

Lecture : 04.03.2021

- Fundamental Standards

Fundamental quantities for which independent standards are defined:
length, time, mass & temperature.

Some of the things we take care of while choosing the standards :

- i) Should not be rare and should be ~~expensive~~ easy to reproduce. Labs all over the world need to have access to them.
- ii) Should be insensitive to modifying and interfering inputs.
- iii) Should be well known and should be accepted throughout the world. Shouldn't be controversial. Should give the same result every time it's reproduced.

Hierarchy in fundamental standards : For traceability

- Global Reference

- i) National Reference standards
- ii) Working standards
- iii) Interlaboratory standards

- Vernier Caliper

A common length measuring instrument. It gives additional accuracy by implementing two scales: The main scale and the vernier scale. Unequal divisions in these scales give way to extra accuracy.

Let, a = Size of main scale division

p = No. of main scale divisions

q = No. of vernier scale divisions

} for same interval

$$\text{Scale readability} = \frac{a}{q}$$

q

- Screw Gauge

Also used to measure length. It has a main scale and a circular scale. We define pitch as the distance moved on the main scale for one screw rotation.

Let, n = No. of divisions on the circular scale

$$p = \text{Pitch}$$

$$\text{Scale readability} = \frac{p}{n}$$

Lecture : 08.03.2021

- Gauge blocks:

Physical objects whose dimensions we can rely on for checking various instruments.

Small blocks of hard & dimensionally stable steel.

Made in sets of wide range and small steps, such that any arbitrary figure can be reproduced using the limited blocks at our disposal.

Extremely important to control all interfering and modifying inputs for greater accuracy.

Similarly, we have angle blocks.

Angular displacement is essentially based on length measurements so a fundamental standard isn't necessary.

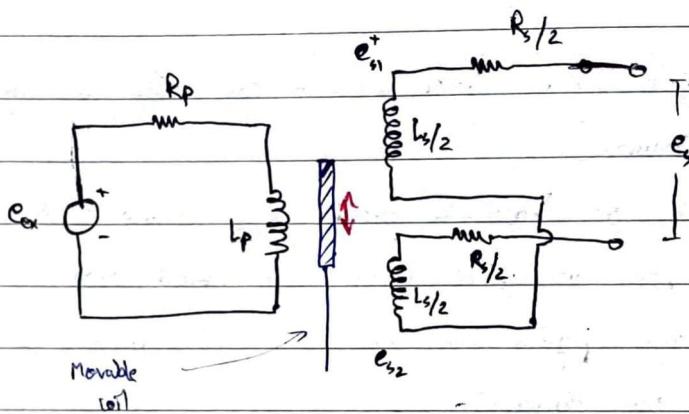
We still have working standards for ease.

Again, angle blocks have a wide range & several steps.

- Magnetoresistive displacement transducer : Refer the slides

Lecture : 09.03.2021

- Variable inductance pickup \rightarrow Differential transformer



$$e_o = i_p R_p + L_p \frac{di_p}{dt} \quad \text{--- (1)}$$

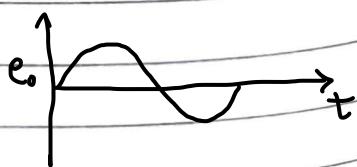
$$\left. \begin{aligned} e_{s1} &= M_1 \frac{di_p}{dt} \\ e_{s2} &= M_2 \frac{di_p}{dt} \end{aligned} \right\} \begin{matrix} M_1 \text{ & } M_2 \text{ are mutual} \\ \text{inductances.} \end{matrix}$$

$$e_s = e_{s1} - e_{s2} = (M_1 - M_2) \frac{di_p}{dt} \quad \text{--- (2)}$$

For symmetric configuration,

$$M_1 = M_2$$

$$e_s = 0$$



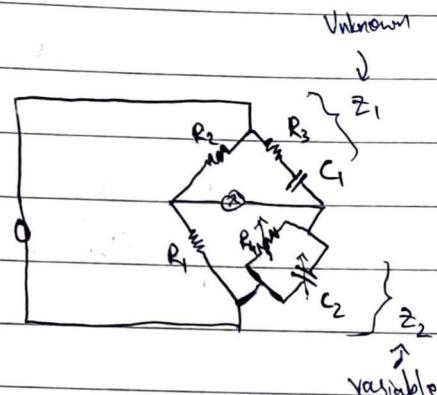
Nearly linear near the null point.

- AC bridges

i) Wien bridge.

Used to measure an unknown capacitance.

When the bridge is balanced:



$$\frac{R_2}{R_1} = \frac{Z_1}{Z_2}$$

$$\text{Impedance } Z_1 = R_3 + j \frac{1}{j\omega C_1} = R_3 - \frac{j}{\omega C_1}$$

$$Z_2 = \frac{j\omega C_2 \times R_4}{R_4 + j\omega C_2} = \frac{R_4}{1 + j\omega C_2 R_4}$$

Substituting,

$$R_2 \left(\frac{R_4}{1 + j\omega C_2 R_4} \right) = R_1 \left(R_3 - \frac{j}{\omega C_1} \right)$$

$$\begin{aligned} \omega C_1 R_2 R_4 &= (R_1 R_3 + R_1 j\omega C_2 R_3 R_4) \omega C_1 \\ &= jR_1 + \omega C_2 R_3 R_4 \end{aligned}$$

Comparing complex & real parts,

$$\omega^2 = \frac{1}{R_3 R_4 C_1 C_2} \quad \text{--- ①}$$

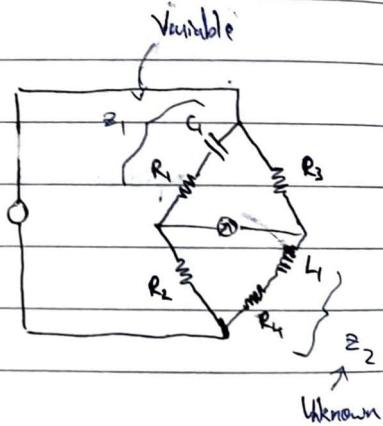
$$wC_1 R_2 R_4 = wR_1 (C_1 R_3 + C_2 R_4)$$

$$\frac{R_2}{R_1} = \frac{C_2}{C_1} + \frac{R_3}{R_4} \quad \text{--- (2)}$$

By changing the variable elements R_3 & C_2 , we can find the value of R_3 & C_1 .

ii) Hay bridge

Used to measure an unknown inductance & resistance.



$$\frac{Z_1}{R_2} = \frac{R_3}{Z_2}$$

$$\frac{R_1 - jwC_1}{R_2} = \frac{R_3}{R_4 + jwL_1}$$

$$\frac{R_1 R_4 + jwL_1 R_1 - jR_4}{wC_1 C_1} + \frac{L_1}{C_1} = R_3 R_2$$

Comparing complex & real parts,

$$\omega^2 = \frac{R_4}{L_1 R_1 C_1} \quad \text{--- (3)}$$

And, $R_1 R_4 + \frac{L_1}{C_1} = R_3 R_2$

Finally, the eqn for unknown are:

$$L_1 = \frac{R_2 R_3 C_1}{1 + \omega^2 R_1^2 C_1^2}$$

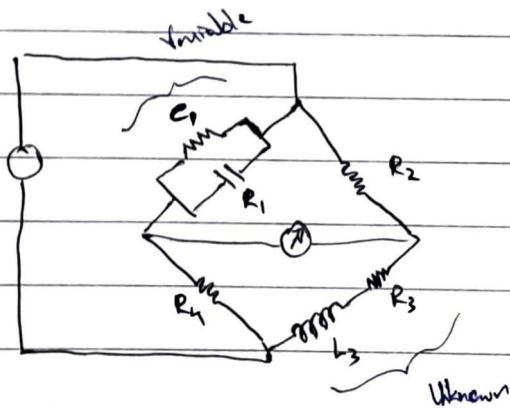
Where ~~ω~~ R, C ,
is set to balance
the bridge

$$R_4 = \frac{\omega^2 C_1^2 R_2 R_3 R_1}{1 + \omega^2 R_1^2 C_1^2}$$

iii) Maxwell - Wien's bridge

Used to measure unknown
resistance & inductance.

Using the same analysis,
we finally get :



$$L_3 = C_1 R_2 R_4$$

$$R_3 = \frac{R_2 R_4}{R_1}$$

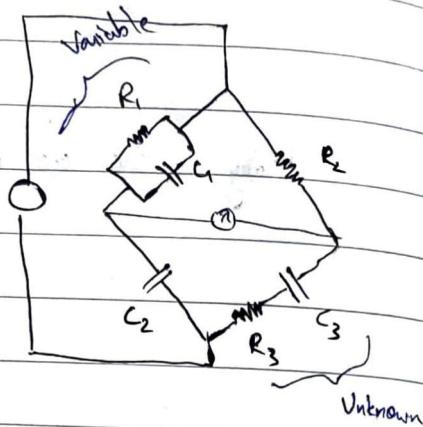
} For a balanced
bridge

Note: The unknowns are independent of frequency ω .

v) Schering bridge .

Used to measure unknown capacitance & resistance.

Using the same analysis,
we get :



$$C_3 = \frac{R_1 C_2}{R_2} \quad \left. \begin{array}{l} \\ \end{array} \right\} \text{For a balanced bridge}$$

$$R_3 = \frac{C_1 R_2}{C_2}$$

Note: The unknowns are independent of frequency ω .

Lecture : 11.03.2021

- Capacitance pick-up

Concept : By changing the overlapping area of the two plates of the capacitor, we change the capacitance which is picked up.

We have seen several bridges that can be used to pick up this capacitance.

