calorimeter* because it is not constructed to actively compensate for the energy lost to these electromagnetic effects and the energy measurements must be corrected offline, known as "jet energy corrections" (JECs). JECs are typically calculated by examining data from collisions producing a boson recoiling against hadronic particles. By accurately measuring the boson energy in the ECAL, the sum of the recoiling hadronic energy in the HCAL is inferred (by momentum conservation) and compared to the detector response. Because the performance of the HCAL can fluctuate with time and run conditions, JECs are regularly recalculated and applied to the raw energy measurements

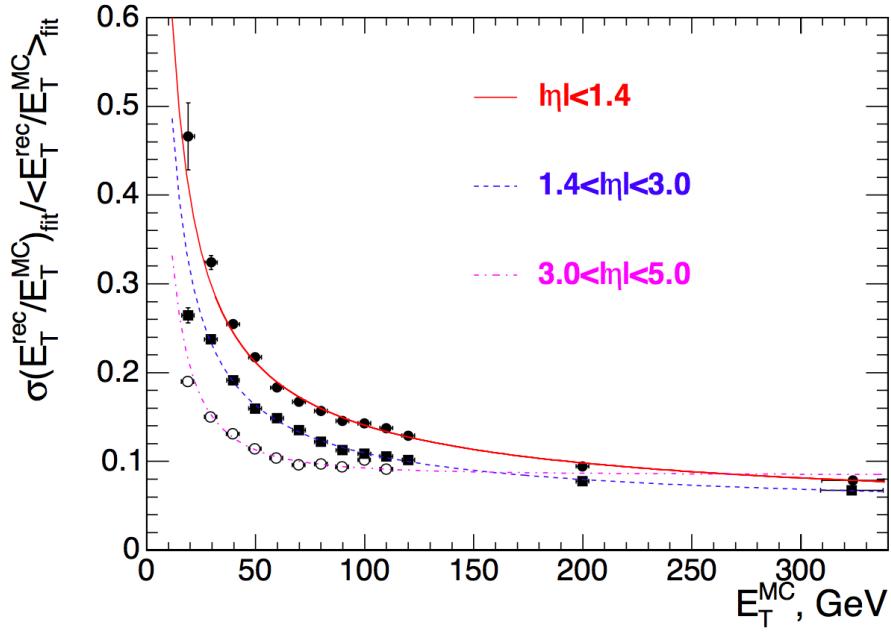


Figure 3.10: The jet transverse energy resolution as a function of the jet transverse energy, in different regions of pseudorapidity. For an explanation of jets, see section 3.3.6

taken by the HCAL to compensate for these effects. Additional information on the corrections contained in JECs is detailed in section 3.3.6. The energy resolution of the HCAL in different regions of pseudorapidity can be seen in figure 3.10.

The geometry of the HCAL can be reduced to four sections. The HCAL barrel (HB) consist of 32 “towers” of alternating absorber/scintillator material spanning the pseudorapidity region $-1.4 < \eta < 1.4$. The HCAL outer (HO) detector lies outside the vacuum tank of the magnetic coil and measures energy from any hadronic particles “leaking” through the HB, covering the pseudorapidity region $-1.26 < \eta < 1.26$. The HCAL endcap (HE) consists of 14 η towers spanning the pseudorapidity range $1.3 < |\eta| < 3.0$. Finally, the HCAL forward (HF) is a different steel/quartz fibre calorimeter spanning the very-forward $3.0 < |\eta| < 5.0$ region. Utilizing Cherenkov radiation generated in the quartz

fibers, the HF preferentially samples neutral hadronic energy and is ideally designed for the hadronic-heavy radiation environment in the forward region.

3.2.4 Muon Detectors

The muon detector is the only subsystem which is constructed outside of the toroidal magnetic field. Interleaved with the magnetic return yoke, different muon detectors are used to aid in the identification and reconstruction of muon tracks. Because muons typically penetrate every other layer of the detector and have a large bending radius, additional measurements in the muon system — combined with measurements in the SVT — can lead to improved resolution of muons with large transverse momentum. The reversed direction of the magnetic field in the muon detector region (beyond the magnetic toroid) causes an s-shaped trajectory and tighter bending radius than in the SVT for the muons, which improves the resolution for particles with transverse momentum above $\sim 200\text{GeV}$ as seen in figure 3.11.

Different types of muon detectors are deployed, sometimes in combination, in different sections of the full muon system. In the barrel region ($-1.2 < \eta < 1.2$) where the muon flux, neutral background, and magnetic field are small, drift tube (DT) chambers are used. In the endcaps where the the muon flux, neutral background, and magnetic field are much greater, cathode strip chambers (CSCs) are used to increase coverage up to $|\eta| < 2.4$. In addition to these technologies, both the barrel and endcap systems are supplemented with resistive plate chambers (RPCs) to provide complementary information to the DT and CSC detectors. The layout of the different muon detector components in the barrel

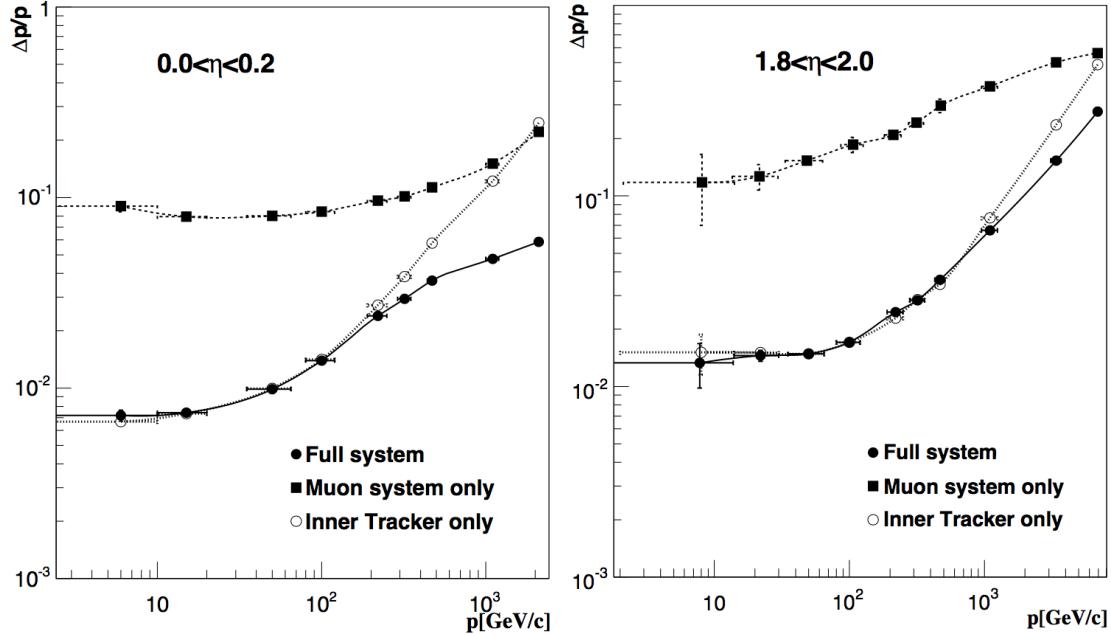


Figure 3.11: Muon momentum resolution as a function of muon momentum using only the inner tracking system, only the muon system, or both combined in the barrel region (left) or endcap region (right).

and endcap can be seen in figure 3.12.

The DT detectors are chambers filled with gas surrounding a wire, with a voltage difference between the wire and outside of the DT. When charged particles pass through the drift chamber, they ionize the gas, and the ionization products will drift across the voltage difference in the tube, resulting in a detectable voltage change in the DT. By measuring both the position along the DT wire where charge is deposited, as well as reconstructing the “drift time” it takes for the ionized particles to reach the wire, DTs provide a 2-dimensional measurement of a particle’s position. In CMS, the DT chambers consist of a dozen layers arranged into 3 groups, each with up to 60 DTs. The layers are arranged in such a way that some measure the direction of the muon parallel to the proton beam and others along the perpendicular coordinate, such that the full muon trajectory

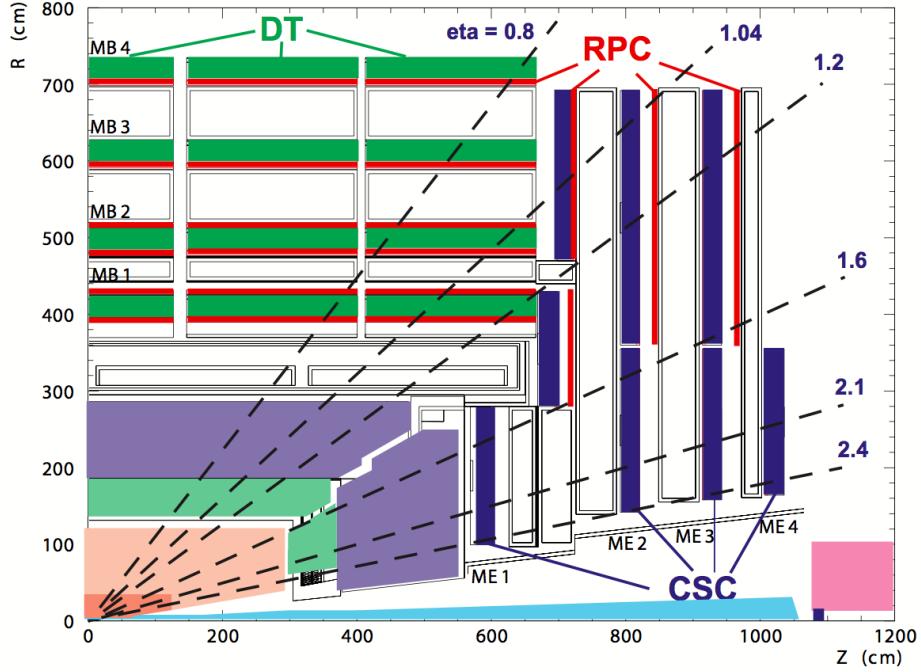


Figure 3.12: Layout of the CMS muon system for initial run configurations. The DT system is used only in the barrel region and CSCs in the endcap region, while RPCs are deployed in both the barrel and endcap.

can be reconstructed by using the DT station information.

The CSCs in each endcap are trapezoidal in shape and overlap in the ϕ coordinate, increasing the fiducial coverage of the muon system. Each CSC consists of 6 gas gaps, where each gap is comprised of radially aligned cathode strips and a plane of anodes wires nearly perpendicular to the strips. When a muon passes through the CSCs, the ensuing ionization and electron avalanche deposits a charge on the anode wire and an image charge on the cathode strips. A measurement of the muon position can be reconstructed by determining the center-of-gravity of the charge distribution on the cathode strips, and the fast readout of the anode wire is used in low-latency trigger decisions as described in section 3.2.5.

The RPCs are complementary systems in both the barrel and endcap used to im-

prove the time resolution of muon measurements. Each RPC consists of two parallel anode and cathode plates, separated by a gas chamber. The ionization of the gas as muons travel through the chamber is quickly deposited on the plates, and used for a precise time measurement of the muon crossing. Though the spatial resolution of the RPCs is poor, in conjunction with the other muon systems they allow for an unambiguous association of muons with different bunch crossings.

3.2.5 Trigger Systems

When operating at design luminosity, the LHC can deliver proton bunches to the collision point at a rate of 40MHz, resulting in an average collision rate on the order of one billion collisions per second. In order to achieve reasonable rates of data collection for offline storage and processing, the detector must suppress the event rate by six orders of magnitude when selecting events of interest to be saved for physics analyses. This is accomplished through a combination of readout electronics and the trigger systems: the Level-1 trigger (L1) processors and online High-Level triggers (HLT) [40].

The L1 trigger system is comprised of specialized hardware processors to rapidly pre-select events of interest based on the calorimeter and muon systems. Based on the beam crossing frequency, the L1 electronics make a readout decision based on buffered readout data collected between each bunch crossing from the front end electronics, and have only a few microseconds to execute logic to select events of interest, such that the total time allotted for L1 trigger calculations is $< 1\mu s$. During this time the bulk of detector data is held in a buffer while L1 trigger decisions are made based on data with reduced granularity

and resolution rapidly collected from the calorimeter and muon systems, where triggers typically check for “trigger primitive” objects (such as photons, muons, electrons, etc.). Trigger primitives must meet certain momentum or energy thresholds, and L1 triggers may also check global data about the event such as the sum of transverse energy or the missing transverse energy (inferred from momentum conservation).

After an L1 trigger tags an event, high-resolution data is read out from buffers for additional data processing before reaching the HLT. Each event of $\sim 1.5\text{MB}$ is transferred to front-end readout buffers which then pipe data to a processor containing the HLT code. HLT code is designed to discard events as soon as possible when making trigger decisions and only a subset of objects or partial reconstruction of events may occur before the final trigger decision is made, though HLT algorithms typically approach the quality of final reconstruction. The HLT reduces the L1 output rate to $\mathcal{O}(1\text{kHz})$ for event storage and full reconstruction.

In this analysis, the primary triggers used for selecting signal events depend on the total hadronic energy (H_T), jet momentum (described in section 3.3.6), or missing energy (E_T^{miss} , described in section 3.3.7). These triggers are seeded by L1 paths which rapidly reconstruct H_T , jet p_T , or E_T^{miss} in an event. If the L1 trigger flags the event, it is passed to HLT triggers that check for higher-level reconstructed quantities of H_T , jet p_T , E_T^{miss} , or a combination thereof. The efficiency for one of the signal trigger paths (HLT PFHT900 OR HLT PFJet450) as a function of the total hadronic energy in the event can be seen in figure 3.13. Additional triggers in this analysis for lepton-rich and QCD-rich control regions are also used as described in chapter 5.

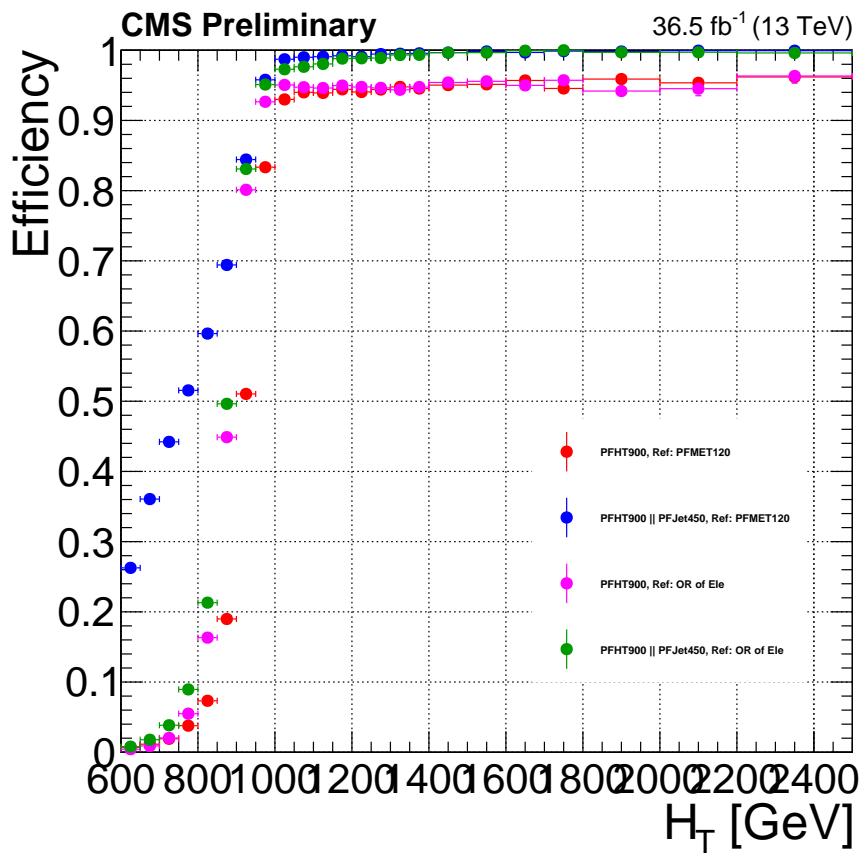


Figure 3.13: The efficiency of the HLT PFHT900 OR HLT PFJet450 triggers as a function of total hadronic energy H_T .

3.3 CMS Physics Objects

When physics events are fully reconstructed, detector data is used to identify *physics objects* representing real particles and event quantities for use in a physics analysis. The physics objects in an event — such as leptons, jets, or missing energy — and their properties are used to select events of interest for physics analyses targeting different final states. The properties of physics objects and global event data are also be used to make analysis level decisions of the quality of different objects. Here we describe some of the physics objects referred to in the $M_{\mathrm{T}2}$ analysis and how they are reconstructed, as well as some global event properties and quality variables used as discriminants for physics objects and events.

3.3.1 Particle Flow

Most of the physics objects described in the following sections are reconstructed and identified in CMS using the particle flow (PF) algorithm [29]. The PF algorithm is a holistic, iterative algorithm which uses all the available data in the detector to classify “PF candidate” particles in an event. PF works iteratively by clustering tracks and calorimeter deposits into a PF candidate, removing all energy and hits associated with the candidate, and repeating the algorithm until all the detector information has been associated to PF objects. First, any muon tracks in the inner tracker associated with muon system hits are associated and removed. Any remaining tracks are then extrapolated into the calorimeters, and the energy deposits consistent with tracks are clustered with the track and removed from further consideration. Once all the tracks have been associated, the remaining energy

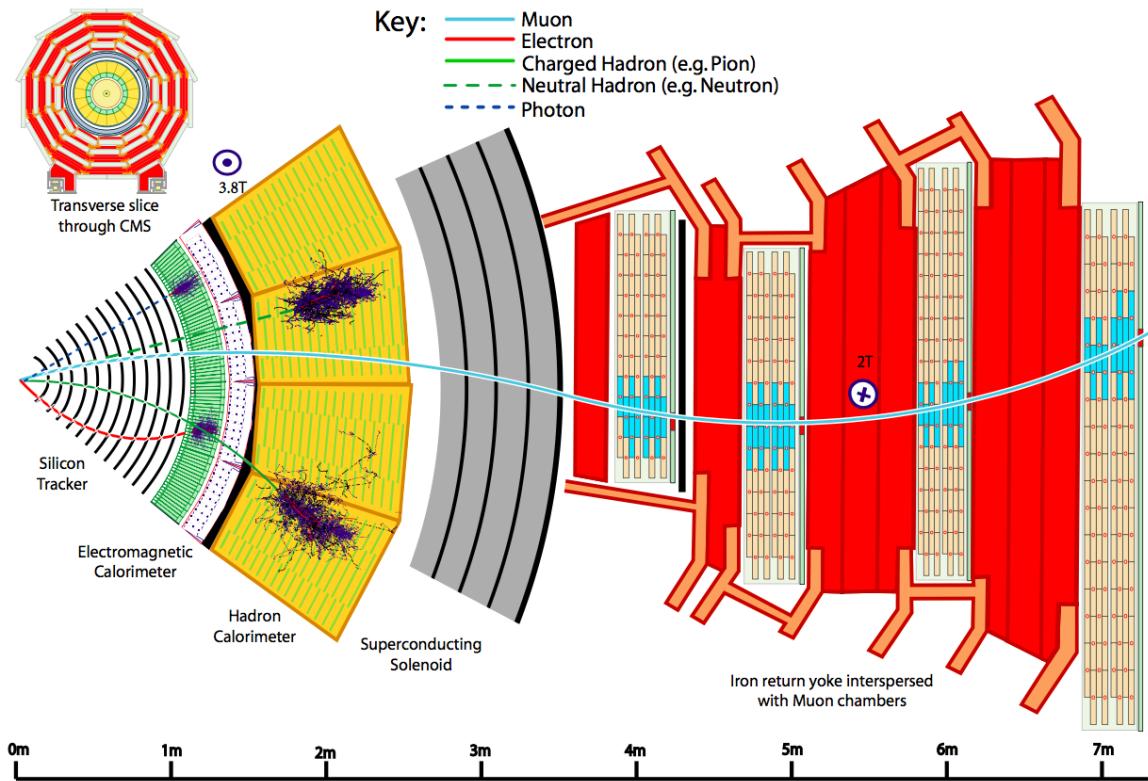


Figure 3.14: A graphic depiction of different particles leaving various signatures in the different CMS detector subsystems. Particles may be detected via tracks hits, energy deposits, or a combination of both.

clusters can be identified with photons and neutral hadrons (depending on their presence in the ECAL or HCAL, respectively). An example of the different tracks and energy deposits associated with various particles can be seen in figure 3.14.

3.3.2 Isolation

Isolation is an important distinguishing variable for physics objects measured in the detector. Simply put, isolation measures the total amount of energy in some proximity to a physics object. Because many physics processes results in particle production outside the primary interaction vertex (such as showering, decays, hadronization, pair production,

etc.), the isolation of a particle parametrizes the amount of other "activity" surrounding the physics object, and consequently proves a useful discriminant when determining if color singlet particles were produced or if particles are part of a hadronizing or showering process. Particles which are produced in the primary physics process in a collision are sometimes referred to as *prompt*, and isolation is a powerful tool in identifying prompt leptons in particular.

3.3.3 Electrons and Photons

The primary leptons considered in this analysis are electrons and muons (discussed in section 3.3.4), and photons are also used as a cross-check in some control regions for background estimates. Prompt electrons and photons are identified in CMS primarily through the use of the ECAL, where they are stopped and deposit all their energy [14, 20]. Prompt electrons are customarily distinguished by both the presence of a track originating at the primary vertex and a sufficiently low isolation.

Electrons in CMS are identified primarily through the existence of a charged particle track terminating in an energy cluster in the ECAL. As described in section 3.2.2, the ECAL crystals are ~ 25 radiation lengths deep, and will stop and contain nearly all of the energy of an electron (except perhaps minuscule losses to scattering in the tracker). The charged track leading to the energy deposit distinguishes the charged particle from a neutral particle, and the direction of curvature of the path in the tracker can be used to identify the charge of an electron (or positron).

Photons are distinguished from electrons by the lack of a track in the SVT. As a

neutral particle, the only sign of a photon will be the energy deposit in the ECAL where the photon is stopped and deposits all its energy. Photons can often be distinguished from other neutral electromagnetic particles (such as a π^0) by the shape of the shower in the ECAL. Whereas photons typically shower through pair production and bremsstrahlung, other hadronic particles may cascade via different physics processes leading to measurable difference in the shape of the energy cluster.

The all-hadronic search described here makes use of electrons and photons in the construction of different control regions to predict background processes (described in detail in chapter 5). While an electron veto is used in the signal region, electrons are selected in the single lepton control region and other dilepton regions. Photons are primarily used as a cross-check in M_{T2} shape studies as described in section 5.4.

3.3.4 Muons

Prompt muons produced in collisions at the LHC are generally minimum ionizing particles, and will penetrate the bulk of the detector. Without any matching calorimeter deposits, they are reconstructed using hits in the muon system matched to inner tracker hits [16]. The muon system flags muon candidates, which are then matched to inner tracker hits for the best fit of the muon track. Given the track hits, the transverse momentum and energy of the muon can be calculated by a fit to the track parameters. As with electrons, muons are customarily distinguished by both the presence of a track originating at the primary vertex and a sufficiently low isolation.

Muons are primarily used in the all-hadronic analysis to select lepton-enriched

control regions for the estimation of different backgrounds (described in detail in chapter 5). A muon veto is applied in the signal region, and muons are used in single lepton and dilepton control regions for different background estimates.

3.3.5 Isolated Tracks

Not all charged tracks in the detector are necessarily associated with leptons or other physics objects. *Isolated tracks* are a collection of objects where a charged track consistent with the primary interaction vertex is also isolated (and thus not likely part of jet hadronization as described in section 3.3.6).

The analysis described in this dissertation makes particular use of isolated tracks in rejecting events with leptons. Because leptons are not always fully reconstructed into higher-level physics objects, an additional isolated track veto as described in chapter 4 removes additional events with prompt leptons. Furthermore, τ leptons (which are the only leptons with sufficient mass to decay into hadrons) can be identified through the presence of isolated tracks, of particular importance in the analysis presented here when implementing a lepton veto.

3.3.6 Jets

Many particles produced at the LHC are created through the strong interaction, which may produce particles carrying color charge in the final state (i.e. quarks or gluons). Colored particles cannot exist individually due to the phenomenon known as *confinement*, and so via the process of *hadronization* they will proliferate in a series of interactions

producing additional particle-antiparticle pairs, cascading in a parton shower to form a multitude of hadronic bound states. When a “bare” quark or gluon is produced in the primary interaction, the ensuing hadronization results in a stream of tightly collimated hadrons aligned with the trajectory of the original bare particle, referred to as a *jet*.

