

# Quantum Materials

## Introduction to Quantum Mechanics

- Historical evolution from classical physics to quantum physics.
- Exploration of wave-particle duality and the concept of quantum states. Superposition and entanglement: their significance in quantum systems and potential applications.
- The uncertainty principle and its impact on measurement in quantum systems.
- Quantum wave function, Schrödinger equation (time independent). Energy quantization: Introduction to discrete energy levels in atoms and their effects on material properties.

## Dual nature of matter

de Broglie wavelength,  $\lambda = h/p$

For a cricket ball

mass = 160 g & velocity = 150 km/h

$\lambda = 0.98 \times 10^{-34} \text{ m} \rightarrow$  so small, not appreciable

For an electron

accelerated by applied potential of 1 V

$$\lambda = 1.23 \text{ nm} \left( \lambda = \frac{h}{p} = \frac{h}{\sqrt{2meV}} \right)$$

## Heisenberg Uncertainty Principle

It is impossible to simultaneously describe with absolute accuracy the position and momentum of a particle.

$$\Delta x \cdot \Delta p_x \geq \hbar/2$$

Predictions from uncertainty principle:

- Nonexistence of free electron in nucleus
- Estimate of radius of Bohr's first orbit
- Zero point energy of simple harmonic oscillators

## How do we describe a system in Quantum Mechanics?

- Described by a wave function:  $\psi(r,t)$
- All information about the system is contained in this  $\psi(r,t)$
- Any variation in property/behaviour of the system will be reflected in  $\psi(r,t)$ .
- Solving the problem in quantum mechanics is **to determine  $\psi(r,t)$**  considering all parameters.

## Requirements for $\psi$

- $\psi$  can be any function, real or complex, akin to the behavior of the real physical system.
- While  $\psi$  itself has no physical interpretation, its square  $\psi^2$  evaluated at a particular place at any time gives the **probability of finding** the particle there at that time.

## Mathematical requirements

- $\psi$  and  $\frac{d\psi}{dr}$  must be finite, continuous and single valued.
- $\psi \rightarrow 0$  as  $r \rightarrow \infty$ .
- $\int_{-\infty}^{+\infty} \psi^2 dV = 1$

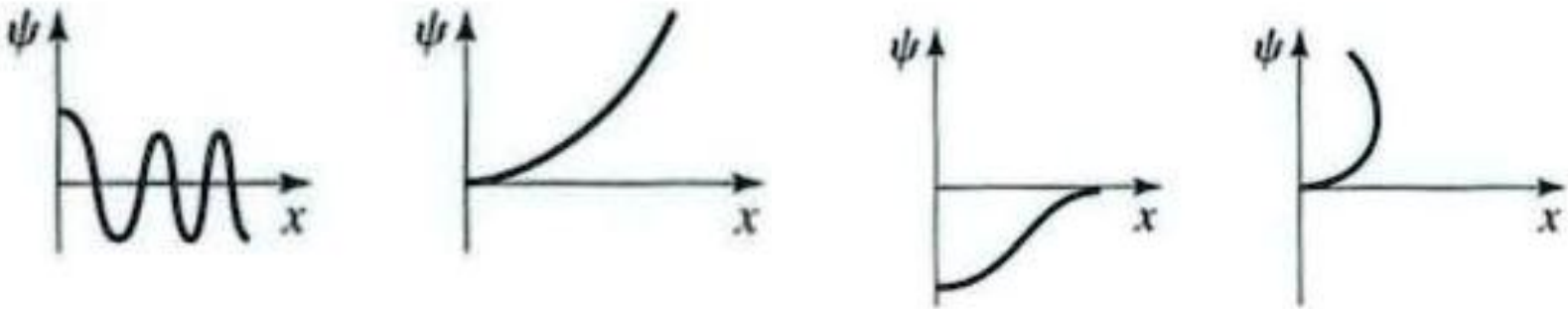
Quantum physics  
is probabilistic in  
nature.

# Wave function $\psi$

## Mathematical requirements

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Which is the acceptable wave form?