Similarly, instead of changing the overlapping area, we can also change the distance between the plates.

We can also use these concepts to measure the angular displacement. Like rotation to change overlap area.

Sensitivity is adjusted based on whether it's used for long range or short range measurements.

$$C = \frac{\epsilon_0 \epsilon_r A}{x}$$

$8.854 \times 10^{-12} \text{ F/m}$

$$\text{Non linearity } \eta = \frac{\Delta h}{h + \Delta h}$$

Sensitivity

- i) When area is changed, i.e., the plate is moved parallel to itself

$$\frac{\Delta C}{\Delta A} = \frac{\epsilon_0 \epsilon_r A}{x} \quad \text{where, } \Delta A = w \Delta l$$

- ii) When the distance b/w the plates ~~is~~ is changed

$$\frac{\Delta C}{\Delta x} = \frac{\epsilon_0 \epsilon_r A}{x^2} \quad (\text{mod for +ve value})$$

- Piezoelectric Transducers.

- Generating electric charge within them upon deformation.
This is called piezoelectric effect.
- The effect is linear & reversible.
- So it is direction sensitive. Polarity switches for compression and tension. It helps us detect the direction.
- Used for measuring force, acceleration & pressure.
- Low response time.

Refer the slides for more.

- Integration route or the differentiation route?

When we have a system with a lot of error, differentiation tends to amplify this noise, while integration averages out the noise and actually give near to perfect solutions.

Lecture : 15.03.2021 & 16.03.2021

- Strain gauge

Concept: Change in resistance with elongation or compression of the strain gauge. This change is picked up using the Wheatstone bridge.

We know,

$$R = \frac{SL}{A}$$

$$\frac{dR}{R} = \left[\frac{L}{A} \frac{ds}{s} + \frac{s}{A} \frac{dl}{l} - \frac{sl}{A^2} \frac{da}{a} \right] \times \frac{A}{sl}$$

$$\frac{dR}{R} = \frac{ds}{s} + \frac{dl}{l} - \frac{da}{a}$$

Also,

$$\frac{A+da}{A} = (1 - \epsilon v)^2 \approx 1 - 2\epsilon v$$

$$\frac{da}{a} = -2\epsilon v = -2v \frac{dl}{l}$$

$$\frac{dR}{R} = \frac{ds}{s} + (1+2v) \frac{dl}{l}$$

$\underbrace{\hspace{2cm}}$
Change in
resistance due to
piezoelectric effect

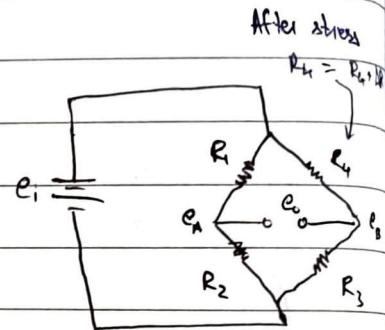
$\underbrace{\hspace{2cm}}$
Change in resistance
due to the change
in length

Change in mechanical property (Resistivity) due to mechanical stress/Force : Piezo electric effect

$$\text{Gauge factor} = \frac{\Delta R/R}{\Delta L/L} = \frac{\Delta S/S}{\Delta L/L} + (1+2\nu)$$

$$e_A - e_B = e_0$$

$$e_0 = \left[\frac{R_1}{R_1 + R_2} - \frac{R_4}{R_3 + R_4} \right] e_i$$



$$\text{Let } R_1 = R_2$$

$$R_3 = R_4 = R_{eff}$$

$$e_0 = e_i \left[\frac{1}{2} - \frac{R_4 + \Delta R}{R_{eff} + R_4 + \Delta R} \right]$$

$$\text{let } S_g = \frac{\Delta R/R_4}{\Delta L/L} \Rightarrow \Delta R = S_g \varepsilon R_4$$

$$\begin{aligned} e_0 &= e_i \left[\frac{1}{2} - \frac{R_4 + S_g \varepsilon R_4}{2R_4 + S_g \varepsilon R_4} \right] \\ &= e_i \left[\frac{1}{2} - \frac{1 + \varepsilon S_g}{2 + \varepsilon S_g} \right] \end{aligned}$$

$$e_0 = \frac{e_i (-\varepsilon_g)}{2(2 + \varepsilon_g)}$$

} Non linear

Here,, input = Strain (ε)
output = Voltage (e_0)

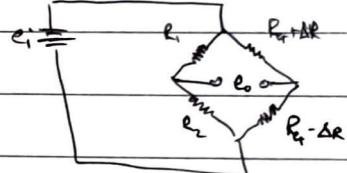
How to make the input output relation linear?

Case I Use two strain gauges



$$e_0 = e_i \left\{ \frac{1}{2} - \frac{R_0 + \Delta R}{2R_0} \right\}$$

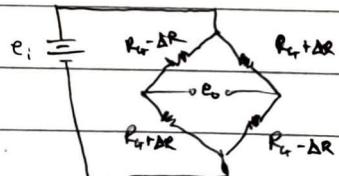
$$= e_i \left(-\frac{\Delta R}{2R_0} \right)$$



$$= -e_i \varepsilon_g \varepsilon / 2$$

Case II Use four strain gauges

$$e_0 = e_i \left\{ \frac{R_0 - \Delta R}{2R_0} - \frac{R_0 + \Delta R}{2R_0} \right\}$$



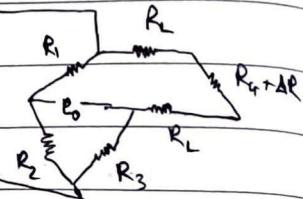
$$= \frac{e_i (-2\Delta R)}{2R_0}$$

$$= -e_i \varepsilon_g \varepsilon$$

$$\frac{e_0}{e_i} = \left[\frac{\frac{1}{R_1 + R_2}}{\frac{2R_L + R_g + \Delta R}{R_g + 2R_L + R_g + \Delta R}} \right] e_i$$

$$= \left[\frac{\frac{1}{2} - \frac{2R_L/R_{eq} + 1 + \Delta R/R_{eq}}{R_g + 2R_L/R_{eq} + 1 + \Delta R/R_L}}{\frac{R_g + 2R_L/R_{eq} + 1 + \Delta R/R_L}{R_L}} \right]$$

$$e_i = \frac{1}{R_1 + R_2 + R_g + \Delta R}$$



Lead wire resistance

Lecture : 18.03.2021

* Second order systems

$$\text{General eqn} : a_2 \frac{d^2q_i}{dt^2} + a_1 \frac{dq_i}{dt} + a_0 q_i = b_0 g_i$$

Let's define :

$$k = \frac{b_0}{a_0}$$

Static sensitivity
Unitless

$$\omega_n = \sqrt{\frac{a_0}{a_2}}$$

Undamped natural frequency
Unit : Hz

$$\xi = \frac{a_1}{2\sqrt{a_0 a_2}}$$

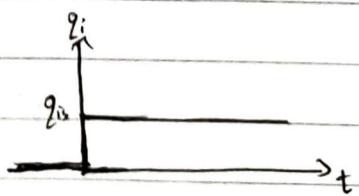
Damping ratio
Unitless

$$\text{Equivalent eqn} : \frac{1}{\omega_n^2} \frac{d^2q_i}{dt^2} + 2\xi \frac{dq_i}{dt} + q_i = kg_i$$

I Step response to a 2nd order system

$$\text{IC} : q_i(0) = 0$$

$$q'_i(0) = 0$$



$$\frac{1}{w_n^2} q_0'' + \frac{2\xi}{w_n} q_0' + q_0 = k q_0$$

Taking Laplace transform:

$$\frac{1}{w_n^2} [s^2 Q_0 - s q_0(0) - q_0'(0)] + \frac{2\xi}{w_n} [s Q_0 - q_0(0)] + Q_0 = \frac{k q_0}{s}$$

$$\left[\frac{s^2}{w_n^2} + \frac{2\xi}{w_n} + 1 \right] Q_0 = \frac{k q_0}{s}$$

$$Q_0 = \frac{k q_0 w_n^2}{s(s^2 + 2\xi w_n + w_n^2)}$$

$$= \frac{k q_0 w_n^2}{[(s + \xi w_n)^2 + (w_n^2 - \xi^2 w_n^2)]s}$$

$$= \frac{k q_0 w_n}{\sqrt{1-\xi^2}} \times \frac{1}{s} \times \frac{\omega \sqrt{1-\xi^2}}{(s+\xi w_n)^2 + ((w_n^2 - \xi^2 w_n^2)s)}$$

Now proceed using the convolution theorem.