Jets present themselves in CMS as collimated energy deposits in both the ECAL and HCAL, as well as tracks originating at the primary vertex [48]. Between the two calorimeters, all the electromagnetic and hadronic components of the jet are stopped and measured by the calorimeters, and any leptons (including muons detected in the muon system) which originate in the cone of the jet are assigned to the total jet energy. The jets are reconstructed using the *anti- k_T algorithm*, which clusters jet activity in cones of some radius [4]. The jet is treated as a single physics object with a total energy equal to the sum of its electromagnetic, hadronic, and leptonic components, which by momentum conservation dictates the energy of the prompt quark or gluon that was produced in the primary interaction.

Furthermore, the substructure of jets and their content can be analyzed to determine the flavor of the parent parton, known as “tagging”. Jets are typically tagged to distinguish between those originating from heavy-flavor quarks (bottom or charmed) and light-flavor quarks. Top quarks produced in primary interactions are unique because of their short lifetime. A top quark will decay before hadronization occurs, instead producing a W boson and down-type (down, bottom, or strange) quark which will subsequently hadronize. Searches involving top quarks in the final state typically employ “top-taggers” to search for these signatures of the top quark.

In practice, jets are complicated objects consisting of many constituent particles, and must be clustered and calibrated so that the jet energy closely matches that of the parent parton which produced the jet. Jet energy corrections (JECs) are applied to raw jet energies to compensate for different experimental deviations from the parent parton energy [28, 21]. In particular, JECs include corrections to compensate for pileup energy as described in section 3.1.1 [3], detector effects (as a function of η), energy scale as a function of jet p_T , and residual corrections to account for differences between data and simulation. The JECs are calculated by collecting data events with a Z boson (decaying to electrons or muons) or photon recoiling against a jet. By precisely measuring the bosons energy with the ECAL, scale factors are derived for jets as a function of pseudorapidity η and momentum p_T .

3.3.7 Missing Energy

Missing energy refers to the sum of all energy which has escaped the detector, and is inferred from a momentum-imbalance in the physics objects measured by the detector [18]. Because the momentum of the incident protons is zero perpendicular to the direction of the beam, the momentum of physics objects produced in any collision must sum to zero in the direction transverse to the beamline. The missing transverse energy (E_T^{miss}), is calculated by taking the negative vector sum of all PF candidates as in equation 3.5. E_T^{miss} is of particular importance to analyses targeting BSM physics, as some BSM models which include dark matter candidates characteristically contain invisible particles in the

final state which escape detection.

$$\vec{E}_T^{\text{miss}} = - \sum_i \vec{p}_T^i \quad (3.5)$$

The missing energy in an event is inferred from all other measured quantities in an event, and there are many sources of E_T^{miss} which are unphysical in nature but rather dependent on experimental effects. Particles from the primary interaction with a sufficiently large pseudorapidity may escape the fiducial region of the detector subsystems, and resolution effects or intrinsic noise in the detector may lead to fluctuations in measured energies. These experimental effects must be suppressed or distinguished from “real” E_T^{miss} due to physics processes creating particles which escape the detector (e.g. neutrinos). Analyses sensitive to final states with E_T^{miss} often employ robust data-driven methods to predict or suppress backgrounds which might generate experimental sources of E_T^{miss} or physics processes which can contribute real E_T^{miss} (such as $Z \rightarrow \nu\bar{\nu}$).

3.3.8 Monte Carlo Simulation

Monte Carlo (MC) simulation of events within the CMS detector is frequently generated to understand the expected physics, design analyses, and help predict certain backgrounds. The simulation begins with the underlying physics process of interest, and subsequently models the trajectories of particles in the final state through the detector volume, their interactions in those volumes and with the detector material, the effects of pileup, and the detector response.

The process of generating MC proceeds in multiple steps interfacing different soft-

ware packages. In this analysis, many MC samples are generated using the POWHEG [44] and MadGraph [2] generator packages to calculate the matrix elements for the underlying physics process. The calculations performed by these generators are then interfaced with Pythia [49] to model parton showering and hadronization effects. The material interaction in the detector volume are modeled with the Geant4 toolkit [5], and the detector’s response is digitized to model the final readout from the event.

Because of the many computationally expensive steps involved in creating MC and the large amount of events required to suppress statistical uncertainties, simulation is occasionally generated with the Fastsim package [17] to decrease computational burden at the expense of some modeling quality. When Fastsim is employed in this analysis for the purpose of generating signal MC, effects due to potential mismodeling introduced by the simulation are accounted for in studies of the systematic uncertainty (as shown in section 6.2.1).

This chapter makes use of figures from the CMS Technical Design Report and Particle Flow reconstruction paper to illustrate the experimental design. The work presented in this dissertation would not be possible without the many contributions from CMS members who designed the detector and algorithms used in this analysis.

Chapter 4

The M_{T2} Variable Search

4.1 Analysis Strategy

Searches for new physics targeting all-hadronic final states present unique challenges and opportunities at the LHC, and have previously been conducted by both the ATLAS collaboration [9, 8, 11, 10, 7, 27] and the CMS collaboration [24, 25, 26]. While such searches typically implement stringent vetoes on lepton candidates and thus suppress the need to correctly identify “real” leptons, the high rate of QCD processes in proton-proton collisions generates large amounts of SM events with all-hadronic final states. Designing a search targeting signatures with all-hadronic final states requires a mechanism to distinguish and suppress the selection of multi-jet QCD events from new physics signatures, as well as robust background estimation methods to predict the yield of SM events which may generate true missing energy (such as $Z \rightarrow \nu\bar{\nu}$).

The M_{T2} analysis harnesses the discriminating power of the M_{T2} variable, some-

times referred as stransverse mass, to distinguish SM events from possible signatures of new physics. By first requiring a significant amount of missing energy in the event, multi-jet QCD processes are greatly suppressed. Additional requirements on the topology of the event implemented using M_{T2} further suppress QCD-like processes and favor events with real missing energy anti-aligned with the hadronic energy deposits in the detector. After estimating the minimal QCD contribution remaining by extrapolating from a region orthogonal to the signal selection, the only remaining backgrounds are leptonic events where the lepton failed reconstruction or identification (or “lost-lepton” events), and SM events creating real E_T^{miss} in the form of neutrinos from a decaying Z boson recoiling against jets (or “invisible Z” events).

4.2 The M_{T2} Variable

M_{T2} is a particularly useful kinematic mass variable for final states where two particles decay (possibly in a chain) to a final state containing an invisible particle X of mass m_X [40]. The typical transverse mass M_T is defined in equation 4.1 for particles $i = 1, 2$, where the mass $m^{\text{vis}(i)}$, transverse momentum $\vec{p}_t^{\text{vis}(i)}$, and transverse energy $E_T^{\text{vis}(i)}$ characterize the visible kinematics of the decay chain, and $\vec{p}_t^{X(i)}$ and $E_T^{X(i)}$ characterize the unknown kinematics of the invisible particle X.

$$(M_T^{(i)})^2 = (m^{\text{vis}(i)})^2 + m_X^2 + 2 \left(E_T^{\text{vis}(i)} \cdot E_T^{X(i)} - \vec{p}_t^{\text{vis}(i)} \cdot \vec{p}_t^{X(i)} \right) \quad (4.1)$$

In principle, if the correct values of m_X and $\vec{p}_t^{X(i)}$ were accessible, then the transverse mass would have a kinematic endpoint and not exceed the mass of the parent particles (disre-

garding any resolution effects). However, the individual momenta of the invisible particles in the two decay chains cannot be measured; the only quantity experimentally accessible is the total missing momentum \vec{p}_T^{miss} . With this in mind, the generalized transverse mass variable $M_{\text{T}2}$ is defined in equation 4.2, where the unknown mass m_X is a free parameter and a minimization is performed over the sum of invisible momenta $\vec{p}_t^{\text{X}(i)}$ that satisfy the measured \vec{p}_T^{miss} constraint.

$$M_{\text{T}2}(m_X) = \min_{\vec{p}_t^{\text{X}(1)} + \vec{p}_t^{\text{X}(2)} = \vec{p}_T^{\text{miss}}} \left[\max \left(M_T^{(1)}, M_T^{(2)} \right) \right] \quad (4.2)$$

Because this analysis selects final states with multiple jets in the final state, the calculation of $M_{\text{T}2}$ first requires grouping the hadronic jet activity into two large *pseudojets* to act as the visible components in the $M_{\text{T}2}$ equation. The jet activity in each event is divided into two hemispheres, and the jets in each hemisphere are summed together to create the pseudojets. The hemisphere algorithm proceeds as follows:

- The direction of the two jets with largest invariant mass is chosen as the initial seed for the two axes.
- Jets are associated to one of the two axes according to the minimal Lund distance [13], such that jet k is associated to hemisphere i instead of j if the condition in equation 4.3 is true.
- After each jet is associated to one of the two axes, the axes are recalculated by summing the momenta of all jets associated to an axis.
- The association algorithm iterates using the new axes, and continues until no jets

are associated to a different axis after an iteration.

$$(E_i - p_i \cos \theta_{ik}) \frac{E_i}{(E_i + E_k)^2} \leq (E_j - p_j \cos \theta_{jk}) \frac{E_j}{(E_j + E_k)^2} \quad (4.3)$$

When clustered using this pseudojet algorithm, QCD multijet events may yield high M_{T2} if the pseudojets have high jet masses, thus the visible masses $m^{\text{vis}(i)}$ are set to zero to suppress such SM events. Since the kinetic components of M_{T2} will be large for signal events, this suppression does not significantly impact sensitivity to many BSM signatures, thus M_{T2} is calculated in this analysis using only E_T^{miss} and the two pseudojets with $m_X \equiv 0$.

4.3 Event Selection Criteria

The general strategy for the event selection is to first apply baseline cuts motivated by hardware and software-level triggers (discussed in section 3.2.5) and reducing the QCD multi-jet background to negligible levels. Events are further categorized using transverse mass (M_{T2}), the scalar sum of the transverse momenta p_T of all selected jets (H_T), the total number of jets in the event (N_{jets}), and the total number of b-tagged jets in the event ($N_{\text{b-jets}}$). A summary of the event preselections can be found in table 4.1.

4.4 Search Regions

There are 213 individual search regions are defined by categorizing events in bins of H_T , N_{jets} , $N_{\text{b-jets}}$, and M_{T2} (in addition to the baseline selection described in section

Table 4.1: Summary of physics objects and preselection for signal events. Here R is the distance parameter of the anti- k_T algorithm, and for veto leptons and tracks, the transverse mass M_T is determined using the veto object and the \vec{E}_T^{miss} , while p_T^{sum} denotes the sum of the transverse momenta of all the PF candidates in a cone around the lepton or track. The size of the cone, in units of $\Delta R \equiv \sqrt{(\Delta\phi)^2 + (\Delta\eta)^2}$ is given in the table.

Trigger	$E_T^{\text{miss}} > 120 \text{ GeV}$ and $H_T^{\text{miss}} > 120 \text{ GeV}$ or $H_T > 300 \text{ GeV}$ and $E_T^{\text{miss}} > 110 \text{ GeV}$ or $H_T > 900 \text{ GeV}$ or jet $p_T > 450 \text{ GeV}$
Jet selection	$R = 0.4$, $p_T > 30 \text{ GeV}$, $ \eta < 2.4$
b tag selection	$p_T > 20 \text{ GeV}$, $ \eta < 2.4$
E_T^{miss}	$E_T^{\text{miss}} > 250 \text{ GeV}$ for $H_T < 1000 \text{ GeV}$, else $E_T^{\text{miss}} > 30 \text{ GeV}$ $\Delta\phi_{\min} = \Delta\phi(p_T^{\text{miss}}, j_{1,2,3,4}) > 0.3$ $ \vec{E}_T^{\text{miss}} - \vec{H}_T^{\text{miss}} /E_T^{\text{miss}} < 0.5$
M_{T2}	$M_{T2} > 200 \text{ GeV}$ for $H_T < 1500 \text{ GeV}$, else $M_{T2} > 400 \text{ GeV}$
Veto muon	$p_T > 10 \text{ GeV}$, $ \eta < 2.4$, $p_T^{\text{sum}} < 0.2 \times p_T^{\text{lep}}$ or $p_T > 5 \text{ GeV}$, $ \eta < 2.4$, $M_T < 100 \text{ GeV}$, $p_T^{\text{sum}} < 0.2 \times p_T^{\text{lep}}$
Veto electron	$p_T > 10 \text{ GeV}$, $ \eta < 2.4$, $p_T^{\text{sum}} < 0.1 \times p_T^{\text{lep}}$ or $p_T > 5 \text{ GeV}$, $ \eta < 2.4$, $M_T < 100 \text{ GeV}$, $p_T^{\text{sum}} < 0.2 \times p_T^{\text{lep}}$
Veto track	$p_T > 10 \text{ GeV}$, $ \eta < 2.4$, $M_T < 100 \text{ GeV}$, $p_T^{\text{sum}} < 0.1 \times p_T^{\text{track}}$
p_T^{sum} cone	Veto e or μ : $\Delta R = \min(0.2, \max(10 \text{ GeV}/p_T^{\text{lep}}, 0.05))$ Veto track: $\Delta R = 0.3$

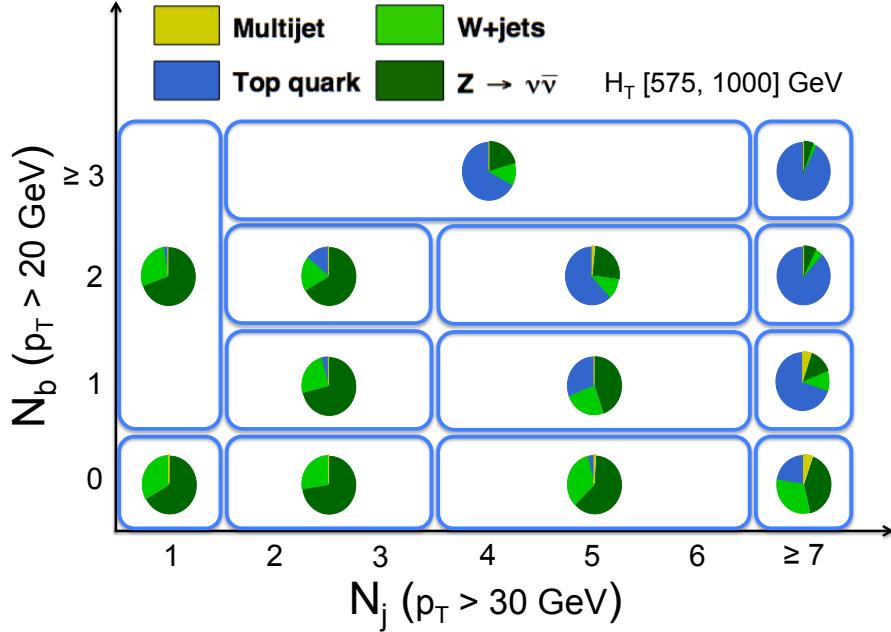


Figure 4.1: Topological regions in N_{jets} and $N_{\text{b-jets}}$ for the $[575, 1000]$ H_{T} region. Within each region, the relative fraction of background events from different SM processes is shown based on simulation.

4.3). First events are categorized into “topological regions” according to H_{T} , N_{jets} , and $N_{\text{b-jets}}$:

- H_{T} (GeV): $[250, 450]$ (Very Low), $[450, 575]$ (Low), $[575, 1000]$ (Medium), $[1000, 1500]$ (High), $[1500, \infty]$ (Extreme)
- $N_{\text{jets}} \& N_{\text{b-jets}}$: 2-3j 0b, 2-3j 1b, 2-3j 2b, 4-6j 0b, 4-6j 1b, 4-6j 2b, $\geq 7j$ 0b, $\geq 7j$ 1b, $\geq 7j$ 2b, 2-6j $\geq 3b$, and $\geq 7j$ $\geq 3b$ (except in the region with $250 < H_{\text{T}} < 450$ GeV, where bins $\geq 7j$ are merged with 4-6j bins due to lack of events).

The different topological regions for one H_{T} region and their background composition are depicted in figure 4.1. Each topological region is further divided into M_{T2} bins. The M_{T2} binning is constructed such that the low edge of the first bin is 400 GeV in regions with H_{T}

$> 1500 \text{ GeV}$ and 200 GeV everywhere else, and the low edge of the final bin is constructed to contain approximately one background event based on simulation and not exceeding the maximum H_T value in that topological region (since an upper limit on H_T places an upper limit on M_{T2}). The detailed M_{T2} binning is as follows:

- Very Low H_T : $[200,300]$, $[300,400]$, $[400,\infty]$
- Low H_T : $[200,300]$, $[300,400]$, $[400,500]$, $[500,\infty]$
- Medium H_T : $[200,300]$, $[300,400]$, $[400,600]$, $[600,800]$, $[800,\infty]$
- High H_T : $[200,400]$, $[400,600]$, $[600,800]$, $[800, 1000]$, $[1000, 1200]$, $[1200,\infty]$
- Extreme H_T : $[400,600]$, $[600,800]$, $[800,1000]$, $[1000,1400]$, $[1400,\infty]$

The various H_T bins and associated M_{T2} binning can be seen in figure 4.2, and the full breakdown of signal regions (including M_{T2} binning) is listed in tables 4.2 and 4.3. In addition to multijet search regions, this analysis also considers monojet events. Because there is only a single jet (and M_{T2} is ill-defined without multiple jets), binning in these regions is defined using $N_{\text{b-jets}}$ and H_T as follows:

- $N_{\text{b-jets}}$: 0b , $\geq 1\text{b}$
- H_T : $[250,350]$, $[350,450]$, $[450,575]$, $[575,700]$, $[700,1000]$, $[1000,1200]$, $[1200,\infty]$

As with the multijet regions, monojet H_T bins with less than one simulated background event in the final bin are merged with the penultimate bin.

Table 4.2: M_{T2} binning in the Very Low, Low, and Medium H_T topological regions.

H_T Range [GeV]	Jet Multiplicities	Binning [GeV]
[250, 450]	2 – 3j, 0b	[200, 300, 400, ∞]
	2 – 3j, 1b	[200, 300, 400, ∞]
	2 – 3j, 2b	[200, 300, 400, ∞]
	\geq 4j, 0b	[200, 300, 400, ∞]
	\geq 4j, 1b	[200, 300, 400, ∞]
	\geq 4j, 2b	[200, 300, 400, ∞]
	\geq 2j, \geq 3b	[200, 300, 400, ∞]
[450, 575]	2 – 3j, 0b	[200, 300, 400, 500, ∞]
	2 – 3j, 1b	[200, 300, 400, 500, ∞]
	2 – 3j, 2b	[200, 300, 400, 500, ∞]
	4 – 6j, 0b	[200, 300, 400, 500, ∞]
	4 – 6j, 1b	[200, 300, 400, 500, ∞]
	4 – 6j, 2b	[200, 300, 400, 500, ∞]
	\geq 7j, 0b	[200, 300, 400, ∞]
	\geq 7j, 1b	[200, 300, 400, ∞]
	\geq 7j, 2b	[200, 300, 400, ∞]
	2 – 6j, \geq 3b	[200, 300, 400, 500, ∞]
[575, 1000]	\geq 7j, \geq 3b	[200, 300, 400, ∞]
	2 – 3j, 0b	[200, 300, 400, 600, 800, ∞]
	2 – 3j, 1b	[200, 300, 400, 600, 800, ∞]
	2 – 3j, 2b	[200, 300, 400, 600, 800, ∞]
	4 – 6j, 0b	[200, 300, 400, 600, 800, ∞]
	4 – 6j, 1b	[200, 300, 400, 600, 800, ∞]
	4 – 6j, 2b	[200, 300, 400, 600, 800, ∞]
	\geq 7j, 0b	[200, 300, 400, 600, 800, ∞]
	\geq 7j, 1b	[200, 300, 400, 600, ∞]
	\geq 7j, 2b	[200, 300, 400, 600, ∞]
	2 – 6j, \geq 3b	[200, 300, 400, 600, ∞]
	\geq 7j, \geq 3b	[200, 300, 400, 600, ∞]

Table 4.3: M_{T2} binning in the High and Extreme H_T topological regions.

H_T Range [GeV]	Jet Multiplicities	Binning [GeV]
[1000, 1500]	2 – 3j, 0b	[200, 400, 600, 800, 1000, 1200, ∞]
	2 – 3j, 1b	[200, 400, 600, 800, 1000, 1200, ∞]
	2 – 3j, 2b	[200, 400, 600, 800, 1000, ∞]
	4 – 6j, 0b	[200, 400, 600, 800, 1000, 1200, ∞]
	4 – 6j, 1b	[200, 400, 600, 800, 1000, 1200, ∞]
	4 – 6j, 2b	[200, 400, 600, 800, 1000, ∞]
	$\geq 7j$, 0b	[200, 400, 600, 800, 1000, ∞]
	$\geq 7j$, 1b	[200, 400, 600, 800, ∞]
	$\geq 7j$, 2b	[200, 400, 600, 800, ∞]
	2 – 6j, $\geq 3b$	[200, 400, 600, ∞]
	$\geq 7j$, $\geq 3b$	[200, 400, 600, ∞]
[1500, ∞]	2 – 3j, 0b	[400, 600, 800, 1000, 1400, ∞]
	2 – 3j, 1b	[400, 600, 800, 1000, ∞]
	2 – 3j, 2b	[400, ∞]
	4 – 6j, 0b	[400, 600, 800, 1000, 1400, ∞]
	4 – 6j, 1b	[400, 600, 800, 1000, 1400, ∞]
	4 – 6j, 2b	[400, 600, 800, ∞]
	$\geq 7j$, 0b	[400, 600, 800, 1000, ∞]
	$\geq 7j$, 1b	[400, 600, 800, ∞]
	$\geq 7j$, 2b	[400, 600, 800, ∞]
	2 – 6j, $\geq 3b$	[400, 600, ∞]
	$\geq 7j$, $\geq 3b$	[400, ∞]

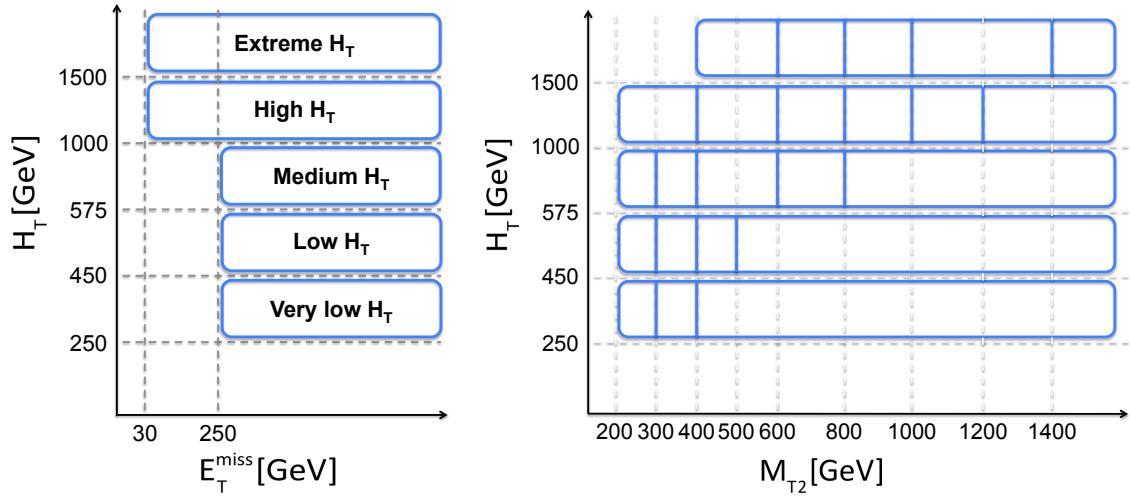


Figure 4.2: Signal region bins in H_T and E_T^{miss} (left) and M_{T2} binning within each H_T region (right). If simulation predicts less than one background event in the greatest M_{T2} bin within a region, it is merged with the previous bin.

