

P 1.3. a) Assuming a flat $\Omega_M = 1$ universe, show that the probability that a light beam from a point source at “ z ” crossing an intervening (foreground) galaxy is given by:

$$P = \frac{2}{3} \frac{c}{H_0} \pi R_{\text{gal}}^2 n_0 [(1+z)^{3/2} - 1], \quad (3.1)$$

where, $R_{\text{gal}} \approx 10 \text{ kpc}$ ⁽¹⁾ is the average effective radius of a galaxy and $n_0 \approx 0.02 \text{ Mpc}^{-3}$ is the comoving number density of galaxies at $z = 0$ ($n(z) = n_0(1+z)^3$).

b) At which redshift does this expression become unity? Assume a Hubble constant value of $H_0 = 68 \text{ km s}^{-1} \text{ Mpc}^{-1}$.

Hint: The probability to intersect a galaxy with redshift in the range $[z, z + dz]$ is

$$dP = \pi R_{\text{gal}}^2 n dl, \quad (3.2)$$

where dl is the distance traveled by the light ray in the redshift interval dz .

Solution

Part a) The distance dl traveled by the light ray in the redshift interval dz is actually the proper distance covered by the light in the same redshift interval along the path which corresponds to the lookback time. Hence, dl in a world time dt is simply (cf. Eq. (4.77) in the book),

$$dl = c dt = c \frac{da}{a} \frac{a}{da} dt = c \frac{da}{a} \frac{a}{\dot{a}} = \frac{cdz}{1+z} \frac{1}{H(z)} = H_0^{-1} \frac{cdz}{(1+z)E(z)}, \quad (3.3)$$

where the expansion rate is given by,

$$E(z) = \frac{H_0}{H(z)} = \sqrt{\Omega_\gamma(1+z)^4 + \Omega_M(1+z)^3 + \Omega_K(1+z)^2 + \Omega_\Lambda}. \quad (3.4)$$

Now, assuming a flat $\Omega_M = 1$ universe, the last expression reduces to $E(z) = (1+z)^{3/2}$, which, after substituted back into Eq. (3.3), gives an expression for the distance traveled by the light ray solely in terms of the redshift z ,

$$dl = \frac{cdz}{H_0(1+z)} \cdot \frac{1}{(1+z)^{3/2}} = \frac{cdz}{H_0(1+z)^{5/2}}. \quad (3.5)$$

Thus, the intersection probability per unit redshift interval, given by Eq (3.2), becomes,

$$dP = \pi R_{\text{gal}}^2 n(z) \frac{cdz}{H_0(1+z)^{5/2}} = \frac{c \pi R_{\text{gal}}^2 n_0 (1+z)^3}{H_0 (1+z)^{5/2}} dz = \frac{c \pi R_{\text{gal}}^2 n_0}{H_0} (1+z)^{1/2} dz, \quad (3.6)$$

and the probability to intersect a galaxy with the redshift z is consequently,

$$P = \frac{c \pi R_{\text{gal}}^2 n_0}{H_0} \int_0^z (1+z')^{1/2} dz' = \frac{c \pi R_{\text{gal}}^2 n_0}{H_0} \left(\frac{2}{3} (1+z')^{3/2} \right) \Big|_0^z, \quad (3.7)$$

or, finally,

$$P = \frac{2\pi c R_{\text{gal}}^2 n_0}{3H_0} \left((1+z)^{3/2} - 1 \right). \quad (3.8)$$

¹ 1 pc = 3.261 light years = $3.086 \times 10^{16} \text{ m}$

Part b) In case of the full coverage of the sky, when every light beam ends up on some galaxy, the probability becomes unity,

$$\frac{2\pi c R_{\text{gal}}^2 n_0}{3H_0} \left((1+z)^{3/2} - 1 \right) = 1. \quad (3.9)$$

From the last equation we can easily obtain the redshift z for which the effective temperature of remote light is redshifted down (out of the visible region, for example),

$$(1+z)^{3/2} - 1 = \frac{3H_0}{2\pi c R_{\text{gal}}^2 n_0}, \quad (3.10)$$

$$z = \left(1 + \frac{3H_0}{2\pi c R_{\text{gal}}^2 n_0} \right)^{2/3} - 1, \quad (3.11)$$

After substituting the given numerical values: $H_0 = 68 \text{ km s}^{-1} \text{ Mpc}^{-1}$, $R_{\text{gal}} \approx 10 \text{ kpc}$, $n_0 \approx 0.02 \text{ Mpc}^{-3}$ and $c = 299\,792\,458 \text{ m s}^{-1}$, we obtain,

$$z = \left(1 + \frac{3 \times (68 \times 10^3 \text{ m s}^{-1} \text{ Mpc}^{-1})}{2\pi \times (299\,792\,458 \text{ m s}^{-1}) \times (10 \times 10^{-3} \text{ Mpc})^2 \times (0.02 \text{ Mpc}^{-3})} \right)^{2/3} - 1, \quad (3.12)$$

$$z = \left(1 + \frac{3 \times 68 \times 10^3}{2\pi \times 299\,792\,458 \times (10 \times 10^{-3})^2 \times 0.02} \cdot \frac{\text{m s}^{-1} \text{ Mpc}^{-1}}{\text{m s}^{-1} \text{ Mpc}^2 \text{ Mpc}^{-3}} \right)^{2/3} - 1, \quad (3.13)$$

