

# A Continuous Interior Penalty Method Framework for Sixth Order Cahn-Hilliard-type Equations with applications to microstructure evolution and microemulsions

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Joint work with Amanda Diegel (Mississippi State University)



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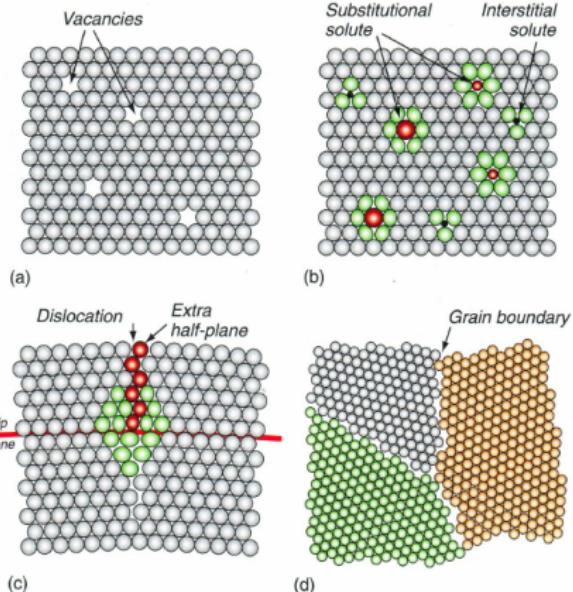
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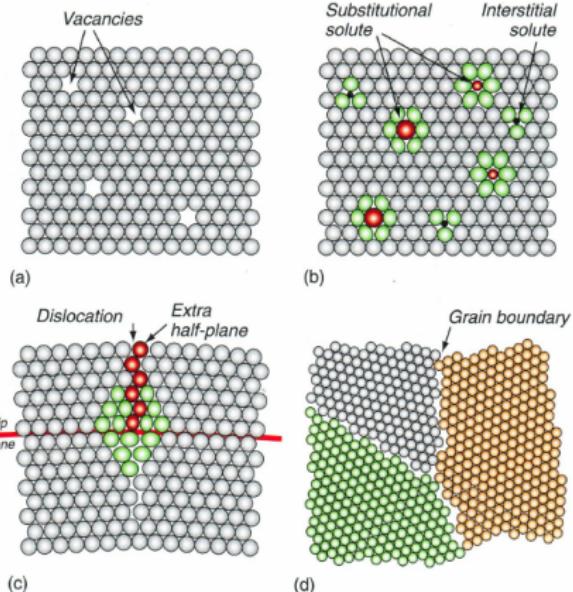
- Microstructure evolution specifically formation of defects within microstructures.
- Microemulsions capturing the dynamics of phase transitions in a oil-water-surfactant mixtures with applications to designing drug delivery systems.

Goal: Develop a Continuous Interior Penalty Method framework to solve these Sixth-Order Phase Field Models.

# Defects in Crystalline Materials

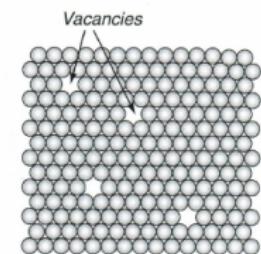


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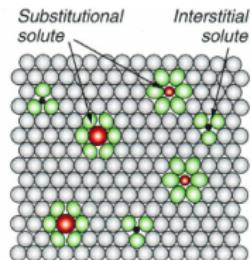


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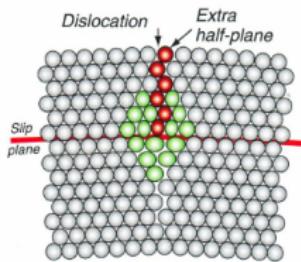
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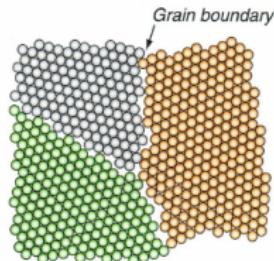
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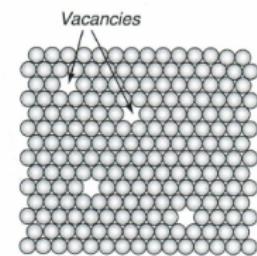


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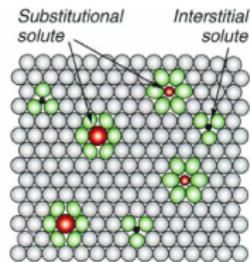
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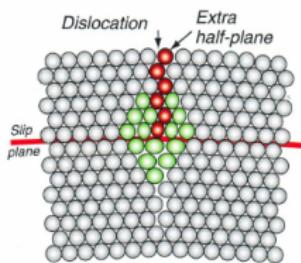
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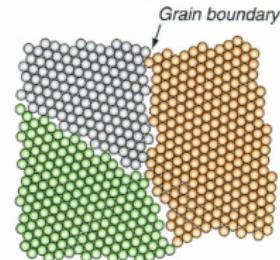
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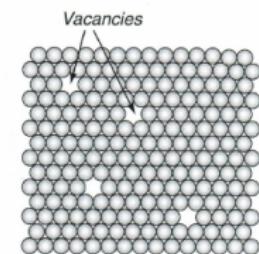


