Probing photonic content of the proton using photon-induced dilepton production in $p + \mathrm{Pb}$ collisions at the LHC

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Abstract

We propose a new experimental method to validate the photon parton distribution function inside the proton (photon PDF) at LHC energies. The method is based on the measurement of dilepton production from the $\gamma p \to \ell^+\ell^- + X$ reaction in proton–lead collisions. These experimental conditions guarantee clean environment, both in terms of reconstruction of the final state and in terms of possible background. We firstly calculate the cross sections for this process with collinear photon PDFs, where we identify optimal choice of the scale, in analogy to deep inelastic scattering kinematics. We then include the virtuality of probed photon in the calculations, based on modern parameterizations of deep inelastic structure functions. Finally we estimate rates of the process for the existing LHC data samples.

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I. INTRODUCTION

Precise calculations of various electroweak reactions in pp collisions at the LHC need to account for, on top of the higher-order corrections, the effects of photon-induced subprocesses. The relevant examples are the production of lepton pairs [1–5] or pairs of electroweak bosons [6–13].

Recently, a precise photon distribution inside the proton has been evaluated in Ref. [14]. This approach provides a model-independent determination of the photon PDF (embedded in so-called LUXqed distribution) and it is based on proton structure function and elastic form factor fits in electron–proton scattering.

Up to date, there are no experimentally clean processes identified that would allow to either strongly constrain or verify the calculations. For example, the extraction of photon PDF from isolated photon production in deep inelastic scattering (DIS) [15] or from inclusive $pp \to \ell^+\ell^- + X$ reaction [2, 16, 17] is limited due to large QCD background. On contrary, the elastic part of the photon PDF is verified via exclusive $\gamma\gamma \to \ell^+\ell^-$ process, measured in pp collisions by ATLAS [18, 19], CMS [20, 21] and recently by CMS+TOTEM [22] collaborations.

We therefore propose a new experimental method to constrain photonic content of the proton. Thanks to the large fluxes of quasi-real photons from the lead ion (Pb) at the LHC, the photon-induced dilepton production in $p + $ Pb collision configuration (where Pb serves as a source of elastic photons) is a very clean way to probe photon PDF inside a proton. This process is shown schematically in Fig. 1, where by analogy to DIS, two leading-order diagrams can be identified. Since the photon flux from the ion scales with Z^2 and QCD-induced cross-sections scale approximately with A, the amount of QCD background is greatly reduced comparing to pp case.

Moreover, as this process does not involve the exchange of color with the photon-emitting nucleus, no significant particle production is expected in the rapidity region between the dilepton system and the nucleus. The photon-emitting nucleus is also expected to produce no neutrons because the photons couple to the entire nucleus. Thus a combination of requirements on rapidity gap and zero neutrons in the same direction provide straightforward criteria to identify these events experimentally.

FIG. 1: Schematic graphs for deep inelastic scattering, $\ell^{\pm} p \to \ell^{\pm} + X$ (a) and photon-induced dilepton prodcution, $\gamma p \to \ell^+ \ell^- + X$, in $p + $Pb collisions for t-channel (b) and u-channel (c) lepton exchange.

II. FORMALISM

A. Elastic photon fluxes

To get the distribution of the elastic photons from the proton, one can express the equivalent photon flux through the electric and magnetic form factors $G_E(Q^2)$ and $G_M(Q^2)$ of the proton. This contribution is obtained as

$$\gamma_{el}^p(x, Q^2) = \frac{\alpha_{\text{em}}}{\pi} \left[\left(1 - \frac{x}{2}\right)^2 \frac{4m_p^2 G_E^2(Q^2) + Q^2 G_M^2(Q^2)}{4m_p^2 + Q^2} + \frac{x^2}{4} G_M^2(Q^2) \right] , \tag{1}$$

where x is the momentum fraction of the proton taken by the photon, Q^2 is the photon virtuality, α_{em} is the electromagnetic structure constant and m_p is the proton mass.

To express the elastic photon flux for the nucleus (γ_{el}^{Pb}), we follow Ref. [23] and replace

$$\frac{4m_p^2 G_E^2(Q^2) + Q^2 G_M^2(Q^2)}{4m_p^2 + Q^2} \longrightarrow Z^2 F_{\text{em}}^2(Q^2) , \tag{2}$$

where $F_{\text{em}}(Q^2)$ is the electromagnetic formfactor of the nucleus and Z is its charge. We also neglect the magnetic formfactor of the ion in the following.

For the Pb nucleus, we use the formfactor parameterization from the STARlight MC generator [24]:

$$F_{\text{em}}(Q^2) = \frac{3}{(QR_A)^3} \left[\sin(QR_A) - QR_A \cos(QR_A) \right] \frac{1}{1 + a^2 Q^2} , \tag{3}$$

where $R_A = 1.1 A^{1/3}$ fm, $a = 0.7$ fm and $Q = \sqrt{Q^2}$.