In addition to these signal regions used to interpret results in the context of various BSM physics models, the analysis also provides results in "super signal regions" (SSRs) as defined in table 4.4. These regions provide a simpler set of selections than the nominal signal regions so that phenomenologists may easily reinterpret results in the context of different signal models. Results obtained using the SSRs are not as sensitive as the nominal binning — finely binned regions have a higher signal-to-background ratio and the global background fit reduces the background uncertainties — but are much easier to use for reinterpretation than the many correlated bins of the full analysis.

4.5 Control Regions

In order to anchor the data-driven background estimates used in this analysis, *control regions* (CR) orthogonal to the signal region selection are defined for various processes.

Table 4.4: Definition of "super signal regions" used in reinterpretations of the analysis.

Region	N_{jets}	$N_{\text{b-jets}}$	H_T [GeV]	M_{T2} [GeV]
2j loose	≥ 2	-	> 1000	> 1200
2j tight	≥ 2	-	> 1500	> 1400
4j loose	≥ 4	-	> 1000	> 1000
4j tight	≥ 4	-	> 1500	> 1400
7j loose	≥ 7	-	> 1000	> 600
7j tight	≥ 7	-	> 1500	> 800
2b loose	≥ 2	≥ 2	> 1000	> 600
2b tight	≥ 2	≥ 2	> 1500	> 600
3b loose	≥ 2	≥ 3	> 1000	> 400
3b tight	≥ 2	≥ 3	> 1500	> 400
7j3b loose	≥ 7	≥ 3	> 1000	> 400
7j3b tight	≥ 7	≥ 3	> 1500	> 400

In particular, there are control regions corresponding to enriched samples of single lepton events, $Z \rightarrow l^+l^-$ events, and QCD multijet events.

4.5.1 Single Lepton Control Region

The single lepton CR is constructed to select a sample enriched with single lepton events, the most dominant contributions being from $t\bar{t}$ and $W + \text{jets}$ production. The same baseline selections described in section 4.3 are applied with the exception of the following:

- In lieu of the lepton veto, exactly one lepton candidate passing the reco or PF lepton selections is required. In order to avoid double counting (for leptons which are reconstructed both as a reco lepton and PF candidate), PF leptons within $\Delta R < 0.1$ of a reco lepton are not considered.
- The transverse mass M_T between the lepton and E_T^{miss} must be less than 100 GeV to reduce possible signal contamination.

Since non-isolated leptons in the fiducial region of the detector are usually successfully reconstructed, the closest jet within $\Delta R < 0.4$ of the lepton is removed and the lepton instead counted as a visible object for the purposes of computing H_T , H_T^{miss} , $\Delta\phi(j_{1234}, E_T^{\text{miss}})$, $|\vec{H}_T^{\text{miss}} - \vec{E}_T^{\text{miss}}|/E_T^{\text{miss}}$, and M_{T2} (as well as the hemispheres used to calculate M_{T2}). Events are further subdivided into the topological regions described in section 4.4 using the modified H_T and N_{jets} and $N_{\text{b-jets}}$, but not in M_{T2} to increase the statistical power of the CR. The signal regions with $\geq 7j, \geq 1b$ are all predicted using CRs with identical H_T bins but $\geq 7j, 1\text{-}2b$ to increase the statistical power of those CRs (and to avoid signal contamination

in regions with $\geq 7j, \geq 3b$). In addition, for regions with $H_T > 1500\text{ GeV}$, the minimum M_{T2} threshold is set to 200 GeV to increase available statistics. The monojet CR is binned identically to the signal region.

4.5.2 Dilepton Control Region

Control regions corresponding to opposite-sign same-flavor leptons (OSSF) from $Z \rightarrow l^+l^-$ events are used to estimate the $Z \rightarrow \nu\bar{\nu}$ background, with corresponding sets of control regions requiring an opposite-sign opposite-flavor (OSOF) pair to estimate the flavor-symmetric background component in the former dilepton CR. The same baseline selections as described in section 4.3 are applied with the exception of the follow:

- In lieu of the lepton veto, exactly 2 leptons (ee , $e\mu$, or $\mu\mu$) passing loose lepton selections are required.
- There is no requirement on E_T^{miss} . Instead, the dilepton system must have a transverse momentum $p_T(\ell\ell) > 200\text{ GeV}$ to mimic the kinematics of the $Z \rightarrow \nu\bar{\nu}$ background and suppress the $t\bar{t}$ contribution.
- Without a missing energy requirement, events are selected in data using leptonic trigger paths. Dimuon events are selected using a combination of dimuon and high- p_T single muon triggers, dielectron events using a combination of dielectron and high- p_T single photon paths (which do not require isolation and recover efficiency for high- p_T electrons or those highly co-linear high- p_T Z bosons), and opposite-flavor

events are selected using a combination of $e\mu$ triggers and higher threshold single-lepton paths (to again recover efficiency for some events).

- To improve trigger efficiency for these regions, the leading lepton is required to have a minimum momentum $p_T > 100 \text{ GeV}$ and the sub-leading lepton $p_T > 30 \text{ GeV}$.
- When selecting Z boson candidates for the $Z \rightarrow \nu\bar{\nu}$ estimate, the leptons are also required to be OSSF with an invariant mass $|m_{\ell\ell} - m_Z| < 20 \text{ GeV}$, where m_Z is the nominal Z boson mass.

Similar to the single lepton control region, the closest jet within $\Delta R < 0.4$ of each lepton is removed and the leptons added to the E_T^{miss} vector for the purposes of computing H_T , H_T^{miss} , $\Delta\phi(j_{1234}, E_T^{\text{miss}})$, $|\vec{H}_T^{\text{miss}} - \vec{E}_T^{\text{miss}}|/E_T^{\text{miss}}$, and M_{T2} (as well as the hemispheres used to calculate M_{T2}). Additional information on the M_{T2} binning for these regions is detailed in section 5.4.

4.5.3 Multijet Control Region

The multijet control region is designed to select a sample enriched in QCD events to estimate the multijet background. The same baseline selections described in section 4.3 are applied with the exception of the $\Delta\phi(j_{1234}, E_T^{\text{miss}})$ requirement, which is inverted to select a sample dominated by QCD events with large fake E_T^{miss} due to jet energy mismeasurements.

The transfer factor which is used to extrapolate the control region yield to the signal region is measured in a separate QCD-dominated region with $M_{T2} < 200 \text{ GeV}$, described

in detail in section 5.2. Because of the lower E_T^{miss} requirement, different trigger paths must be used. For regions with $H_T > 1000 \text{ GeV}$, trigger paths seeded by a single high- p_T jet are used. For other H_T regions, similar trigger paths with lower jet p_T thresholds are used, but due to the rate of QCD events creating low- p_T jets these paths are prescaled to suppress the data acquisition rate.

This chapter makes use of figures and tables from the M_{T2} paper and internal analysis note to illustrate the analysis design, methodology, and results. This work was made possible by contributions from the rest of the Surf & Turf group, our collaborators at ETH Zurich, and the many other CMS members in the SUSY group and beyond.

Chapter 5

Background Estimates

5.1 Types of Backgrounds

The M_{T2} all-hadronic search is primarily a search for missing energy, and the backgrounds present in the signal region can be roughly categorized into two groups: events with “fake” missing energy due to mismeasurement of energies or other experimental effects, and those with “real” missing energy due to the underlying physics process which escapes the detector. In this analysis, the primary backgrounds from SM processes with E_T^{miss} and hadronic activity in the final state can be divided into three categories:

- *Multijet events*, with no true missing energy generating fake E_T^{miss} . Such all-hadronic events usually enter analysis signal regions due to mismeasurement of jet energy or sporadic noise.

- *Lost lepton events*, with a lepton in the final state that is either not identified, reconstructed, in the fiducial region of the detector, or sufficiently isolated to pass the lepton selections. The bulk of these events are due to leptonic W boson or top quark decays (which may also produce significant real E_T^{miss} with an associated neutrino) recoiling against jets.
- *Invisible Z events*, where a Z boson produced in association with jets decays into neutrinos. Neutrinos are weakly interacting particles which escape the detector without leaving any signature, and so this background is fundamentally very similar to many of the BSM signals which may appear in each search region. It is a dominant background in many signal regions, though the relative fraction decreases in regions with more b-tagged jets ($N_{\text{b-jets}}$).

In order to provide a robust prediction of each background that is not heavily dependent on simulation alone, each background is estimated using data-driven techniques as described in the following sections. The control regions for each background can be found in chapter 4.

5.2 Multijet Estimate

The estimate of the SM multijet background utilizes two different techniques depending on the number of jets in a given signal region. For regions with two or more jets, the background is estimated by extrapolating to regions with high $\Delta\phi_{\min}$ after inverting the $\Delta\phi_{\min}$ requirement. In the monojet signal regions, a sample enriched in unbalanced

dijet events is used to extrapolate to regions with low sub-leading jet momentum.

5.2.1 Multijet Signal Region Prediction

As outlined in section 4.5.3, the multijet background in control regions with two or more jets is estimated using a QCD-enriched sample where H_T triggers are used to select events with an inverted $\Delta\phi_{\min}$ cut. The ratio of events with high to low $\Delta\phi_{\min}$ (r_ϕ) is modeled as a power law function in M_{T2} , as shown in equation 5.1.

$$r_\phi(M_{T2}) = \frac{N(\Delta\phi_{\min} > 0.3)}{N(\Delta\phi_{\min} < 0.3)} = a \cdot M_{T2}^b \quad (5.1)$$

The power-law dependence of r_ϕ on M_{T2} is verified in simulation for events with $M_{T2} > 60 \text{ GeV}$, as illustrated in figure 5.1. Because the dominant source of E_T^{miss} in low M_{T2} events is not necessarily due to jet mismeasurement, the fit is performed in the window $60 < M_{T2} < 100 \text{ GeV}$ except in extreme H_T regions where a lower bound of 70 GeV is used (as a conservative measure since these regions have high statistics). The upper bound of 100 GeV is chosen such that the contamination of electroweak and top-quark events is small compared to the QCD multijet yield. A systematic uncertainty in the r_ϕ value is assigned based on variations of the fit window, by shifting the lower boundary of the window in either direction while preserving the fit statistics by shifting the upper boundary of the window. The systematic uncertainty is then taken as the maximal deviation of all such variation with respect to the nominal window.

Due to the total integrated luminosity available and the deliberate suppression rate with which some H_T triggers save events (known as the trigger *prescale*), the statistics

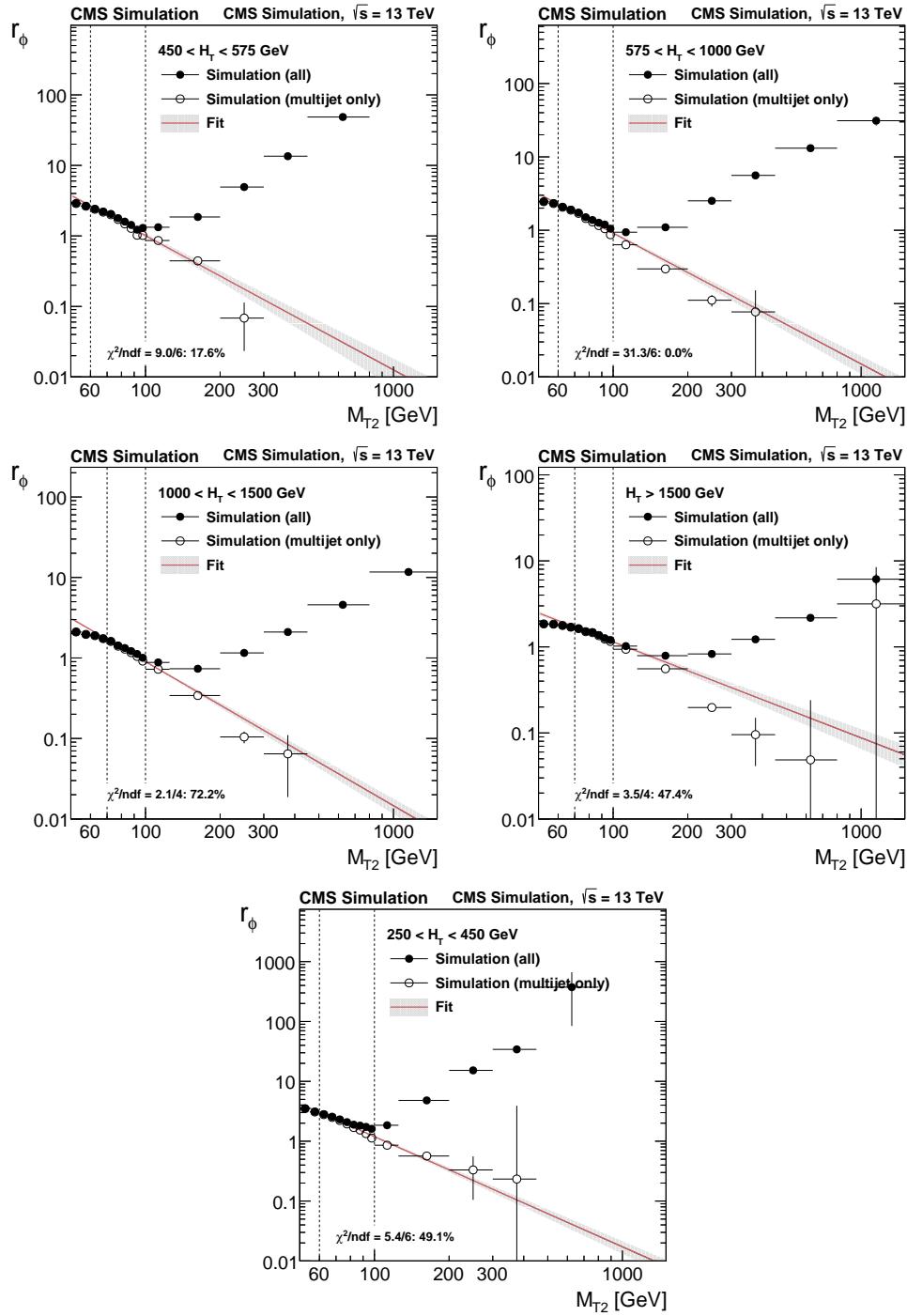


Figure 5.1: Simulated distribution of the r_ϕ ratio as a function of M_{T2} for the low (top left), medium (top right), high (medium left), extreme (medium right), and very low (bottom) H_T regions. Solid points represent the total simulated background, and hollow points show the QCD multijet contribution only. The red line and grey error band illustrates the best fit to a power law function performed in the dashed line fit window

Table 5.1: Relative uncertainty of f_j and r_b associated with the assumed invariance with respect to M_{T2} and $\Delta\phi$ (and H_T for r_b).

Observable	f_{23}	f_{46}	f_{7+}	r_0	r_1	r_2	r_{3+}
Syst. Error	25%	7%	20%	8%	20%	35%	70%

are sufficient to perform fits in H_T regions inclusive in N_{jets} and $N_{\text{b-jets}}$. The inclusive multijet background estimate can be determined using r_ϕ as a function of M_{T2} $N_{\text{inc}}^{SR} = r_\phi(M_{T2}) \cdot N_{\text{inc}}^{CR}(M_{T2})$. The final estimate in a given $N_{\text{jets}}\text{-}N_{\text{b-jets}}$ bin can be determined as in equation 5.2, where f_j is the fraction of multijet events in bin N_{jets} in a given H_T bin, and r_b is the ratio of events with $N_{\text{b-jets}}$ b-tags in a given N_{jets} bin.

$$N_{j,b}^{SR}(M_{T2}) = r_\phi(M_{T2}) \cdot N_{\text{inc}}^{CR}(M_{T2}) \cdot f_j(H_T) \cdot r_b(N_{\text{jets}}) \quad (5.2)$$

The values of f_j and r_b are measured in data using QCD-enriched regions with an inverted $\Delta\phi$ requirement and $100 < M_{T2} < 200$. Based on simulation, f_j and r_b do not significantly depend on M_{T2} , and r_b is independent of H_T . Since f_j is dependent on H_T and r_b on N_{jets} , different values of f_j are measured in each inclusive H_T region, and different values of r_b are measured in inclusive N_{jets} regions. Figures 5.2 and 5.3 illustrates the data-driven values of f_j and r_b respectively, along with a comparison to simulation. The robustness of invariance with respect to M_{T2} and $\Delta\phi$ (and H_T for r_b) is calculated by making several variations of the aforementioned variables and measuring consequent variations in f_j and r_b , and is summarized in table 5.1.

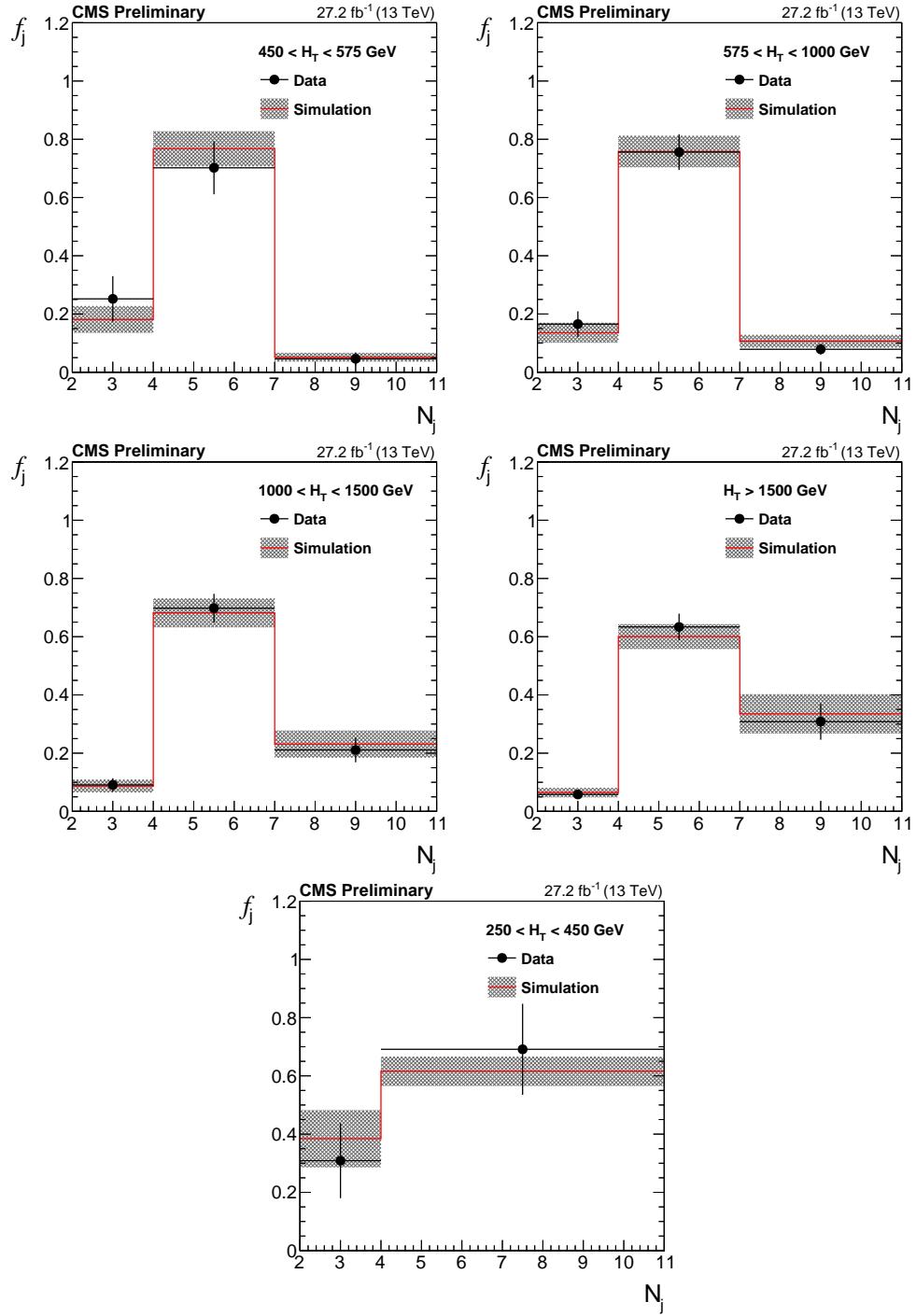


Figure 5.2: The values of f_j as measured in data in different H_T regions, compared to simulation. The uncertainties include both the statistical error and the systematic sources as listed in table 5.1

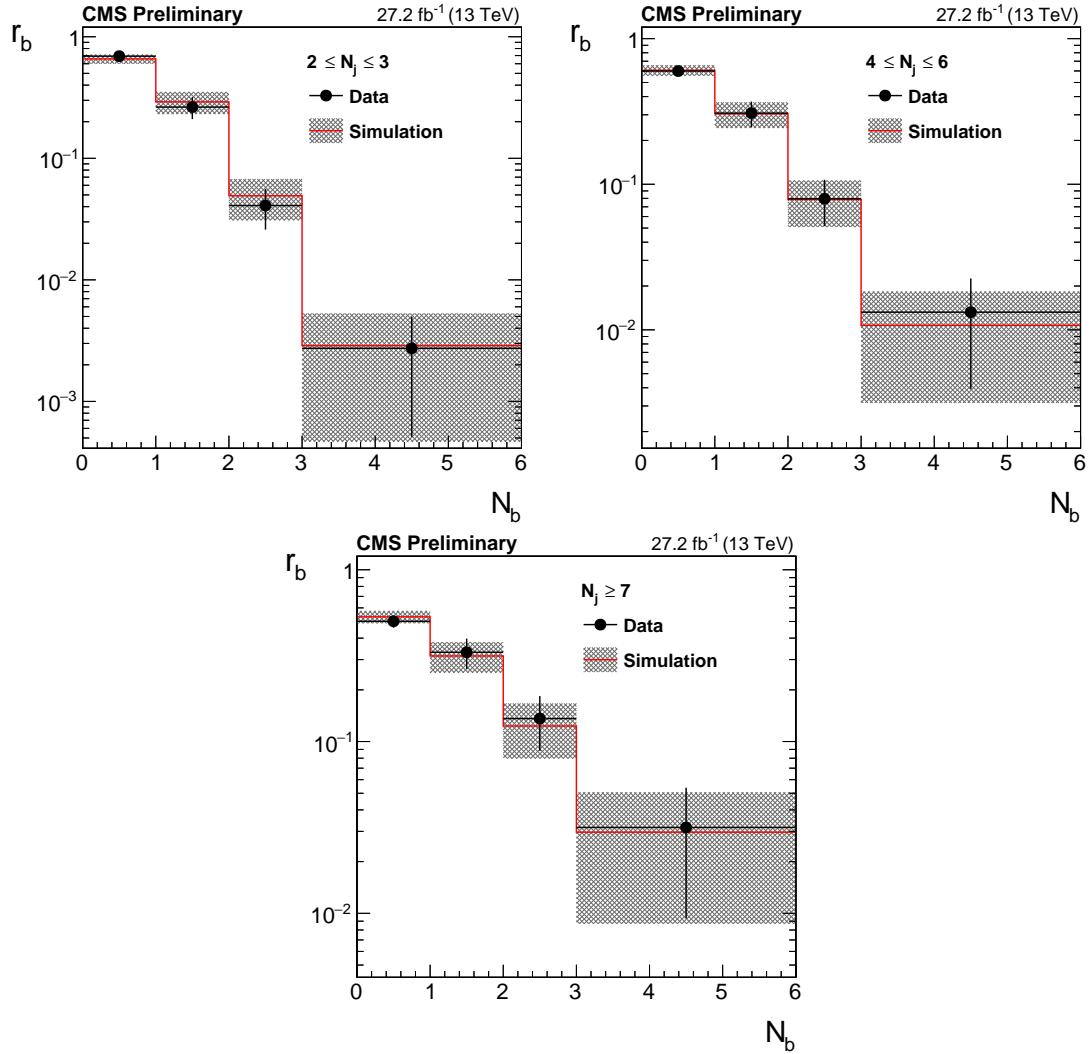


Figure 5.3: The values of r_b as measured in data in different N_{jets} regions, compared to simulation. The uncertainties include both the statistical error and the systematic sources as listed in table 5.1

5.2.2 Monojet Signal Region Prediction

The multijet background in control regions with a single jet cannot be estimated using the $\Delta\phi$ technique since E_T^{miss} is usually very similar to the H_T in these events and typically anti-aligned with the jet. As outlined in section 4.5.3, a separate control region is devised which instead selects dijet events that are orthogonal to the multijet signal regions because of an inverted $\Delta\phi(j_{1234}, E_T^{\text{miss}})$ cut (and orthogonal to the monojet SR with the presence of multiple jets).