Final solution:

$$\text{Case 1: } \xi = 0$$

$$q_0(t) = [1 - \cos(\omega_n t)] k q_0$$

Case 2 : $0 < \xi < 1$

$$q_0(t) = \left[1 - \frac{e^{-\xi w_n t}}{\sqrt{1-\xi^2}} \sin(\sqrt{1-\xi^2} w_n t + \phi) \right] k g_{13}$$

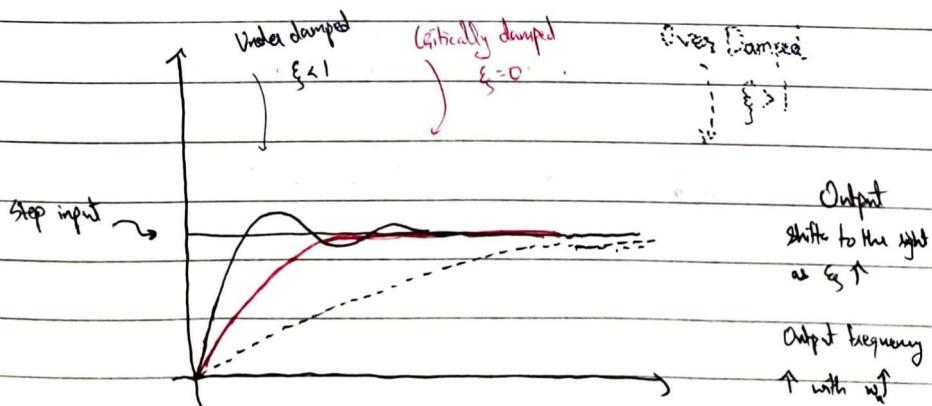
Where, $\theta = \sin^{-1}(\sqrt{1-\xi^2})$

Case 3 : $\xi = 1$

$$q_0(t) = k g_{13} (1 - e^{-w_n t} - w_n t e^{-w_n t})$$

Case 4 : $\xi > 1$

$$\begin{aligned} q_0(t) = & \frac{1}{k g_{13}} \left(\frac{\xi + \sqrt{\xi^2 - 1}}{2\sqrt{\xi^2 - 1}} \right) \exp \left[(-\xi + \cancel{\sqrt{\xi^2 - 1}}) w_n t \right] \\ & + \left(\frac{\xi - \sqrt{\xi^2 - 1}}{2\sqrt{\xi^2 - 1}} \right) \exp \left[(-\xi - \cancel{\sqrt{\xi^2 - 1}}) w_n t \right] \end{aligned}$$

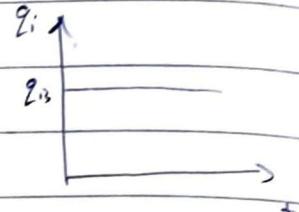


Homework: ME226

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Q1

Step response:



Initial condition: $q_0(0) = 0$

$$\frac{dq_0}{dt} = 0 \quad |_{t=0}$$

$$\frac{1}{\omega_n^2} \frac{d^2 q_0}{dt^2} + \frac{2\xi}{\omega_n} \frac{dq_0}{dt} + q_0 = kq_B$$

Laplace
transform

$$\frac{1}{\omega_n^2} [s^2 Q_0 - s q_0(0) - q'(0)] + \frac{2\xi}{\omega_n} [sq_0 - q_0(0)] + Q_0 = \frac{kq_B}{s}$$

$$Q_0 = \frac{kq_B \omega_n^2}{s(s^2 + 2\xi\omega_n s + \omega_n^2)}$$

$$= \frac{kq_B \omega_n}{\sqrt{1-\xi^2}} \times \frac{1}{s} \times \frac{\sqrt{1-\xi^2} \omega_n}{s^2 + 2\xi\omega_n s + \omega_n^2}$$

$f(s)$

$(g(s))$

$$\frac{b}{(sa)^2 + b^2} \Rightarrow b = \omega_n \sqrt{1-\xi^2}$$

$$a = -\omega_n \xi$$

Inverse Laplace using convolution theorem,

$$q_0 = \frac{kq_B \omega_n}{\sqrt{1-\xi^2}} \times \int u(\tau) e^{-\xi\omega_n(t-\tau)} \sin[\sqrt{1-\xi^2} \omega_n(t-\tau)] d\tau$$

Look at the denominator: $s^2 + 2\xi w_n s + w_n^2$

Case I : $\xi > 1 \Rightarrow$ ~~2 real roots~~

$$\zeta_1, \zeta_2 = -w_n \xi \pm w_n \sqrt{\xi^2 - 1}$$

$$Q_0 = \frac{k g_{13} \times \zeta_1 \zeta_2}{\cancel{s(s-\zeta_1)(s-\zeta_2)}} = k g_{13} \frac{1}{s} \times \frac{\zeta_1 \zeta_2}{\cancel{(s-\zeta_1)(s-\zeta_2)}} \underbrace{F(s)}_{G(s)}$$

~~$q_0 = k g_{13} \left[1 - \left| \frac{s e^{\zeta_1 t} - \zeta_1 e^{\zeta_1 t}}{s - \zeta_1} \right| \right]$~~

$$q_0 = k g_{13} \cancel{\int_0^t} H(t-\tau) \left(\frac{e^{\zeta_1 t} - e^{\zeta_2 t}}{\zeta_2 - \zeta_1} \right) d\tau$$

$$= k g_{13} \cancel{\left[\frac{s_1 e^{\zeta_1 t} - s_2 e^{\zeta_2 t}}{\zeta_2 - \zeta_1} \right]}^t$$

$$= k g_{13} \left[\frac{s_1 e^{\zeta_1 t} - s_2 e^{\zeta_2 t}}{\zeta_2 - \zeta_1} + 1 \right]$$

$$= k g_{13} \left[\frac{w_n (-\xi + \sqrt{\xi^2 - 1}) e^{w_n t (-\xi + \sqrt{\xi^2 - 1})}}{-2 w_n \sqrt{\xi^2 - 1}} - \frac{w_n (-\xi - \sqrt{\xi^2 - 1}) e^{w_n t (\xi - \sqrt{\xi^2 - 1})}}{+1} \right]$$

Case II : $\xi = 1 \Rightarrow \zeta_1 = \zeta_2 = -w_n \xi = -w_n$

$$Q_0 = k g_{13} w_n^2 \times \frac{1}{\cancel{s(s+w_n)^2}} \underbrace{F(s)}_{G(s)}$$

$$q_0 = k g_{13} w_n^2 \int_0^t H(t-\tau) \times T e^{-w_n \tau} d\tau$$

$$= k g_{13} w_n^2 \left[\gamma(-\lambda_{1n}) e^{-w_n t} - \frac{1}{2} w_n^2 e^{-w_n t} \right]_0^t$$

$$= k g_{13} \left[-t w_n e^{-w_n t} - e^{-w_n t} + 1 \right]$$

Case III : $\xi < 1$

$$Q_0 = \frac{k g_{13} w_n}{\sqrt{1-\xi^2}} \times \underbrace{\frac{1}{s}}_{F(s)} \times \frac{\sqrt{1-\xi^2} w_n}{\underbrace{s^2 + 2\xi w_n + w_n^2}_{G(s)}}$$

$$q_0 = \frac{k g_{13} w_n}{\sqrt{1-\xi^2}} \int_0^t H(\frac{t-\tau}{\xi}) e^{-\xi w_n (\frac{t-\tau}{\xi})} \sin(\sqrt{1-\xi^2} w_n \frac{t-\tau}{\xi}) d\tau$$

$$= \frac{k g_{13} w_n}{\sqrt{1-\xi^2}} \int_0^t e^{-\xi w_n \tau} \sin(\sqrt{1-\xi^2} w_n \tau) d\tau$$

$$= \frac{k g_{13} w_n}{\sqrt{1-\xi^2}} \int_0^t e^{a\tau} \sin(b\tau) d\tau$$

$$a = -\xi w_n, \quad b = \sqrt{1-\xi^2} w_n$$

$$q_u = \frac{k g_{13} w_n}{\sqrt{1-\xi^2}} \left\{ \frac{e^{a\tau}}{a^2+b^2} (a \sin b\tau - b \cos b\tau) \right\}_0^t$$

$$\frac{q_0}{k g_{13}} = 1 - \frac{e^{\xi w_n t}}{\sqrt{1-\xi^2}} \left(\xi \sin(\sqrt{1-\xi^2} w_n t) - \sqrt{1-\xi^2} \cos(\sqrt{1-\xi^2} w_n t) \right)$$

Homework : ME 226

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Q2.

Partial fractions :

$$Q_0 = \frac{k q_{13} w_n^2}{s(s^2 + 2\xi w_n s + w_n^2)}$$

$\underbrace{\qquad\qquad\qquad}_{\downarrow}$

Roots for this quadratic are :

$$\begin{aligned}s_1, s_2 &= -2\xi w_n \pm \frac{\sqrt{4\xi^2 w_n^2 - 4 w_n^2}}{2} \\&= -\xi w_n \pm w_n \sqrt{\xi^2 - 1} \\&= w_n (\xi \pm \sqrt{\xi^2 - 1})\end{aligned}$$

So,

$$s_1 = w_n (\xi + \sqrt{\xi^2 - 1})$$

$$s_2 = w_n (\xi - \sqrt{\xi^2 - 1})$$

Case I : $\xi > 1$

$$\frac{1}{s(s^2 + 2\xi w_n s + w_n^2)} = \frac{A}{s} + \frac{B}{s-s_1} + \frac{C}{s-s_2}$$

$$\Rightarrow A(s-s_1)(s-s_2) + B(s-s_2) + C(s-s_1) = 0$$

$$A + B + C = 0 \quad \text{--- (1)}$$

$$A(s_1 + s_2) + B s_2 + C s_1 = 0 \quad \text{--- (2)}$$

$$A s_1 s_2 = 1 \quad \text{--- (3)}$$

$$A = \frac{1}{s_1 s_2} = \frac{1}{\omega_n^2}$$

$$\text{Eq } [(2) - s_1 \times (1)] \Rightarrow A s_2 + B (s_2 - s_1) = 0$$

$$B = -\frac{As_2}{s_2 - s_1}$$

$$= \frac{1}{s_1 (s_1 - s_2)}$$

$$B = \frac{1}{s_1 (s_1 - s_2)} = \frac{1}{\omega_n (\xi + \sqrt{\xi^2 - 1})} \times \frac{1}{2\omega_n \sqrt{\xi^2 - 1}}$$

$$= \frac{1}{2\omega_n^2 (\xi + \sqrt{\xi^2 - 1}) \sqrt{\xi^2 - 1}} = -\frac{1}{2\omega_n^2} \frac{(\xi + \sqrt{\xi^2 - 1})}{\sqrt{\xi^2 - 1}}$$

$$C = -A - B$$

$$= \frac{1}{2\omega_n^2} \frac{(\xi - \sqrt{\xi^2 - 1})}{\sqrt{\xi^2 - 1}}$$

Inverse Laplace

$$\text{So, } Q_o = k g_o \omega_n^2 \left[\frac{A}{s} + \frac{B}{s - s_1} + \frac{C}{s - s_2} \right]$$

$$q_o = k g_o \omega_n^2 (A + B e^{s_1 t} + C e^{s_2 t})$$

$$q_0 = k g_{12} \left| \begin{array}{l} 1 - \frac{(\xi + \sqrt{\xi^2 - 1}) e^{wt(-\xi + \sqrt{\xi^2 - 1})}}{2\sqrt{\xi^2 - 1}} \\ + \frac{\xi - \sqrt{\xi^2 - 1}}{2\sqrt{\xi^2 - 1}} e^{wt(-\xi - \sqrt{\xi^2 - 1})} \end{array} \right|$$