B. Collinear-factorization approach and choice of the scale

The inelastic processes, with breakup of a proton, can be also considered. At LO and at a given scale μ^2, the photon parton distribution $\gamma_{inel}^p(x, \mu^2)$ of photons carrying a fraction x of the proton's momentum, obeys the DGLAP equation:

$$\frac{d\gamma_{inel}^p(x, \mu^2)}{d\log \mu^2} = \frac{\alpha_{\rm em}}{2\pi} \int_x^1 \frac{dy}{y} \left[\sum_q P_{\gamma \leftarrow q}(y) q(\frac{x}{y}, \mu^2) + P_{\gamma \leftarrow \gamma}(y) \gamma_{inel}^p(\frac{x}{y}, \mu^2) \right] , \qquad (4)$$

where $q(x, \mu^2)$ is the quark PDF, e_q is the quark charge, $P_{\gamma \leftarrow q}$ is the $q \to \gamma$ splitting function, and $P_{\gamma \leftarrow \gamma}$ corresponds to the virtual self-energy correction to the photon propagator. This is the basis for collinear photon-PDF determination in all initial [25, 26] and more recent [14, 16, 17, 27–30] analyses.

In these studies the choice of the scale μ^2, which is needed to calculate the cross section for the process of interest, is often not well-motivated. The usual choice for μ is the mass of the system (motivated by the s-channel quark–antiquark annihilation process) or the transverse momentum of the leading object. These choices are however not optimal for the s- and t-channel initiated photon-induced process. By analogy to DIS (Fig. 1), where the scale is associated with the virtuality of the exchanged photon, it is possible to define the scale in case of the $\gamma\gamma \to \ell^+\ell^-$ process. This is achieved by taking the virtuality of the massive t- or u-channel propagator (Fig. 1b or c). Hence, $\mu^2 = -(p^{\gamma^{Pb}} - p^{\ell^-})^2$ for the t-channel diagram and $\mu^2 = -(p^{\gamma^{Pb}} - p^{\ell^+})^2$ for the u-channel exchange, repsectively, where $p^{\gamma^{Pb}}$ is the four momentum of the photon emitted by lead and $p^{\ell^\pm}$ is the four momentum of the lepton of the corresponding charge. Note that the u- and t- channel diagrams have vanishing interference in the zero lepton mass limit. Therefore, they can be separated while convoluting PDFs with the partonic cross section.

In the collinear approach, the $p + {\rm Pb} \to {\rm Pb} + \ell^+\ell^- + X$ production cross section can be written as

$$\sigma = S^2 \int dx_p dx_{\rm Pb} \left[\left(\gamma_{el}^p(x_p) + \gamma_{inel}^p(x_p, \mu^2) \right) \gamma_{el}^{\rm Pb}(x_{\rm Pb}) \sigma_{\gamma\gamma \to \ell^+\ell^-}(x_p, x_{\rm Pb}) \right] , \qquad (5)$$

where $\sigma_{\gamma\gamma \to \ell^+\ell^-}$ is the elementary cross section for the $\gamma\gamma \to \ell^+\ell^-$ subprocess and S^2 is the so-called survival factor which takes into account the requirement that there be no hadronic interactions between the proton and the ion.

C. k_T-factorization approach

At lowest order, the calculations with collinear photons produce leptons that are back-to-back in transverse kinematics. Therefore, to take the effect of transverse momentum smearing into account, a dedicated parton shower algorithms are usually used.

In the k_T factorization approach, one can parametrize the $\gamma^* p \to X$ vertices in terms of the proton structure functions. The photons from inelastic production have transverse momenta and non-zero virtualities Q^2 and the unintegrated photon distributions are used, in contrast to collinear distributions. In the DIS limit, the unintegrated inelastic photon flux can be obtained using the following equation [4, 31]:

$$\gamma_{inel}^p(x, \vec{q}_T) = \frac{1}{x} \frac{1}{\pi \vec{q}_T^2} \int_{M_{\text{thr}}^2} dM_X^2 \mathcal{F}_{\gamma^* \leftarrow p}^{\text{in}}(x, \vec{q}_T, M_X^2), \tag{6}$$

and we use the functions $\mathcal{F}_{\gamma^* \leftarrow p}^{\text{in}}$ from [13, 23]:

$$\begin{aligned}
\mathcal{F}_{\gamma^* \leftarrow p}^{\text{in}}(x, \vec{q}_T) &= \frac{\alpha_{\text{em}}}{\pi} \left\{ (1-x) \left(\frac{\vec{q}_T^2}{\vec{q}_T^2 + x(M_X^2 - m_p^2) + x^2 m_p^2} \right)^2 \frac{F_2(x_{\text{Bj}}, Q^2)}{Q^2 + M_X^2 - m_p^2} \right. \\
&\quad + \left. \frac{x^2}{4 x_{\text{Bj}}^2} \frac{\vec{q}_T^2}{\vec{q}_T^2 + x(M_X^2 - m_p^2) + x^2 m_p^2} \frac{2 x_{\text{Bj}} F_1(x_{\text{Bj}}, Q^2)}{Q^2 + M_X^2 - m_p^2} \right\}.
\end{aligned} \tag{7}$$

The virtuality Q^2 of the photon depends on the photon transverse momentum ($\vec{q}_T^2$) and the proton remnant mass (M_X):

$$Q^2 = \frac{\vec{q}_T^2 + x(M_X^2 - m_p^2) + x^2 m_p^2}{(1-x)}. \tag{8}$$

Moreover, the proton structure functions $F_1(x_{\text{Bj}}, Q^2)$ and $F_2(x_{\text{Bj}}, Q^2)$ require the argument

$$x_{\text{Bj}} = \frac{Q^2}{Q^2 + M_X^2 - m_p^2}. \tag{9}$$

Note that in Eq. 7 instead of using $F_2(x_{\text{Bj}}, Q^2), F_1(x_{\text{Bj}}, Q^2)$, we in practise use the pair $F_2(x_{\text{Bj}}, Q^2), F_L(x_{\text{Bj}}, Q^2)$, where

$$F_L(x_{\text{Bj}}, Q^2) = \left(1 + \frac{4 x_{\text{Bj}}^2 m_p^2}{Q^2} \right) F_2(x_{\text{Bj}}, Q^2) - 2 x_{\text{Bj}} F_1(x_{\text{Bj}}, Q^2) \tag{10}$$

is the longitudinal structure function of the proton.