Is  $\psi$  a well behaved wavefunction if

- (a)  $\psi = A \sin x$ ,      (b)  $\psi = B \tan x$ ,      (c)  $\psi = C e^{kx}$ ,      (d)  $\psi = D e^{ikx}$ ,  
(e)  $\psi = A \sin x + D e^{ikx}$

## Normalization

The probability that its position  $x$  will be in the interval  $a \leq x \leq b$  is the **integral of the density** over this interval:

$$P(t) = \int_a^b |\psi(x, t)|^2 dx \quad \text{where } t \text{ is the time at which the particle was measured.}$$

This leads to the **normalization condition**:

$$\int_{-\infty}^{+\infty} |\psi(x, t)|^2 dx = 1 \quad \text{because if the particle is measured, there is 100% probability that it will be somewhere..}$$

A particle limited to the  $x$  axis has the wave function

$$\psi = Ax \text{ for } 0 \leq x \leq 1;$$

$$\psi = 0 \text{ elsewhere.}$$

- (a) Normalize the wave function (find the normalization constant  $A$ ).
- (b) Probability to find the particle between  $x = 0.45$  and  $x = 0.55$ .
- (c) Expectation value  $x$  of the particle's position.

- Applying certain operations/instruction on wave function  $\psi$ , information/ observable can be extracted.
- These information specific operations are called **operator** for a given observable.

For example momentum operator,  $\hat{p}_x \equiv \left(-i\hbar \frac{\partial}{\partial x}\right)$ ,

Energy operator  $\hat{E} \equiv \left(i\hbar \frac{\partial}{\partial t}\right)$

## How to use an operator?

$$\hat{A} \psi(r, t) = \lambda \psi(r, t)$$

Operator  $\hat{A}$  operates on the wave function  $\psi$  to give the Eigen value  $\lambda$ .

## Expectation value (classically, average value)

$$\int_{-\infty}^{+\infty} \psi^* \hat{A} \psi dV = \langle \hat{A} \rangle$$

# Extraction of information



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- Is the function  $e^{3x+5}$  an Eigen-function of the operator  $\frac{d^2}{dx^2}$ ? If so, what is the corresponding Eigen-value?
- Obtain expressions for the following operator  $\left(\frac{d}{dx} + x\right)^2$  and  $\left(x \frac{d}{dx}\right)^2$

1. State of a physical system: Wave function
2. Operator corresponding to a classical observable
3. Measurement of physical quantity: eigen value equation
4. Probabilistic measurement: expectation value
5. Evolution of wave function (The Schrodinger equation)

## Schrodinger equation

$$\left( -\frac{\hbar^2}{2m} \frac{\partial^2}{\partial x^2} + V \right) \psi(x) = E\psi(x)$$
$$\Rightarrow \frac{\partial^2 \psi(x)}{\partial x^2} + \frac{2m(E - V)}{\hbar^2} \psi(x) = 0$$

Whole of quantum mechanics is based on finding solution to this equation.



# Schrodinger equation

$$\frac{\partial^2 \psi(x)}{\partial x^2} + \frac{2m(E - V)}{\hbar^2} \psi(x) = 0$$

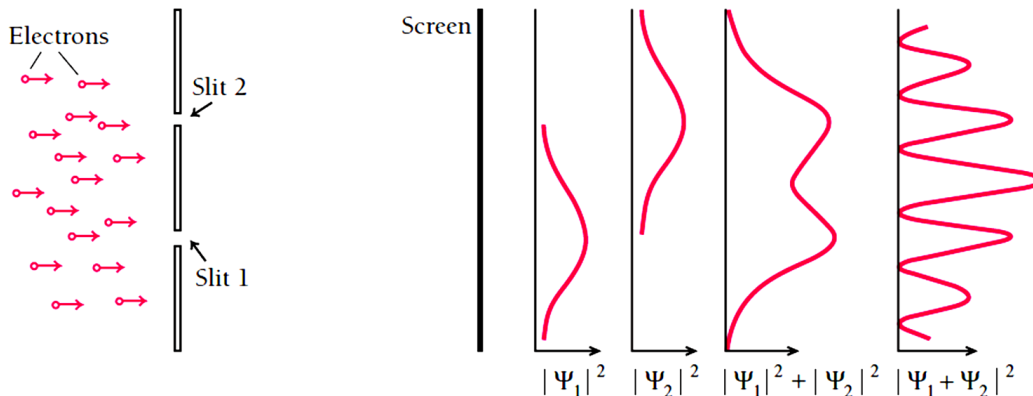
This equation may have multiple solutions!