Case II : $\xi = 1$

Equal roots

~~Product of roots~~

$$g_1 = g_2 = -w_n$$

$$Q_0 = \frac{k g_{12} w_n^2}{s(s+w_n)^2}$$

$$\frac{1}{s(s+w_n)^2} = \frac{A}{s} + \frac{B}{s+w_n} + \frac{C}{(s+w_n)^2}$$

$$A(s+w_n)^2 + Bs(s+w_n) + Cs = 1$$

$$A + B = 0$$

$$A = \lambda w_n^2$$

$$B = -\lambda w_n^2$$

$$2Aw_ns + Bw_ns + Cs = 0$$

$$C = -\lambda w_n$$

$$Q_o = k g_{12} w_n^2 \left[\frac{1}{w_n^2 s} - \frac{1}{w_n^2(s+w_n)} - \frac{1}{w_n(s+w_n)^2} \right]$$

$$= k g_{12} \left[1 - \frac{1}{s+w_n} - \frac{w_n}{(s+w_n)^2} \right]$$

Inverse Laplace :

$$q_o = k g_{12} \left[1 - e^{-w_n t} - w_n t e^{-w_n t} \right] \cancel{/}$$

Case III : $\xi < 1$

Imaginary roots :

$$Q_o = \frac{k g_{12} w_n^2}{s(s^2 + 2w_n \xi s + w_n^2)}$$

$$\frac{1}{s(s^2 + 2w_n \xi s + w_n^2)} = \frac{A}{s} + \frac{Bs + C}{(s^2 + 2w_n \xi s + w_n^2)}$$

$$A(s^2 + 2w_n \xi s + w_n^2) + (Bs + C)s = 1$$

$$A = \lambda w_n^2$$

$$B = -A = -\lambda w_n^2$$

$$C = -2A\xi w_n = -\frac{2\xi}{w_n}$$

$$Q_0 = k g_{13} \left\{ \frac{1}{s} - \frac{(s + 2\xi w_n)}{s^2 + 2\xi w_n s + w_n^2} \right\}$$

$$= k g_{13} \left\{ \frac{1}{s} - \frac{(s + 2\xi w_n)}{(s + w_n \xi)^2 + w_n^2(1 - \xi^2)} \right\}$$

$$= k g_{13} \left\{ \frac{1}{s} - \frac{(s + \xi w_n)}{(s + w_n \xi)^2 + w_n^2(1 + \xi^2)} - \frac{\xi w_n}{(s + w_n \xi)^2 + w_n^2(1 + \xi^2)} \right\}$$

Let $\Rightarrow w_d = w_n \sqrt{1 - \xi^2}$

$$Q_0 = k g_{13} \left\{ \frac{1}{s} - \frac{(s + \xi w_n)}{(s + w_n \xi)^2 + w_d^2} - \frac{\xi w_d}{(s + w_n \xi)^2 + w_d^2} \times \frac{1}{\sqrt{1 - \xi^2}} \right\}$$

Inverse
laplace

$$\frac{s-a}{(sa)^2 + b^2}$$

$$\frac{u}{(s-a)^2 + b^2}$$

$$q_0 = k g_{13} \left\{ 1 - e^{-\xi w_n t} \cos(w_n t) - \frac{\xi}{\sqrt{1-\xi^2}} e^{-\xi w_n t} \sin(w_n t) \right\}$$

Let, $\phi = \sin^{-1}(\sqrt{1-\xi^2})$

$$\sin\phi = 1 - \xi^2$$

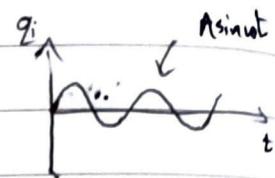
$$\sqrt{1-\xi^2} \cos(w_n t) + \xi \sin(w_n t) = \sin(w_n t + \phi)$$

$$q_0 = k g_{13} \left\{ 1 - \frac{e^{-\xi w_n t}}{\sqrt{1-\xi^2}} \times \sin(\sqrt{1-\xi^2} w_n t + \phi) \right\}$$

Lecture : 23.03.2021

- Frequency response of a 2nd order system

$$\frac{1}{w_n^2} \frac{d^2 q_o}{dt^2} + \frac{2\xi}{w_n} \frac{dq_o}{dt} + q_o = k A \sin \omega t$$



- For the homogeneous eqn:

$$\frac{1}{w_n^2} \frac{d^2 q_{oh}}{dt^2} + \frac{2\xi}{w_n} \frac{dq_{oh}}{dt} + q_{oh} = 0$$

Let us assume, $q_{oh} = e^{\lambda t}$, Substituting

$$\frac{\lambda^2}{w_n^2} + \frac{2\xi\lambda}{w_n} + 1 = 0$$

$$\lambda^2 + 2\xi w_n \lambda + w_n^2 = 0$$

$$\text{So, } \lambda_1, \lambda_2 = -\xi w_n \pm w_n \sqrt{\xi^2 - 1}$$

$$\text{And, } q_{oh} = c_1 e^{\lambda_1 t} + c_2 e^{\lambda_2 t}$$

For the particular solⁿ: $q_{op} = B \cos \omega t + C \sin \omega t$

$$\frac{-\omega^2}{w_n^2} (B \cos \omega t + C \sin \omega t) + \frac{2\xi \omega}{w_n} (-B \sin \omega t + C \cos \omega t) + (B \cos \omega t + C \sin \omega t) = kA \sin \omega t$$

Solve for B & C in this eqn.

Finally, we get

$$q_o = q_{oh} + q_{op}$$

$$q_o = c_1 e^{\lambda_1 t} + c_2 e^{\lambda_2 t} + B \cos \omega t + C \sin \omega t$$

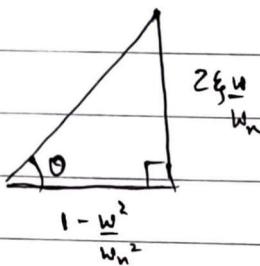
$$q_o = c_1 e^{w_n t(-\xi + \sqrt{\xi^2 - 1})} + c_2 e^{-w_n t(\xi + \sqrt{\xi^2 - 1})} + \frac{2\xi (w_n) kA}{(1 - w_n^2)^2 + 4\xi^2 w_n^2} \cos \omega t + \frac{(1 - w_n^2)}{(1 - w_n^2)^2 + 4\xi^2 w_n^2} \sin \omega t$$

} Transient : gets affected by initial condⁿ
} Steady state

So, the steady state output will always be the same irrespective of the initial condⁿ.

For steady state,

$$\frac{q_0 \text{ steady}}{kA} = \frac{2\zeta_p (\omega/\omega_n) \cos \omega t + (1 - \frac{\omega^2}{\omega_n^2}) \sin \omega t}{(1 - \frac{\omega^2}{\omega_n^2})^2 + 4\zeta_p^2 \frac{\omega^2}{\omega_n^2}}$$



$$\sin \theta = \frac{2\zeta_p \omega / \omega_n}{\sqrt{(1 - \frac{\omega^2}{\omega_n^2})^2 + 4\zeta_p^2 \frac{\omega^2}{\omega_n^2}}}$$

$$\frac{q_0}{kA} = \frac{\sin \theta \cos \omega t + \cos \theta \sin \omega t}{\sqrt{(1 - \frac{\omega^2}{\omega_n^2})^2 + 4\zeta_p^2 \frac{\omega^2}{\omega_n^2}}}$$

$$= \frac{\sin(\omega t + \theta)}{\sqrt{(1 - \frac{\omega^2}{\omega_n^2})^2 + 4\zeta_p^2 \frac{\omega^2}{\omega_n^2}}}$$

- Frequency Response (Without transient state)

$$\frac{1}{\omega_n^2} \frac{d^2 q_0}{dt^2} + \frac{2\zeta_p}{\omega_n} \frac{dq_0}{dt} + q_0 = kq_0$$

Let $D = \frac{d}{dt}$

$$\left(\frac{D^2}{w_n^2} + \frac{2\xi D + 1}{w_n} \right) q_0 = k q_{13}$$

$$\frac{q_0}{k q_{13}} = \frac{1}{\left(\frac{D^2}{w_n^2} + \frac{2\xi D + 1}{w_n} \right)}$$

For $D = jw$,

$$\frac{q_0}{k q_{13}} = \frac{1}{\left[\frac{j^2 w^2}{w_n^2} + \frac{2\xi jw + 1}{w_n} \right]}$$