These unintegrated photon fluxes enter the $p + \text{Pb} \to \text{Pb} + \ell^+ \ell^- + X$ production cross section as

$$\sigma = S^2 \int dx_p dx_{\mathrm{Pb}} d\vec{q}_T \Big[\left(\gamma_{el}^p(x_p, \vec{q}_T) + \gamma_{inel}^p(x_p, \vec{q}_T) \right) \gamma_{el}^{\mathrm{Pb}}(x_{\mathrm{Pb}}) \sigma_{\gamma^*\gamma \to \ell^+\ell^-}(x_p, x_{\mathrm{Pb}}, \vec{q}_T) \Big] \, , \quad (11)$$

where $\sigma_{\gamma^*\gamma \to \ell^+\ell^-}$ is the off-shell elementary cross-section [31] and for $x_p \ll 1$ we have $Q^2 \approx \vec{q}_T^2$ (see Eq. 8). One should note that while the fluxes do not depend on the direction of $\vec{q}_T$, averaging over directions of $\vec{q}_T$ in the off-shell cross section replaces the average over photon polarizations in the collinear case.

III.   EXAMPLE EXPERIMENTAL CONFIGURATION AND POSSIBLE BACK-GROUND SOURCES

We assume collision setup from recent $p + \mathrm{Pb}$ run at the LHC, carried out at the centre-of-mass energy per nucleon pair $\sqrt{s_{NN}} = 8.16$ TeV. Since the energy per nucleon in the proton beam is larger than in the lead beam, the nucleon–nucleon centre-of-mass system has a rapidity in the laboratory frame of $+0.465$.

As an example of method's applicability, we will use the geometry of ATLAS [32] and CMS [33] detectors in the following. We also consider dimuon-channel only, however the integrated results for ee and $\mu\mu$ channels can be obtained by simply multiplying the cross-sections by a factor of two.

We start with applying minimum transverse momentum requirement of 4 GeV to both muons. This requirement is imposed to ensure high lepton reconstruction and triggering efficiency. Moreover, due to limited acceptance of the detectors, each muon is required to have a pseudorapidity (η^ℓ) that satisfies $|\eta^\ell| < 2.4$ condition. Our calculations are carried out for a minimum dilepton invariant mass of $m_{\ell^+\ell^-} = 10$ GeV. Such a choice is due to removal of possible contamination from $\Upsilon(\to \ell^+\ell^-)$ photoproduction process. Summary of all selection requirements is presented in Table I.

Possible background for this process can arise from inclusive lepton-pair production, e.g. from Drell–Yan process [34–37]. This processes would lead to disintegration of the incoming ion, and zero-degree calorimeters (ZDC) [38, 39] can be used to veto very-forward-going neutral fragments which would allow to fully reduce this background. Another background can arise from diffractive interactions, hence possibly mimicking signal topology. However, since the Pb nucleus is a fragile object (with the nucleon binding energy of just 8 MeV) even

Variable	Requirement
lepton transverse momentum, p_T^ℓ	> 4 GeV
lepton pseudorapidity, $|\eta^\ell|$	< 2.4
dilepton invariant mass, $m_{\ell^+\ell^-}$	> 10 GeV

TABLE I: Definition of the fiducial region used in the studies.

the softest diffractive interaction will likely result in the emission of a few nucleons from the ion, detectable in the ZDC.

Another background category is the photon-induced process with resolved photon, i.e. $\gamma p \to Z/\gamma^* + X$ reaction. Here, the rapidity gap is expected to be smaller than in the signal process due to the additional particle production associated with the "photon remnant". Any other residual contamination of this process can be controlled using dedicated region, with a dilepton invariant mass around the Z-boson mass.

IV.   RESULTS WITH COLLINEAR PHOTON-PDFS

We start with the calculation of the elastic contribution, $p + \mathrm{Pb} \to p + \mathrm{Pb} + \ell^+\ell^-$. In this case the photon flux becomes:

$$\gamma_{el}^p(x, Q^2) = \gamma_{el}^p(x) \tag{12}$$

and the following parameterization is used [23]:

$$\gamma_{el}^p(x) = \frac{\alpha_{\mathrm{em}}}{\pi}\left(\frac{1 - x + 0.5x^2}{x}\right)\left(\frac{F+3}{F-1}\log F - \frac{17}{6} - \frac{4}{3F} + \frac{1}{6F^2}\right), \tag{13}$$

where $F = 1 + \frac{Q_0^2(1-x)}{x^2 m_p^2}$ and $Q_0^2 = 0.71$ GeV^2. This parameterization is a good analytical approximation of Eq. 1 integrated over Q^2. The results for the elastic case are cross-checked with the calculation from STARlight MC and a good agreement between the fiducial cross-sections is found: $\sigma_{fid}^{\mathrm{el}} = 17.5$ nb, whereas $\sigma_{fid}^{\mathrm{STARlight}} = 17.0$ nb. Both calculations are also corrected by a factor $S^2 = 0.96$ which is calculated using STARlight, where the hard-sphere proton–nucleus requirement [24] is used.

