

# A Tutorial on Hybrid Finite Element Formulation

## Assumed Stress-Displacement Field

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# Outline

- 1 Conventional FEM
- 2 Hybrid Formulation
- 3 Getting the matrix equations
- 4 The Stress Field Matrices
- 5 Element Stiffness Matrix - Pseudo Code

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# Conventional Formulation for FEM

$$\int_{\Omega} \mathbf{v} \cdot (\nabla \cdot \boldsymbol{\tau} + \mathbf{b}) d\Omega + \int_{\Gamma_t} \mathbf{v} \cdot (\mathbf{t} - \bar{\mathbf{t}}) d\Gamma = 0 \quad \forall \mathbf{v} \in V_u$$

- $\Omega$  - domain
- $\Gamma$  - boundary.
- $\Gamma_t$  - traction specified boundary.
- $\mathbf{v}$  - variation of the displacement.
- $\mathbf{t}$  - tractions on the boundary
- $\bar{\mathbf{t}}$  - specified tractions
- $\boldsymbol{\tau}$  - stress tensor
  - $\boldsymbol{\tau} = \mathbf{C} : \boldsymbol{\varepsilon}(\mathbf{u})$
  - $\mathbf{C}$  fourth order constitutive tensor.
  - $\boldsymbol{\varepsilon}(\mathbf{u})$  linear strain tensor
    - $\frac{1}{2} (\nabla \mathbf{u} + \nabla \mathbf{u}^T)$
  - $\mathbf{u}$  Displacements

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- Find  $\mathbf{u}$  for all possible  $\mathbf{v}$
- Use the identity
  - $\mathbf{v} \cdot (\nabla \cdot \boldsymbol{\tau}) = \nabla \cdot (\boldsymbol{\tau} \mathbf{v}) - \boldsymbol{\tau} : \frac{1}{2} (\nabla \mathbf{v} + \nabla \mathbf{v}^T)$ 
    - $\frac{1}{2} (\nabla \mathbf{v} + \nabla \mathbf{v}^T)$  denoted by  $\boldsymbol{\varepsilon}(\mathbf{v})$
  - Used symmetry of  $\boldsymbol{\tau}$
- Use Divergence theorem
  - cancel out the term  $\int_{\Gamma_t} \mathbf{v} \cdot \mathbf{t} d\Gamma$
- This results in

$$\int_{\Omega} \mathbf{v} \cdot \mathbf{b} d\Omega + \int_{\Gamma_t} \mathbf{v} \cdot \bar{\mathbf{t}} d\Gamma = \int_{\Omega} \boldsymbol{\varepsilon}(\mathbf{v}) : \boldsymbol{\tau} d\Omega \quad \forall \mathbf{v} \in V_u$$

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# Conventional Formulation for FEM

$$\int_{\Omega} \mathbf{v} \cdot \mathbf{b} d\Omega + \int_{\Gamma_t} \mathbf{v} \cdot \bar{\mathbf{t}} d\Gamma = \int_{\Omega} \boldsymbol{\varepsilon}(\mathbf{v}) : \boldsymbol{\tau} d\Omega \quad \forall \mathbf{v} \in V_u$$

- $\boldsymbol{\tau}$  is replaced by  $\mathbf{C} : \boldsymbol{\varepsilon}(\mathbf{u})$
- $\mathbf{v}$  and  $\mathbf{u}$  are interpolated

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# Hybrid Formulation for FEM

$$\int_{\Omega} \mathbf{v} \cdot (\nabla \cdot \boldsymbol{\tau} + \mathbf{b}) d\Omega + \int_{\Gamma_t} \mathbf{v} \cdot (\mathbf{t} - \bar{\mathbf{t}}) d\Gamma + \int_{\Omega} \boldsymbol{\sigma} : (\boldsymbol{\varepsilon}(\mathbf{u}) - \mathbf{C}^{-1} : \boldsymbol{\tau}) d\Omega = 0$$

$$\forall (\mathbf{v}, \boldsymbol{\sigma}) \in (V_u \times V_{\tau})$$

- $\boldsymbol{\sigma}$  Variation of the stress field.
- Results in two sets of equations
  - set variation in each to zero

$$\int_{\Omega} \boldsymbol{\varepsilon}(\mathbf{v}) : \boldsymbol{\tau} d\Omega = \int_{\Omega} \mathbf{v} \cdot \mathbf{b} d\Omega + \int_{\Gamma_t} \mathbf{v} \cdot \bar{\mathbf{t}} d\Gamma$$