**Which solution is of our use?**

## Superposition

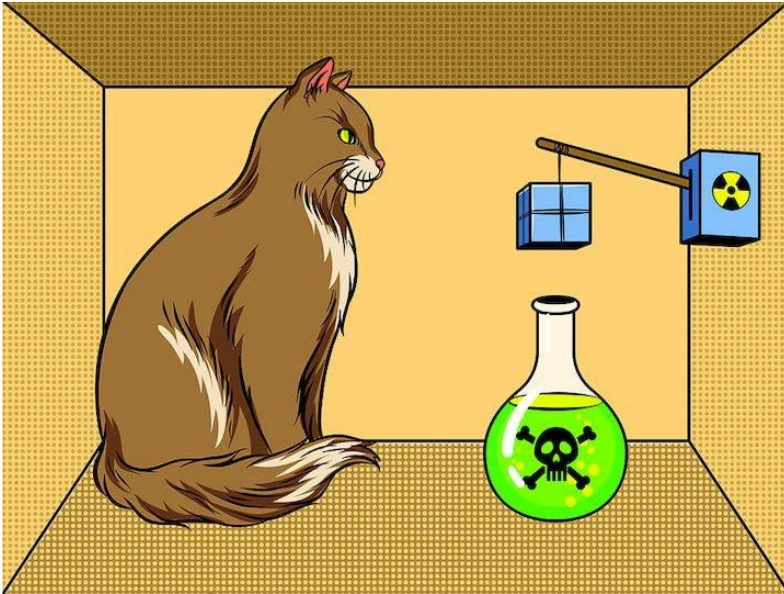
- It is the characteristic property of the waves.
- A linear combination of solutions of Schrödinger's equation is also a solution!

$$\psi = a\psi_1 + b\psi_2$$



It is wave function, not the probability, which are added.

# Schrodinger equation



$$\psi_{cat} = a\psi_{cat}^{alive} + b\psi_{cat}^{dead}$$

On opening the box  $\psi_{cat}$  will collapse to either  $\psi_{cat}^{alive}$  or  $\psi_{cat}^{dead}$ .

# Superposition & Entanglement

For an electron, there are two quantum states, i.e., spin-up and spin-down states represented by  $\psi^\downarrow$  and  $\psi^\uparrow$ .

For electron A, Quantum states are:  $\psi_A^\uparrow$  and  $\psi_A^\downarrow$

For electron B, Quantum states are:  $\psi_B^\uparrow$  and  $\psi_B^\downarrow$

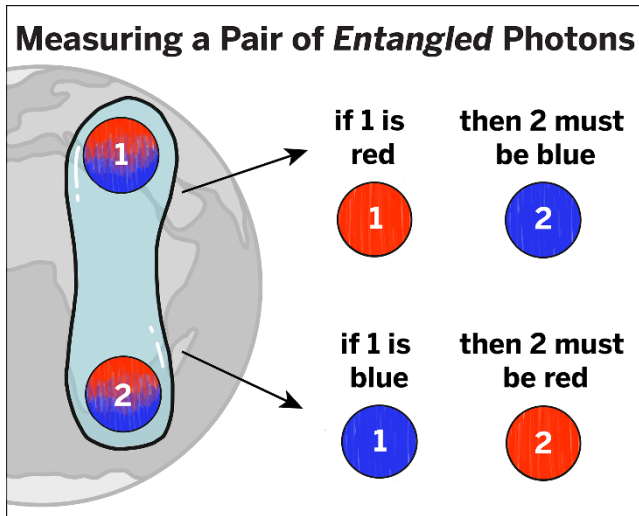
For the combination the quantum state will be the linear combination:

$$\psi = \frac{1}{\sqrt{2}} (\psi_A^\uparrow \psi_B^\downarrow \pm \psi_A^\downarrow \psi_B^\uparrow)$$

→ requirement of Pauli exclusion principle.

