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## Homework 2

1. Starting with  $C_V = dU/dT = \frac{3}{2}Nk_B$  and  $C_P = dU/dT + PdV/dT = \frac{5}{2}Nk_B$ , we have:

$$C_P/C_V = \frac{\frac{5}{2}Nk_B}{\frac{3}{2}Nk_B} \tag{1}$$

$$=\frac{5}{3}\tag{2}$$

$$=\gamma$$
 (3)

$$C_P - C_V = \frac{5}{2}Nk_B - \frac{3}{2}Nk_B \tag{4}$$

$$=Nk_{B} \tag{5}$$

2. (a) To find the speed at which the fastest electron should be moving, we want to set nV times the integral of the Maxwell-Boltzmann distribution equal to one and solve for the velocity. Given that  $\int_x^\infty y^2 e^{-y^2} dy \approx x e^{-x^2}/2$  for  $x \gg 1$ , we set  $x = \sqrt{mv^2/2k_BT}$  and find:

$$1 = nV \frac{1}{2}xe^{-x^2} \tag{6}$$

$$= \frac{nV}{2} \sqrt{\frac{mv^2}{2k_B T}} e^{-\frac{mv^2}{2k_B T}}$$
 (7)

Solving this numerically using scipy.optimize.fsolve, we find that the maximum expected velocity is 0.224c, which is much less than the cosmic ray velocity. Therefore cosmic rays cannot be of thermal origin.

3. (a) For a particle cross section of  $\sigma$ , the number density at which the mean free path equals the scale height is:

$$\frac{1}{\sigma n_{\rm esc}} = \frac{k_B T}{mg} \tag{8}$$

$$n_{\rm esc} = \frac{mg}{k_B T \sigma} \tag{9}$$

(b) The upwards flux of particles is:

$$\phi(v)dv = \frac{f(v)dv}{4\pi} \int_0^{2\pi} d\phi \int_0^{\pi/2} v \cos\theta \sin\theta d\theta \tag{10}$$

$$= \frac{f(v)dv}{2} \int_0^{\pi/2} v \cos \theta \sin \theta d\theta \tag{11}$$

$$=\frac{vf(v)dv}{4}\tag{12}$$

(c) The total flux of escaping particles is:

$$\phi = \int_{v_{\rm esc}}^{\infty} \phi(v) dv \tag{13}$$

$$= \frac{1}{4} n_{\rm esc} \int_{v_{\rm esc}}^{\infty} v f(v) dv \tag{14}$$

$$= \frac{1}{4} n_{\rm esc} \int_{v_{\rm esc}}^{\infty} 4\pi v^3 \left(\frac{m}{2\pi k_B T}\right)^{3/2} e^{-\frac{mv^2}{2k_B T}} dv \tag{15}$$

$$= \pi n_{\rm esc} \left(\frac{m}{2\pi k_B T}\right)^{3/2} \int_{v_{\rm esc}}^{\infty} v^3 e^{-\frac{mv^2}{2k_B T}} dv$$
 (16)

$$= \pi n_{\rm esc} \left(\frac{m}{2\pi k_B T}\right)^{3/2} \frac{k_B T}{m^2} e^{\frac{-mv_{\rm esc}^2}{2k_B T}} \left(mv_{\rm esc}^2 + 2k_B T\right)$$
 (17)

$$= \pi n_{\rm esc} \left(\frac{m}{2\pi k_B T}\right)^{3/2} e^{\frac{-mv_{\rm esc}^2}{2k_B T}} \left(\frac{k_B T}{m} v_{\rm esc}^2 + \frac{2(k_B T)^2}{m^2}\right)$$
(18)

Substituting in  $v_s = \sqrt{2k_BT/m}$  and  $\lambda_{\rm esc} = (v_{\rm esc}/v_s)^2 = mv_{\rm esc}^2/2k_BT$ , we have:

$$\phi = \pi n_{\text{esc}} \left( \frac{m}{2\pi k_B T} \right)^{3/2} e^{\frac{-mv_{\text{esc}}^2}{2k_B T}} \left( \frac{k_B T}{m} v_{\text{esc}}^2 + \frac{2(k_B T)^2}{m^2} \right)$$
(19)

$$= \frac{n_{\rm esc}}{2\sqrt{\pi}} v_s^{-3} e^{-\lambda_{\rm esc}} \left( v_s^2 v_{\rm esc}^2 + v_s^4 \right) \tag{20}$$

$$= \frac{v_s n_{\rm esc}}{2\sqrt{\pi}} e^{-\lambda_{\rm esc}} \left(\lambda_{\rm esc} + 1\right) \tag{21}$$

- (d) I find that the rate of H2 loss over the entire exosphere is  $1.33 \times 10^{26}$  particles per second, which means that over 1 Gyr  $4.20 \times 10^{42}$  particles will escape. This is about half of the current hydrogen content of Earth's atmosphere.
- (e) For oxygen, I find that the loss rate is on the order of  $10^{-70}$  particles per second, which comes out to much less than a single particle being lost over 1 Gyr. This implies a stable oxygen abundance over time.
- (f) For deuterium, I find a loss rate of  $9.6 \times 10^{36}$  particles per second, and a total loss over 1 Gyr of  $3 \times 10^{36}$  particles. This is a slower loss rate than H2, but not slow enough to be stable.