0.1

q_is) phase

1

89495 0.0201992

76362 0.0416425

52023 0.0658387

11366 0.0949517

.16372 0.1325515

73779 0.1853479

083558 0.2679104

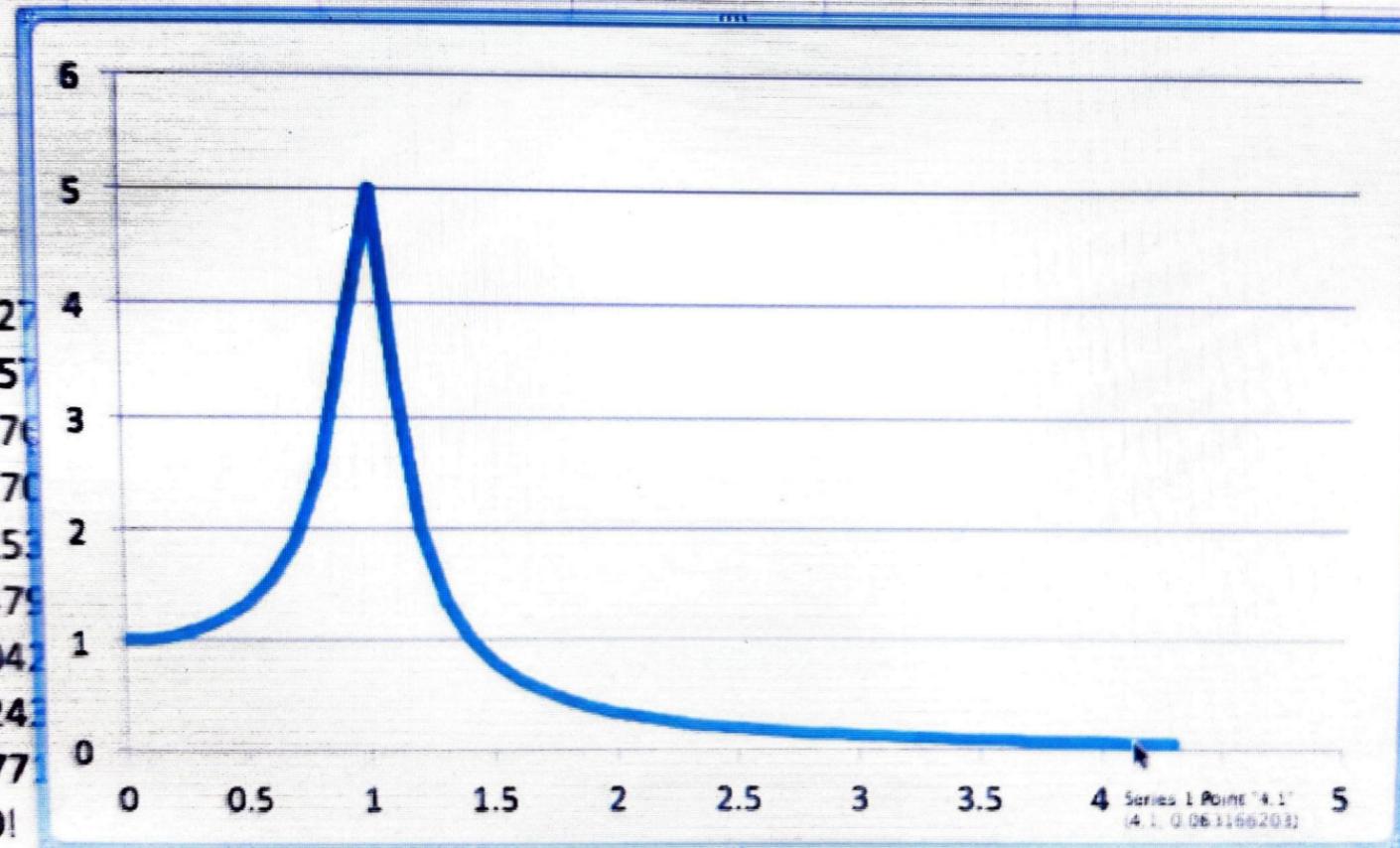
336541 0.4182241

208036 0.7583777

5 #DIV/0!

797975 -0.80864979

521721 -0.49934672



0.2

_ls) phase

1 0

0482 0.020047439

7753 0.040382076

0615 0.061304161

0685 0.083141232

1981 0.106264863

6942 0.131111648

5839 0.158210768

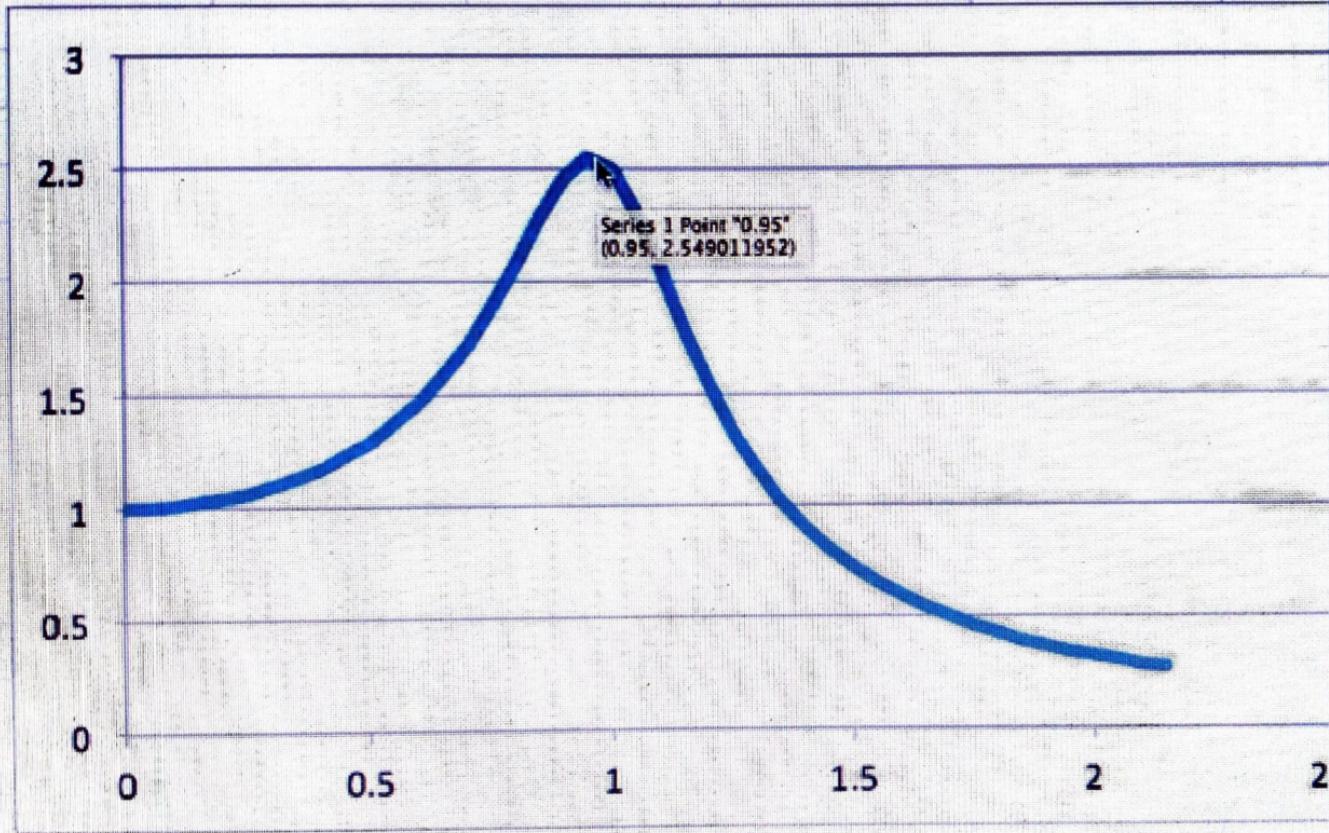
5067 0.188221505

5008 0.221985683

1325 0.260602392

9166 0.305535759

1434 0.35877067



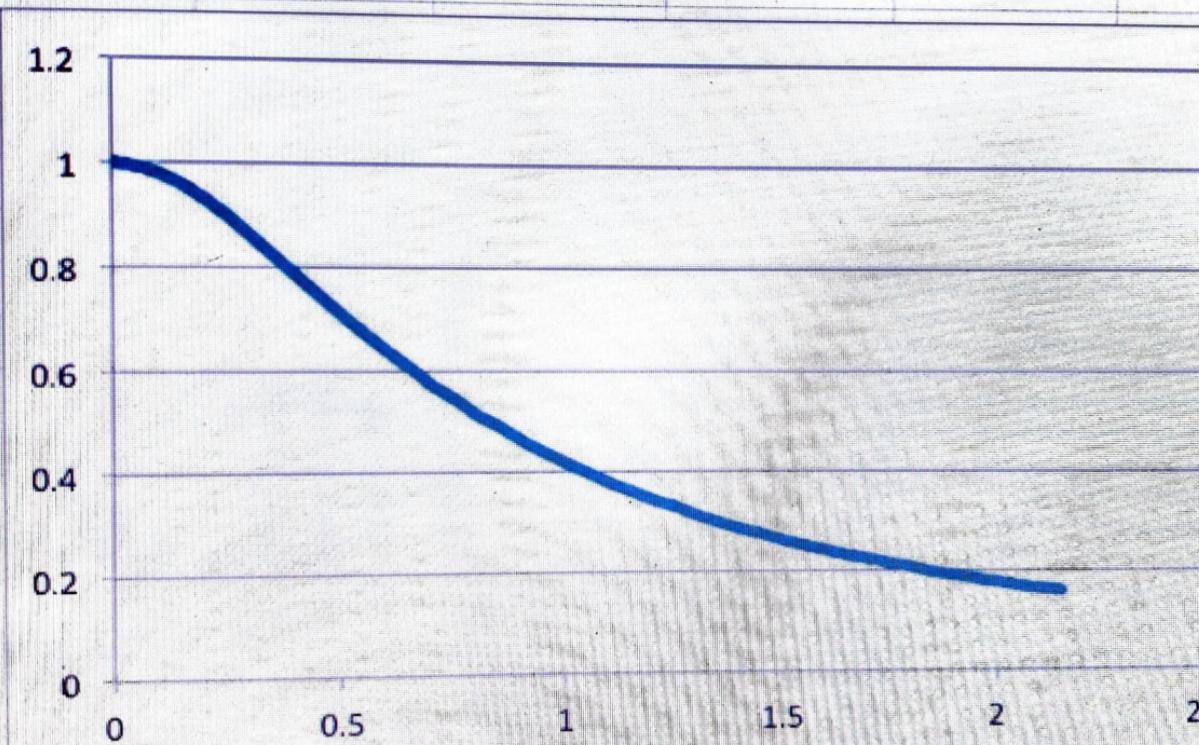
New Scale Other Line Column Way Loss Select Switch Plot

f= q_0/(K q_is)