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$$\int_{\Omega} \boldsymbol{\sigma} : (\varepsilon(\mathbf{u}) - \mathbf{C}^{-1} : \boldsymbol{\tau}) d\Omega = 0$$

- $\mathbf{v}$  and  $\mathbf{u}$  interpolated
- $\boldsymbol{\sigma}$  and  $\boldsymbol{\tau}$  interpolated

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## Getting the matrix equations

- $\mathbf{u} = N\hat{\mathbf{u}}$  and  $\mathbf{v} = N\hat{\mathbf{v}}$ 
  - $\hat{\mathbf{u}}$  and  $\hat{\mathbf{v}}$  are nodal displacements and its variations.
  - viewed as constants helping to describe the displacement field within.
- $\boldsymbol{\tau} = P\hat{\boldsymbol{\beta}}$  and  $\boldsymbol{\sigma} = P\hat{\boldsymbol{\gamma}}$ 
  - $\hat{\boldsymbol{\beta}}$  and  $\hat{\boldsymbol{\gamma}}$  are constants in an element that help in description of the stress field within.
- This results in  $\boldsymbol{\varepsilon}(\mathbf{u}) = B\hat{\mathbf{u}}$  and  $\boldsymbol{\varepsilon}(\mathbf{v}) = B\hat{\mathbf{v}}$ 
  - $B$  is the strain-displacement matrix

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# Getting the matrix equations

$$\int_{\Omega} \varepsilon(\mathbf{v}) : \boldsymbol{\tau} d\Omega = \int_{\Omega} \mathbf{v} \cdot \mathbf{b} d\Omega + \int_{\Gamma_t} \mathbf{v} \cdot \bar{\mathbf{t}} d\Gamma$$

$$\int_{\Omega} \boldsymbol{\sigma} : (\boldsymbol{\varepsilon}(\mathbf{u}) - \mathbf{C}^{-1} : \boldsymbol{\tau}) d\Omega = 0$$

# Getting the matrix equations

$$\boldsymbol{\Phi}^T \left( \int_{\Omega} B^T P d\Omega \right) \hat{\boldsymbol{\beta}} = \boldsymbol{\Phi}^T \left( \int_{\Omega} N^T \mathbf{b} d\Omega \right) + \boldsymbol{\Phi}^T \left( \int_{\Gamma_t} N^t \bar{\mathbf{t}} d\Gamma \right)$$

$$\hat{\boldsymbol{\gamma}}^T \int_{\Omega} P^T \left( B \hat{\mathbf{u}} - \mathbf{C}^{-1} P \hat{\boldsymbol{\beta}} \right) d\Omega = 0$$

## Getting the matrix equations

$$\mathbf{\hat{v}}^T \left( \int_{\Omega} B^T P d\Omega \right) \mathbf{\hat{\beta}} = \mathbf{\hat{v}}^T \left( \int_{\Omega} N^T \mathbf{b} d\Omega \right) + \mathbf{\hat{v}}^T \left( \int_{\Gamma_t} N^t \bar{\mathbf{t}} d\Gamma \right)$$

$$\hat{\gamma}^T \left( \int_{\Omega} P^T B d\Omega \right) \mathbf{\hat{u}} - \hat{\gamma}^T \left( \int_{\Omega} P^T \mathbf{C}^{-1} P d\Omega \right) \mathbf{\hat{\beta}} = 0$$