Next, for the inelastic case ($\gamma p \to \ell^+\ell^- + X$), several recent parameterizations of the photon parton distributions are studied: CT14qed [15], HKR16qed [29], LUXqed17 [40]

and NNPDF3.1luxQED [30]. All predictions are scaled by $S^2 = 0.95$, again derived from STARlight. This value of S^2 is lower than for the purely elastic case, due to slightly smaller average impact parameter between the proton and the ion in the inelastic reaction. One should note that all of these PDF sets include both elastic and inelastic parts of the photon spectrum. We keep the elastic part now (as provided by each group), but we subtract it later in Sec. VI for the comparison with k_T-factorization results.

The integrated fiducial cross-sections for $p + \text{Pb} \rightarrow \text{Pb} + \ell^+\ell^- + X$ production at $\sqrt{s_{NN}} = 8.16$ TeV for different collinear photon PDF sets are summarized in Tab. II. Comparison of several lepton kinematic distributions between different photon-PDFs is shown in Fig. 2, including invariant mass and rapidity of lepton pair, and single-lepton transverse momentum/pseudorapidity distributions. The asymmetry visible on pair rapidity and single-lepton pseudorapidity distributions is due to expermental setup, which assumes a difference in the energy per nucleon between the proton beam and the lead beam (see Sec. III). All photon PDF parameterizations agree within 10% with each other. The small differences are mainly due to overall PDF normalization, as no change in the shape of various kinematic distributions is observed.

Contribution	$p_T^\ell > 4$ GeV	$p_T^\ell > 4$ GeV, $|\eta^\ell| < 2.4$, $m_{\ell^+\ell^-} > 10$ GeV
γ_{el}^p	44.9(1) nb	17.5(1) nb
$\gamma_{\text{el}}^p + \gamma_{\text{inel}}^p$ [CT14qed_inc]	98.0(1) nb	39.7(1) nb
$\gamma_{\text{el}}^p + \gamma_{\text{inel}}^p$ [LUXqed17]	105.8(1) nb	44.1(1) nb
$\gamma_{\text{el}}^p + \gamma_{\text{inel}}^p$ [NNPDF3.1luxQED]	115.6(1) nb	45.9(1) nb
$\gamma_{\text{el}}^p + \gamma_{\text{inel}}^p$ [HKR16qed]	121.6(1) nb	49.4(1) nb

TABLE II: Integrated fiducial cross sections for $p + \text{Pb} \rightarrow \text{Pb} + \ell^+\ell^- + X$ production at $\sqrt{s_{NN}} = 8.16$ TeV for different collinear photon PDF sets. The effect of applying only p_T^ℓ requirement is shown in second column. For comparison, the cross section for purely elastic contribution is also shown.
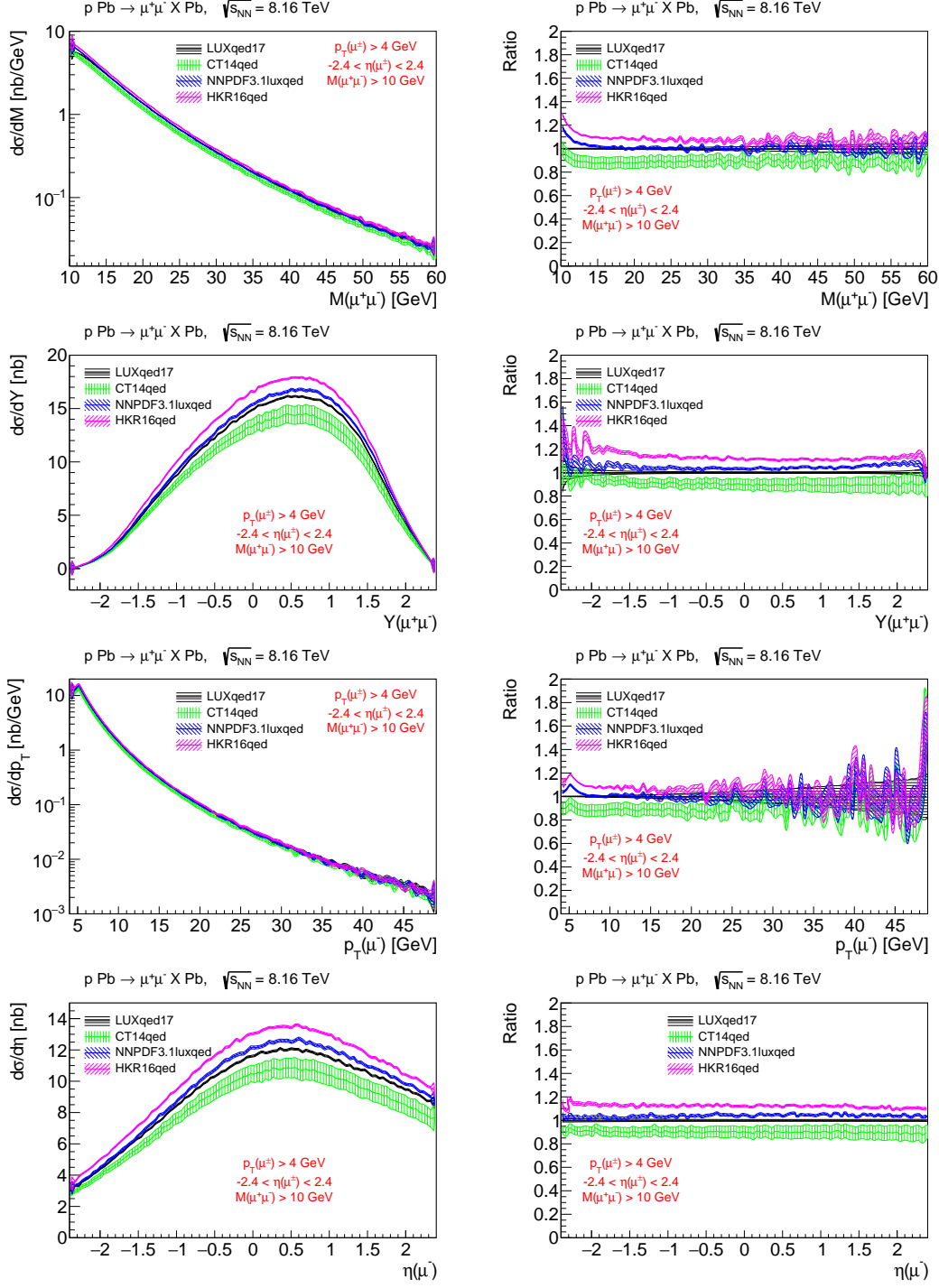
FIG. 2: Differential cross sections in the fiducial region for $p + \text{Pb} \rightarrow \text{Pb} + \ell^+\ell^- + X$ production at $\sqrt{s_{NN}} = 8.16$ TeV for different collinear photon PDF sets. Four differential distributions are shown (from top to bottom): invariant mass of lepton pair, pair rapidity, transverse momentum of negatively-charged lepton and its pseudorapidity. Figures on the right show the ratios to LUXqed17 PDF. The bands denote the PDF uncertainties calculated at 68% CL.