## Entangled-state of two quantum particles

These two electrons are related with each other with some condition, they are entangled



## Quantum entanglement

Phenomenon of a group of particles being generated, interacting, or sharing spatial proximity in such a way that the quantum state of each particle of the group cannot be described independently, even when the particles are separated by a large distance.

### Key Characteristics:

**Non-locality:** Particles seem to be "connected" over large distances.

**Quantum Correlation:** Measurement of one particle affects the state of the other.

### Examples

- **Superconductors:** Cooper pair.
- **Superconducting qubits:** IBM's or Google's quantum computers.
- **Trapped ions/atoms:** Ions or atoms can be entangled by controlled laser interactions.
- **Beam Splitters:** Entangled photon pairs can be generated using beam splitters and polarizers.

# Particle inside a box – application of Schrödinger's equation

**Free particle** (No restriction – no force – no potential !)

Particle is free to move on x-axis.

**Time independent Schrodinger Equation:**

$$\frac{\partial^2 \psi(x)}{\partial x^2} + \frac{2mE}{\hbar^2} \psi(x) = 0$$

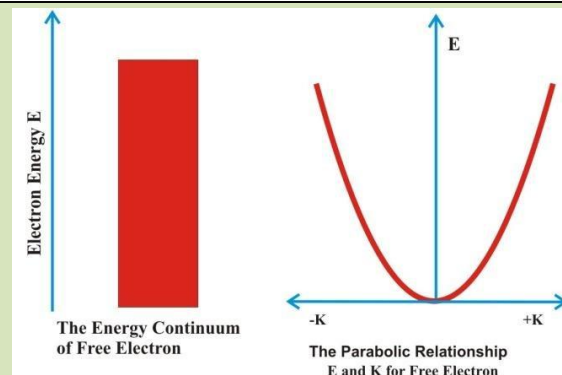
$$\frac{\partial^2 \psi(x)}{\partial x^2} + k^2 \psi(x) = 0 \quad \text{when } k = \sqrt{\frac{2mE}{\hbar^2}}$$

Solution of this equation can be

$$\psi(x) = Ae^{i(kx)} \text{ or } Be^{-i(kx)}$$

$$E = \frac{\hbar^2 k^2}{2m}$$

Since k may have any value, particle may have *any* energy!



# Particle inside a box – application of Schrödinger's equation

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**Bound particle** (A restriction is in place!)

Particle is free to move on x-axis only in a permitted region.

# Particle inside a box – application of Schrödinger's equation

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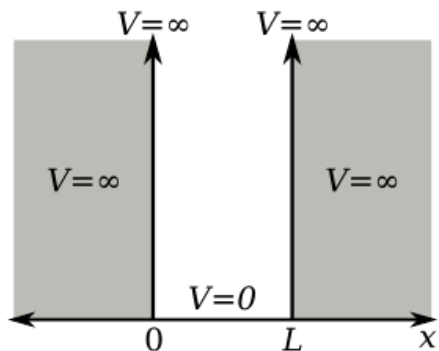
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# Particle inside a box – application of Schrödinger's equation