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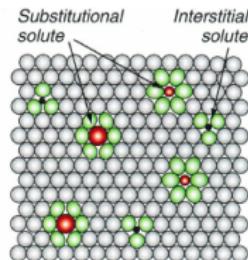
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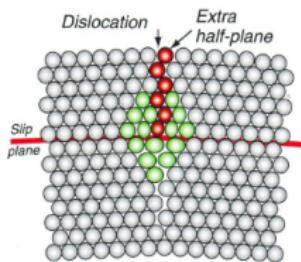
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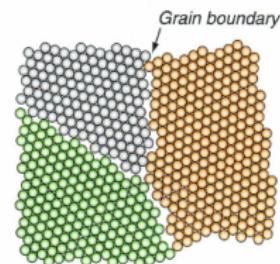
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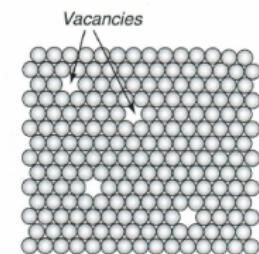


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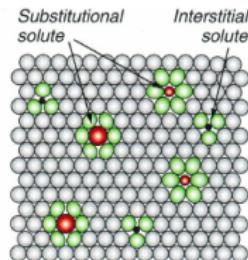
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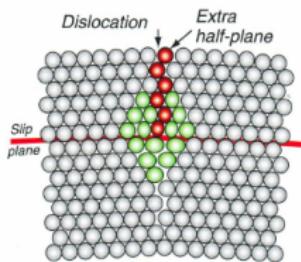
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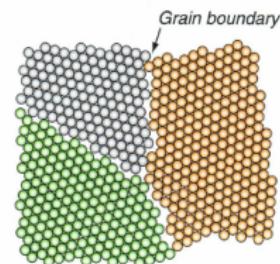
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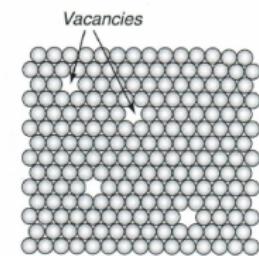


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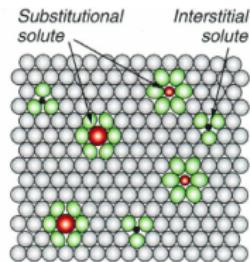
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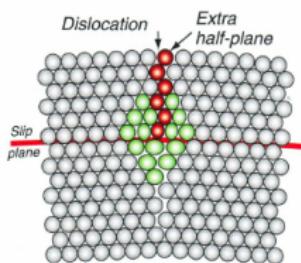
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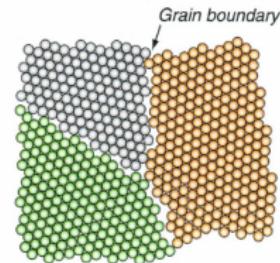
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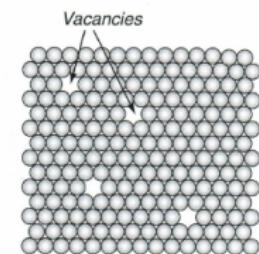
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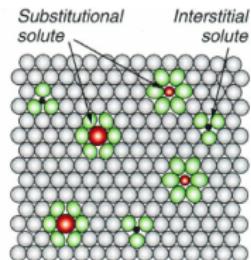
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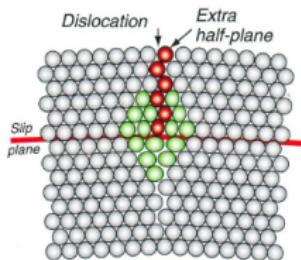
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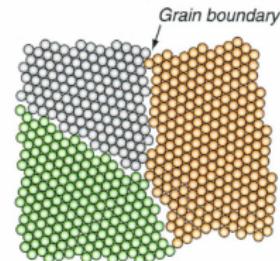
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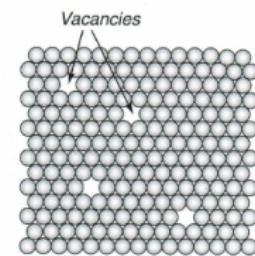
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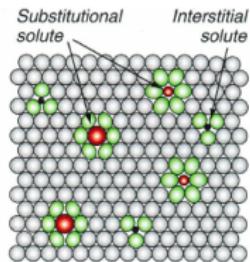
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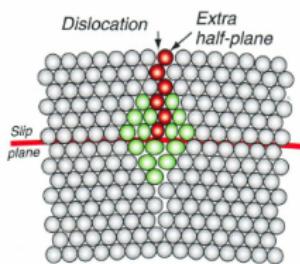
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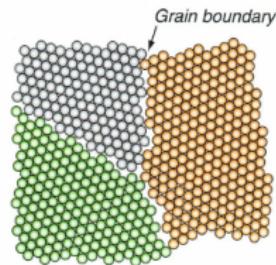
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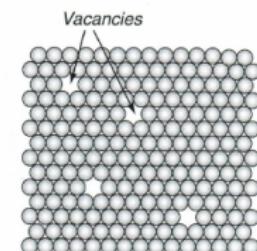
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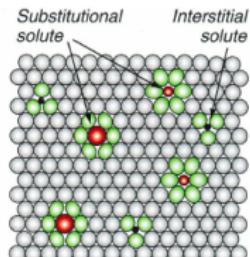
**Approach:**