V. RESULTS USING k_T-FACTORIZATION APPROACH

Several different parametrizations of proton strucure functions are used. Those are labeled as:

- ALLM [41, 42]: This parametrization gives a good fit to F_2 in most of the measured regions.

- SY [43]: This parameterization of Suri and Yennie from the early 1970's does not include DGLAP evolution. It is still used as one of the defaults in the LPAIR event generator [44].

- SU [45]: A parametrization which concentrates to give a good description at smal and intermediate Q^2 at not too small x. At large Q^2, it is complemented by the NNLO calculation of F_2 and F_L from NNLO MSTW 2008 PDF analysis [46].

- LUX-like: a recently constructed parametrization, described in details in Ref. [13]. This setup closely follows the LUXqed work from Ref.[40].

To model $\gamma_{el}^p(x, Q^2)$ we use Eq. 1 with so-called dipole parametrization of the proton formfactors:

$$G_E(Q^2) = \left(1 + \frac{Q^2}{Q_0^2}\right)^{-2} \tag{14}$$

$$G_M(Q^2) = \mu_p G_E(Q^2) , \tag{15}$$

where μ_p is the proton magnetic moment.

Table III shows the comparison of integrated fiducial cross sections for inelastic $p + \mathrm{Pb} \rightarrow \mathrm{Pb} + \ell^+\ell^- + X$ production at $\sqrt{s_{NN}} = 8.16$ TeV for different proton structure functions used. All structure functions provide similar fiducial cross-section, at the level of 16–18 nb. These inelastic cross-sections are also similar in size to the elastic contribution (18 nb) and are slightly lower than the numbers from collinear analysis, subtracted for elastic part (see Table II). A comparison is also made with LUX-like parametrization when the longitudinal structure function (F_L) is explicitly considered. This leads to the decrease of the cross section by 2%, similarly as in Ref. [13].

Figure 3 presents differential cross sections for several lepton kinematic distributions: invariant mass of lepton pair, leading lepton transverse momentum, lepton pseudorapidity difference and leading lepton pseudorapidity. The shapes of the distributions obtained

Contribution	$p_T^\ell > 4$ GeV	$p_T^\ell > 4$ GeV, $|\eta^\ell| < 2.4$, $m_{\ell^+\ell^-} > 10$ GeV
γ_{el}^p	47.9 nb	18.3 nb
γ_{inel}^p [LUX-like F_2]	43.6 nb	17.4 nb
γ_{inel}^p [LUX-like $F_2 + F_L$]	42.6 nb	17.1 nb
γ_{inel}^p [ALLM97 F_2]	41.7 nb	16.4 nb
γ_{inel}^p [SU F_2]	41.7 nb	16.7 nb
γ_{inel}^p [SY F_2]	40.4 nb	16.0 nb

TABLE III: Integrated fiducial cross sections for inelastic $p + \mathrm{Pb} \to \mathrm{Pb} + \ell^+\ell^- + X$ production at $\sqrt{s_{NN}} = 8.16$ TeV for different proton structure functions. The effect of applying only p_T^ℓ requirement is shown in second column.

with various proton structure functions are very similar. For completeness, differential cross sections as a function of lepton pair transverse momentum and azimuthal angle difference between the pair are shown in Fig. 4. Quite large (small) transverse momenta (angle differences) are possible, in contrast to leading-order calculations with collinear photons where the corresponding distributions are just a Dirac delta functions. The k_T-factorization approach should be considered more appropriate here. It is also visible that the SY parametrization gives lower predictions at larger pair-p_T, comparing to the other parametrizations used. This is because SY parametrization does not include explicit DGLAP evolution terms, which are relevant for large photon virtualities.