F G H I J K L

I 1.2
_0/(K q_is) phase

1	0
0.9953298	0.119725398
0.98166677	0.237835927
0.95998372	0.352871862
0.93169499	0.463647609
0.89842283	0.569313191
0.86178081	0.669357414
0.82321781	0.763567571
0.78393406	0.851966327
0.74485829	0.93474387
0.70666525	1.012197011
0.66981381	1.084680589



II Impulse response to a 2nd order system : $q_i = AS$

Case 1 : $\xi \neq 0$

$$\underbrace{q_o(t)}_{KA} = w_n \sin(w_n t)$$

Case 2 : $0 < \xi_g < 1$

$$\frac{q_0(t)}{kA} = \frac{w_n e^{-\xi_g w_n t}}{\sqrt{1-\xi_g^2}} \sin(w_n t \sqrt{1-\xi_g^2})$$

Case 3 : $\xi_g = 1$

$$\frac{q_0(t)}{kA} = \frac{w_n^2 t}{e^{-w_n t}}$$

Case 4 : $\xi_g > 1$

$$\frac{q_0(t)}{kA} = \frac{w_n}{2\xi_g^2 - 1} \left[e^{-w_n t (\xi_g - \sqrt{\xi_g^2 - 1})} - e^{-w_n t (\xi_g + \sqrt{\xi_g^2 - 1})} \right]$$

Q1. Impulse function : AS

$$\frac{1}{w_n^2} \frac{d^2 q_0}{dt^2} + \frac{2\xi}{w_n} \frac{dq_0}{dt} + q_0 = kAS$$

IC:
$q_0(0) = 0$
$\dot{q}_0(0) = 0$

Taking Laplace,

$$\frac{1}{w_n^2} [s^2 Q_0 - s q_0(0) - \dot{q}_0(0)] + \frac{2\xi}{w_n} [s Q_0 - q_0(0)] + Q_0 = kA$$

$$Q_0 \left[\frac{s^2}{w_n^2} + \frac{2\xi s}{w_n} + 1 \right] = kA$$

$$Q_0 = \frac{kAw_n^2}{s^2 + 2\xi w_n s + w_n^2}$$

$$= \frac{kAw_n^2}{(s + w_n\xi)^2 + (w_n^2 - \xi^2 w_n^2)}$$

Case 1 : $w_n^2 - \xi^2 w_n^2 > 0$

$$Q_0 = \frac{kAw_n^2}{\sqrt{w_n^2 - \xi^2 w_n^2}} \frac{\sqrt{w_n^2 - \xi^2 w_n^2}}{(s + w_n\xi)^2 + (\sqrt{w_n^2 - \xi^2 w_n^2})^2}$$

Taking inverse Laplace,

$$q_0(t) = \frac{kAw_n^2}{\sqrt{w_n^2 - \xi^2 w_n^2}} e^{-w_n\xi t} \sin(w_n t \sqrt{1 - \xi^2})$$

Case 2 : $w_n^2 - \xi^2 w_n^2 = 0$

$$Q_0 = \frac{kAw_n^2}{(s + w_n\xi)^2}$$

Taking inverse Laplace,

$$q_0(t) = \frac{kAw_n^2}{w_n^2} t e^{-w_n t}$$

Case 3 : $w_n^2 - \xi^2 w_n^2 < 0$

$$Q_0 = \frac{kAw_n^2}{(s + w_n\xi)^2 - (\xi^2 w_n^2 - w_n^2)}$$

$$Q_0 = \frac{kA w_n^2}{\sqrt{\xi^2 w_n^2 - w_n^2}} \times \frac{\frac{1}{2} \xi w_n^2 - w_n^2}{(\xi w_n)^2 - (\sqrt{\xi^2 w_n^2 - w_n^2})^2}$$

Taking inverse Laplace:

$$\begin{aligned} q(t) &= \frac{kA w_n^2}{\sqrt{\xi^2 w_n^2 - w_n^2}} \times e^{-w_n t} \left[\frac{e^{w_n \sqrt{\xi^2 - 1} t} + e^{-w_n \sqrt{\xi^2 - 1} t}}{2} \right] \\ &= \frac{kA w_n}{2 \sqrt{\xi^2 - 1}} e^{-w_n t} \left(e^{w_n \sqrt{\xi^2 - 1} t} + e^{-w_n \sqrt{\xi^2 - 1} t} \right) \end{aligned}$$

Q2.

Ramp ~~initial~~ response: q_{it}

$$\frac{1}{w_n^2} \frac{d^2 q_0}{dt^2} + \frac{2\xi}{w_n} \frac{dq_0}{dt} + q_0 = k q_i t$$

Solving for homogeneous eqn: $q_p = e^{xt}$

$$\frac{1}{w_n^2} \lambda^2 e^{xt} + \frac{2\xi \lambda}{w_n} e^{xt} + e^{xt} = 0$$

$$\frac{\lambda^2}{w_n^2} + \frac{2\xi \lambda}{w_n} + 1 = 0$$

$$\lambda^2 + 2\xi w_n \lambda + w_n^2 = 0$$

$$\lambda = -\xi w_n \pm w_n \sqrt{\xi^2 - 1}$$

$$\therefore q_{inh} = c_1 e^{\lambda_1 t} + c_2 e^{\lambda_2 t}$$

$$\text{Where, } \lambda_1 = (-\xi + \sqrt{\xi^2 - 1}) w_n$$

$$\lambda_2 = -(\xi + \sqrt{\xi^2 - 1}) w_n$$

Solving for particular solⁿ: $q_{opp} = At + B$

$$0 + \frac{2\xi}{w_n} A + At + B = k q_i t$$

$$A = k q_i$$

$$B = -\frac{2\xi}{w_n} A = -\frac{2\xi}{w_n} k q_i$$

$$q_i(t) = q_{0,n} + q_{0,p}$$

$$= c_1 e^{\lambda_1 t} + c_2 t e^{\lambda_2 t} + kq_i t - \frac{2\xi}{\omega_n} kq_i$$

Impose initial condition:

$$q_i(0) = c_1 + c_2 - \frac{2\xi}{\omega_n} kq_i = 0$$

$$q'_i(0) = \lambda_1 c_1 + \lambda_2 c_2 + kq_i = 0$$

We get,

$$c_1 = \frac{\lambda_2 \left(-\frac{2\xi}{\omega_n} kq_i \right) - kq_i}{\lambda_2 - \lambda_1} = \frac{kq_i}{2\omega_n} \left[2\xi + \frac{2\xi^2 - 1}{\sqrt{\xi^2 - 1}} \right]$$

$$c_2 = \frac{\lambda_1 \left(-\frac{2\xi}{\omega_n} kq_i \right) - kq_i}{\lambda_2 - \lambda_1} = \frac{kq_i}{2\omega_n} \left[2\xi - \frac{2\xi^2 - 1}{\sqrt{\xi^2 - 1}} \right]$$

So finally, we get

$$q_i(t) = \frac{kq_i}{2\omega_n} \left[2\xi + \frac{2\xi^2 - 1}{\sqrt{\xi^2 - 1}} \right] e^{w_n t (\xi + \sqrt{\xi^2 - 1})} + \frac{kq_i}{2\omega_n} \left[2\xi - \frac{2\xi^2 - 1}{\sqrt{\xi^2 - 1}} \right] e^{-w_n t (\xi + \sqrt{\xi^2 - 1})}$$

$$+ kq_i \left(t - \frac{2\xi}{\omega_n} \right)$$

Case 2: $\xi = 1$

$$\lambda_1 = \lambda_2 = -\xi \omega_n$$

$$q_{0,n} = c_1 e^{\lambda t} + c_2 t e^{\lambda t}$$

$$q_i(t) = c_1 e^{\lambda t} + c_2 t e^{\lambda t} + kq_i \left(t - \frac{2\xi}{\omega_n} \right)$$

Impose IC:

$$q_i(0) = c_1 - \frac{2kq_i \xi}{\omega_n} = 0 \rightarrow c_1 = kq_i \times \frac{2\xi}{\omega_n} = \frac{2kq_i}{\omega_n}$$

$$q'_i(0) = \lambda c_1 + c_2 + kq_i = 0$$

$$-\frac{\omega_n \xi}{\omega_n} 2kq_i + c_2 + kq_i = 0$$

$$c_2 = kq_i$$

$$q_0(t) = \left(\frac{2k_{gi}}{\omega_n} + t k_{gi} \right) e^{-\omega_n t} + k_{gi} \left(t - \frac{2}{\omega_n} \right)$$

