

- ① Using a second-order Adams-Bashforth scheme for the advection terms, and a second order Crank-Nicholson scheme for the viscous term, the predictor step (Wesseling, 2001):

$$\frac{\underline{u}^* - \underline{u}^n}{\Delta t} = -\frac{3}{2} \underline{A}_h(\underline{u}^n) + \frac{1}{2} \underline{A}_h(\underline{u}^{n-1}) + \frac{\nu}{2} \left( \underline{D}_h(\underline{u}^n) + \underline{D}_h(\underline{u}^*) \right) + \text{force term} \quad (2.9)$$

and the correction step is:

$$\frac{\underline{u}^{n+1} - \underline{u}^*}{\Delta t} = -\nabla_h \Phi^{n+1} \quad (2.10)$$

$\Phi$  is not exactly pressure





②

It can be shown:

$$-\nabla \phi^{n+1} = -\nabla p^{n+1} + \frac{\nu}{2} \left( \nabla_h(\underline{u}^{n+1}) - \nabla_h(\underline{u}^*) \right) \quad (2.11)$$

$\underline{u}^*$  is not div.-free, so pseudo  
-pressure can be found as:

$$\nabla_h^2 \phi^{n+1} = \frac{\nabla_h \cdot \underline{u}^*}{\Delta^+} \quad (2.12)$$

(2.12) must be solved before  $\underline{u}^{n+1}$  is  
obtained.



③

Re-arrange (2.9); to get a  
Helmholtz eq. for  $u^n$ :

$$D_h(u^n) - \frac{2}{\nu \Delta t} u^n = \frac{3}{\nu} A_h(u^n) - \frac{1}{\nu} A_h(u^{n-1})$$

$$- D_h(u^n) - \frac{2}{\nu \Delta t} u^n \equiv \text{R.H.S.} \quad (2.13)$$





④ First find the average over each control volume:

$$\underline{A}(\underline{u}^n) = \frac{1}{\Delta V} \int_V \underline{\nabla} \cdot (\underline{u}^n \underline{u}^n) dV$$

$$= \frac{1}{\Delta V} \oint_S \underline{u}^n (\underline{u}^n \cdot \underline{n}) dS \quad (2.14)$$

and  $\underline{D}(\underline{u}^n) = \frac{1}{\Delta V} \int_V \nabla^2 \underline{u}^n dV \quad (2.15)$

$\Delta V$  is the volume of C.V.

$S$  is the surface of C.V.

Div. Theorem:  $\iiint_V (\underline{\nabla} \cdot \underline{F}) dV = \oint_S (\underline{F} \cdot \underline{n}) dS$





⑤ Continuity eq., (2.2) can be numerically approximated,  $\nabla_h \cdot \underline{u}^{n+1} = 0$  through interpretation over volume & converting to surface integral:

$$\frac{1}{\Delta V} \int_V \nabla \cdot \underline{u}^{n+1} dV = \frac{1}{\Delta V} \oint_S \underline{u}^{n+1} \cdot \underline{n} dS \quad (2.16)$$

Incomp:

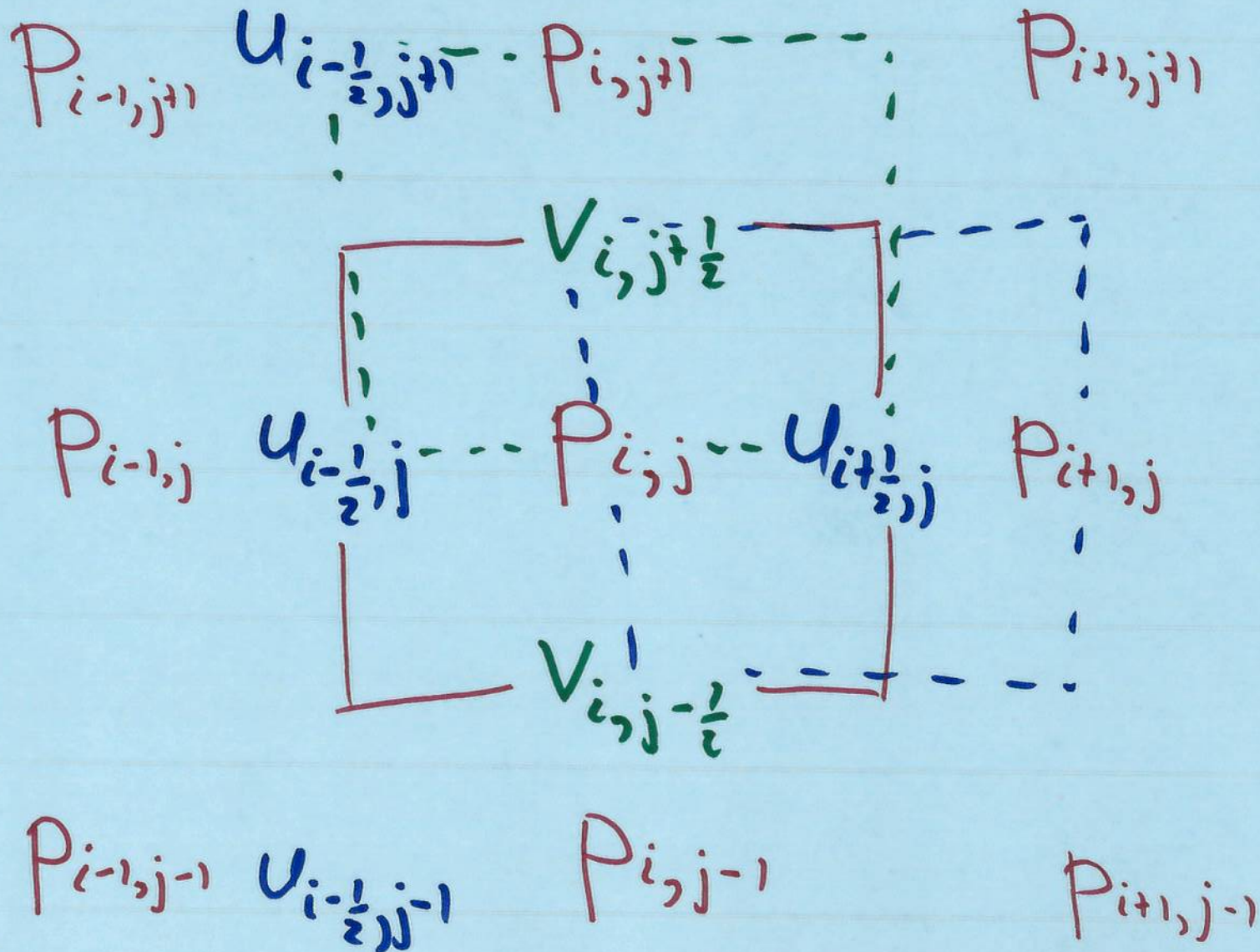
$$u_{i+\frac{1}{2},j}^{n+1} - u_{i-\frac{1}{2},j}^{n+1} + v_{i,j+\frac{1}{2}}^{n+1} - v_{i,j-\frac{1}{2}}^{n+1} = 0 \quad (2.17)$$

Assumes same grid spacing in both directions.





⑥



⑦ Discrete forms of eq. (2.3) & (2.4) for  $u$  vel. in a C.V. centered at  $(i+\frac{1}{2}, j)$  and  $v$  vel. centered @  $(i, j+\frac{1}{2})$  are:

$$u_{i+\frac{1}{2}, j}^* = u_{i+\frac{1}{2}, j}^n + \Delta t \left( -(A_x)_{i+\frac{1}{2}, j}^n + v(D_x)_{i+\frac{1}{2}, j}^n + (f_x)_{i+\frac{1}{2}, j} \right)$$

$$v_{i, j+\frac{1}{2}}^* =$$

(2.18)

- predictor step.

