

# Columns

2013

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The most critical mode of failure for such aircraft structures as thin webs stiffened by slender longerons or stringers is **buckling**.

Extremely important to predict the buckling loads of the following structures: 1) columns, 2) thin plates, 3) stiffened panels.

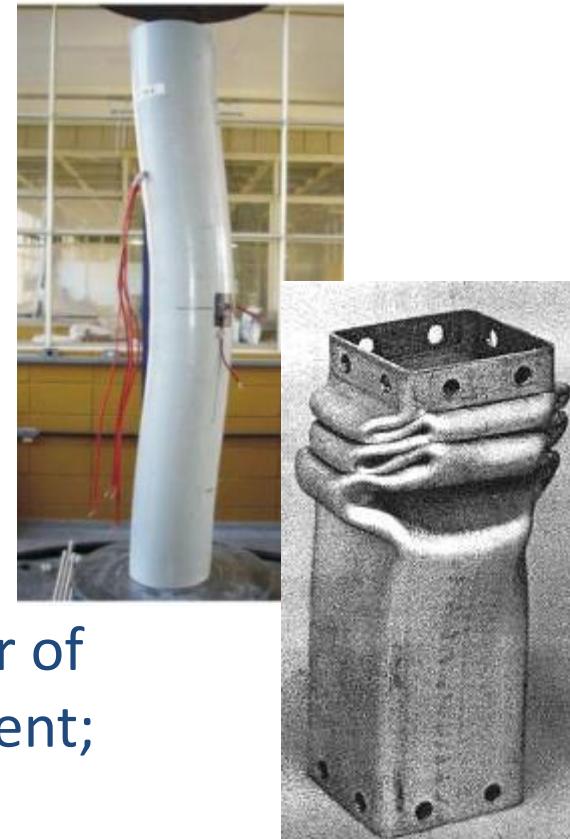
Two types of structural instability arise:

## Primary:

- no change in cross-sectional area;
- the wavelength of the buckle is of the same order as the length of the element;
- solid and thick-walled columns.

## Secondary:

- changes in cross-sectional area occur;
- the wavelength of the buckle is of the order of the cross-sectional dimensions of the element;
- thin-walled columns and stiffened plates.



# Euler buckling of columns

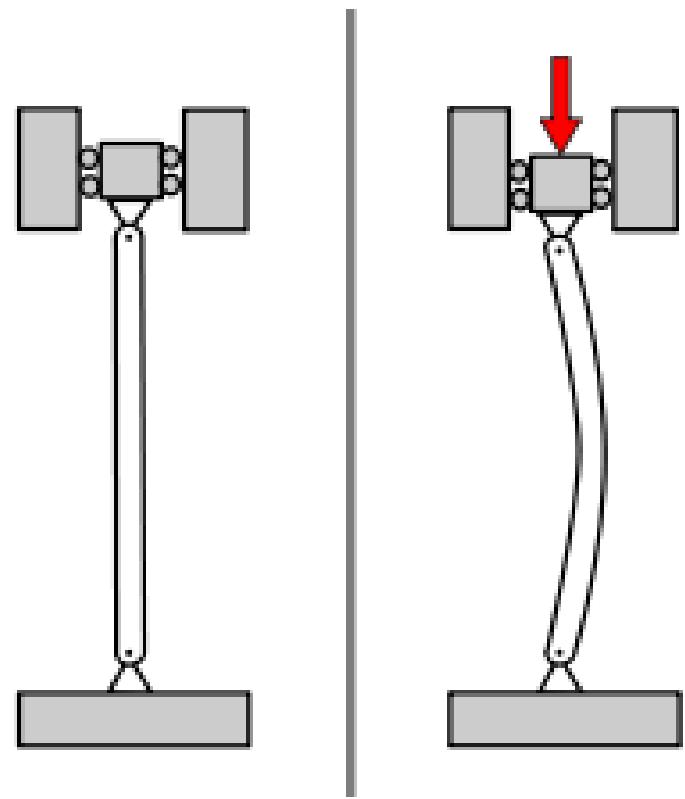
If an increasing axial compressive load is applied to a slender column there is a value of the load called **buckling load** at which the column will suddenly bow or buckle in some unpredetermined direction.

Direction of the buckling may depend:

- degree of asymmetry;
- geometrical imperfection;
- asymmetrical load application;
- material imperfection.