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Schrodinger Equation for  $0 \leq x \leq L$

$$\frac{\partial^2 \psi(x)}{\partial x^2} + \frac{2m(E - \cancel{V})}{\hbar^2} \psi(x) = 0, \quad k = \sqrt{\frac{2mE}{\hbar^2}}$$

$$\int_{-\infty}^{+\infty} |\psi(x)|^2 dx = 1$$

$$\int_{-\infty}^0 |\cancel{\psi(x)}|^2 dx + \int_0^L |\psi(x)|^2 dx + \int_L^{+\infty} |\cancel{\psi(x)}|^2 dx = 1$$

$$\psi(x) = A \sin(kx) + B \cos(kx)$$

Boundary conditions:  $\psi_{x=0} = 0$  and  $\psi_{x=L} = 0$

$$\psi_{x=0} = 0 = 0 + B$$

$$\psi_{x=L} = 0 = A \sin(kL)$$

$$k = \frac{n\pi}{L}, \quad n = 0, \pm 1, \pm 2, \dots$$

$$\psi_n(x) = A \sin(k_n x) = A \sin\left(\frac{n\pi}{L} x\right)$$

$$E_n = \frac{\hbar^2 k_n^2}{2m} = \frac{\hbar^2 n^2 \pi^2}{2mL^2}, \quad n = 1, 2, 3, \dots$$

# Particle inside a box – application of Schrödinger's equation

$$\int_{x=0}^{x=L} |\psi(x)|^2 dx = 1 \quad \psi_n(x) = A \sin\left(\frac{n\pi x}{L}\right)$$

$$\int_{x=0}^{x=L} \left| A \sin\left(\frac{n\pi x}{L}\right) \right|^2 dx = 1$$

$$\frac{A^2}{2} \int_{x=0}^{x=L} \left[ 1 - \cos\left(\frac{2n\pi x}{L}\right) \right] dx = 1$$

$$\frac{A^2}{2} \left[ x - \frac{L}{2n\pi} \sin\left(\frac{2n\pi x}{L}\right) \right]_0^L = 1$$

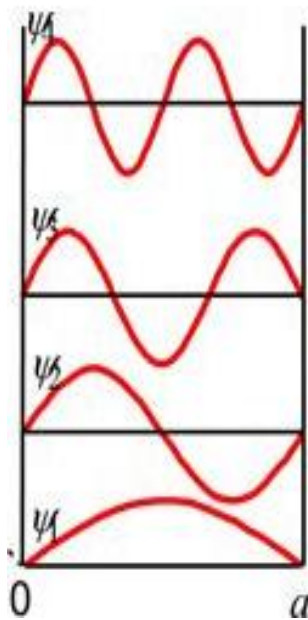
$$\frac{A^2}{2} L = 1 \quad A = \sqrt{\frac{2}{L}}$$

$$\psi_n(x) = \sqrt{\frac{2}{L}} \sin\left(\frac{n\pi x}{L}\right)$$

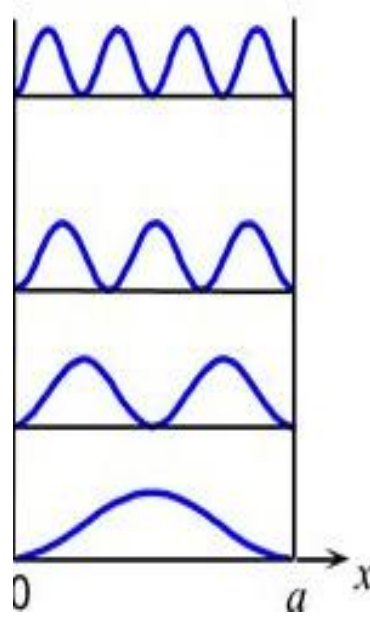
$$E = \frac{\hbar^2 \pi^2}{2mL^2} n^2$$

# Particle inside a box – application of Schrödinger's equation

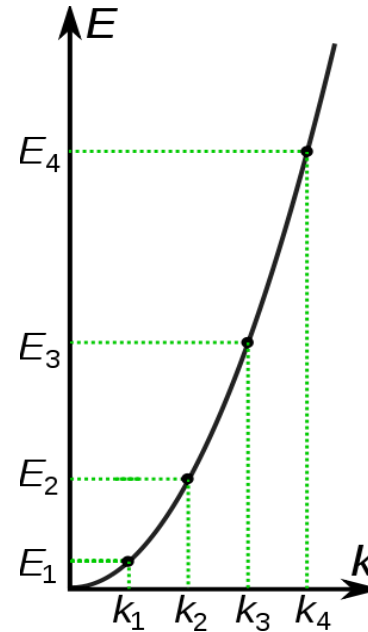
$$\psi_n(x) = A \sin\left(\frac{n\pi x}{L}\right) = \sqrt{\frac{2}{L}} \sin\left(\frac{n\pi x}{L}\right) \quad E = \frac{\hbar^2 \pi^2}{2mL^2} n^2$$