The sub-leading jet momentum in this CR can be seen in figure 5.4. Because jets with p_T below 30 GeV are not considered in this analysis, monojet events can be classified as those with $p_T^{\text{jet}2} < 30 \text{ GeV}$, and the CR is used to extrapolate into the regime where $p_T^{\text{jet}2}$ is small. With decreasing $p_T^{\text{jet}2}$, events appear more imbalanced and approximate the topology of a true monojet event, as depicted in figure 5.5.

The predicted yield of multijet background in a monojet $p_T^{\text{jet}1}$ bin is determined according to equation 5.3, where f_{QCD} is the fraction of QCD events as measured in the region with $30 < p_T^{\text{jet}2} < 60 \text{ GeV}$ and N_{data} is the yield in data of dijet events with $30 < p_T^{\text{jet}2} < 60 \text{ GeV}$. Assuming $N_{\text{data}}(0 - 30) < N_{\text{data}}(30 - 60)$, this estimate provides an upper bound on the total multijet background contribution in each monojet CR. A systematic uncertainty of 50% on the QCD fraction f_{QCD} is assigned as a conservative estimate.

$$N_{\text{QCD}}(p_T^{\text{jet}1}) = N_{\text{data}}(30 - 60, p_T^{\text{jet}1}) \cdot f_{\text{QCD}}(30 - 60, p_T^{\text{jet}1}) \quad (5.3)$$

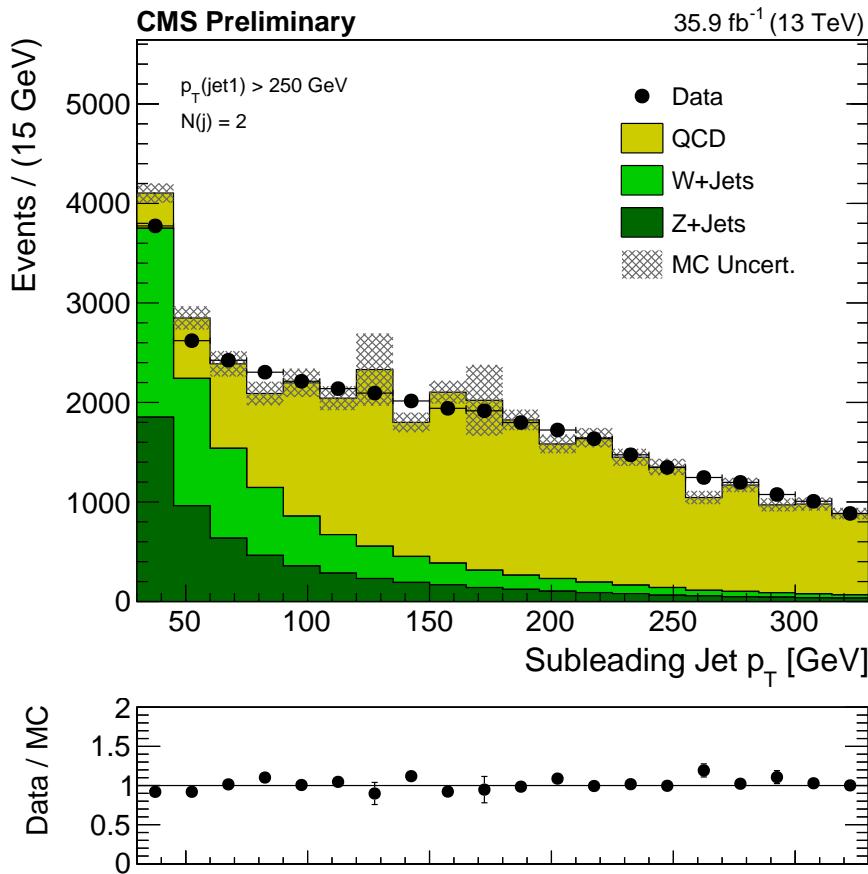


Figure 5.4: The transverse momentum of the sub-leading jet for dijet events in the monojet QCD background control region. The total yield of the simulation is normalized to the overall yield in data.

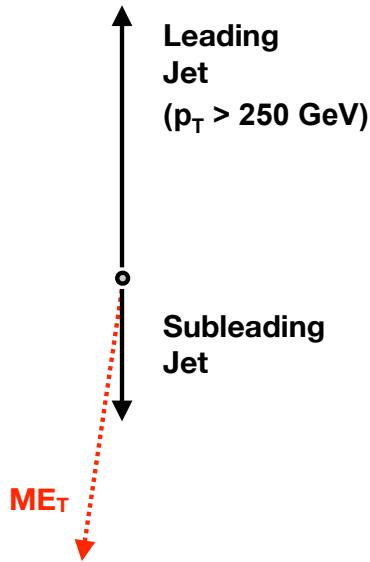


Figure 5.5: An illustration of “unbalanced” dijet events. As the momentum of the sub-leading jet decreases, E_T^{miss} is more anti-aligned with the primary jet and approaches the topology of a monojet event.

5.3 Lost Lepton Estimate

The lost lepton background is predicted by taking the yield of single lepton events in similar kinematic regions, measuring the ratio of events where the lepton is lost to those where it is found, and extrapolating into the lost lepton regime. The control region for the lost lepton estimate is described in detail in section 4.5.1, and is comprised of events in data with the same signal triggers and preselection, with the exception of an inverted lepton veto and an additional requirement on lepton M_T (to reduce signal contamination).

5.3.1 Background Prediction

The estimate in a given signal region is obtained from a control region using transfer factors as described in equation 5.4, where N_{1L}^{CR} is the number of events in the corresponding CR, $R_{MC}^{0l/1l}$ is the ratio of zero-lepton to one-lepton events derived from simulation, and $k(M_{T2})$ is the transfer factor into bins of M_{T2} in a given topological region. The ratio $R_{MC}^{0l/1l}$ is measured in Monte Carlo after normalizing the yield to data in each CR and accounting for differences in lepton efficiency between data and simulation, and also factors in lepton acceptance, efficiency of reconstruction and identification, as well as contributions from W bosons decaying through τ leptons to a hadronic final state. The M_{T2} transfer factor $k(M_{T2})$ is taken from data where statistics permit, and uses information from simulation to project events into bins of M_{T2} in the low-statistics regime.

$$N_{LL}^{SR}(H_T, N_j, N_b, M_{T2}) = N_{1L}^{CR}(H_T, N_j, N_b, M_{T2}) \cdot R_{MC}^{0l/1l}(H_T, N_j, N_b, M_{T2}) \cdot k(M_{T2}) \quad (5.4)$$

To reduce the dependence of the estimate on the M_{T2} shape modeling in simulation, the transfer factor $k(M_{T2})$ uses a combination of data and simulation information. In each topological region, the greatest M_{T2} bin is iteratively combined with the next-greatest bin until the total expected SM background yield in simulation is at least 50 events. These combined bins together form the CR for a range of M_{T2} values, where the fraction of events falling in a particular M_{T2} bin, $k(M_{T2})$, is determined from simulation. In all the other M_{T2} bins in the topological region, statistics are sufficient for a direct measurement in data and $k(M_{T2}) = 1$. The extrapolation point for each topological region can be found in table 5.2. The shape modeling in simulation is verified in data by selecting an inclusive

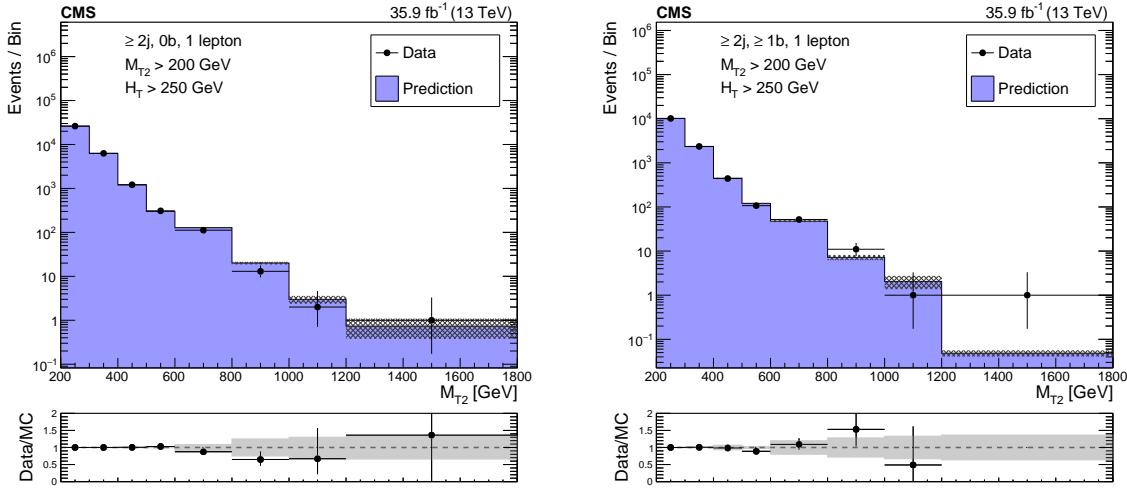


Figure 5.6: The M_{T2} shape in data and simulation using the single lepton CR selection, for events with zero b-tagged jets (left) or at least one b-tagged jet (right). The simulation is normalized to data in each topological region before summing to create the inclusive region. The hatched bands in each upper plot show the MC statistical uncertainty, while the shaded bands in each lower plot represent the systematic shape uncertainty.

sample enriched in either W boson or $t\bar{t}$ production (using the number of b-tags in the event), and predicting the M_{T2} distribution using simulation after normalizing simulation yield to data and summing all the control regions together, as seen in figure 5.6.

5.3.2 Systematic Uncertainties

Several sources of uncertainty are assessed for the lost lepton estimate, including those associated with the topological transfer factor and the M_{T2} shape modeling in simulation. The full list of systematic uncertainties is as follows:

- *Control region statistical error:* the Poisson error on the number of observed events in data is taken as a correlated uncertainty in each signal region using the same control region. The error is uncorrelated amongst M_{T2} bins, except in regions sharing merged

Table 5.2: The last M_{T2} bin and the M_{T2} extrapolation point for each topological region, beyond which shape data from simulation is used to extrapolate the lost lepton estimate into M_{T2} bins.

Control Region H_T [GeV]	N_j	N_b	Extrapolation Point M_{T2} [GeV]	Last M_{T2} bin M_{T2} [GeV]
[250, 450]	2-3	0	> 300	> 400
[250, 450]	2-3	1	> 300	> 400
[250, 450]	2-3	2	> 300	> 400
[250, 450]	4+	0	> 300	> 400
[250, 450]	4+	1	> 300	> 400
[250, 450]	4+	2	> 200	> 400
[250, 450]	2-3	3+	> 200	> 400
[450, 575]	2-3	0	> 400	> 500
[450, 575]	2-3	1	> 300	> 500
[450, 575]	2-3	2	> 200	> 500
[450, 575]	4-6	0	> 400	> 500
[450, 575]	4-6	1	> 300	> 500
[450, 575]	4-6	2	> 300	> 500
[450, 575]	7+	0	> 200	> 400
[450, 575]	7+	1-2	> 200	> 400
[450, 575]	2-6	3+	> 200	> 500
[575, 1000]	2-3	0	> 600	> 800
[575, 1000]	2-3	1	> 400	> 800
[575, 1000]	2-3	2	> 200	> 800
[575, 1000]	4-6	0	> 400	> 800
[575, 1000]	4-6	1	> 400	> 800
[575, 1000]	4-6	2	> 400	> 800
[575, 1000]	7+	0	> 300	> 800
[575, 1000]	7+	1-2	> 300	> 600
[575, 1000]	2-6	3+	> 200	> 600
[1000, 1500]	2-3	0	> 400	> 1200
[1000, 1500]	2-3	1	> 200	> 1200
[1000, 1500]	2-3	2	> 200	> 1000
[1000, 1500]	4-6	0	> 400	> 1200
[1000, 1500]	4-6	1	> 400	> 1200
[1000, 1500]	4-6	2	> 200	> 1000
[1000, 1500]	7+	0	> 200	> 1000
[1000, 1500]	7+	1-2	> 200	> 800
[1000, 1500]	2-6	3+	> 200	> 600
[1500, ∞]	2-3	0	> 200	> 1400
[1500, ∞]	2-3	1	> 200	> 1000
[1500, ∞]	2-3	2	> 200	> 400
[1500, ∞]	4-6	0	> 400	> 1400
[1500, ∞]	4-6	1	> 200	> 1400
[1500, ∞]	4-6	2	> 200	> 800
[1500, ∞]	7+	0	> 200	> 1000
[1500, ∞]	7+	1-2	> 200	> 800
[1500, ∞]	2-6	3+	> 200	> 600

M_{T2} bins for the estimate.

- *Monte Carlo statistical error*: where the simulation is used to compute the transfer factor (and in some cases, the M_{T2} shape), the MC statistical uncertainty ranges from 1-50%, taken as uncorrelated amongst all bins.
- *Electron and muon selection efficiency*: the reconstruction and identification efficiency of electrons and muons is computed using a procedure known as *tag and probe*, where leptons from $Z \rightarrow \ell\ell$ decays are used to evaluate the lepton id efficiencies as a function of lepton p_T and η and lepton isolation efficiencies as a function of p_T and nearby activity. The scale factors applied to simulation to compensate for such effects approach unity with uncertainties on the order of a few percent, and are correlated across all bins. The maximum effect of this uncertainty is up to 7% in some signal regions.
- *Tau selection efficiency*: the efficiency for hadronically decaying taus is classified according to the number of charged particles in the τ decay, whether a *1-prong* tau leaving a single charged track, or a *3-prong* tau leaving three charged tracks in the final state. This efficiency is measured in simulation by measuring the isolation efficiency as a function of candidate p_T for electrons, muons, and taus in various decay modes. Based on half the difference in efficiency between 1-prong taus and muons, an uncertainty of 10% is taken for 1-prong taus which also covers any differences in veto efficiency as a function of the primary kinematic variables. For 3-prong taus, the PF hadron veto is very inefficient since most fail the isolation requirement (the typical

selection efficiency is 8%), and a 100% relative uncertainty is taken to cover any differences as a function of kinematics. These uncertainties are correlated amongst all bins.

- *M_T cut efficiency:* the use of simulation to compute the CR-to-SR transfer factor also relies on the satisfactory modeling of the M_T cut in Monte Carlo. A sample of $Z \rightarrow \ell\ell$ events is selected in data with one of the leptons “deleted” from the event to mimic a leptonic W boson decay, and compared with simulation. Based on data-MC agreement, a correlated error of 3% is taken across all bins.
- *b-tagging efficiency:* the effect of varying the b-tag scale factor efficiency is calculated in each bin, and taken as a correlated error amongst all bins. The maximum effect of this variation is about 4% in some bins.
- *Jet energy corrections:* by varying the jet energy scales across all bins, a maximum deviation of about 5% is observed in regions with sufficient statistics, and a correlated error of 5% is taken across all bins.
- *MC renormalization and factorization scales:* the overall effect of varying the simulation renormalization and factorization scales of the underlying physics processes (and subsequent effect on event kinematics) is computed separately in each bin. Taken as correlated across all bins, it is typically on the order of a few percent, but ranges up to 10% in some regions.

- M_{T2} shape uncertainty: in regions where the simulation is used to model the M_{T2} distribution, additional variations of the renormalization and factorization scales, parton distribution functions, b-tagging scale factor uncertainties, and jet energy scale uncertainties are performed to measure their effect on the M_{T2} shape modeling. The most significant impact is seen in the highest M_{T2} bins from theoretical uncertainties ($\sim 15\%$), and up to 40% in low statistics bins due to jet energy scale variations. With this in mind, the shape uncertainty (in regions where MC M_{T2} shape modeling is used) is assigned as a linear morphing of the M_{T2} shape starting in the first bin from which MC extrapolation is used, growing to 40% in the final bin. The shape morphing in every distinct topological region is taken as an uncorrelated error.

5.3.3 Signal Contamination

Nearly every control region in this analysis (including those for multijet and invisible Z backgrounds) is not only orthogonal to the signal region selection, but also crafted such that any potential BSM signal contribution to any CR is negligible and will not bias the background estimate. However, certain SUSY simplified models yield final states with prompt lepton decays — in some cases kinematically similar to SM $t\bar{t}$ decays — which may be non-negligible in the lost lepton CR. To account for this effect when calculating limits on such models (described in detail in section 6.2), the amount by which the lost lepton background is overestimated is modeled as a loss in signal efficiency. The modified signal yield $N_{\text{sig}}^{SR'}$ is defined in equation 5.5, where N_{sig}^{SR} and N_{sig}^{CR} represent the simulated

signal yield in the signal and control regions respectively, and TF the total transfer factor used in the lost lepton estimate for a given signal region.

$$N_{\text{sig}}^{SR'} = N_{\text{sig}}^{SR} - TF \cdot N_{\text{sig}}^{CR} \quad (5.5)$$

The correction is most significant for simplified models with stop decays, where the stop-neutralino mass splitting is close to the SM top quark mass. In such cases, the signal contamination is maximally 5% of the expected background yields in each CR for mass points near the expected exclusion limits at high mass.

5.4 Invisible Z Estimate

The invisible Z background is the dominant SM contribution in many signal regions due to the irreducible nature of the underlying physics. Because the primary interaction produces a massive particle decaying to an invisible final state ($Z \rightarrow \nu\bar{\nu}$) recoiling against hadronic activity, the signature is fundamentally similar to that of the BSM physics the search is designed to target and difficult to reduce by conventional cuts removing SM background contributions. A robust method leveraging the well-understood Drell-Yan process ($Z \rightarrow l^+l^-$) — an interaction kinematically similar to the invisible Z background — is used to predict the expected SM contribution based on data. However the ratio the branching fraction of $Z \rightarrow \nu\bar{\nu}$ decays (20%) to $Z \rightarrow l^+l^-$ decays (7%) presents a challenge, requiring a robust extrapolation method that can overcome the statistical limitations of the control region while minimizing the reliance on Monte Carlo modeling of the kinematics.

5.4.1 Background Prediction

The invisible Z background is estimated using dilepton events selected in data. The control region consists of events selected with dilepton triggers, with addition requirements that the leptons are of the same flavor and opposite sign. The momentum of the leading and trailing lepton must also be at least 100 GeV and 30 GeV, respectively, and the invariant mass of the lepton system $m_{\ell\ell}$ must be withing 20 GeV of the Z boson mass. The individual CRs are then constructed by removing the dilepton system from the event and applying the baseline preselection requirements as for the signal regions. A detailed description of the CR can be found in section 4.5.2.

The invisible Z prediction is computed as described by equation 5.6, where $N_{\ell\ell}^{\text{CR(SF)}}$ is the number of events in the dilepton same-flavor (SF) region, $N_{\ell\ell}^{\text{CR(OF)}}$ the number of events in the dilepton opposite-flavor (OF) region, $R^{\text{SF/OF}}$ the SF-OF transfer factor, $R_{\text{MC}}^{Z \rightarrow \nu\bar{\nu}/Z \rightarrow l^+l^-}$ the transfer factor from $Z \rightarrow l^+l^-$ to $Z \rightarrow \nu\bar{\nu}$ events, and $k(M_{\text{T2}})$ the transfer factor into bins of M_{T2} . Each factor is explained in detail below.

$$N_{Z \rightarrow \nu\bar{\nu}}^{\text{SR}}(H_{\text{T}}, N_{\text{j}}, N_{\text{b}}, M_{\text{T2}}) = \left[N_{\ell\ell}^{\text{CR(SF)}}(H_{\text{T}}, N_{\text{j}}, N_{\text{b}}) - N_{\ell\ell}^{\text{CR(OF)}}(H_{\text{T}}, N_{\text{j}}, N_{\text{b}}) \cdot R^{\text{SF/OF}} \right] \times R_{\text{MC}}^{Z \rightarrow \nu\bar{\nu}/Z \rightarrow l^+l^-}(H_{\text{T}}, N_{\text{j}}, N_{\text{b}}) \cdot k(M_{\text{T2}}) \quad (5.6)$$

The second term in equation 5.6 ($N_{\ell\ell}^{\text{CR(OF)}} \cdot R^{\text{SF/OF}}$) is a correction factor applied to the control region yield to correct for the contribution from processes producing SF and OF event, or *flavor-symmetric processes* such as $t\bar{t}$ production. To compensate for this contribution, a separate control region enriched in $t\bar{t}$ events is selected from data by using the same selections as the invisible Z CR, except for an inverted selection on the

dilepton p_T and $m_{\ell\ell}$ requirements. The ratio of $t\bar{t}$ events is then measured directly from data by counting the yield of SF (ee or $\mu\mu$) and OF ($e\mu$ or μe) events in this CR. The ratio $R^{\text{SF}/\text{OF}}$ is expected to be close to unity based on the underlying physics process, but due to varying acceptance and efficiency effects for different flavor leptons is measured as $R^{\text{SF}/\text{OF}} = 1.13 \pm 0.15$, and is stable with respect to event kinematics as shown in figure 5.7. The Drell-Yan yield in each control region, as well as the SF yield, OF yield, and transfer factors can be found in tables 5.3 and 5.4.

The ratio of invisible Z events to Drell-Yan events, $R_{\text{MC}}^{Z \rightarrow \nu\bar{\nu}/Z \rightarrow l^+l^-}$, is measured in simulation for each CR. The ratio accounts for the branching fraction differences between $Z \rightarrow l^+l^-$ and $Z \rightarrow \nu\bar{\nu}$ decays, as well as differences in lepton acceptance and efficiency for dilepton pairs in the CR (including corrections for differences in lepton efficiency between data and simulation).

The transfer factor $k(M_{\text{T2}})$ uses a combination of data and simulation information to reduce the dependence of the prediction on the M_{T2} shape modeling in simulation. Based on simulation, the M_{T2} shape in ever H_{T} region is independent of $N_{\text{b-jets}}$. In addition, in the extreme H_{T} region $H_{\text{T}} > 1500 \text{ GeV}$, the shape is also independent of N_{jets} . Based on these findings, M_{T2} shape templates are constructed for each H_{T} and N_{jets} region (inclusive in $N_{\text{b-jets}}$ ¹), with one single template for the extreme H_{T} region (which is also inclusive in N_{jets}).

The M_{T2} distribution, $k(M_{\text{T2}})$, for each topological region is constructed from

¹The only exception is for regions with more than 2 b-tags (2j3b and 2-6j3b), where at least 3 jets are required so as to avoid biasing the N_{jets} distribution when requiring more b-tags than jets.