Based on Fig. 4, it is also possible to separate experimentally the elastic part ($p + \mathrm{Pb} \to p + \mathrm{Pb} + \ell^+\ell^-$) with striking back-to-back topology, out of the inelastic contribution. With k_T-factorization, one can also calculate the mass of the proton remnants (M_X). This is shown in Fig. 5, where in contrary to the elastic case ($M_X = m_p$) quite large masses of the remnant system can be achieved.

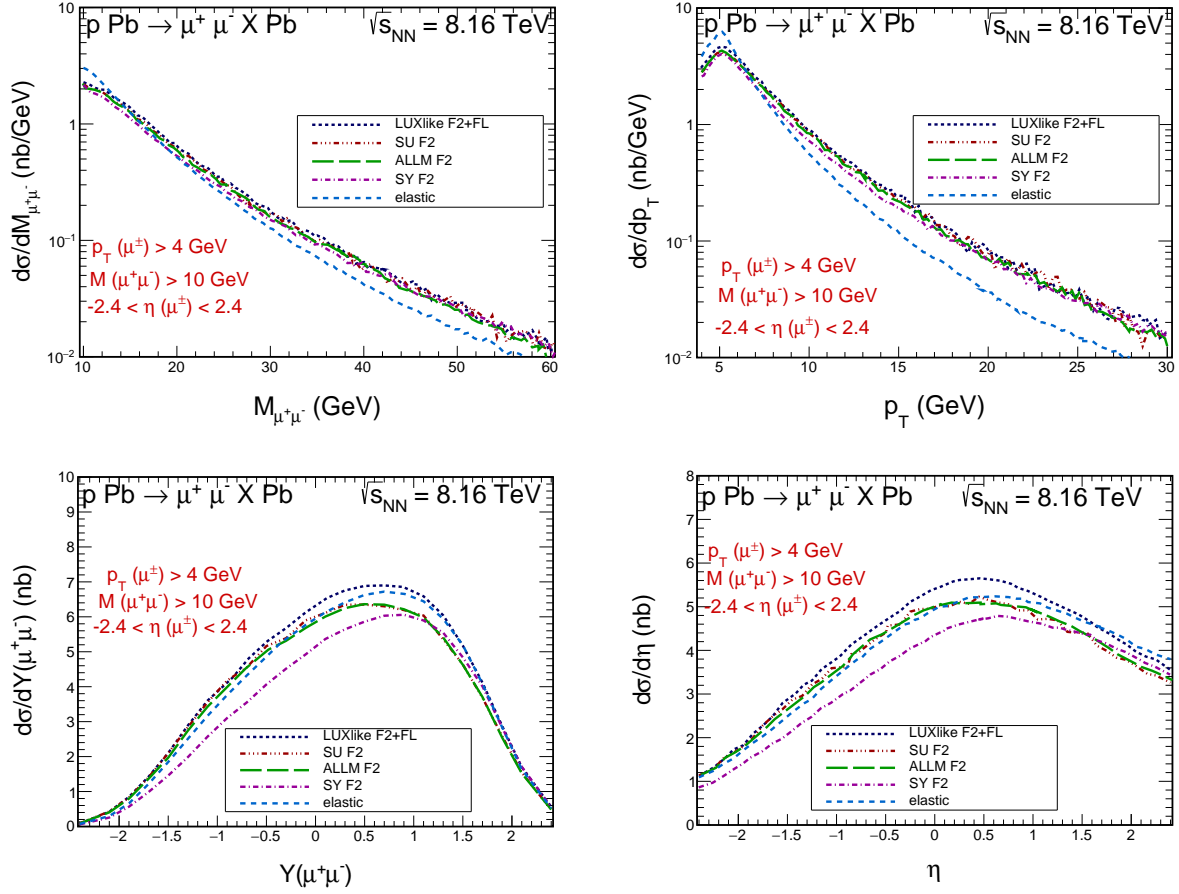FIG. 3: Differential cross sections in the fiducial region for $p + \mathrm{Pb} \to \mathrm{Pb} + \ell^+\ell^- + X$ production at $\sqrt{s_{NN}} = 8.16$ TeV in k_T factorization approach for several proton structure functions. Four differential distributions are shown: invariant mass of lepton pair (top left), leading lepton transverse momentum (top right), dilepton rapidity (bottom left) and leading lepton pseudorapidity (bottom right). For comparison, the elastic contribution ($p + \mathrm{Pb} \to p + \mathrm{Pb} + \ell^+\ell^-$) is also shown.
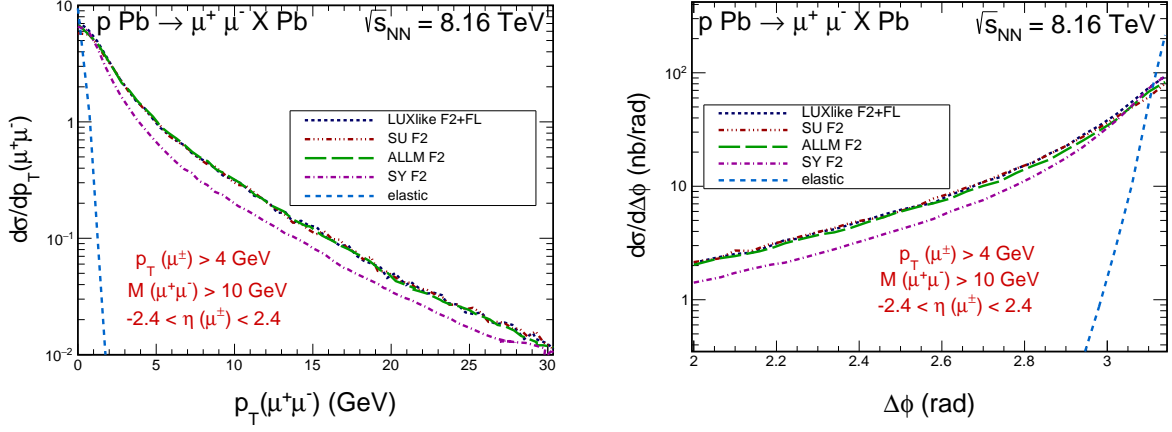
FIG. 4: Differential cross sections in the fiducial region for $p + \mathrm{Pb} \to \mathrm{Pb} + \ell^+\ell^- + X$ production at $\sqrt{s_{NN}} = 8.16$ TeV in k_T factorization approach for several proton structure functions. Two differential distributions are shown: transverse momentum of lepton pair (left) and azimuthal angle difference between the pair (right). For comparison, the elastic contribution ($p + \mathrm{Pb} \to p + \mathrm{Pb} + \ell^+\ell^-$) is also shown.
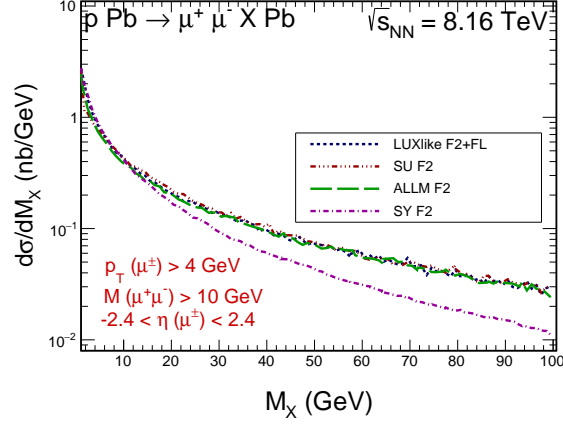