Consider a perfect column in which

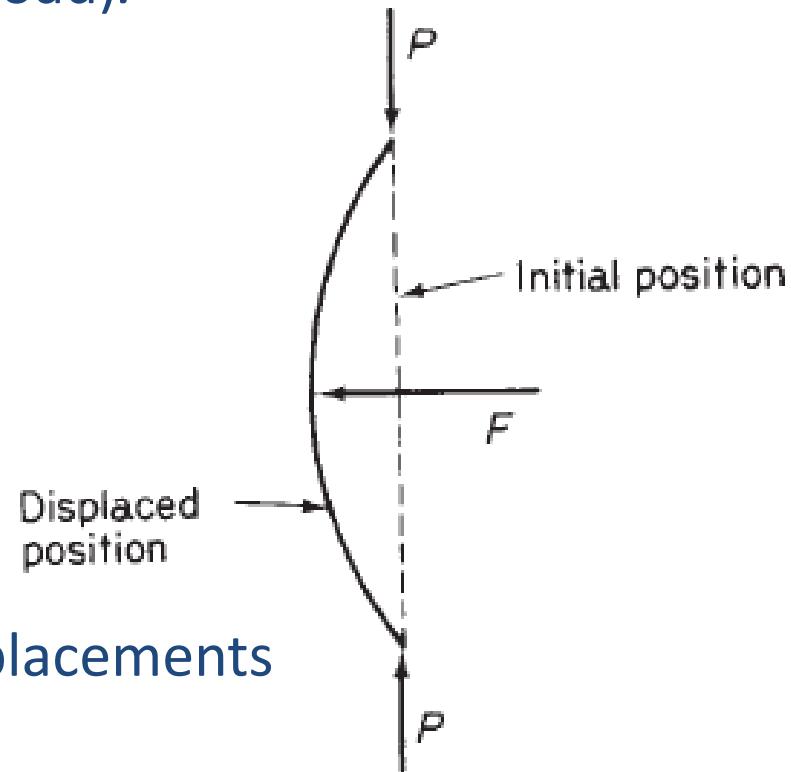
- the load is applied precisely along the perfectly straight centroidal axis,
- there is perfect symmetry so that, theoretically, there can be no sudden bowing or buckling.



# Euler buckling of columns

## Properties of the “perfect” column

1. Only shortening of the perfect column occurs no matter what the value of compressive load  $P$ .
2. The column is displaced a small amount by a lateral load  $F$  at values of  $P < P_{CR}$  ( $P_{CR}$  – critical or buckling load).
3. Removal of  $F$  results in a return of the column to its undisturbed position (state of *stable equilibrium*).
4. At the  $P_{CR}$  the displacement does not disappear and the column will remain in any displaced position as long as the displacement is small (state of a *neutral equilibrium*).
5. For  $P > P_{CR}$  enforced lateral displacements increase and the column is *unstable*.

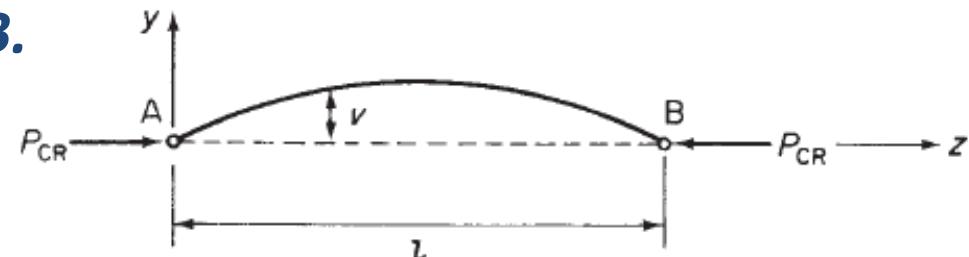


# Euler buckling of columns

Consider the pin-ended column **AB**.

Assume that it is in the displaced state of neutral equilibrium ( $P_{CR}$ ).

Simple bending theory gives:



$$EI \frac{d^2v}{dz^2} = -M \quad \rightarrow \quad EI \frac{d^2v}{dz^2} = -P_{CR} v \quad \rightarrow \quad \frac{d^2v}{dz^2} + \frac{P_{CR}}{EI} v = 0$$

The well-known solution  $v = A \cos \mu z + B \sin \mu z$

where  $\mu^2 = P_{CR}/EI$ ,  $A$  and  $B$  – unknown coefficients

The boundary conditions for this particular case are

$v = 0$  at  $z = 0$  and  $l$ , thus  $A = 0$  and  $B \sin \mu l = 0$

For a non-trivial solution (i.e.  $v \neq 0$ ) then

$\sin \mu l = 0$  or  $\mu l = n\pi$  where  $n = 1, 2, 3, \dots$

giving  $\frac{P_{CR} l^2}{EI} = n^2 \pi^2$  or  $P_{CR} = \frac{n^2 \pi^2 EI}{l^2}$

# Euler buckling of columns

Note that solution of differential equation  $v = A \cos \mu z + B \sin \mu z$  cannot be solved for  $v$  no matter how many of the available boundary conditions are inserted, since neutral state of equilibrium means that  $v$  is indeterminate.