Case 3 : $\xi_i < 1$

$$\lambda = -\xi_i \omega \pm i \sqrt{1-\xi_i^2}$$

$$q_0(t) = e^{-\xi_i \omega t} \left[c_1 \cos(\omega_n t \sqrt{1-\xi_i^2}) + c_2 \sin(\omega_n t \sqrt{1-\xi_i^2}) \right] + k_{gi} \left(t - \frac{2\xi_i}{\omega_n} \right)$$

Imposing IC :

$$q_0(0) = c_1 - k_{gi} \times \frac{2\xi_i}{\omega_n} = 0 \Rightarrow c_1 = k_{gi} \frac{2\xi_i}{\omega_n}$$

$$q'_0(0) = -\xi_i \omega_n c_1 + c_2 \omega_n \sqrt{1-\xi_i^2} + k_{gi} = 0 \Rightarrow c_2 = \frac{(2\xi_i^2 - 1) k_{gi}}{\omega_n \sqrt{1-\xi_i^2}}$$

So finally,

$$\frac{q_0(t)}{k_{gi}} = e^{-\xi_i \omega_n t} \left[\frac{2\xi_i}{\omega_n} \cos(\omega_n t \sqrt{1-\xi_i^2}) + \frac{(2\xi_i^2 - 1)}{\omega_n \sqrt{1-\xi_i^2}} \sin(\omega_n t \sqrt{1-\xi_i^2}) \right] + t - \frac{2\xi_i}{\omega_n}$$

G

Lecture

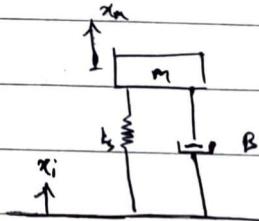
30.03.2021

classmate

Tutor
Maya

- Seismic Pick-up Transducer

A second order instrument to detect earthquakes.



$$k_s(x_i - x_m) + B(x_i - x_m) + m\ddot{x}_i = m\ddot{x}_m + m\ddot{x}_i$$

$$m(\ddot{x}_i - \ddot{x}_m) + B(\ddot{x}_i - \ddot{x}_m) + k_s(x_i - x_m) = m\ddot{x}_i$$

$$m\ddot{x}_i + B\ddot{x}_i + k_s x_0 = m\ddot{x}_i$$

$$\frac{m}{k_s} \ddot{x}_i + \frac{B}{k_s} \dot{x}_i + x_0 = \frac{m}{k_s} \ddot{x}_i$$

$$\left(\frac{D^2}{w_n^2} + \frac{2\xi}{w_n} + 1 \right) x_0 = \frac{D^2}{w_n^2} x_i$$

Note : Assuming that the input sinusoidal

$$w_n = 2\pi f_n$$

$$\frac{m}{k_s} = \frac{1}{w_n^2} = \frac{1}{(2\pi f_n)^2}; \quad \frac{B}{k_s} = \frac{2\xi}{w_n}$$

$$\text{Also, } x_0 = x_i - x_m$$

$$\frac{x_0}{x_i} = \frac{\frac{D^2}{w_n^2}}{\left(\frac{D^2}{w_n^2} + \frac{2\zeta_1 D + 1}{w_n^2} \right)}$$

$$D = j\omega$$

$$\frac{x_0}{x_i} = \frac{-\omega^2/w_n^2}{-\omega^2/w_n^2 + \frac{2\zeta_1(j\omega)}{w_n}}$$

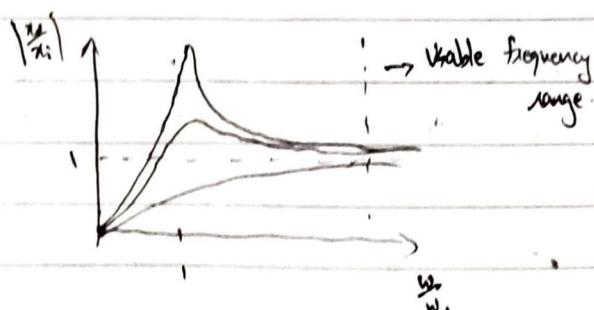
$$= \frac{-\omega^2/w_n^2}{(1 - \omega^2/w_n^2) + \frac{2\zeta_1 \omega}{w_n} j}$$

$$\left| \frac{x_0}{x_i} \right| = \frac{\rho^2}{\sqrt{(1-\rho^2)^2 + (2\zeta_1 \rho)^2}} \quad (\rho = \frac{\omega}{w_n})$$

Analyzing, when $\rho \rightarrow 0$, $\left| \frac{x_0}{x_i} \right| \rightarrow 0$

$$\rho \rightarrow 1, \quad \left| \frac{x_0}{x_i} \right| \gg 1$$

$$\rho \rightarrow \infty, \quad \left| \frac{x_0}{x_i} \right| = 1$$



for large frequency, $\omega \rightarrow \infty$

so we compress the spring, but by the time it reacts,
we elongate it.

High frequency doesn't allow the spring to react &
so $\Delta x = 0$.

Since seismic transducers are made for high ω values,
we can design the system to choose the natural
frequency ω_n .

$$\omega_n = \sqrt{\frac{k_s}{m}}$$

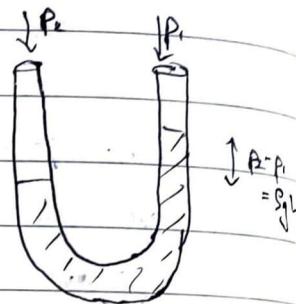
For example, to measure $\omega = 30 \text{ Hz}$

We can set $\omega_n = 5 \text{ Hz}$

Lecture : 01.04.2021

- Pressure manometer

- Manometer is a deflection type gauge.
- Continuous output
- Pressure $\uparrow \rightarrow$ Height \uparrow



Refinements / Corrections required

For high accuracy :

- Thermal expansion of scale needs to be considered
- Variation of manometer fluid density with temperature
- Local value of g should be determined.

Additional error sources?

- Non verticality of tubes
- Difficulty in reading h due to meniscus formed by capillarity.

Lecture : 05.04.2021

- McLeod Gauge

Refer the slides for setup :

Used for measuring very small pressures p_i .

We enclose gas in one of the arms to produce pressure $P > p_i$.

$$\text{So now, } P - p_i = \rho gh$$

Using Boyle's law, $PV = \text{constant}$

$$P \nabla = \rho h A_t$$

$$P = \frac{P \nabla}{h A_t}$$

So,

$$P - p_i = \rho gh$$

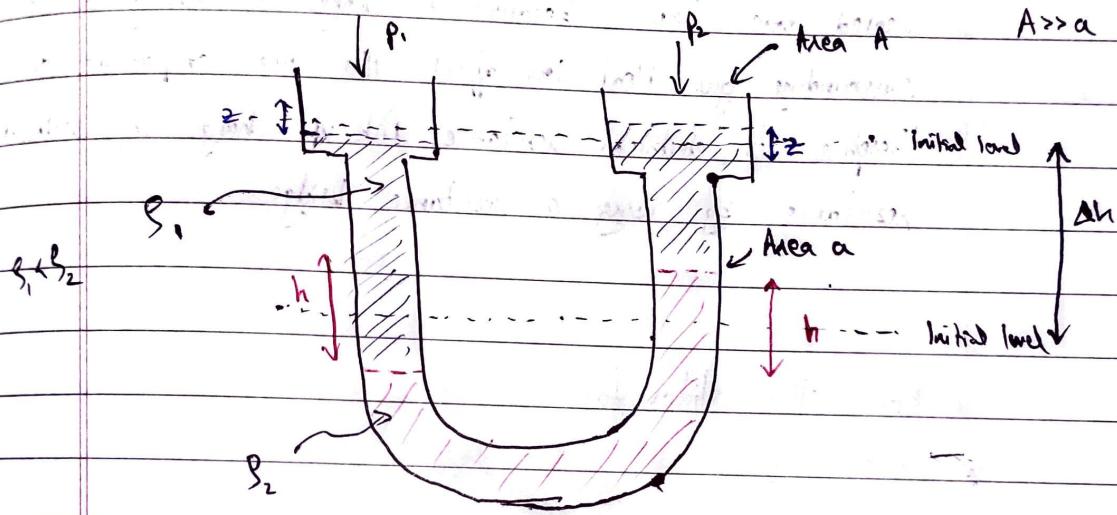
$$\frac{P \nabla}{h A_t} - p_i = \rho gh$$

$$p_i = \frac{\rho gh^2 A_t}{\nabla - h A_t}$$

If $h A_t \ll \nabla$,

$$P_1 = \frac{\rho g h^2 A_1}{A}$$

Micromanometer



Substituting volume, $A \frac{z}{2} = a \frac{h}{2}$

$$A_2 = ah$$

$$P_1 + \rho_1 g \left(\frac{\Delta h - z + h}{2} \right) = P_2 + \rho_2 g \left(\frac{\Delta h + z - h}{2} \right) + \rho_2 gh$$

$$P_1 - P_2 = \rho_1 g (z - h) + \rho_2 g h \approx (\rho_2 - \rho_1) gh$$

Since $A \gg a$, $z \ll h$

Lecture : 08.04.2021

- Pizani Gauge

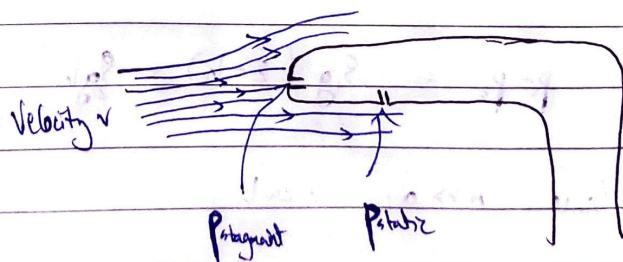
Change in gas pressure will affect heat loss from a heated wire. This is because of change in density of the surrounding gas. Heat loss affects the wire temperature and therefore its electrical resistance. Pick up change in electrical resistance by using a resistance bridge.