FIG. 5: Differential cross section as a function of the mass of the proton remnants in the fiducial region for $p + \mathrm{Pb} \to \mathrm{Pb} + \ell^+\ell^- + X$ production at $\sqrt{s_{NN}} = 8.16$ TeV in k_T factorization approach for several proton structure functions.

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VI. DISCUSSION

Figure 6 compares several differential distribution computed using the two approaches. For the collinear approach pure inelastic contribution is estimated by subtracting elastic part computed following Eq. 13. For the invariant mass distribution and lepton pseudorapidity the shapes are similar and the main difference between the two predictions is observed in the normalization. For the distribution of the lepton pair rapidity the two predictions agree at larger rapidities while disagreement concentrates in the central region. The biggest difference is observed for the transverse momentum distribution of the lepton where at low p_T collinear approximation exceeds the estimate from k_T-factorization approach while at high p_T the ordering is reversed. This suggests that at low p_T (close to the boundary of the fiducial region) the difference is due to the smearing of dilepton transverse momentum introduced by the k_T-factorization approach.

We also take the opportunity to calculate expected number of events for realistic assumption on total integrated luminosity. Based on the previous $p+$Pb runs at the LHC, we assume $\int L dt = 200$ nb^{-1}. We also assume possible experimental efficiencies, mainly due to trigger and reconstruction of leptons, which we embed in a single correction factor $C = 0.7$.