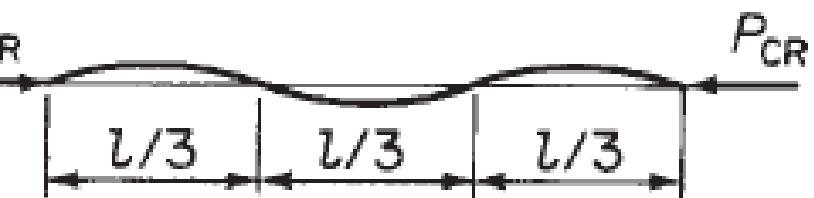
The smallest value of critical buckling load at  $n = 1$   $P_{CR} = \frac{\pi^2 EI}{l^2}$

Other values of  $P_{CR}$  (at  $n = 2, 3, \dots$ )  $P_{CR} = \frac{4\pi^2 EI}{l^2}, \frac{9\pi^2 EI}{l^2}, \dots$

These values of buckling load cause more **complex modes** of buckling:



$$P_{CR} = 4\pi^2 EI/l^2$$



$$P_{CR} = 9\pi^2 EI/l^2$$

Can be produced by applying external restraints to a very slender column at the points of contraflexure to prevent lateral movement.

# Euler buckling of columns

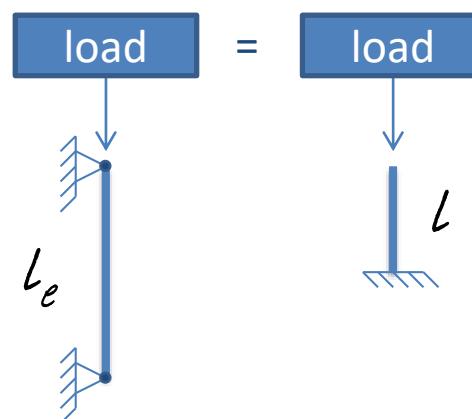
$$P_{\text{CR}} = \frac{\pi^2 EI}{l^2} \rightarrow \sigma_{\text{CR}} = \frac{\pi^2 E}{(l/r)^2}$$
 here  $r = \sqrt{I/A}$  is the **radius of gyration** of the cross-sectional area of the column.

The term  $l/r$  is known as the **slenderness ratio** of the column.

- for unsymmetrical column  $r$  is the least (since the column will bend about an axis about which the flexural rigidity  $EI$  is least);
- for column with restraints  $EI$  is the flexural rigidity in unrestrained plane.

Can rewrite foregoing equations:  $P_{\text{CR}} = \frac{\pi^2 EI}{l_e^2}$      $\sigma_{\text{CR}} = \frac{\pi^2 E}{(l_e/r)^2}$

where  $l_e$  is the **effective length of the column**, i.e. the length of a pinned-ended column that would have the same critical load as that of a column of length  $l$ , but with different end conditions.



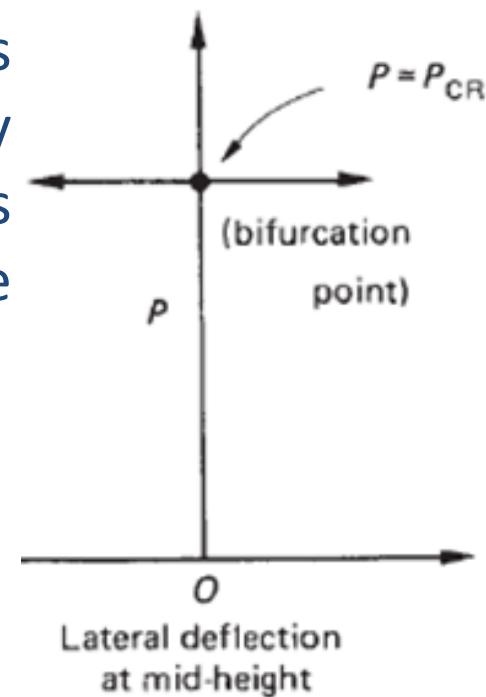
# Euler buckling of columns

Ends	$l_e/l$	Boundary conditions
Both pinned	1.0	$v = 0$ at $z = 0$ and $l$
Both fixed	0.5	$v = 0$ at $z = 0$ and $z = l$ , $dv/dz = 0$ at $z = l$
One fixed, the other free	2.0	$v = 0$ and $dv/dz = 0$ at $z = 0$
One fixed, the other pinned	0.6998	$dv/dz = 0$ at $z = 0$ , $v = 0$ at $z = l$ and $z = 0$

If the lateral load  $F$  is removed the column (which is perfectly straight, homogeneous and loaded exactly along its axis) will suffer only axial compression as  $P$  is increased. This situation, theoretically, would continue until yielding of the material of the column occurred.