Wave functions



Probabilities



# Particle inside a box – application of Schrödinger's equation

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- What is the probability that a particle, trapped in a box of width  $L$ , is found to be between  $0.4L$  and  $0.5L$  in 1<sup>st</sup> excited state?
- Find the expectation value of position of a particle, trapped in a box of width  $L$  in ground state and 1<sup>st</sup> excited state.

# Particle inside a box – application of Schrödinger's equation

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A small 0.40-kg cart is moving back and forth along an air track between two bumpers 2.0 m apart. Assuming no friction; collisions with the bumpers are perfectly elastic so that between the bumpers, the cart maintains a constant speed of 0.50 m/s. Treating the cart as a quantum particle, estimate the value of the principal quantum number that corresponds to its classical energy.

$$K = \frac{1}{2}mv^2 = \frac{(0.4)(0.5)^2}{2} = 0.05 \text{ J}$$

If the cart as a quantum particle then energy is

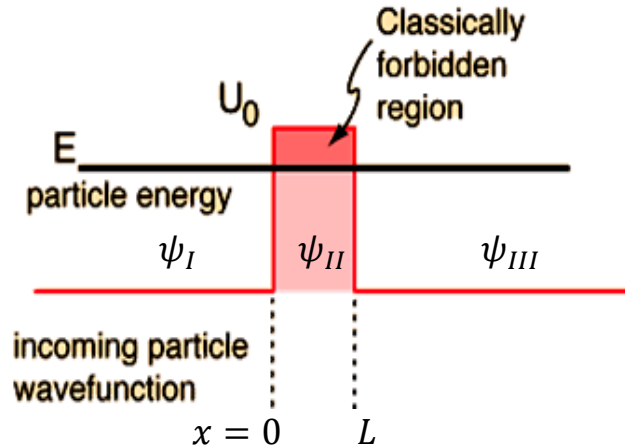
$$E_1 = \frac{\pi^2 \hbar^2}{2mL^2} n^2 = \frac{\pi^2 (1.05 \times 10^{-34})^2}{2(0.4)(2)^2} = 1.7 \times 10^{-68} n^2 \text{ J}$$

$$0.05 \text{ J} = 1.7 \times 10^{-68} n^2 \text{ J}$$

$$n = (K/E_1)^{1/2} = (0.050/1.700 \times 10^{-68})^{1/2} = 1.2 \times 10^{33}$$

In the limit of high quantum numbers, there is no advantage in using quantum formalism because we can obtain the same results using classical mechanics.

# Problem of finite potential barrier



$$V = \begin{cases} 0 & x \leq 0 \\ U_0 & 0 < x < L \\ 0 & x \geq L \end{cases}$$

$$\frac{\partial^2 \psi(x)}{\partial x^2} + \frac{2m(E - V)}{\hbar^2} \psi(x) = 0$$

$$\frac{\partial^2 \psi_I(x)}{\partial x^2} + \frac{2mE}{\hbar^2} \psi_I(x) = 0$$

$$\frac{\partial^2 \psi_{II}(x)}{\partial x^2} + \frac{2m(E - U_0)}{\hbar^2} \psi_{II}(x) = 0$$

$$\frac{\partial^2 \psi_{III}(x)}{\partial x^2} + \frac{2mE}{\hbar^2} \psi_{III}(x) = 0$$

$$k_1^2 = \frac{2mE}{\hbar^2}, \quad k_2^2 = \frac{2m(U_0 - E)}{\hbar^2}$$



# Problem of finite potential barrier

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$$\frac{\partial^2 \psi_I(x)}{\partial x^2} + k_1^2 \psi_I(x) = 0$$