Read the slides for:

- Pizani Gauge
- Ionization gauges
- Diaphragm based pressure gauges

* Flow Measurement

- Velocity measurement using Pitot Static Tube



$$P_{atm} + \frac{1}{2} \rho v^2 = P_{stagnant}$$

$$v = \sqrt{\frac{2(P_{stagnant} - P_{atm})}{\rho}}$$

Used in aircrafts and missiles to determine the speed and altitude.

Sources of error in P_{atm}

- i Misalignment of tube axis: This exposes the tube to a component of velocity.
- ii Non zero tube diameter: Flow accelerates around the tube. This causes a decrease in static pressure.
- iii Influence of tube support: Flow stagnates on tube support leading edge. Pressure increases at the stagnation point. This effect propagates upstream.

So, non zero tube diameter accelerates flow.

Influence of tube ~~support~~ de-accelerates flow.

These effects are made to cancel out in Pitot tube
Refer slides for set up.

Sources of error in P_{stop} :

- i) Misalignment: Prevents formation of a true stagnation point at the measuring hole.
 - ii) 2D & 3D velocity fields: Due to finite probe size and non uniform velocity field, ~~there~~ there is actually a range of velocity at the tube tip.
 - iii) Effect of viscosity: Due to finite viscosity fluid, correction in $(P_{\text{stagn}} - P_{\text{total}})$ is required

$$P_{\text{drag}} - P_{\text{total}} = \frac{C_D V^2}{2}$$

$$C = 1 + \frac{h}{k_e}$$

Where, $10 < \text{Re} \left(= \frac{\rho V R}{\mu} \right) < 100$

R is the outer radius of the tube

- Variation of pressure with altitude in our atmosphere

Lecture : 15.04.2021

* Trace particles & properties : PIV

PIV works on the principle that the trace particles are located at two points of time and the distance travelled by them is recorded.

This info gives us the flow velocity.

Synchronizer controls the flow of time of these events.

The two points of time can be long if the flow velocity is low.

Otherwise, the interval should be short if the flow velocity is high.

We can say that the distance travelled by the trace particle could be almost constant.

Now, what particles could be used for tracing?

* Mechanical constraints :

Sp. gravity should be near to that of water.

Shouldn't be too high that it settles down nor too low.

Stokes no. : $St \ll 1$

A particle with low Stokes number follows fluid streamlines.

$$St = \frac{T_p}{T_f} \quad \text{where} \quad T_p = \frac{\rho \cdot d_p^2}{16 \mu}$$

\therefore Particle size should be low.

* Optical constraint :

Should be able to scatter and reflect light.

Larger particle size required to scatter light.

Trade off.

* Chemical constraint

Should not react with fluid.

* Magnetic flowmeter

- Uses Faraday's law of Electromagnetic induction.
- Magnetic field is generated and channeled into the fluid flowing through the pipe.
- Flow of conductive fluid through the magnetic field will produce emf.
- Voltage signals sensed by electrodes located on the flow tube walls.
- Faster the flow, more the voltage generated. Proportional

Major advantage : Linear relation

Constraint : Flow has to be uniform.

Should have sufficient length upstream for a uniform fully developed flow.

Minimum 3-5 diameter length upstream.

Flow meter does not obstruct flow. Can be applied to any kind of dirty or abrasive liquids.

Liquid should be conductive. Constraint : Hydrocarbons & gases cannot be measured.

* Ultrasonic Flowmeter

Uses sound waves to determine the velocity of a fluid flowing in a pipe.

Under flow condition, the frequency of the reflected wave is different due to the Doppler effect.

When the fluid flow is faster, frequency shift increases linearly.

Do not obstruct flow & can be applied to dirty & abrasive liquids.

Advantages :

- Handles high pressure
- High turndown
- Handles extreme temperature
- Low maintenance
- Highly reliable
- Clamped outside the pipe without penetration

Disadvantages :

- High cost
- Sensitive to high stray vibrations
- Problems with pipe diameter change due to buildup

Lecture : 19.04.2021

* Hot wire Anemometry

Principle : Amount of cooling experienced by a heated wire can be related to the local flow velocity.

Current is passed through the wire to maintain a certain temperature.

More the flow, faster the cooling and more is the current required to maintain the temperature.

We need a high response rate. We need a quick response. For this we need a low thermal inertia.

For that, we need the wire to be small and mass to be less.

$$\frac{dQ}{dt} = P - F$$

Q is the internal energy of sensor

P is the electrical input power to sensor

F is the total rate of heat transferred from the sensor

$$F = q_c + q_r + q_u + q_s \quad \leftarrow$$

Convection from
sensor to flow

Conduction to
supporting porous

Radiation from
sensor to surrounding

transfer to quartz
substrate

- q_p (Heat loss to prong support) can be neglected if the ratio ℓ/d of the wire is large.
- q_x (Radiation loss) can be neglected.

Free convection is without an air flow but the air around the heated body getting hot and rising. Replaced by colder air.

Forced convection is due to blown air over the surface to cool it.

$$P = I^2 R, \quad \text{where } R_f = R_f [1 + \alpha (T_s - T_f)]$$

$$F = h \times \text{Area} \times \Delta T = [C_o + C_i \sqrt{U}] \times \text{Area} \times \Delta T$$

$$\frac{dQ}{dt} = P(I, T_s) - F(V, T_s)$$

$$\frac{m C_p dT_s}{dt} = P(I, T_s) - F(V, T_s)$$

We want to find velocity V .

So, we can keep I or T_s constant and relate V to the variable term.

Constant temperature mode is preferred so that the thermal inertia ($m(C_p dT_s/dt)$ term) does not come in the picture.

Lecture : 20.04.2021

* Uncertainty

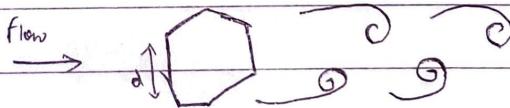
$$\text{let, } y = f(x_1, x_2, x_3, \dots)$$

$$\Delta y = \frac{\partial y}{\partial x_1} \Delta x_1 + \frac{\partial y}{\partial x_2} \Delta x_2 + \dots$$

$$\Delta y = \sqrt{\left[\frac{\partial y}{\partial x_1} \Delta x_1 \right]^2 + \left[\frac{\partial y}{\partial x_2} \Delta x_2 \right]^2 + \dots}$$

* Vortex Flowmeter

Bluff body

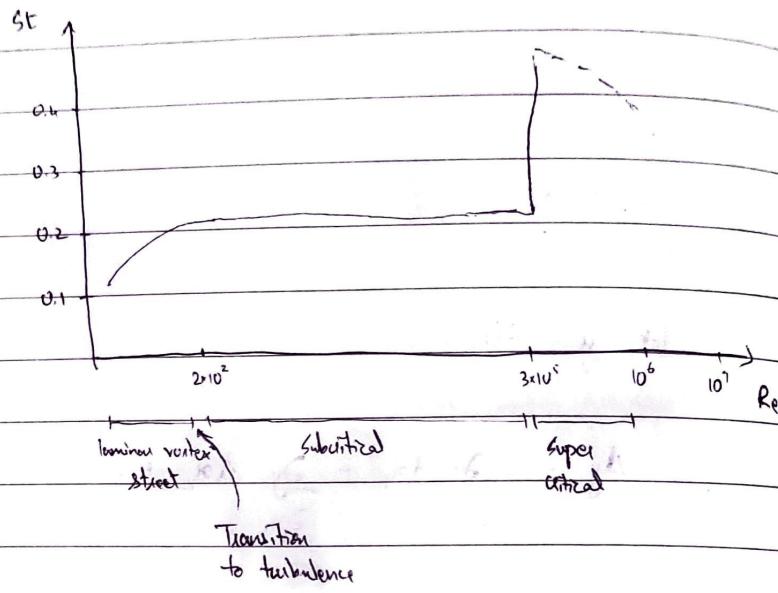


$$\text{Strouhal number (St)} : \frac{d \times f}{V_{av}} \quad (\text{Dimensionless})$$

d - Diameter of bluff body

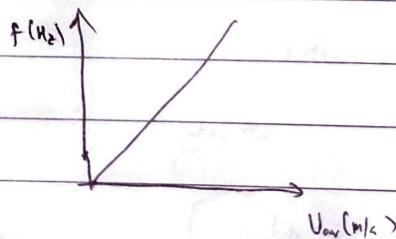
f - Frequency of vortex

V_{av} - Average flow velocity



We can see that St -no. is constant for a range of Reynolds no.

For a given bluff body, we can say that : $f \propto V_{\infty}$



Advantages :

- Low pressure loss
- Large turn down ratio
- Low cost : compared to Coriolis & Ultrasonic flowmeter
- Wide operating temp & pressure range.
- Suitable for various fluids : Air, water, steam
- Linear output : $f \propto V_{\infty}$

Disadvantage :

- Sensitive to upstream flow conditions.

Concept :

- A bluff body intercepts the flow, generating wakes behind it.
- We define Strouhal number and operate ~~it~~ in the range where it is constant. So we have : $f \propto V_{\infty}$
- To find the frequency, we use some sensors, generally piezoelectric, to detect fall & rise in pressure.
- We can optimize the location of the sensor for accurate readings and to minimize permanent pressure loss.