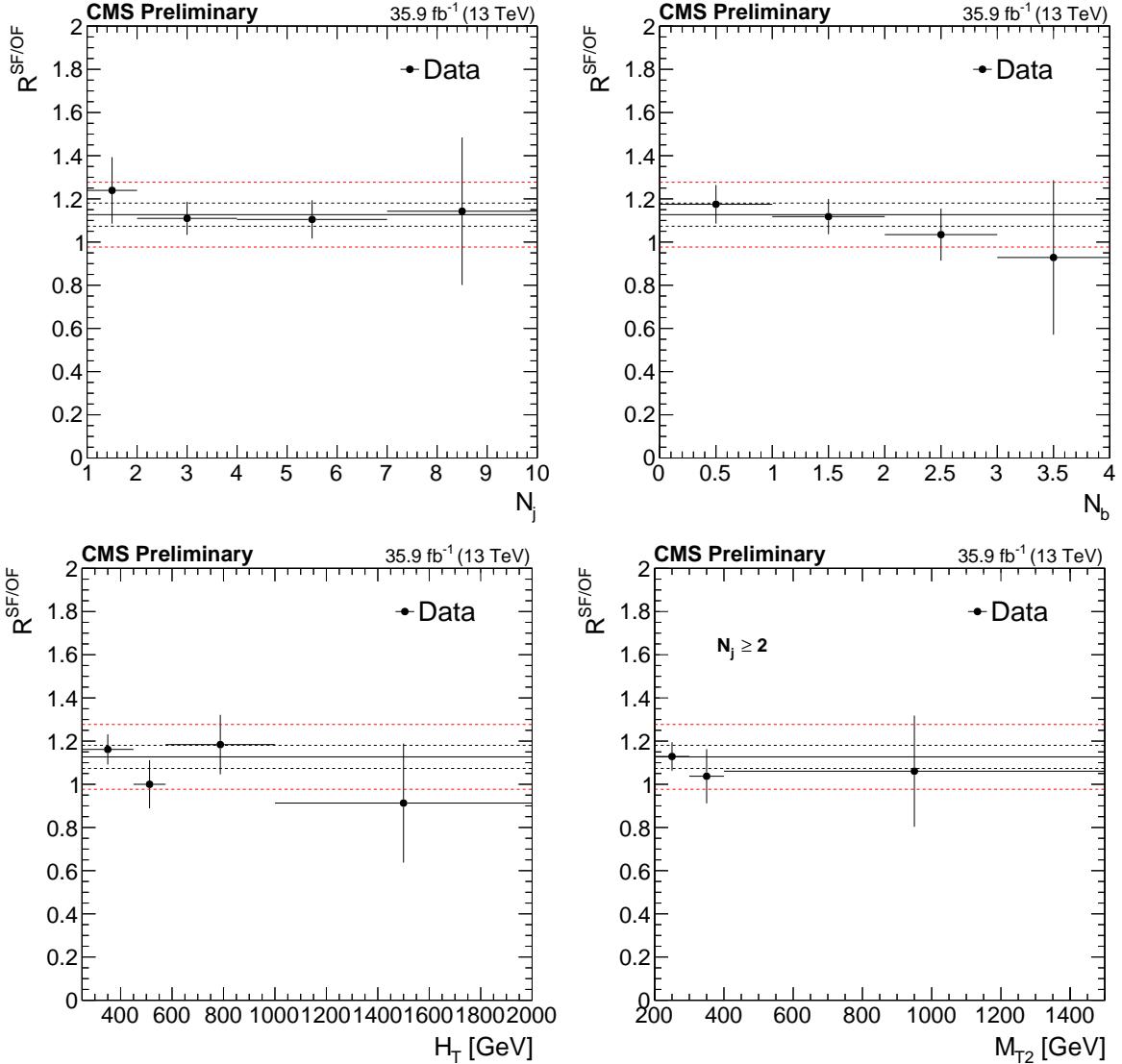


Figure 5.7: The ratio of same-flavor to opposite-flavor events in a $t\bar{t}$ enriched control region, as a function of N_{jets} (top left), $N_{\text{b-jets}}$ (top right), H_{T} (bottom left), and M_{T2} (bottom left). The solid black line corresponds to a constant value of 1.13 ± 0.15 , while the dashed black lines correspond to the statistical uncertainty and the dashed red lines the total systematic uncertainty.

Table 5.3: The control region predicted Drell-Yan (DY) yield, SF yield, OF yield, purity (the rfraction of $Z \rightarrow l^+l^-$ events), and ratio $R_{\text{MC}}^{Z \rightarrow \nu\bar{\nu}/Z \rightarrow l^+l^-}$ for the very low, low, and medium H_T topological regions. Note that the 7+ jet regions with b-tags (marked with an asterisk) share the same CR, and the fraction of events with different numbers of b-tags is folded into the ratio.

Invisible Z							
Region			DY Yield	SF Yield	OF Yield	Purity	Ratio
H_T [GeV]	N_j	N_b					
[250, 450]	2-3	0	8707.6	8749	37	0.995 ± 0.001	5.61
[250, 450]	2-3	1	986.0	1005	17	0.981 ± 0.004	5.47
[250, 450]	2-3	2	122.8	125	2	0.982 ± 0.012	5.31
[250, 450]	4+	0	1124.5	1129	4	0.996 ± 0.002	5.88
[250, 450]	4+	1	215.2	223	7	0.965 ± 0.012	5.70
[250, 450]	4+	2	35.9	37	1	0.970 ± 0.028	5.52
[250, 450]	2-3	3+	7.0	7	0	1.000 ± 0.000	6.29
[450, 575]	2-3	0	1862.7	1875	11	0.993 ± 0.002	5.08
[450, 575]	2-3	1	273.2	281	7	0.972 ± 0.010	4.89
[450, 575]	2-3	2	19.8	22	2	0.898 ± 0.064	4.75
[450, 575]	4-6	0	650.9	661	9	0.985 ± 0.005	5.49
[450, 575]	4-6	1	143.4	149	5	0.962 ± 0.016	5.58
[450, 575]	4-6	2	27.6	31	3	0.892 ± 0.056	5.21
[450, 575]	7+	0	7.0	7	0	1.000 ± 0.000	5.56
[450, 575]	7+	1	4.0	4	0	1.000 ± 0.000	5.18 ^(*)
[450, 575]	7+	2	4.0	4	0	1.000 ± 0.000	1.87 ^(*)
[450, 575]	2-6	3+	7.0	7	0	1.000 ± 0.000	5.78
[450, 575]	7+	3+	4.0	4	0	1.000 ± 0.000	0.48 ^(*)
[575, 1000]	2-3	0	1347.0	1356	8	0.993 ± 0.002	4.76
[575, 1000]	2-3	1	157.3	164	6	0.959 ± 0.015	4.55
[575, 1000]	2-3	2	25.8	28	2	0.920 ± 0.051	4.68
[575, 1000]	4-6	0	812.4	818	5	0.993 ± 0.003	5.16
[575, 1000]	4-6	1	180.0	189	8	0.953 ± 0.015	5.09
[575, 1000]	4-6	2	26.3	33	6	0.796 ± 0.070	5.29
[575, 1000]	7+	0	32.9	34	1	0.967 ± 0.031	5.84
[575, 1000]	7+	1	15.6	19	3	0.823 ± 0.088	4.41 ^(*)
[575, 1000]	7+	2	15.6	19	3	0.823 ± 0.088	1.10 ^(*)
[575, 1000]	2-6	3+	5.8	8	2	0.720 ± 0.159	5.82
[575, 1000]	7+	3+	15.6	19	3	0.823 ± 0.088	0.22 ^(*)

Table 5.4: The control region predicted Drell-Yan (DY) yield, SF yield, OF yield, purity (the rfraction of $Z \rightarrow l^+l^-$ events), and ratio $R_{\text{MC}}^{Z \rightarrow \nu\bar{\nu}/Z \rightarrow l^+l^-}$ for the high and extreme H_T topological regions. Note that the 7+ jet regions with b-tags (marked with an asterisk) share the same CR, and the fraction of events with different numbers of b-tags is folded into the ratio.

Invisible Z							
Region			DY Yield	SF Yield	OF Yield	Purity	Ratio
H_T [GeV]	N_j	N_b					
[1000, 1500]	2-3	0	129.0	129	0	1.000 ± 0.000	4.76
[1000, 1500]	2-3	1	25.9	27	1	0.959 ± 0.038	4.63
[1000, 1500]	2-3	2	1.9	3	1	0.627 ± 0.279	4.73
[1000, 1500]	4-6	0	154.0	154	0	1.000 ± 0.000	5.10
[1000, 1500]	4-6	1	42.0	42	0	1.000 ± 0.000	4.97
[1000, 1500]	4-6	2	11.0	11	0	1.000 ± 0.000	5.07
[1000, 1500]	7+	0	19.0	19	0	1.000 ± 0.000	5.63
[1000, 1500]	7+	1	10.0	10	0	1.000 ± 0.000	4.63 ^(*)
[1000, 1500]	7+	2	10.0	10	0	1.000 ± 0.000	1.22 ^(*)
[1000, 1500]	2-6	3+	1.0	1	1	1.000 ± 1.000	4.25
[1000, 1500]	7+	3+	10.0	10	0	1.000 ± 0.000	0.20 ^(*)
[1500, ∞]	2-3	0	29.0	29	0	1.000 ± 0.000	5.00
[1500, ∞]	2-3	1	8.0	8	0	1.000 ± 0.000	4.66
[1500, ∞]	4-6	0	28.9	30	1	0.963 ± 0.035	5.09
[1500, ∞]	4-6	1	14.0	14	0	1.000 ± 0.000	5.25
[1500, ∞]	4-6	2	2.9	4	1	0.720 ± 0.224	4.80
[1500, ∞]	7+	0	5.0	5	0	1.000 ± 0.000	5.16
[1500, ∞]	7+	1	1.9	3	1	0.627 ± 0.279	3.97 ^(*)
[1500, ∞]	7+	2	1.9	3	1	0.627 ± 0.279	1.18 ^(*)
[1500, ∞]	2-6	3+	1.0	1	0	1.000 ± 0.000	6.27
[1500, ∞]	7+	3+	1.9	3	1	0.627 ± 0.279	0.18 ^(*)

dilepton events in data where statistics allows, and $Z \rightarrow \nu\bar{\nu}$ simulation where statistics are sparse. In each template region, the greatest M_{T2} bin is iteratively combined with the next-greatest bin until the total expected SM background yield in simulation is at least 50 events. For uncombined bins (where statistics in data is sufficient), the M_{T2} shape is taken directly from dilepton data, corrected for the ratio $R_{\text{MC}}^{Z \rightarrow \nu\bar{\nu}/Z \rightarrow l^+l^-}$. In the low-statistics combined regime, the M_{T2} shape in $Z \rightarrow \nu\bar{\nu}$ simulation is used to determine the fraction of events in each M_{T2} bin after normalizing the simulation yield to data in the combined bins. The extrapolation point after which the M_{T2} shape is based on simulation in each signal region can be found in tables 5.5 and 5.6.

The accuracy of the M_{T2} shape modeling in simulation is verified using other control samples enriched in γ , $W \rightarrow \ell\nu$, and $Z \rightarrow l^+l^-$ events in each H_T bin, as shown in figure 5.8. The γ -enriched sample is selected using photon triggers and requiring $p_T^\gamma > 180 \text{ GeV}$, with corrections applied for multijet background contributions and the ratio of M_{T2} distributions for photon to Z boson events, $R_{\text{MC}}^{Z/\gamma}$. The W and Z boson samples are selected in data using leptonic triggers, corrected to compensate for the contribution from top quark production, as well as the ratio of M_{T2} distributions, $R_{\text{MC}}^{Z/W}$ and $R_{\text{MC}}^{Z \rightarrow \nu\bar{\nu}/Z \rightarrow l^+l^-}$ respectively.

5.4.2 Systematic Uncertainties

Several sources of uncertainty are assessed for the invisible Z estimate, including those associated with the various transfer factors and the M_{T2} shape modeling in simulation. The dominant uncertainty in regions using an M_{T2} template constructed from data

Table 5.5: The M_{T2} extrapolation point for the very low, low, and medium H_T topological regions, beyond which shape data from simulation is used to extrapolate the invisible Z estimate into M_{T2} bins. “NA” indicates regions where the simulation shape is not used at all since dilepton statistics in data are sufficiently large to perform the estimate bin-by-bin.

Invisible Z				
Template Region			Extrapolation Point	
	H_T [GeV]	N_j	N_b	M_{T2} [GeV]
[250, 450]	2-3	0		NA
[250, 450]	2-3	1		NA
[250, 450]	2-3	2		NA
[250, 450]	4+	0		300
[250, 450]	4+	1		300
[250, 450]	4+	2		300
[250, 450]	2+	3+		NA
[450, 575]	2-3	0		400
[450, 575]	2-3	1		400
[450, 575]	2-3	2		400
[450, 575]	4-6	0		400
[450, 575]	4-6	1		400
[450, 575]	4-6	2		400
[450, 575]	7+	0		200
[450, 575]	7+	1		200
[450, 575]	7+	2		200
[450, 575]	2-6	3+		NA
[450, 575]	7+	3+		200
[575, 1000]	2-3	0		600
[575, 1000]	2-3	1		600
[575, 1000]	2-3	2		600
[575, 1000]	4-6	0		600
[575, 1000]	4-6	1		600
[575, 1000]	4-6	2		600
[575, 1000]	7+	0		200
[575, 1000]	7+	1		200
[575, 1000]	7+	2		200
[575, 1000]	2-6	3+		NA
[575, 1000]	7+	3+		200

Table 5.6: The M_{T2} extrapolation point for the high and extreme H_T topological regions, beyond which shape data from simulation is used to extrapolate the invisible Z estimate into M_{T2} bins. “NA” indicates regions where the simulation shape is not used at all since dilepton statistics in data are sufficiently large to perform the estimate bin-by-bin.

Invisible Z			
Template H_T [GeV]	Region N_j	N_b	Extrapolation Point M_{T2} [GeV]
[1000, 1500]	2-3	0	400
[1000, 1500]	2-3	1	400
[1000, 1500]	2-3	2	400
[1000, 1500]	4-6	0	400
[1000, 1500]	4-6	1	400
[1000, 1500]	4-6	2	400
[1000, 1500]	7+	0	200
[1000, 1500]	7+	1	200
[1000, 1500]	7+	2	200
[1000, 1500]	2-6	3+	NA
[1000, 1500]	7+	3+	200
[1500, ∞]	2-3	0	400
[1500, ∞]	2-3	1	400
[1500, ∞]	4-6	0	400
[1500, ∞]	4-6	1	400
[1500, ∞]	4-6	2	400
[1500, ∞]	7+	0	400
[1500, ∞]	7+	1	400
[1500, ∞]	7+	2	400
[1500, ∞]	2-6	3+	400
[1500, ∞]	7+	3+	400

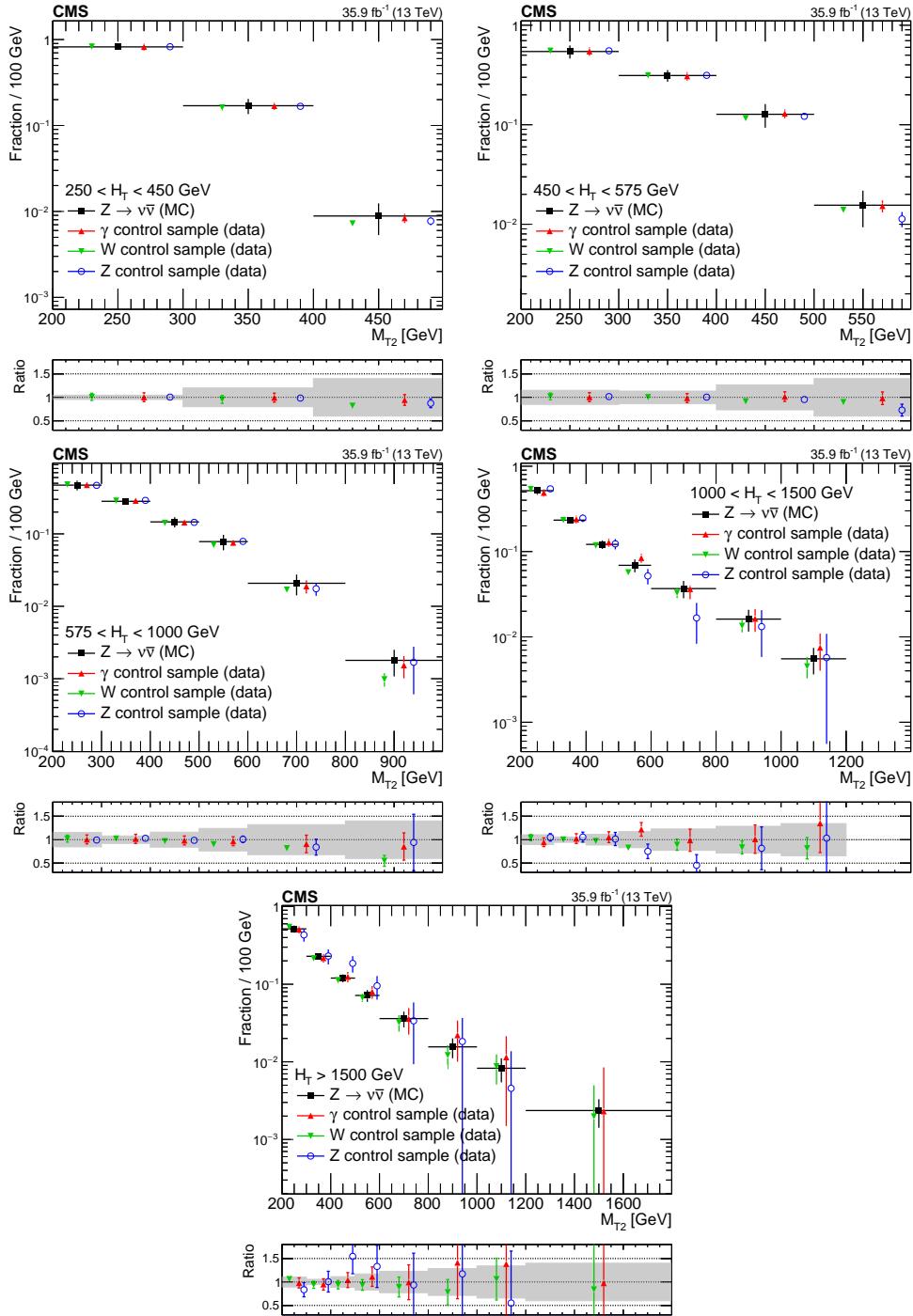


Figure 5.8: The M_{T2} shape distribution in $Z \rightarrow \nu\bar{\nu}$ simulation compared to γ , $W \rightarrow \ell\nu$, and $Z \rightarrow l^+l^-$ enriched samples in data, for each H_T region. The solid grey band indicates the systematic uncertainty associated with the M_{T2} shape modeling.

is the statistics of the template. The full list of systematic uncertainties is as follows:

- *Control region statistical error*: the Poisson error on the number of observed events in $Z \rightarrow l^+l^-$ data is taken as an uncorrelated uncertainty across all signal regions.
- $R_{\text{MC}}^{Z \rightarrow \nu\bar{\nu}/Z \rightarrow l^+l^-}$ *statistical error*: the statistical error associated with the number of MC events generated factors into the transfer factor.
- $R_{\text{MC}}^{Z \rightarrow \nu\bar{\nu}/Z \rightarrow l^+l^-}$ *systematic error*: a 5.5% uncertainty based on variations of lepton efficiency uncertainties as well as other modeling parameters (jet energy scales, factorization and renormalization scales, etc.) is applied as a correlated error in each topological region.
- *Flavor-symmetric subtraction statistical error*: a Poisson error based on the number of observed opposite-flavor events is assigned to the purity correction.
- *Flavor-symmetric subtraction systematic error*: a 15% uncertainty on the $t\bar{t}$ contamination is taken based on the $R^{\text{SF/OF}}$ uncertainty.
- *M_{T2} shape uncertainty*: in regions where the simulation is used to model the M_{T2} distribution, additional variations of the renormalization and factorization scales, parton distribution functions, b-tagging scale factor uncertainties, and jet energy scale uncertainties are performed to measure their effect on the M_{T2} shape modeling. The most significant impact is seen in the highest M_{T2} bins of up to 20%. In order to cover the uncertainty from additional electro-weak corrections not present in simulation (and possibly not covered by the above variations), a conservative

upper threshold of 40% is used, and the shape uncertainty (in regions where MC M_{T2} shape modeling is used) is assigned as a linear morphing of the M_{T2} shape starting in the first bin from which MC extrapolation is used, growing to 40% in the final bin. The shape morphing in every distinct topological region is taken as an uncorrelated error.

This chapter makes use of figures and tables from the M_{T2} paper and internal analysis note to illustrate the analysis design, methodology, and results. This work was made possible by contributions from the rest of the Surf & Turf group, our collaborators at ETH Zurich, and the many other CMS members in the SUSY group and beyond.

Chapter 6

Results

The techniques described in this analysis are applied to 35.9fb^{-1} of proton-proton collision data gathered at the LHC and recorded by the CMS detector. Observed yields in each signal region are further interpreted in the context of simplified SUSY models to establish new limits on the masses of hypothesized BSM particles.

6.1 Yields and Background Fits

The observed yield in the search regions is statistically compatible with the predicted background from SM processes. A summary of the total yield in each signal region and predicted background contribution relying only on the techniques described in chapter 5, referred to as *pre-fit* results, is illustrated for each topological region in figure 6.1, and the individual yield in each M_{T2} bin can be found in figures 6.2, 6.3, and 6.4. Approximately 80% of signal bins are consistent (based on the poisson error) with the nominal background

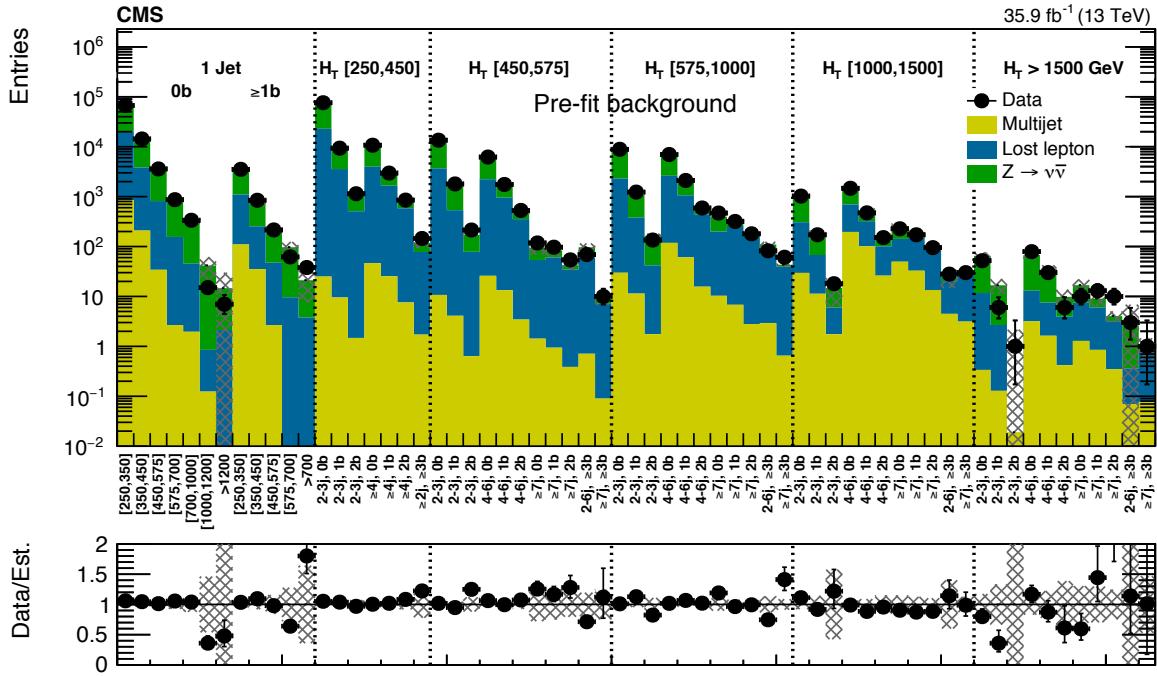


Figure 6.1: The data yield in each topological region compared to the pre-fit background prediction. The hatched bands illustrate the total uncertainty in the background estimate. Results in the monojet regions are binned in jet p_T , while those in the multijet regions are labeled according to N_{jets} and $N_{\text{b-jets}}$.

prediction, and approximately 95% of bins are consistent with the nominal background prediction including the total uncertainty on the background estimate (a full global background fit is also performed as described in the following paragraph). No significant excess above the predicted SM background is observed.