$$\psi_I = A e^{ik_1 x} + B e^{-ik_1 x}$$

$$\frac{\partial^2 \psi_{II}(x)}{\partial x^2} - k_2^2 \psi_{II}(x) = 0$$

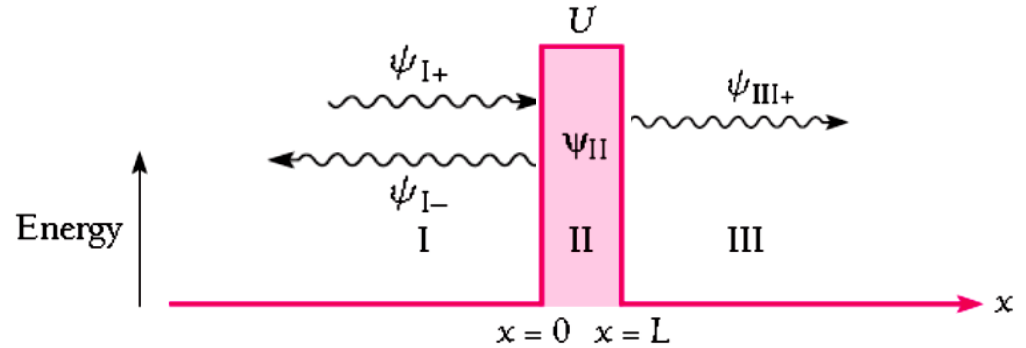
$$\psi_{II} = C e^{k_2 x} + D e^{-k_2 x}$$

$$\frac{\partial^2 \psi_{III}(x)}{\partial x^2} + k_1^2 \psi_{III}(x) = 0$$

$$\psi_{III} = F e^{ik_1 x} + G e^{-ik_1 x}$$

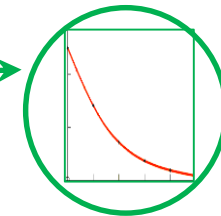
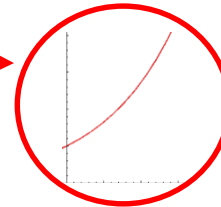
# Problem of finite potential barrier

$$\psi_I = \underbrace{Ae^{ik_1x}}_{\psi_{I+}} + \underbrace{Be^{-ik_1x}}_{\psi_{I-}}$$



$$\psi_{II} = \underbrace{Ce^{k_2x}}_{\text{red dot}} + \underbrace{De^{-k_2x}}_{\text{green dot}}$$

$$\psi_{III} = \underbrace{Fe^{ik_1x}}_{\psi_{III+}} + \underbrace{Ge^{-k_1x}}_{\psi_{III-}}$$