However

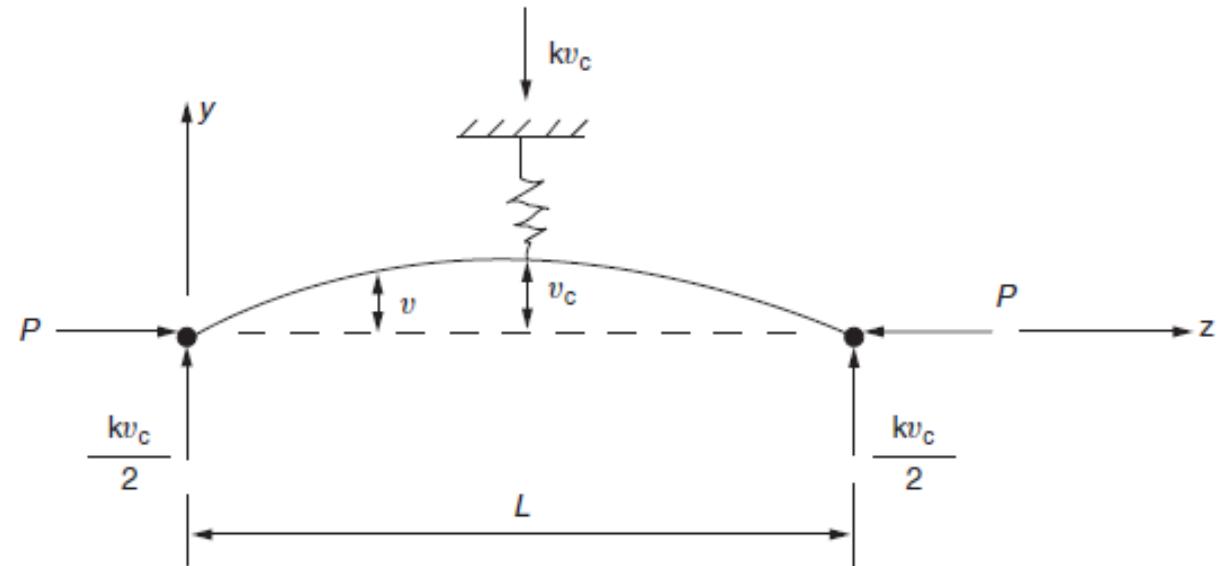
- for values of  $P < P_{CR}$  the column is in stable equilibrium;
- whereas for  $P > P_{CR}$  the column is unstable.



# Euler buckling of columns

A uniform column of

- length  $L$
- flexural stiffness  $EI$
- is simply supported
- has an elastic support at midspan.



This support is such that if a lateral displacement  $v_c$  occurs at this point a restoring force  $kv_c$  is generated at the point.

Derive

- an equation giving the buckling load of the column.
- If the buckling load is  $4\pi^2EI/L^2$  find the value of  $k$ .
- Also if the elastic support is infinitely stiff show that the buckling load is given by the equation  $\tan \lambda L/2 = \lambda L/2$  where  $\lambda = \sqrt{P/EI}$ .

# Euler buckling of columns

The bending moment at any section of the column  $M = Pv - \frac{kv_c}{2}z$

$$EI \frac{d^2v}{dz^2} = -M \quad \xrightarrow{M = Pv - \frac{kv_c}{2}z} \quad EI \frac{d^2v}{dz^2} = -Pv + \frac{kv_c}{2}z \quad \xrightarrow{\frac{d^2v}{dz^2} + \lambda^2 v = \frac{kv_c}{2EI}z}$$

The solution  $v = A \cos \lambda z + B \sin \lambda z + \frac{kv_c}{2P}z$

The constants  $A$  and  $B$  are found using the boundary conditions of the column which are:

- $v = 0$  when  $z = 0$ ,
- $v = v_c$ , when  $z = L/2$  and
- $(dv/dz) = 0$  when  $z = L/2$ .

$$A = 0 \quad B = \frac{v_c}{\sin(\lambda L/2)} \left(1 - \frac{k\lambda}{4P}\right)$$

$$\left(1 - \frac{kL}{4P}\right) \cos \frac{\lambda L}{2} + \frac{k}{2P\lambda} \sin \frac{\lambda L}{2} = 0 \quad \xrightarrow{} \quad P = \frac{kL}{4} \left(1 - \frac{\tan(\lambda L/2)}{\lambda L/2}\right)$$

If  $P = 4\pi^2 EI/L^2$  then  $\lambda L/2 = \pi$  so that  $k = 4P/L$  (as  $\tan \pi = 0$ ).