Assuming that the input data was specified by at least two significant digits (which is generous, as both R_{gal} and n_0 have only one significant digit!), we get,


$$z \approx 13. \quad (3.14)$$

According to a more thorough uncertainty analysis (see the appendix), we get the value,

$$z \approx 13.5 \pm 5.7, \quad (3.15)$$

where the number following the symbol \pm is the numerical value of a combined standard uncertainty with a level of confidence of approximately 68 %.

Appendix: P1.3 Uncertainty Analysis

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$$\left(1 + \frac{3 H_0}{2 \pi c (R_{gal})^2 n_0}\right)^{2/3} - 1 \mapsto \begin{pmatrix} H_0 & 67\,800 \pm 770 & \text{Normal}\mathcal{D} & (\text{m / s}) / \text{Mpc} \\ c & \text{CODATA [c]} & & \\ R_{gal} & 0.010 \pm 0.005 & \text{Uniform}\mathcal{D} & \text{Mpc} \\ n_0 & 0.02 \pm 0.005 & \text{Uniform}\mathcal{D} & \text{Mpc}^{-3} \end{pmatrix}$$

Assumptions

- (1) As of 21 March 2013, the Hubble constant, as measured by the Planck Mission, is 67.80 ± 0.77 (km/s)/Mpc (combined Planck + WP + highL + BAO 68% limits).
- (2) The average effective radius of a galaxy R_{gal} is assumed to be given with 1 significant digit.
- (3) The comoving number density of galaxies n_0 is assumed to be given with 1 significant digit.

Result (1 σ coverage)

$$\begin{aligned} z &= 13.4608 \pm 5.633\,98; \text{ Normal}\mathcal{D}, k = 1\sigma \\ &\in [7.826\,87; 19.0948] \\ &\approx (1.35 \pm 0.57) \times 10^1 = 1.35(57) \times 10^1 \end{aligned}$$

Uncertainty Analysis

$$y = -1 + \left(1 + \frac{3 x_1}{2 \pi x_2 x_3^2 x_4}\right)^{2/3}$$

Quantity		Estimate \pm Uncertainty	Units	Distribution	$ \partial f / \partial x_i $
x_1	H_0	$(6.780 \pm 0.077) \times 10^4$	$\text{m s}^{-1} \text{Mpc}^{-1}$	Normal, 1σ	1.39606×10^{-4}
x_2	c	299 792 458 (exact)	m s^{-1}	-	3.15727×10^{-8}
x_3	R_{gal}	$(1.0 \pm 0.5) \times 10^{-2}$	Mpc	Uniform	1.89305×10^3
x_4	n_0	$(2.0 \pm 0.5) \times 10^{-2}$	Mpc^{-3}	Uniform	4.73263×10^2

y	13.460 848 083 285 334 613 033 274 934 2
u_c	5.633 978 024 840 062 514 352 474 239 73 = 41.9 %
$y \pm u_c$	$(1.35 \pm 0.57) \times 10^1$ = $1.35(57) \times 10^1$
ϵ_{\max}	12.154 055 346 395 523 654 381 490 573 8 = 90.3 %
$y \pm \epsilon_{\max}$	$(1.3 \pm 1.2) \times 10^1$ = $1.3(12) \times 10^1$
y_{\min}	6.255 318 59 = $y - 7.20553$
y_{\max}	41.711 953 26 = $y + 28.2511$

$\Phi\text{UC}[\Phi f\text{Estimate} \pm \Phi f\text{MaximumUncertainty}, 2, 12]$

$$\begin{aligned} &13.460\,848\,083\,3 \pm 12.1541 \\ &\in [1.3068; 25.615] \\ &\approx (1.3 \pm 1.2) \times 10^1 = 1.3(12) \times 10^1 \end{aligned}$$

$\Phi\text{UC}[\Phi f\text{Estimate} \pm \Phi f\text{StandardUncertainty}, 2, 12]$

$$\begin{aligned} &13.460\,848\,083\,3 \pm 5.633\,98 \\ &\in [7.826\,87; 19.0948] \\ &\approx (1.35 \pm 0.57) \times 10^1 = 1.35(57) \times 10^1 \end{aligned}$$

$\Phi\text{UC}[\Phi\text{MonteCarlo}[10^6], 3, 10]$

$$\begin{aligned} z &= 15.928\,576\,29 \pm 7.505\,15; \text{ Normal}\mathcal{D}, k = 1\sigma \\ &\in [8.423\,423; 23.433\,73] \\ &\approx (1.593 \pm 0.751) \times 10^1 = 1.593(751) \times 10^1 \end{aligned}$$