The total background estimate is further refined by performing a maximum likelihood fit to data in each signal region, referred to as *post-fit* results. The fit is performed using both background-only or background-plus-signal hypotheses to set limits on simplified physics models as described in section 6.2. The estimates and uncertainties on each background as described in chapter 5 are used as inputs to the fitting procedure, where the likelihood is constructed as a product of Poisson probability density functions for each

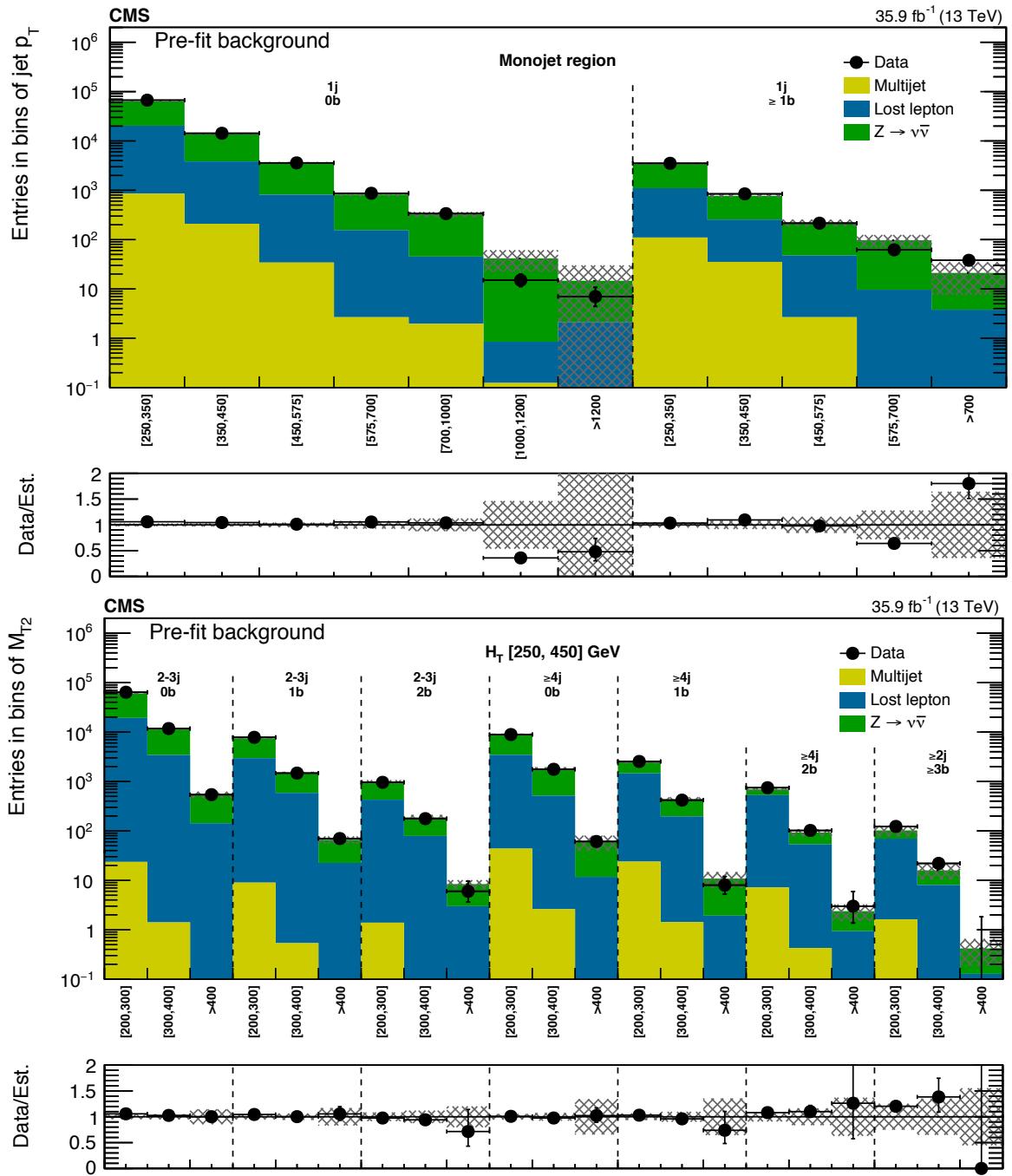


Figure 6.2: The data yield in the monojet and very-low H_T regions compared to the pre-fit background prediction. The hatched bands illustrate the total uncertainty in the background estimate. Results in the monojet regions are binned in jet p_T in units of GeV, while those in the multijet regions are labeled according to M_{T2} bin in units of GeV.

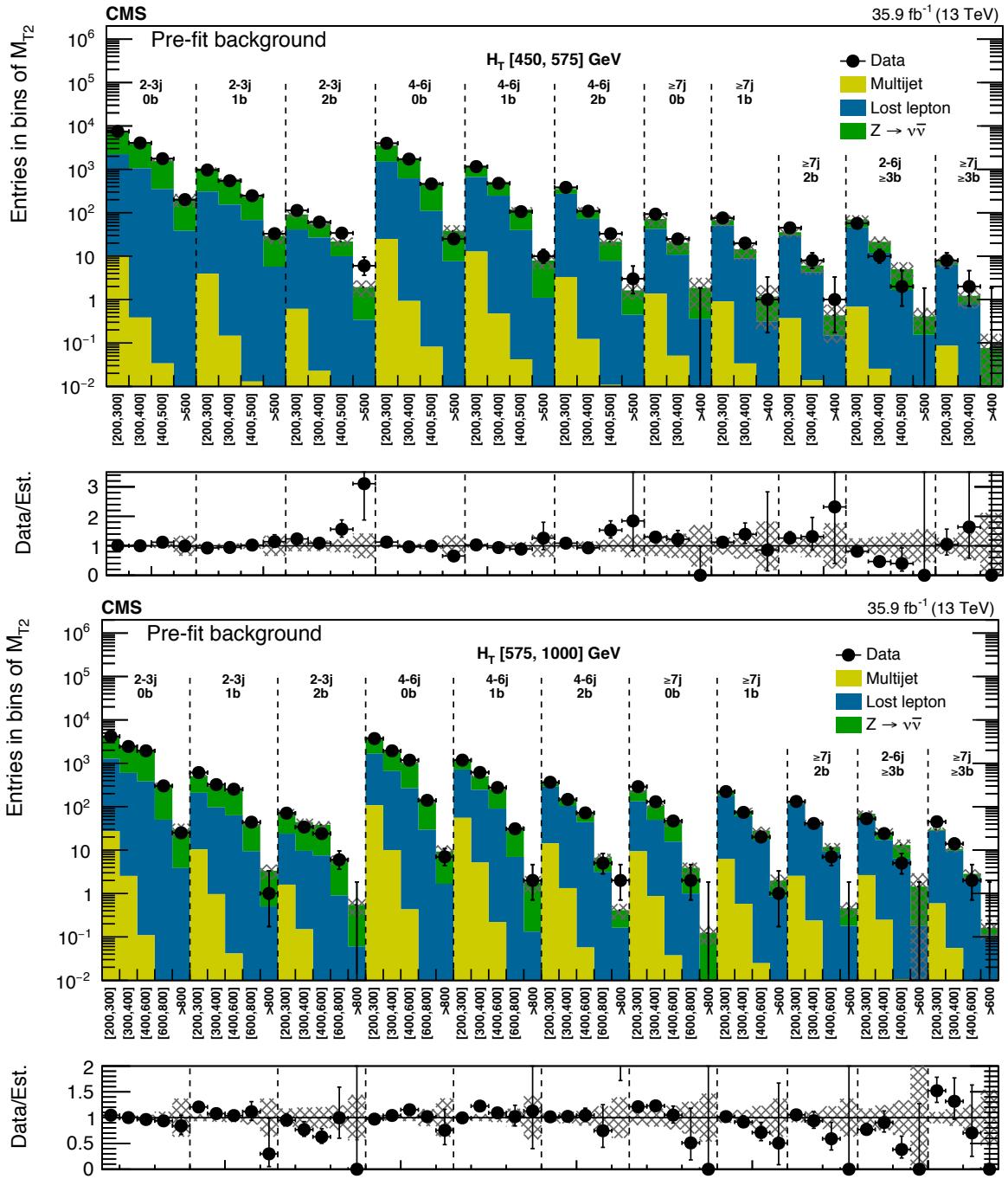


Figure 6.3: The data yield in the low H_T and medium H_T regions compared to the pre-fit background prediction. The hatched bands illustrate the total uncertainty in the background estimate. Results are labeled according to M_{T2} bin in units of GeV.

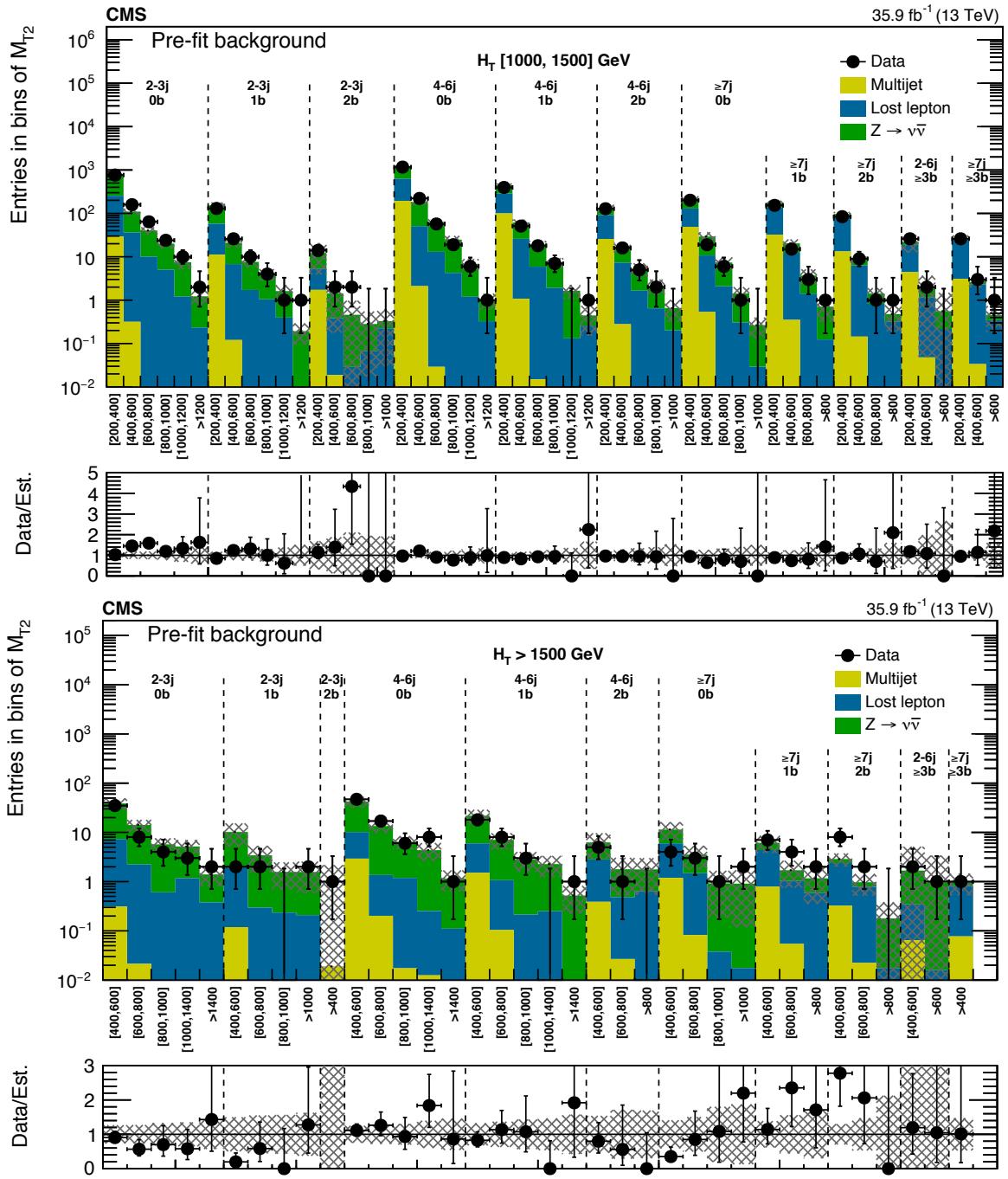


Figure 6.4: The data yield in the high H_T and extreme H_T regions compared to the pre-fit background prediction. The hatched bands illustrate the total uncertainty in the background estimate. Results are labeled according to M_{T2} bin in units of GeV.

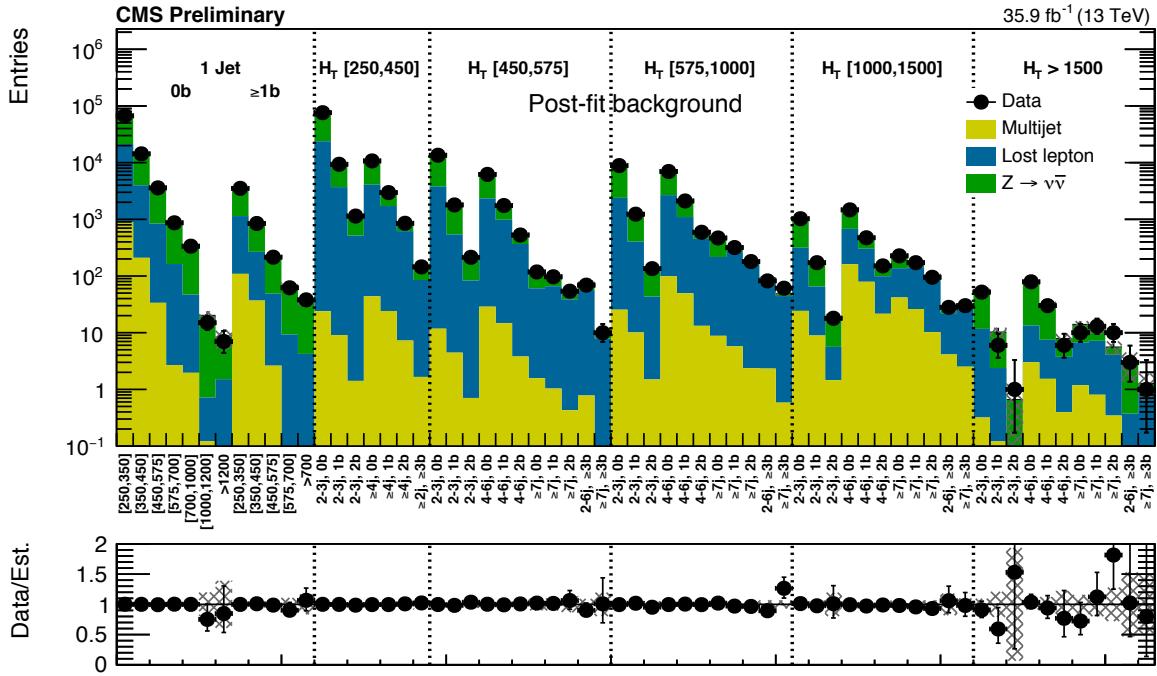


Figure 6.5: The data yield in each topological region compared to the post-fit background prediction. The hatched bands illustrate the total uncertainty in the background estimate. Results in the monojet regions are binned in jet p_T , while those in the multijet regions are labeled according to N_{jets} and $N_{\text{b-jets}}$.

signal region with constraints set according to the background uncertainties and signal uncertainties (described in section 6.2.1). The post-fit yield for each topological region is illustrated in figure 6.5, and the individual yield in each M_{T2} bin can be found in figures 6.6, 6.7, and 6.8. The post-fit procedure has the effect of constraining background and its associated uncertainties when the fitting procedure is applied to data consistent with predictions modeling uncertainties appropriately.

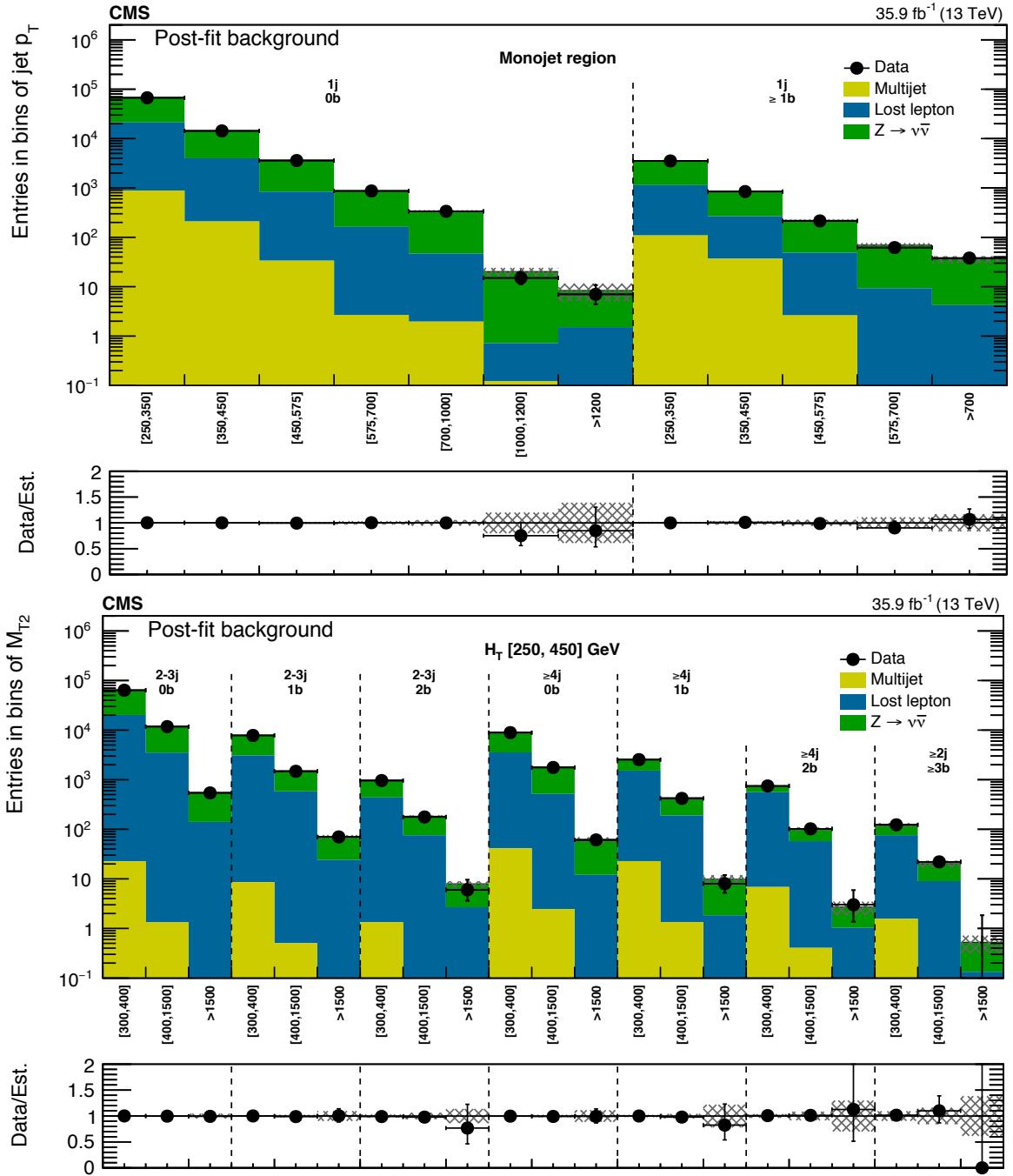


Figure 6.6: The data yield in the monojet and very-low H_T regions compared to the post-fit background prediction. The hatched bands illustrate the total uncertainty in the background estimate. Results in the monojet regions are binned in jet p_T in units of GeV, while those in the multijet regions are labeled according to M_{T2} bin in units of GeV.

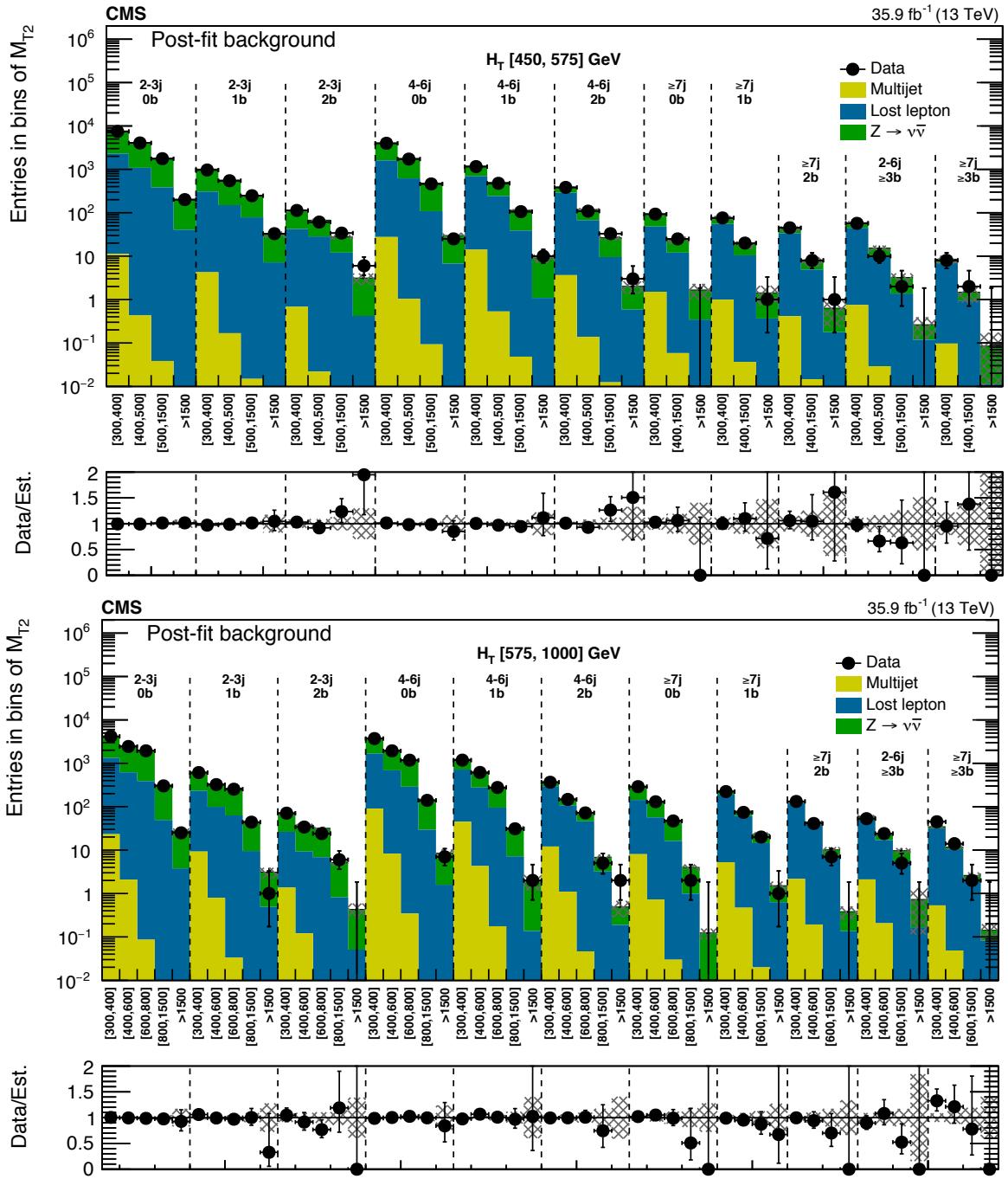


Figure 6.7: The data yield in the low H_T and medium H_T regions compared to the post-fit background prediction. The hatched bands illustrate the total uncertainty in the background estimate. Results are labeled according to M_{T2} bin in units of GeV.