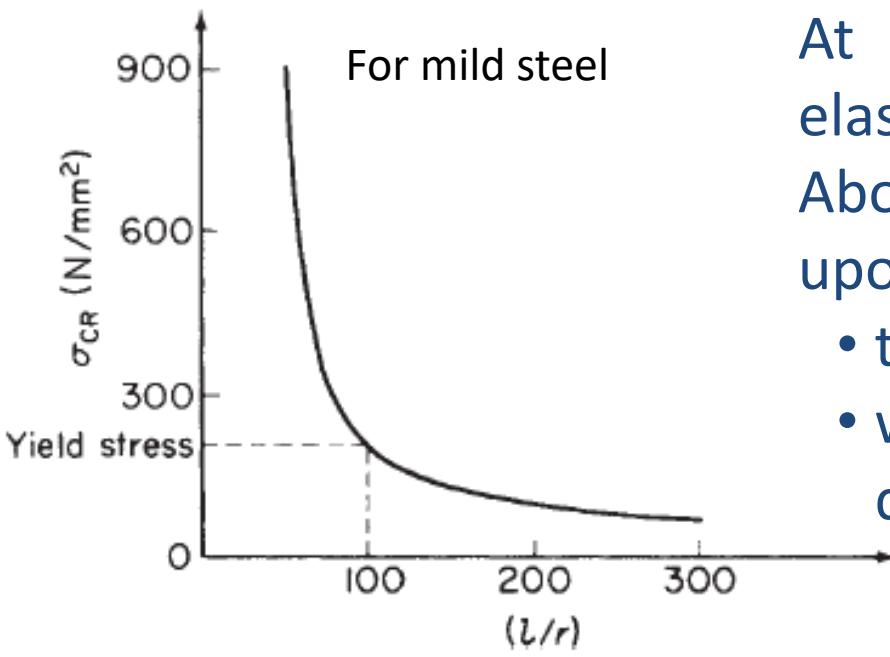
Finally, if  $k \rightarrow \infty$  then  $\tan(\lambda L/2) = (\lambda L/2)$

# Inelastic buckling

The critical stress  $\sigma_{CR} = \frac{\pi^2 E}{(l_e/r)^2}$  (applicable for  $\sigma < \sigma_y$ ) depends only on

- the elastic modulus  $E$  of the material of the column and
- the slenderness ratio  $l/r$ .

The critical stress increases as the slenderness ratio decreases; i.e. as the column becomes shorter and thicker.



At the elastic region the modulus of elasticity  $E$  ( $=d\sigma/d\varepsilon$ ) is constant. Above the elastic limit  $d\sigma/d\varepsilon$  depends upon

- the value of stress and
- whether the stress is increasing or decreasing.

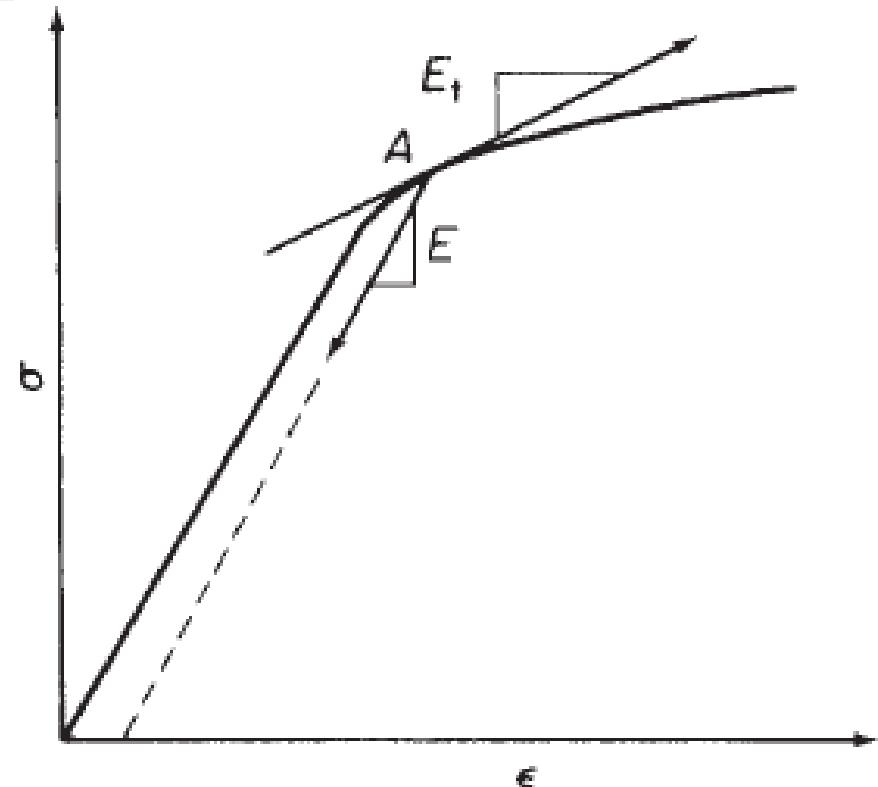
# Inelastic buckling

The elastic modulus at the point **A** is

- the tangent modulus  $E_t$  if the stress is increasing but
- $E$  if the stress is decreasing.

Consider

- a column having a plane of symmetry
- subjected to a compressive load  $P$  (so stress is above the elastic limit).



If the column is given a small deflection,  $v$ , in its plane of symmetry, then the stress on the concave side increases while the stress on the convex side decreases.