Table IV shows the expected number of events for $p+$Pb $\to$ Pb$+\ell^+\ell^-+X$ production at $\sqrt{s_{NN}} = 8.16$ TeV and configuration described above. Approximately 2500 elastic dilepton events are expected. Depending on the calculations, 3400 (collinear with LUXqed17 PDF) or 2400 (k_T-factorization with LUX-like $F_2 + F_L$) reconstructed inelastic events are predicted. The difference between collinear and k_T-factorization can be therefore constrained with large significance, using existing datasets collected by ATLAS and CMS.
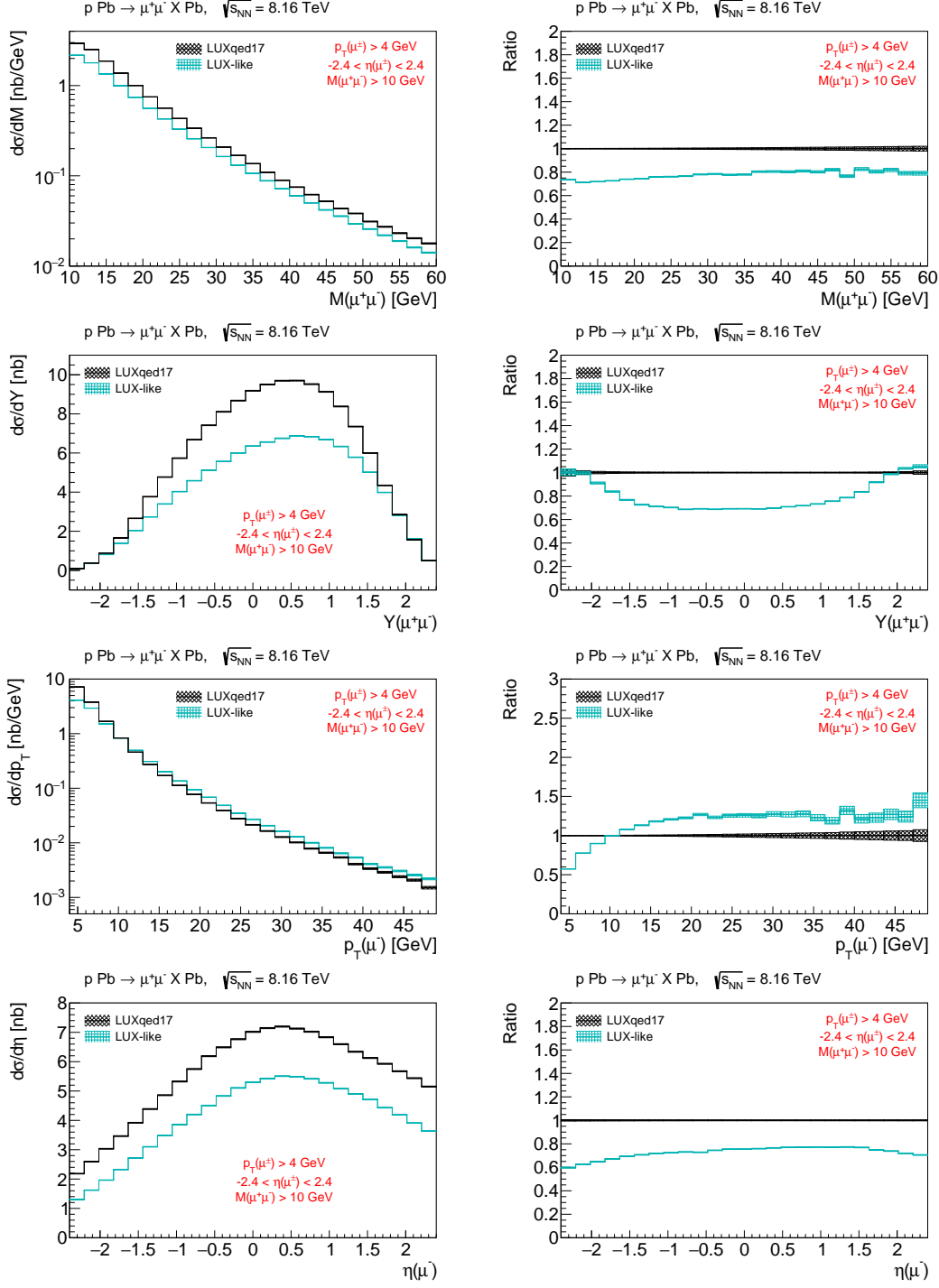
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FIG. 6: Differential cross sections in the fiducial region for $p + \mathrm{Pb} \rightarrow \mathrm{Pb} + \ell^+\ell^- + X$ production at $\sqrt{s_{NN}} = 8.16$ TeV for collinear LUXqed17 photon PDF and for LUX-like $F_2 + F_L$ photon PDF with k_T-factorization. Four differential distributions are shown (from top to bottom): invariant mass of lepton pair, pair rapidity, transverse momentum of negatively-charged lepton and its pseudorapidity. Figures on the right show the ratios to LUXqed17 PDF.

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Contribution	Expected events $(C = 1)$	Expected events $(C = 0.7)$
γ_{el}^{p}	3600	2500
γ_{inel}^{p} [LUXqed17 collinear]	5600	3900
γ_{inel}^{p} [LUX-like $F_2 + F_L$]	3400	2400

TABLE IV: Expected number of events for $p+\text{Pb} \to \text{Pb}+\ell^+\ell^-+X$ production at $\sqrt{s_{NN}} = 8.16$ TeV assuming $\int Ldt = 200$ nb^{-1}. Shown are several contributions: purely elastic, inelastic with collinear LUXqed17 PDF and inelastic with k_T-factorization and LUX-like F_2+F_L proton structure function parameterization. An effect of possible experimental efficiencies is shown in last column.

VII. SUMMARY

In summary, we propose a method that would unambiguously allow to test and constrain the photon parton distribution at LHC energies. This method is based on the measurement of the cross-section for the reaction $p+\text{Pb} \to \text{Pb}+\ell^+\ell^-+X$, where the expected background is small comparing to the analogous process in pp collisions. Results are shown for different choices of collinear photon PDFs, and a comparison is made with unintegrated photon distributions that include non-zero photon transverse momentum. Due to the smearing of dilepton transverse momentum introduced by the k_T-factorization approach, these two approaches lead to the cross sections that differ by about 30%. Moreover, for collinear approach and by analogy to DIS, the correct choice of the scale is identified. Using simple (realistic) experimental requirements on lepton kinematics, it is shown that one can expect O(3000) inelastic events with the existing datasets recorded by ATLAS/CMS at $\sqrt{s_{NN}} = 8.16$ TeV for each lepton flavour.

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