# Problem of finite potential barrier

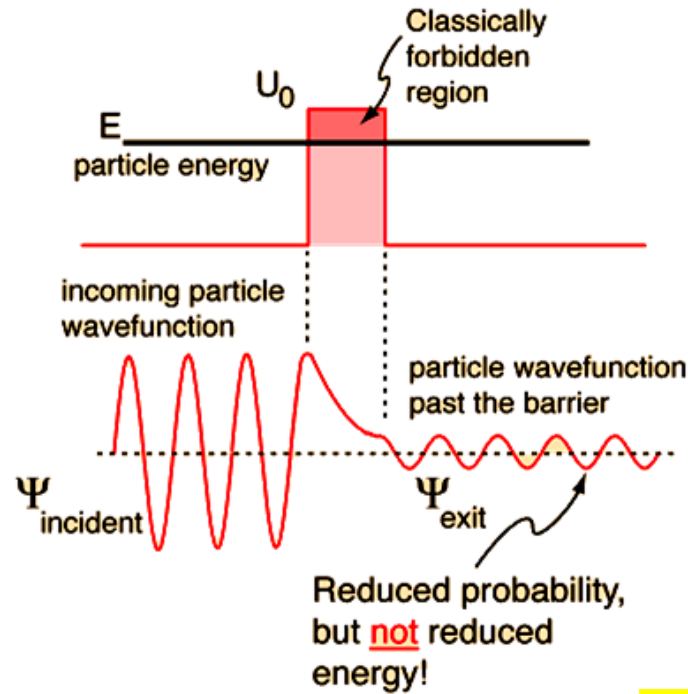
$$\psi_i = Ae^{ik_1x} + Be^{-ik_1x}$$

$$\psi_b = De^{-k_2x}$$

$$\psi_o = Fe^{ik_1x}$$

$$k_1^2 = \frac{2mE}{\hbar^2},$$

$$k_2^2 = \frac{2m(U_0 - E)}{\hbar^2}$$



Quantum mechanical  
tunneling

Transmission probability

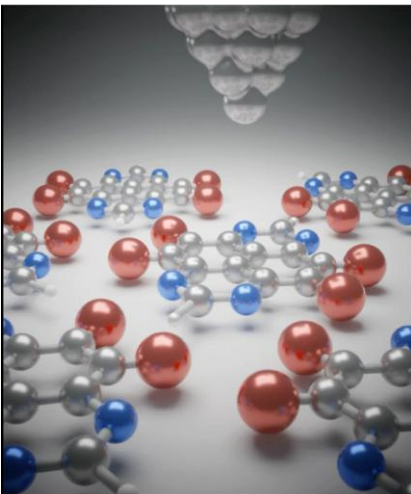
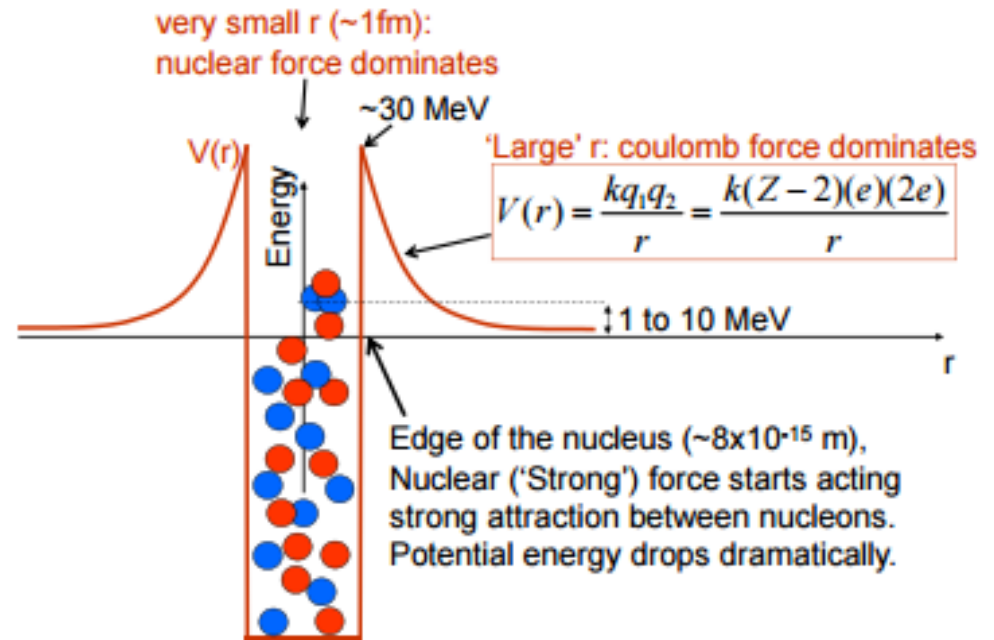
$$T = \left| \frac{F}{A} \right|^2$$

$$T \approx 16 \frac{E}{U_0} \left( 1 - \frac{E}{U_0} \right) e^{-2k_2L}$$

# Applications - Particle encountering a potential barrier

## Radioactive decay:

Probability of nucleons coming outside the nucleus energy well.



- **Semiconductors:** Tunnelling is essential in tunnel diodes and transistors.
- **Nuclear Fusion:** Tunnelling allows particles to overcome repulsive forces in nuclear reactions.
- **Scanning Tunneling Microscopy (STM):** Utilizes tunnelling to image surfaces at the atomic level.