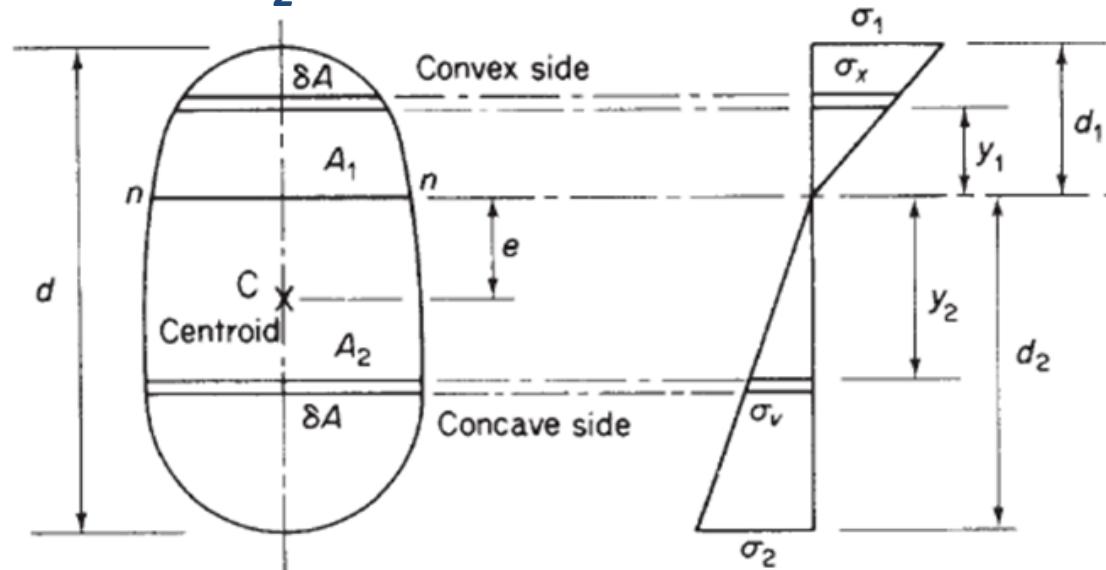
# Inelastic buckling

In the cross-section of the column the compressive stress decreases in the area  $A_1$  and increases in the area  $A_2$ , while the stress on the line  $nn$  is unchanged.

The modulus of elasticity of the material in the area  $A_1$  is  $E$  while that in  $A_2$  is  $E_t$ .

The homogeneous column now behaves as if it were non-homogeneous.

The linearity of the distribution follows from an assumption that plane sections remain plane.



$$\int_0^{d_1} \sigma_x dA = \int_0^{d_2} \sigma_v dA \quad \text{as load } P \text{ is unchanged by the disturbance.}$$

$$\int_0^{d_1} \sigma_x(y_1 + e) dA + \int_0^{d_2} \sigma_v(y_2 - e) dA = -Pv \quad \text{as } P \text{ is applied at C.}$$

# Inelastic buckling

According to given figure (left)

$$\sigma_x = \frac{\sigma_1}{d_1} y_1 \quad \sigma_v = \frac{\sigma_2}{d_2} y_2$$

The change of slope equals to the angle between two close sections of the column, and equal to angle  $\delta\varphi$  in the strain diagram.

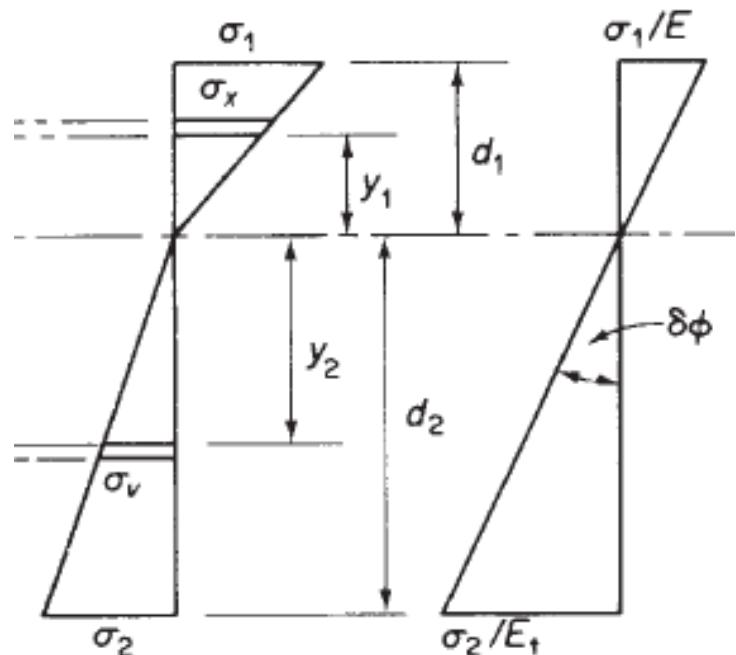
$$\frac{d^2v}{dz^2} = \frac{\sigma_1}{Ed_1} = \frac{\sigma_2}{E_t d_2}$$

$$\int_0^{d_1} \sigma_x dA = \int_0^{d_2} \sigma_v dA$$

$$\frac{d^2v}{dz^2} = \frac{\sigma_1}{Ed_1} = \frac{\sigma_2}{E_t d_2}$$

$$\sigma_x = \frac{\sigma_1}{d_1} y_1 \quad \sigma_v = \frac{\sigma_2}{d_2} y_2$$

$$E \frac{d^2v}{dz^2} \int_0^{d_1} y_1 dA - E_t \frac{d^2v}{dz^2} \int_0^{d_2} y_2 dA = 0$$



