

The "missing step" is as follows

$$V_a(t) = \iint E(l, m, t) dl dm$$

but we will express  $E(l, m, t)$  as a function of its inverse Fourier transform

$$E(l, m, t) = \int E(l, m, \nu) e^{i2\pi\nu t} d\nu$$

$$\rightarrow V_a(t) = \iiint E(l, m, \nu) e^{i2\pi\nu t} dl dm d\nu$$

∴ we'll further write

$$E(l, m, \nu) = A(l, m, \nu) e^{i\phi(l, m, \nu)}$$

Now

$$V_a(t) V_b^*(t) = \iiint \iiint A(l, m, \nu) A^*(l', m', \nu') e^{i\phi(l, m, \nu)} e^{-i\phi(l', m', \nu')} \\ \cdot e^{i2\pi(x_a l + y_a m) - i2\pi(x_b l' + y_b m)} e^{i2\pi\nu t} e^{-i2\pi\nu' t} dl dl' dm dm' d\nu d\nu'$$

But if there is no spatial coherence the phase term  $\phi(l, m, \nu)$  varies very rapidly (i.e., more than  $A(l, m, \nu)$ )

∴ we write

$$\int e^{i\phi(l,m,\nu)} e^{-i\phi(l',m',\nu')} dl' dm' d\nu' = \delta(l-l', m-m') \delta(\nu-\nu')$$

Then

$$V_a(t) V_b^*(\lambda) = \iiint |A(l,m,\nu)|^2 e^{i2\pi\nu(t-\lambda)} e^{i(\Delta x l + \Delta y m)} dl dm d\nu$$

If we set  $\lambda = t - \tau$ , with  $\tau$  the time lag

$$V_a(t) V_b^*(t-\tau) = \iiint |A(l,m,\nu)|^2 e^{i2\pi\nu\tau} e^{i(\Delta x l + \Delta y m)} dl dm d\nu$$

Therefore

$$\begin{aligned} \langle V_a(t) V_b(t-\tau) \rangle &= \lim_{T \rightarrow \infty} \frac{1}{2T} \iiint_{-T}^T |A(l,m,\nu)|^2 e^{i2\pi\nu\tau} e^{i(\Delta x l + \Delta y m)} dl dm d\nu dt \\ &= \iiint |A(l,m,\nu)|^2 e^{i2\pi\nu\tau} e^{i(\Delta x l + \Delta y m)} dl dm d\nu \\ &= \iiint I(l,m,\nu) e^{i2\pi\nu\tau} e^{i(\Delta x l + \Delta y m)} dl dm d\nu \end{aligned}$$

So, we can say that at a given frequency (which is how observations are made)

$$\langle V_a(t) V_b(t-\tau) \rangle_{\nu} = F[I(l,m,\nu)]$$



this subscript means "frequency of observations".