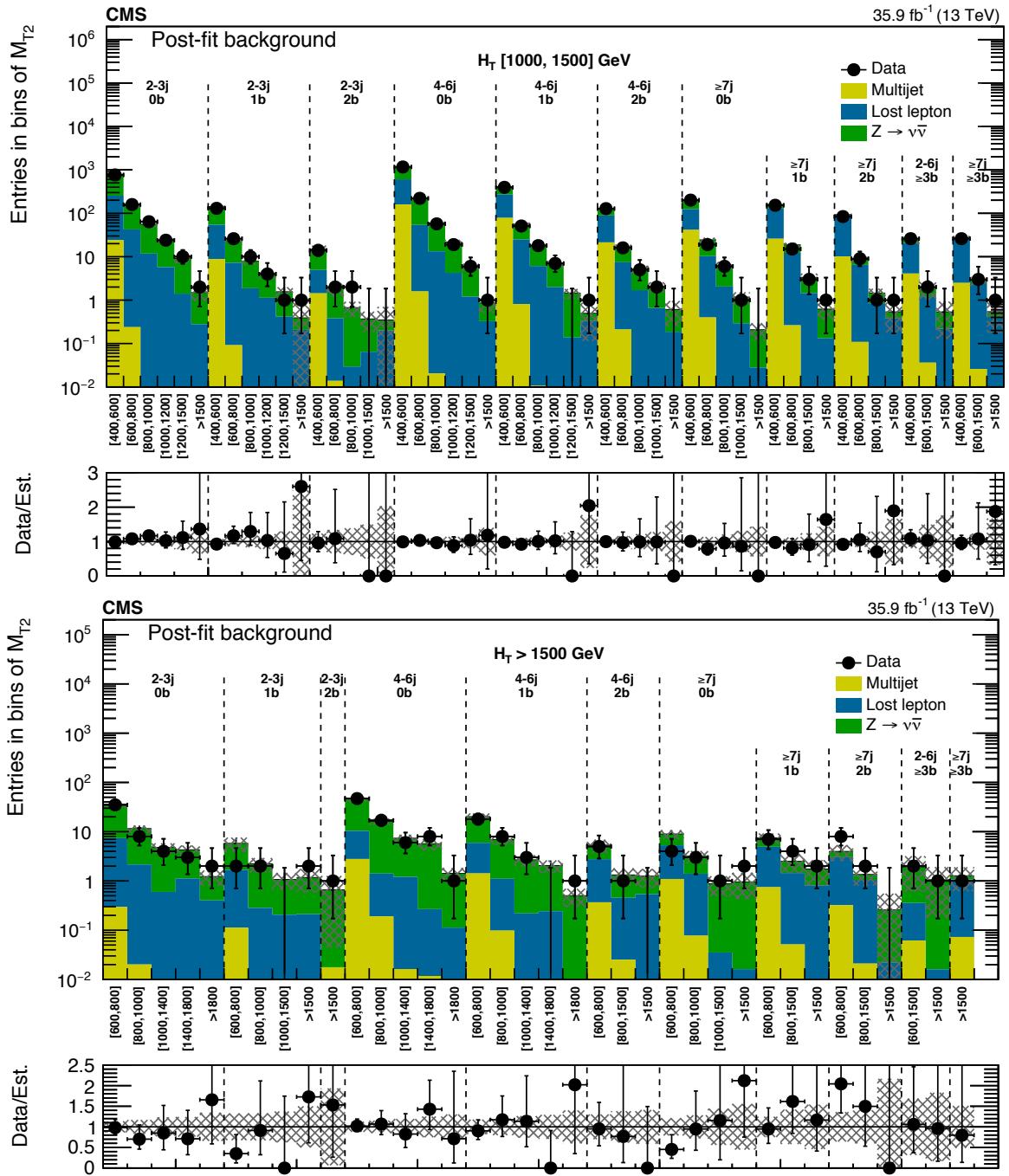


Figure 6.8: The data yield in the high H_T and extreme H_T regions compared to the post-fit background prediction. The hatched bands illustrate the total uncertainty in the background estimate. Results are labeled according to M_{T2} bin in units of GeV.

6.2 Signal Interpretations

The results of the search can be interpreted as a constraint on the mass of hypothesized particles of various BSM models. Based on the simulated production processes of simplified SUSY models, this section describes some of the uncertainties associated with the signal yield, as well as the limits associated with various simplified models of interest.

6.2.1 Signal Yield Systematic Uncertainties

The uncertainties associated with the signal yield are summarized in table 6.1. The various sources of uncertainty are described in detail below:

- *Luminosity*: the uncertainty of the total integrated luminosity delivered by the LHC is evaluated in *Van Der Meer scans*, where the rate of interactions is measured while scanning the colliding proton beams across each other. The uncertainty in the luminosity based on these studies is 2.5%.
- *Simulation statistics*: the limited size of the Monte Carlo samples for each sample affect the statistics in each signal region bin. After applying the signal selection, the statistical uncertainty can range from 1-100% in different bins, though bins with large uncertainty typically correspond to those with low signal acceptance, and thus do not drive the sensitivity of the analysis to that simplified model mass point.
- *Renormalization and factorization scales*: the overall effect of varying the simulation renormalization and factorization scales of the underlying physics processes (and

subsequent effect on event kinematics) is computed separately in each bin. The result of this variation is relatively flat and minor across the signal bins, and is conservatively estimated at 5% for all regions.

- *Initial-state radiation (ISR) recoil*: variations of the boost due to initial-state radiation are performed to test the modeling of ISR effects in simulation. The effect of these variations ranges from 0-30% across signal regions.
- *b-tagging efficiency*: the effect of varying the b-tag scale factor efficiency is calculated in each bin, and taken as a correlated error amongst all bins. The effect of this variation for heavy (light) flavor jets is 0-40% (0-20%) across signal regions.
- *Lepton efficiency*: the effect of varying the electron and muon scale factors applied to simulation is calculated in each bin, and taken as a correlated error amongst all bins. The effect of this variation is 0-20% across signal regions.
- *Jet energy scale*: the effect of varying the jet energy scales is calculated in each bin, and the results are compatible with statistical uncertainty for low-statistics bins. Based on findings in the high-statistics bins, the uncertainty is estimated at 5% for all regions.
- *Fast simulation modeling*: The signal Monte Carlo samples are generated using the FastSim package, which may result in modeling differences compared to Fullsim MC. Studies of the Fastsim kinematics with respect to E_T^{miss} and pile-up modeling indicate differences of up to 5% in some signal regions.

Table 6.1: Typical values of the systematic uncertainties associated with the simplified SUSY model signal yield for each interpretation in this search. Some systematics are taken as correlated amongst all signal regions; the rest are uncorrelated everywhere. Note that the large range of statistical uncertainty is driven by a small number of signal regions with low acceptance (which are not typically sensitive to those model points).

Source	Typical Values	Correlated?
Luminosity	2.6%	✓
Limited size of MC samples	1–100%	-
Renormalization and factorization scales	5%	-
“ISR” recoil	0–30%	✓
B-tagging efficiency, heavy flavor	0–40%	✓
B-tagging efficiency, light flavor	0–20%	✓
Lepton efficiency	0–20%	✓
Jet energy scale	5%	-
Fast simulation E_T^{miss} modeling	0–5%	✓
Fast simulation pileup modeling	4.6%	✓

6.2.2 Exclusion Limits

The final results are interpreted in the context of various simplified SUSY models: gluino-mediated squark pair production, direct production of squarks, and alternative models of top squark production with different decay modes. Each of these models is illustrated in figure 6.9. For each pair producing gluino (squark) scenario, the models assume all SUSY particles other than the gluino (squark) and lightest neutralino are too massive to be produced directly and the gluino (squark) decays promptly. In addition, each model assumes that the gluino (squark) decays with a 100% branching fraction into the decay products depicted in figure 6.9, except for models where the decays of the two squarks differ where a 50% branching fraction for each decay mode is assumed. When considering top squark pair production, the polarization of the top quark is model dependent (and a function of the top-squark and neutralino mixing matrices), and events are generated without polarization.

The cross-section exclusion limits are calculated at 95% confidence level (CL) for each simplified model [46, 39, 32]. The limits are obtained using the background-fitting procedure described in section 6.1 obtained with the higgsCombine tool [50]. The 95% CL exclusion limits for gluino-mediated models is shown in figure 6.10, for direct squark production in figure 6.11, and for alternate top squark decay modes in figure 6.12. The constraints on the masses of SUSY particles excluded by this search is summarized in table 6.2.

This chapter makes use of figures and tables from the M_{T2} paper and internal

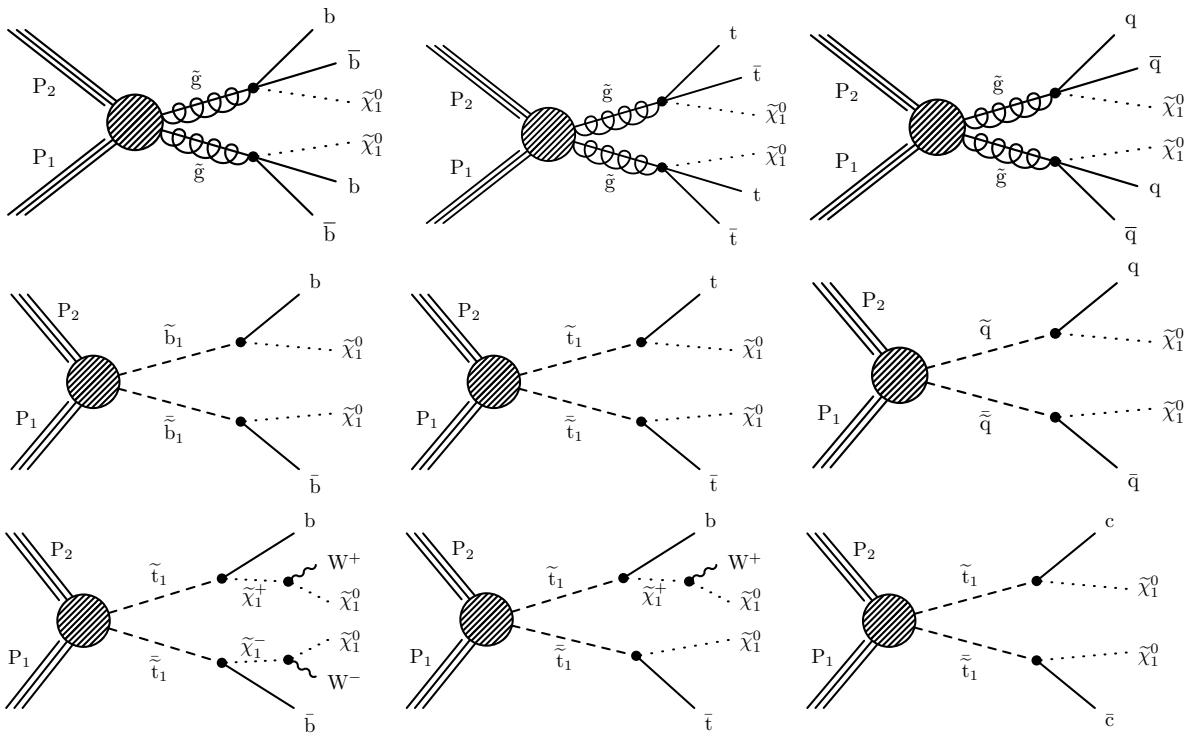


Figure 6.9: Feynman diagrams depicting the different simplified SUSY models considered in this analysis. Results are interpreted in the context of gluino-mediated bottom, top, and light-flavor squark production (top), direct production of bottom, top, and light-flavor squarks (middle), and alternate decay modes of direct top squark production (bottom).

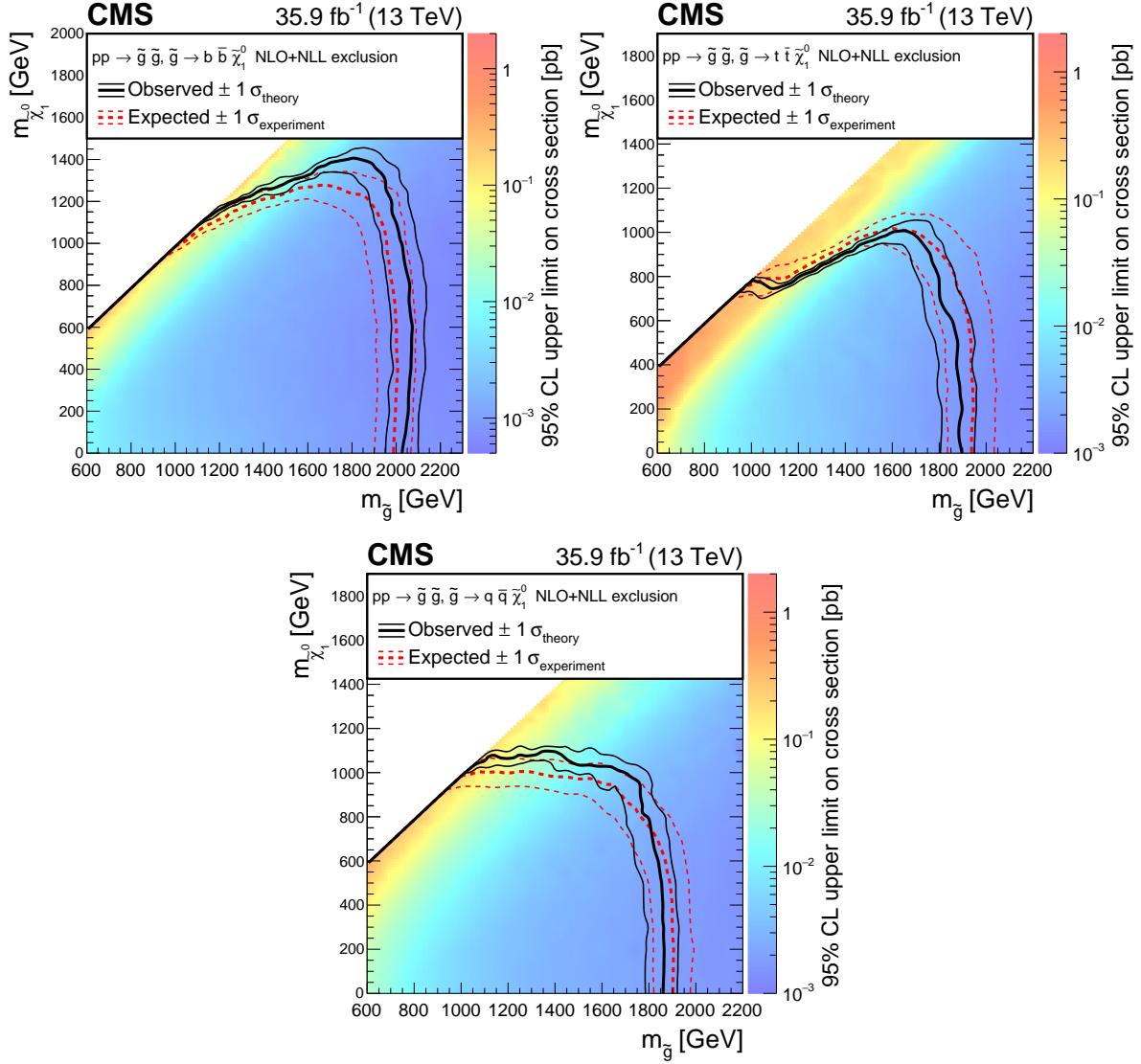


Figure 6.10: 95% confidence level exclusion limits for gluino-mediated squark production of bottom (top left), top (top right), and light-flavor (bottom) squarks. Dashed red lines indicate the expected sensitivity and associated uncertainty, while black lines indicate the observed exclusion limit and its associated theoretical uncertainty based on the signal cross-section.

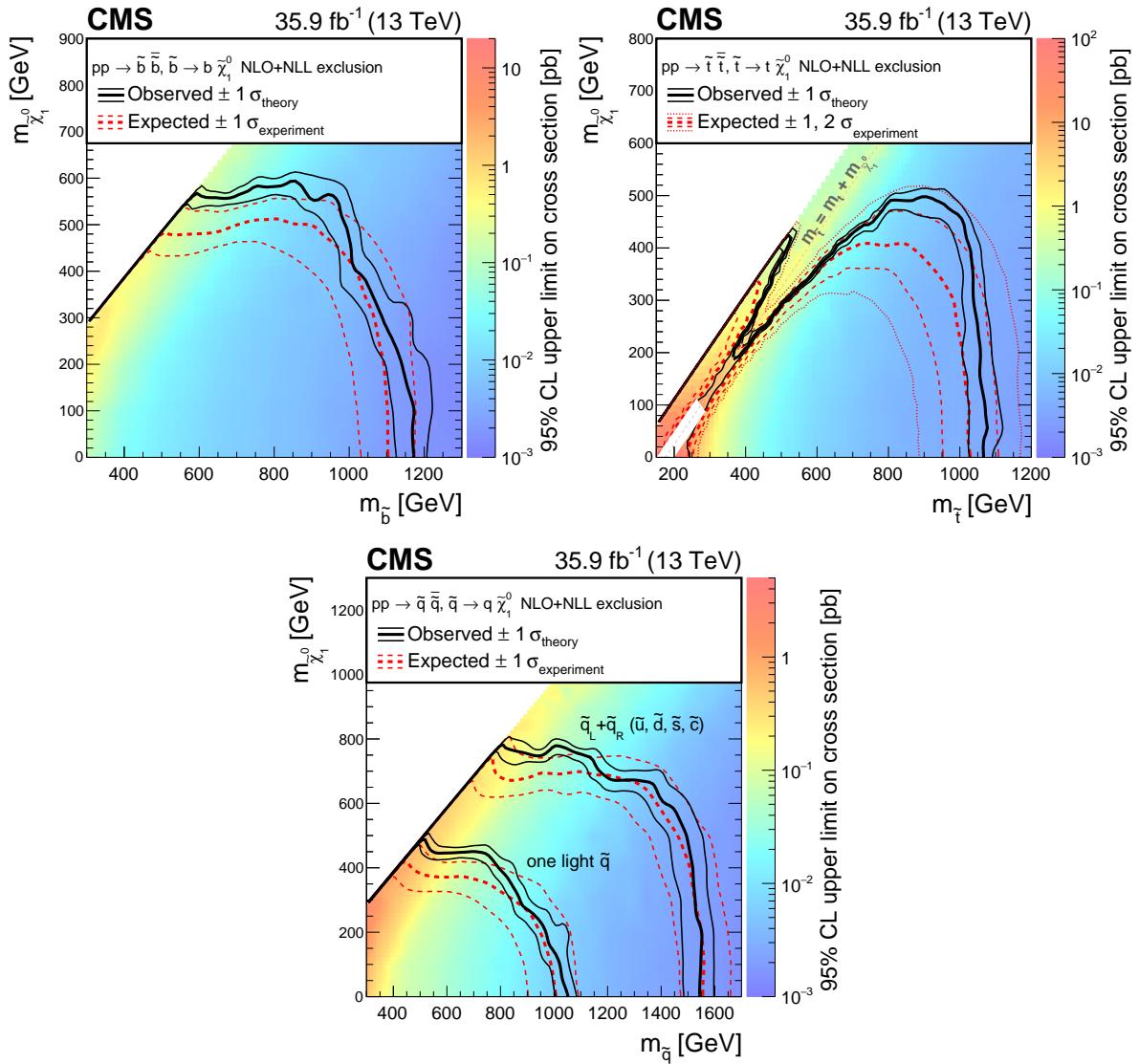


Figure 6.11: 95% confidence level exclusion limits for direct squark production of bottom (top left), top (top right), and light-flavor (bottom) squarks. Dashed red lines indicate the expected sensitivity and associated uncertainty, while black lines indicate the observed exclusion limit and its associated theoretical uncertainty based on the signal cross-section.

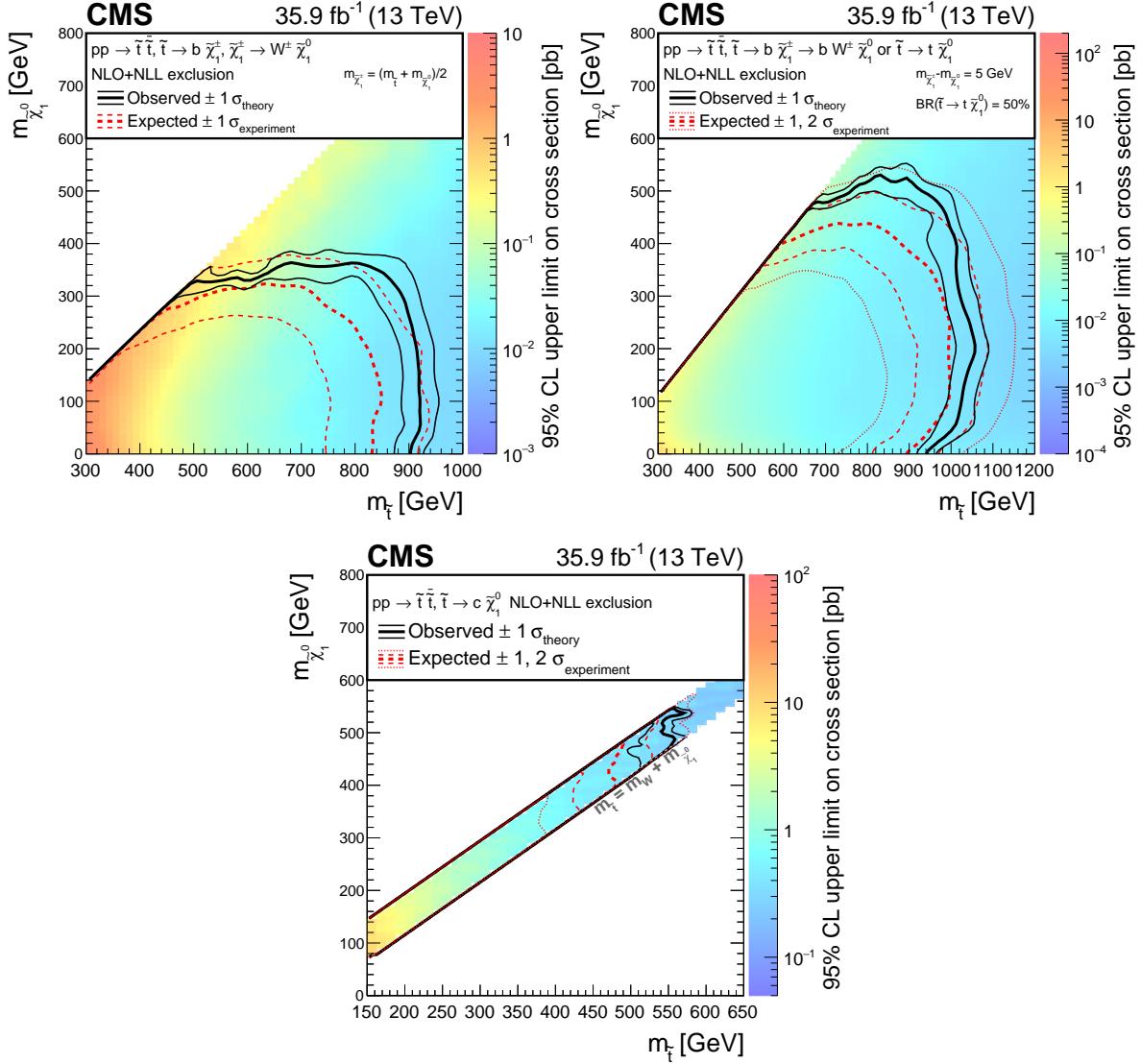


Figure 6.12: 95% confidence level exclusion limits for top squark production, decaying either via charginos (top left), where the chargino mass is taken as halfway-between the top squark and neutralino; a mixed mode with charginos and neutralinos (top right), where the $\tilde{\chi}^\pm$ mass is 5 GeV above the $\tilde{\chi}_1^0$ mass; or a compressed scenario, where the top squark is kinematically constrained to decay to charm quarks (bottom). Dashed red lines indicate the expected sensitivity and associated uncertainty, while black lines indicate the observed exclusion limit and its associated theoretical uncertainty based on the signal cross-section.

Table 6.2: Summary of 95% CL exclusion limits on the masses of SUSY particles (sparticles) produced by various simplified models. The limit on the produced sparticle is listed for a massless $\tilde{\chi}_1^0$, along with the greatest excluded mass of $\tilde{\chi}_1^0$ for any mass of produced sparticle.

Simplified model	Limit on produced sparticle mass [GeV] for $m_{\tilde{\chi}_1^0} = 0 \text{ GeV}$	Highest limit on the $\tilde{\chi}_1^0$ mass [GeV]
Direct squark production:		
Bottom squark	1175	590
Top squark	1070	550
Single light squark	1050	475
Eight degenerate light squarks	1550	775
Gluino-mediated production:		
$\tilde{g} \rightarrow b\bar{b}\tilde{\chi}_1^0$	2025	1400
$\tilde{g} \rightarrow t\bar{t}\tilde{\chi}_1^0$	1900	1010
$\tilde{g} \rightarrow q\bar{q}\tilde{\chi}_1^0$	1860	1100

analysis note to illustrate the analysis design, methodology, and results. This work was made possible by contributions from the rest of the Surf & Turf group, our collaborators at ETH Zurich, and the many other CMS members in the SUSY group and beyond.

Chapter 7

Extending the All-Hadronic Search

The results presented in this analysis thus far exclude significant amounts of SUSY parameter phase space compared to previous searches for similar signatures. However, the spectrum of SUSY models that might be realized in nature is wide and varied. As the LHC continues to deliver an increasing amount of collision data, the ability to probe increasing rare and more subtle signatures increases.

In this final section we discuss extensions to the M_{T2} analysis that modify the search to target additional SUSY models, in particular those which may produce very low- p_T leptons in the final state, otherwise known as *soft leptons*. By requiring a lepton in the final state, the previously irreducible invisible Z background of the all-hadronic analysis is greatly suppressed, and in principle one might hope to probe additional phase space for new physics.