# Inelastic buckling

In similar manner

$$\int_0^{d_1} \sigma_x(y_1 + e) dA + \int_0^{d_2} \sigma_v(y_2 - e) dA = -Pv$$

↓

$$\frac{d^2v}{dz^2} = \frac{\sigma_1}{Ed_1} = \frac{\sigma_2}{E_t d_2}$$
$$\sigma_x = \frac{\sigma_1}{d_1} y_1 \quad \sigma_v = \frac{\sigma_2}{d_2} y_2$$
$$\frac{d^2v}{dz^2} \left( E \int_0^{d_1} y_1^2 dA + E_t \int_0^{d_2} y_2^2 dA \right) + e \frac{d^2v}{dz^2} \left( E \int_0^{d_1} y_1 dA - E_t \int_0^{d_2} y_2 dA \right) = -Pv$$

= 0

Therefore we have  $\frac{d^2v}{dz^2}(EI_1 + E_tI_2) = -Pv$

where  $I_1 = \int_0^{d_1} y_1^2 dA$  and  $I_2 = \int_0^{d_2} y_2^2 dA$

the second moments of area about **nn** of the convex and concave sides of the column respectively.

# Inelastic buckling

Putting  $E_r I = EI_1 + E_t I_2$  or  $E_r = E \frac{I_1}{I} + E_t \frac{I_2}{I}$   $E_r$  - reduced modulus

$$\frac{d^2v}{dz^2}(EI_1 + E_t I_2) = -Pv \xrightarrow{E_r I = EI_1 + E_t I_2} E_r I \frac{d^2v}{dz^2} + Pv = 0$$

Comparing with equation of bending of the column

$$\frac{d^2v}{dz^2} + \frac{P_{CR}}{EI} v = 0$$

we see that if  $P$  is the critical load  $P_{CR}$  then  $P_{CR} = \frac{\pi^2 E_r I}{l_e^2}$   $\sigma_{CR} = \frac{\pi^2 E_r}{(l_e/r)^2}$

The above method for predicting critical loads and stresses outside the elastic range is known as the *reduced modulus theory*

$$E \frac{d^2v}{dz^2} \int_0^{d_1} y_1 dA - E_t \frac{d^2v}{dz^2} \int_0^{d_2} y_2 dA = 0 \xrightarrow{} E \int_0^{d_1} y_1 dA - E_t \int_0^{d_2} y_2 dA = 0$$

The last equation together with the relationship  $d = d_1 + d_2$  enables the position of  $nn$  to be found

# Inelastic buckling

Also there are two possible cases:

the axial load  $P$  is increased at the time of the lateral disturbance of the column:

- no strain reversal on convex side
- the compressive stress therefore increases over the complete section
- the tangent modulus applies over the whole cross-section

reduction in  $P$  could result in a decrease in stress over the whole cross-section:

- strain reversal on both sides
- the compressive stress therefore decreases over the complete section
- usual modulus applies over the whole cross-section

$$P_{CR} = \frac{\pi^2 E_t I}{l_e^2}$$

$$\sigma_{CR} = \frac{\pi^2 E_t}{(l_e/r)^2}$$

$$P_{CR} = \frac{\pi^2 EI}{l_e^2}$$

$$\sigma_{CR} = \frac{\pi^2 E}{(l_e/r)^2}$$

The buckling load of columns is given most accurately for practical purposes by the **tangent modulus theory**

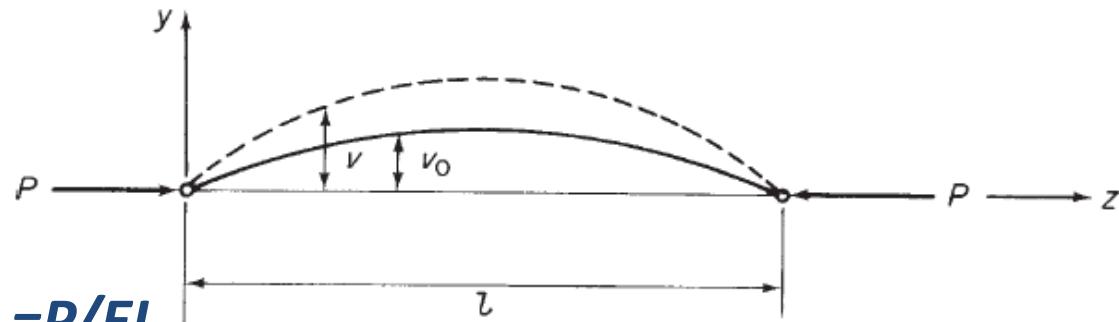
# Effect of initial imperfections

Possible imperfections influence to a large degree the behaviour of the column which, unlike the perfect column, begins to bend immediately the axial load is applied.