## Getting the matrix equations

- Eliminating the Variations we get

$$\overbrace{\left( \int_{\Omega} B^T P d\Omega \right)}^{G^T} \hat{\beta} = \overbrace{\left( \int_{\Omega} N^T \mathbf{b} d\Omega \right) + \left( \int_{\Gamma_t} N^t \bar{\mathbf{t}} d\Gamma \right)}^{\hat{\mathbf{f}}}$$

$$\overbrace{\left( \int_{\Omega} P^T B d\Omega \right)}^G \hat{\mathbf{u}} - \overbrace{\left( \int_{\Omega} P^T \mathbf{C}^{-1} P d\Omega \right)}^H \hat{\beta} = 0$$

# Getting the matrix equations

- The Matrix equations

$$G^T \hat{\beta} = \hat{f}$$

$$G\alpha - H\hat{\beta} = 0$$

## Getting the matrix equations

- Condensing out the  $\hat{\beta}$

$$\left(G^T H^{-1} G\right) \hat{\mathbf{u}} = \hat{\mathbf{f}}$$

- Can be split across elements

$$\sum_e \left(G_e^T H_e^{-1} G_e\right) \hat{\mathbf{u}}_e = \hat{\mathbf{f}}_e$$

- So every element stiffness matrix formulation needs a matrix inversion.

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## Obtaining the ' $P$ ' Matrix

- Governed by the stress shape functions.
- Defined in local co-ordinate system of the element.
  - Transformed to global coordinate system
  - At every point inside the element
  - $\tau = J^T \tau_{loc} J$ 
    - $J$  Jacobian matrix
    - Linear in  $\tau_{loc}$
    - For vectorized  $\tau$  one can get a Transformation matrix  $T$  size  $6 \times 6$  in 3D
- $P_{loc}$  These are of size  $6 \times num \beta s$  in 3D
- $P = T \times P_{loc}$

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## 'P' Matrix an example

- In local coords for a 4 node quad

$$\bullet \begin{bmatrix} \tau_{xx} \\ \tau_{yy} \\ \tau_{xy} \end{bmatrix} = \begin{bmatrix} \beta_1 + \beta_2 y \\ \beta_3 + \beta_4 x \\ \beta_5 \end{bmatrix}$$

$$\bullet \begin{bmatrix} \tau_{xx} \\ \tau_{yy} \\ \tau_{xy} \end{bmatrix} = \underbrace{\begin{bmatrix} 1 & y & 0 & 0 & 0 \\ 0 & 0 & 1 & x & 0 \\ 0 & 0 & 0 & 0 & 1 \end{bmatrix}}_{P_{loc}} \begin{bmatrix} \beta_1 \\ \beta_2 \\ \beta_3 \\ \beta_4 \\ \beta_5 \end{bmatrix}$$

- Transformation matrix

$$\bullet \begin{bmatrix} J_{11}^2 & J_{21}^2 & 2J_{11}J_{21} \\ J_{12}^2 & J_{22}^2 & 2J_{12}J_{22} \\ J_{11}J_{12} & J_{21}J_{22} & J_{11}J_{22} + J_{12}J_{21} \end{bmatrix}$$

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## Stress Shape Functions the rule of the thumb

- $\tau_{xx}$ - differentiation of the displacement field  $u_x$
- $\tau_{yy}$ - differentiation of the displacement field  $u_y$
- $\tau_{zz}$ - differentiation of the displacement field  $u_z$
- $\tau_{xy}$ - intersection of  $\tau_{xx}$  and  $\tau_{yy}$
- $\tau_{xz}$ - intersection of  $\tau_{xx}$  and  $\tau_{zz}$
- $\tau_{yz}$ - intersection of  $\tau_{yy}$  and  $\tau_{zz}$



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## Pseudo - Code

For every gauss point:

displacement shape functions (disp)

displacement derivatives (derdisp)

$P_{local}$

Jacobian  $J$

$J^{-1}$

$Det(J)$

$T$  - using  $J$

$B$  - using  $J^{-1}derdisp$

$P = T \times P_{local}$

$G = G + P^T \times B \times Det(J)$

$H = H + P^T \times C \times P \times Det(J)$

end for

$KE = G^T H^{-1} G$

Thank You

# For Further Reading I



C.S. Jog. A 27-node hybrid brick and a 21-node hybrid wedge element for structural analysis, 2007