Requiring a soft lepton in the final state not only reduces the dominant invisible Z background in this all hadronic search, but can also increase the sensitivity to models

where the mass difference of the parent particles and their decay products are very small, or those with *compressed mass spectra* (for example, in scenarios close to the diagonal as in figures 6.10, 6.11, and 6.12). Many supersymmetric models in particular include a relative mixing of the superpartners of the Higgs and electroweak (EW) gauge bosons into mass eigenstates with a highly-degenerate mass spectrum. Charged EW bosons ($\tilde{\chi}^\pm$) can decay into neutral EW bosons or the lightest supersymmetric particle ($\tilde{\chi}^0$) in cascades with leptons in the final state. In models with compressed mass spectra where the mass difference between the charged and neutral EW bosons is small, the outgoing leptons will be soft, and searches targeting soft leptons in the final state can be a useful tool in searches for these SUSY scenarios.

7.1 The Soft Lepton Search

A preliminary search extending the all-hadronic analysis was conducted on 2.3 fb^{-1} of data collected by the CMS detector through 2015 [22]. The results of this preliminary analysis are presented here and interpreted in the context of two simplified SUSY models with compressed mass spectra.

7.1.1 Event Selection

The soft lepton search defines a baseline selection similar to that used in the M_{T2} analysis, with the notable exception of requiring a soft electron (muon) with $5 < p_{\text{T}} < 20\text{ GeV}$ in the barrel region, $|\eta| < 1.4442$ (1.479). Similar $E_{\text{T}}^{\text{miss}}$ triggers are used, with

a minimum E_T^{miss} requirement of 200 GeV to saturate the trigger efficiency. Identical $\Delta\phi(j_{1234}, E_T^{\text{miss}})$ and $|\vec{H}_T^{\text{miss}} - \vec{E}_T^{\text{miss}}|/E_T^{\text{miss}}$ requirements are used to protect against E_T^{miss} mismeasurements, and to reject dilepton events the same lepton vetoes are applied to additional leptons in the event. There is an additional requirement on transverse mass of the soft leptons such that $m_T > 20$ GeV, in order to reduce the background contribution from $Z \rightarrow \tau\tau$ events. Events must contain at least 1 jet and a minimum of $H_T > 200$ GeV of hadronic activity.

Following the baseline selection above, events are further categorized into topological regions using the H_T , E_T^{miss} , and N_{jets} content, as well as the number of b-tags $N_{\text{b-jets}}^{\text{soft}}$ and $N_{\text{b-jets}}^{\text{hard}}$, where a b-jet is considered soft if $20 < p_T < 50$ GeV, and hard if $p_T > 50$ GeV. There are three “tail” regions corresponding to different kinematic variables:

- $N_{\text{b-jets}}$ (soft or hard) ≥ 3
- $E_T^{\text{miss}} > 500$ GeV and $N_{\text{b-jets}} \leq 3$
- $H_T > 1000$ GeV, $N_{\text{b-jets}} \leq 3$, and $E_T^{\text{miss}} < 500$ GeV

The remaining phase space is divided into topological regions as follows:

- E_T^{miss} : [200,300], [300,500] GeV. These are merged in the monojet regions with b-tags.
- $N_{\text{b-jets}}$: There are four $N_{\text{b-jets}}$ regions, $N_{\text{b-jets}} = 0$, $N_{\text{b-jets}}^{\text{soft}} = [1,2]$ (with $N_{\text{b-jets}}^{\text{hard}} = 0$), and $N_{\text{b-jets}}^{\text{hard}} = [1,2]$, respectively referred to as the 0b, soft b, and hard b regions.

- N_{jets} : In regions with no b-tags, events are binned in regions with 1, 2-3, 4-5, or ≥ 6 jets. In regions with one or more b-tags, events are binned in regions with 1, 2-3, or ≥ 4 jets.

Finally, each topological region is subdivided into three bins of m_T : [20,90], [90,120], and $\geq 120 \text{ GeV}$. In total, there are 21 topological regions each subdivided into 3 m_T bins.

7.1.2 Backgrounds

The background for the soft lepton analysis can be categorized into three separate sources:

- *Single Lepton*: real SM processes can generate events with a single lepton in the final state. Typically such events contain a single leptonically decaying W boson (from either W production, top quark production, or other rare processes), and also generate significant missing energy due to the neutrino associated with the W decay.
- *Dilepton*: SM events with two (or more) leptons in the final state can be misclassified when other leptons are misidentified or not reconstructed. Typically such events are due to two leptonically decaying W bosons (from either $t\bar{t}$ or diboson production), where one of the W-associated leptons is not found. Because there are usually two neutrinos contributing to the E_T^{miss} vector, m_T is less effective in discriminating against this background.
- *Fake Lepton*: events without any lepton in the final state (such as QCD multijet production of $Z \rightarrow \nu\bar{\nu}$) can enter the signal region if a jet is misidentified as a lepton.

ton or a non-prompt lepton passes the lepton selection. This background is highly suppressed and typically negligible in all but some high- m_T regions.

7.1.3 Single Lepton Estimate

The single lepton background is primarily due to highly-asymmetric W decays where the neutrino carries a large fraction of the W momentum, resulting in a large E_T^{miss} and a soft lepton. These events are greatly suppressed by sampling the m_T distribution as values greater than the W mass, and is estimated in all signal regions using a control region with an inverted asymmetry. A single high- p_T control region is created by selecting events with a single lepton above the corresponding E_T^{miss} threshold in a given signal region ($p_T \geq 200, 300, \text{ or } 500 \text{ GeV}$), with a maximum threshold on the missing energy $E_T^{\text{miss}} < 60 \text{ GeV}$ so that the W momentum distribution is kinematically similar to that in the signal regions. The number of control region events in a given topological region (N^{CR1L}) is multiplied by the ratio of SR-to-CR events as calculated in simulation ($R_{\text{MC}}^{\text{SR/CR1L}}$) for the region, and extrapolated into bins of m_T based on simulation ($k_{\text{MC}}(m_T)$) as described in equation 7.1. The extrapolation in bins of m_T is necessary not only to preserve statistics but because the shape of the m_T distribution in each control region is very different from the signal region distribution, due to the low- E_T^{miss} requirement in the control region where resolution effects cause a significant smearing in the m_T shape.

$$N_{1l}^{\text{SR}} = N^{\text{CR1L}} \cdot R_{\text{MC}}^{\text{SR/CR1L}} \cdot k_{\text{MC}}(m_T) \quad (7.1)$$

The uncertainties associated with the single lepton estimate account for statistical

fluctuations in data and simulation, as well as possible sources of systematic error in simulation propagated to the quantities measured in MC. The polarization modeling is studied in an inclusive W sample and found to be well modeled by simulation, and additional studies reweighting simulation based on W polarization result in an uncertainty of 10-20% on the transfer factor $R_{\text{MC}}^{\text{SR/CR1L}}$. Additional uncertainties due to the relative fraction of top and W events, renormalization and factorization scales, lepton efficiency, b-tagging efficiency, jet energy corrections, and uncertainty on top momentum are propagated through simulation, with varying effects between 1-10%. In addition, multiple uncertainties associated with m_T modeling in simulation are studied, with effects due to W-top fraction and E_T^{miss} resolution yielding uncertainties ranging between 5-30% in m_T shape, depending on the relative fraction of top events in a given bin.

7.1.4 Dilepton Estimate

Backgrounds due to dilepton events where one lepton is not found are estimated using a control region instead requiring two leptons, replacing the second-lepton veto with a tight requirement for a second, high- p_T lepton. To suppress the contribution from events with fake leptons, the second lepton is required to have a minimum p_T of at least 25 GeV, and the expected statistics, kinematics, and composition of the control region is similar to that of the signal region. Because of the low rate of this background in the signal region (and corresponding low statistics in each control region), the estimate is performed in a similar manner to the single lepton estimate, but with an additional extrapolation in the E_T^{miss} dimension based on simulation. The number of control region events in a given

region of H_T , N_{jets} , and $N_{\text{b-jets}}$ (N^{CR2L}) is multiplied by the ratio of SR-to-CR events as calculated in simulation ($R_{\text{MC}}^{\text{SR/CR2L}}$) for that region, and extrapolated into bins of E_T^{miss} and m_T based on simulation ($k_{\text{MC}}(E_T^{\text{miss}}, m_T)$) as described in equation 7.2.

$$N_{2l}^{\text{SR}} = N^{\text{CR2L}} \cdot R_{\text{MC}}^{\text{SR/CR1L}} \cdot k_{\text{MC}}(E_T^{\text{miss}}, m_T) \quad (7.2)$$

In addition to statistical fluctuations, the uncertainties associated with the dilepton estimate account for several effects, including lepton acceptance in simulation. Varying the renormalization and factorization scales, lepton efficiency, and jet energy corrections in simulation is propagated to the ratio $R_{\text{MC}}^{\text{SR/CR1L}}$, yielding uncertainties of 1-10%, 5%, and 5-20% respectively (the b-tagging uncertainty is found to be negligible). Because the $Z \rightarrow l^+l^-$ yield in the control region is slightly different than that in the corresponding signal regions, an additional uncertainty of 5-10% is assigned to each region based on the relative fraction of $Z \rightarrow l^+l^-$ events. Possible systematic effects of the E_T^{miss} and m_T modeling in simulation are also studied in dilepton $t\bar{t}$ control regions, yielding uncertainties of 10% (35%) in the 300-500 GeV (≥ 500 GeV) E_T^{miss} regions and 30% (50%) in the medium (high) m_T bins.

7.1.5 Fake Lepton Estimate

The fake lepton background has multiple contributions from different SM processes. Missing energy mismeasurement may yield a contribution from QCD multijet events, but this background is strongly suppressed due to sufficiently large E_T^{miss} requirements in every signal region. Additional contributions from electroweak SM processes with zero or one

lepton in the final state but significant E_T^{miss} can also contribute, when a fake lepton is found in a zero lepton event (such as $Z \rightarrow \nu\bar{\nu}$) or a real lepton lost and a fake lepton found in a one lepton event (such as W production). In most regions, the fake lepton background is negligible (with respect to uncertainties on the other backgrounds), and so the prediction is taken directly from simulation in the low and medium m_T regions with a 100% uncertainty on the yield. However, because the fake lepton background is not constrained by $m_T < m_W$ and falls more slowly than other background distributions in m_T , in the high- m_T regions it is more significant and estimated from data. To calculate the predicted yield, a loose-not-tight (L!T) control region is constructed with identical kinematic requirements to the signal region, except the soft lepton is required to pass the loose lepton identification requirements but fail the tight requirements ($N_{\text{data}}^{\text{L!T}}$). Once the contribution from prompt leptons passing the L!T selection ($N_{\text{prompt}}^{\text{L!T}}$), the yield is weighted by the probability that a fake lepton passing the loose selection also passes the tight selection (ϵ_{TL}). The probability is measured as a function of lepton p_T in a sample of QCD multijet data enriched in fake leptons, selected using a high-statistics sample of QCD events recorded with pure H_T triggers. The fake lepton estimate in each high- m_T signal region is described by equation 7.3.

$$N_{\text{SR}}^{\text{fakes}} = \sum_{p_T} \left(N_{\text{data}}^{\text{L!T}}(p_T) - N_{\text{prompt}}^{\text{L!T}}(p_T) \right) \times \frac{\epsilon_{\text{TL}}(p_T)}{1 - \epsilon_{\text{TL}}(p_T)} \quad (7.3)$$

Based on studies in simulation, the fake-enriched QCD multijet sample has negligible prompt lepton contamination when requiring $E_T^{\text{miss}} < 50 \text{ GeV}$ and $m_T < 40 \text{ GeV}$. The tight-to-loose ratio is $\mathcal{O}(10\%)$, and the expected fake yield in the high- m_T bin is $\mathcal{O}(1)$

events. Thus the estimate is dominated by statistics in the L!T region, ranging between 50-100%.

7.1.6 Results and Interpretations

The predicted background yields are compared with 2.3fb^{-1} of data collected by the CMS detector through 2015. No significant deviations from the expected SM background are observed. As with the all-hadronic analysis, the backgrounds are also estimated using a maximum likelihood fitting procedure with a background-only hypothesis as described in section 6.1. Both the pre-fit and post-fit background predictions compared to data are illustrated in figure 7.1.

Additional fits using a background-plus-signal hypothesis are used to set upper limits on the production cross sections of some simplified SUSY models producing soft leptons. A summary of the uncertainties on the simulated signal yield can be found in table 7.1. The post-fit background yield based on these inputs is used to set 95% confidence level exclusion limits as shown in figure 7.2.

7.2 Future Extensions for an All-Hadronic Search

The extension of the all-hadronic analysis presented in this section illustrates one possible way to broaden the scope of an all-hadronic search to target additional sectors where new physics might reveal itself. However, there are similar analyses which may supersede the results of a naive single soft lepton search, and an additional question re-

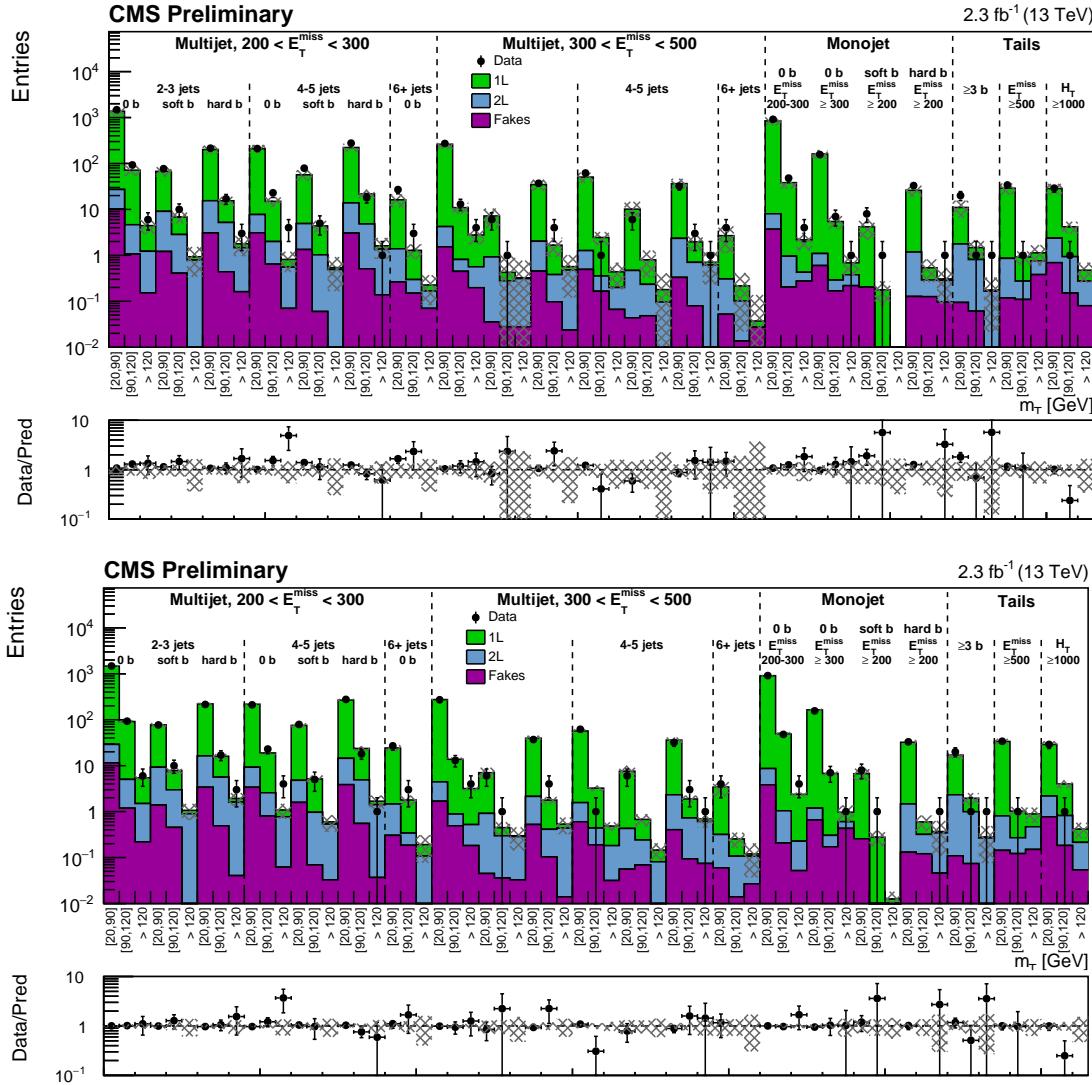


Figure 7.1: The predicted background yields compared to the observed number of events in data for pre-fit (top) and post-fit (bottom) background estimates. The m_T bins are shown on the x -axis, and the grey hatched band represents the total uncertainty on the background yields for the pre-fit background estimates and the fit uncertainty for the post-fit estimates.

Table 7.1: Typical values of the systematic uncertainties associated with the simplified SUSY model signal yield for each interpretation in this search.

Source	Typical values [%]
Integrated luminosity	2.7
Lepton efficiency	10
Jet energy scale	5
b tagging efficiency	0–20
ISR	15–30
Renormalization and factorization scales	5
Limited size of MC samples	1–70

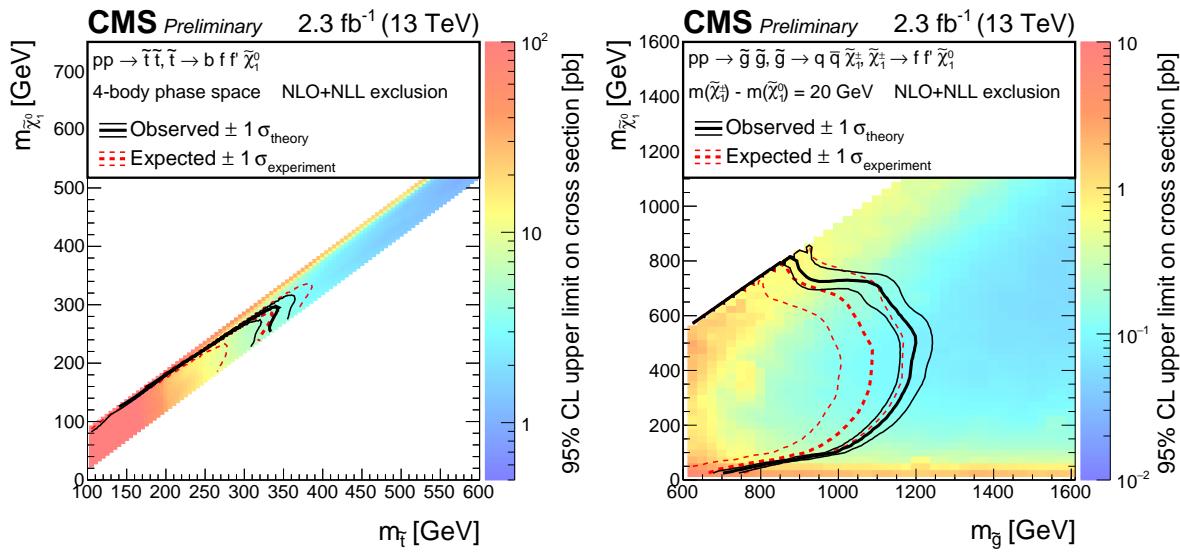


Figure 7.2: 95% confidence level exclusion limits for top squark (left) and gluino (right) production. Dashed red lines indicate the expected sensitivity and associated uncertainty, while black lines indicate the observed exclusion limit and its associated theoretical uncertainty based on the signal cross-section.

mains – how much does the extension requiring a soft lepton outperform a traditional all-hadronic analysis? An all-hadronic search should have some discriminating power even against models which always generate soft leptons in the final state (since these leptons will occasionally be lost or mis-identified), and ironically may even have greatly discriminatory power than a soft lepton search when the leptons generated are ultra-soft ($p_T << 5 \text{ GeV}$) and not reconstructed at all.

To benchmark the possible performance of an improved soft lepton analysis against an all-hadronic search, the performance of the M_{T2} analysis on a squark production model generating soft leptons (T5qqWW) is compared to a soft lepton extension. With additional optimization of signal region binning — in particular a finer binning in the low- E_T^{miss} regime to target compressed spectra — and assuming a typical background estimate uncertainty commensurate with previous analyses, it is possible for a soft lepton analysis to outperform an all-hadronic search as illustrated in figure 7.3.

Near the mass-diagonal where the parent particle mass (in this case, the squark) is very close to the LSP mass, the soft-lepton search can significantly outperform an all-hadronic equivalent, at the cost of performance in the light-LSP regime. The evidence suggests that a soft-lepton search could be used in conjunction with typical all-hadronic analyses to increase the total excluded phase space when targeting compressed models with nearly-degenerate mass splittings.

This chapter makes use of figures and tables from the soft lepton physics analysis summary to illustrate the analysis design, methodology, and results. This work was made possible by contributions from Giovanni Zevi Della Porta, the rest of the Surf & Turf

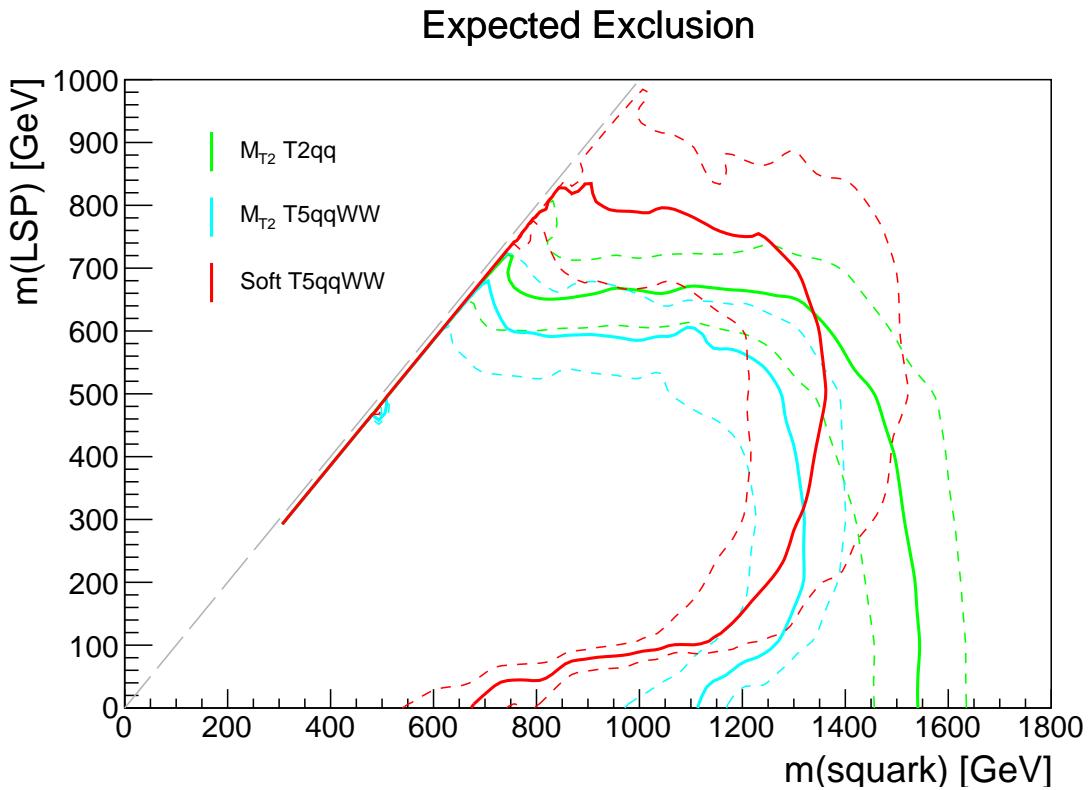


Figure 7.3: The expected limits of the all-hadronic M_{T2} analysis compared with a hypothetical soft lepton extension. The “T2qq” model is a typical squark production model as previously shown in figure 6.11, whereas the “T5qqWW” model is a similar squark production model where the squarks decay in a cascade including charginos, which radiate W bosons that may decay to soft leptons. While the performance of the M_{T2} analysis in the hadronic model (green) is similar to the soft lepton model (blue), a soft lepton search can out-perform the all-hadronic equivalent near the diagonal where the mass splitting between the squarks and lightest supersymmetric particle is very small.

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