Let us suppose that a column, initially bent, is subjected to an increasing axial load  $P$ .

$$EI \frac{d^2v}{dz^2} - EI \frac{d^2v_0}{dz^2} = -Pv$$

$$\frac{d^2v}{dz^2} + \lambda^2 v = \frac{d^2v_0}{dz^2} \quad \text{where } \lambda^2 = P/EI.$$



The final deflected shape,  $v$ , of the column depends upon the form of its unloaded shape,  $v_0$ .

$$v_0 = \sum_{n=1}^{\infty} A_n \sin \frac{n\pi z}{l}$$

$\frac{d^2v}{dz^2} + \lambda^2 v = \frac{d^2v_0}{dz^2} \xrightarrow{\text{blue arrow}} \frac{d^2v}{dz^2} + \lambda^2 v = -\frac{\pi^2}{l^2} \sum_{n=1}^{\infty} n^2 A_n \sin \frac{n\pi z}{l}$

# Effect of initial imperfections

Solution

$$v = B \cos \lambda z + D \sin \lambda z + \sum_{n=1}^{\infty} \frac{n^2 A_n}{n^2 - \alpha} \sin \frac{n\pi z}{l}$$

where  $B$  and  $D$  are constants of integration and  $\alpha = \lambda^2 l^2 / \pi^2$

The boundary conditions

$$v = 0 \text{ at } z = 0 \text{ then } B = 0$$

$$v = 0 \text{ at } z = l \text{ then } D = 0$$

$$v = \sum_{n=1}^{\infty} \frac{n^2 A_n}{n^2 - \alpha} \sin \frac{n\pi z}{l} \quad \alpha = \frac{Pl^2}{\pi^2 EI} = \frac{P}{P_{CR}} \quad \alpha < 1 \rightarrow \text{first term in series dominates}$$

A good approximation when the axial load is close to critical load

$$v = \frac{A_1}{1 - \alpha} \sin \frac{\pi z}{l} \quad \text{or at the centre of the column } z = l/2 \quad v = \frac{A_1}{1 - P/P_{CR}}$$

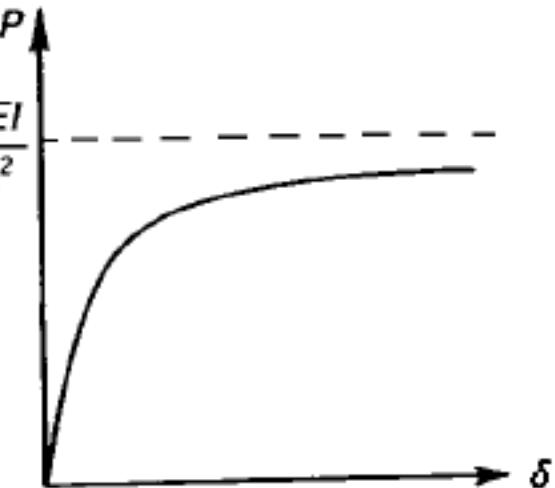
$A_1$  – initial central deflection

# Effect of initial imperfections

Imagine that we measure central deflections from the initially bowed position of the column, i.e.  $\delta = v - A_1$

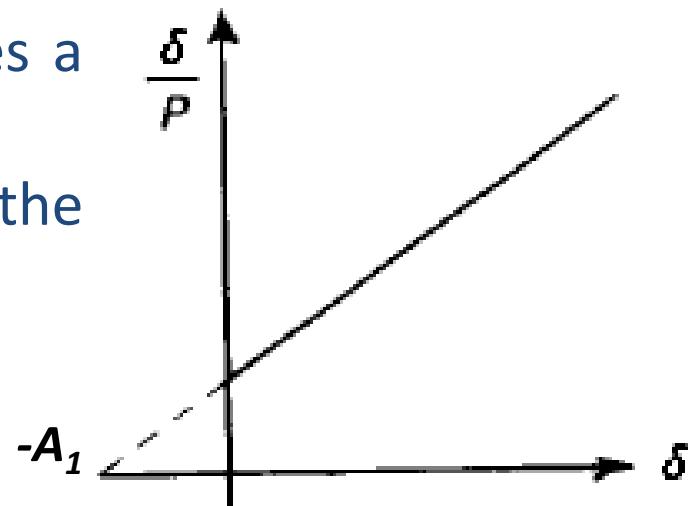
$$\frac{A_1}{1 - P/P_{CR}} - A_1 = \delta \quad \rightarrow \quad \delta = P_{CR} \frac{\delta}{P} - A_1 \quad P_{CR} = \frac{\pi^2 EI}{L^2}$$

Theoretically load  $P_{CR}$  can be attained at an infinitely large deflection (in practice the column breaks before  $P_{CR}$  could be attained).



Plotting  $\delta$  against  $\delta/P$  in a column test gives a value of initial central deflection  $A_1$ .

The slope of this straight line defines  $P_{CR}$  (the buckling load for a perfect-straight column).



Obrigado!