- Use **phase field crystal equation** as our atomistic model.

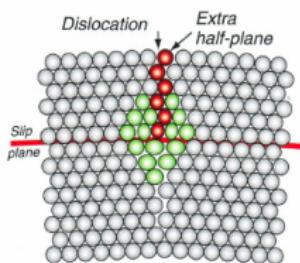
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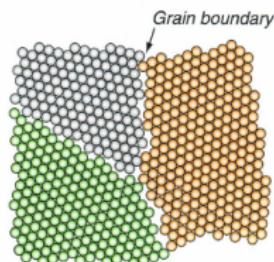
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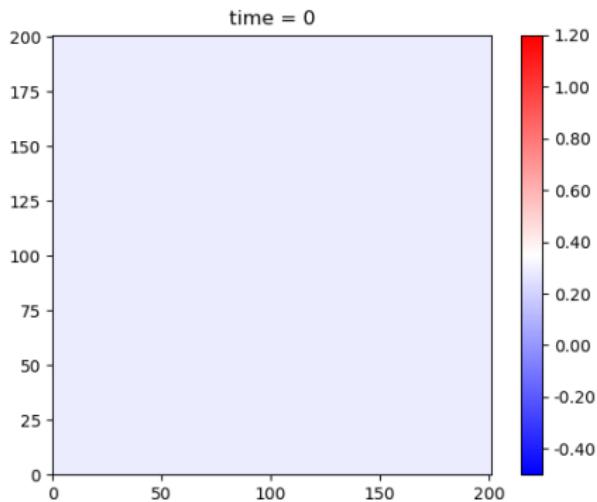
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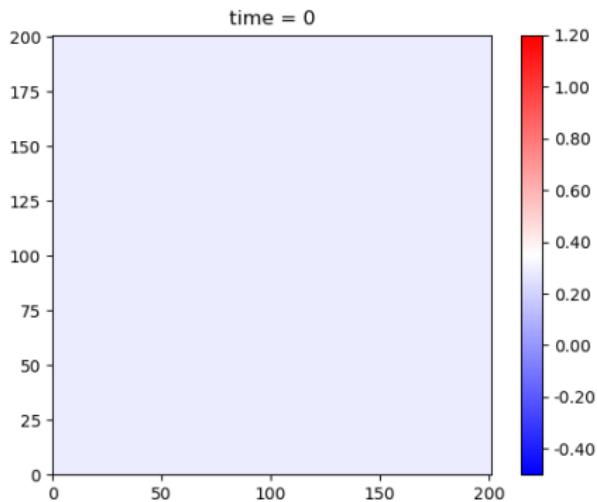
- Use **phase field crystal equation** as our atomistic model.
- Develop an accurate, efficient, easy-to-compute numerical scheme.

# Phase Field Crystal Equation



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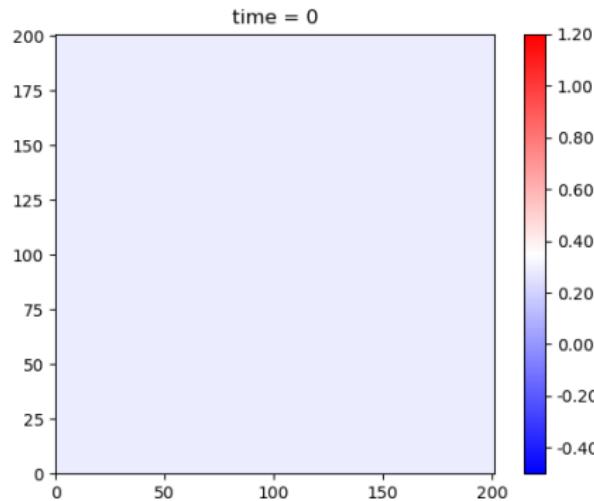
Two-phase system



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$\varphi$ : number density of atoms in the material occupying  $\Omega$  with

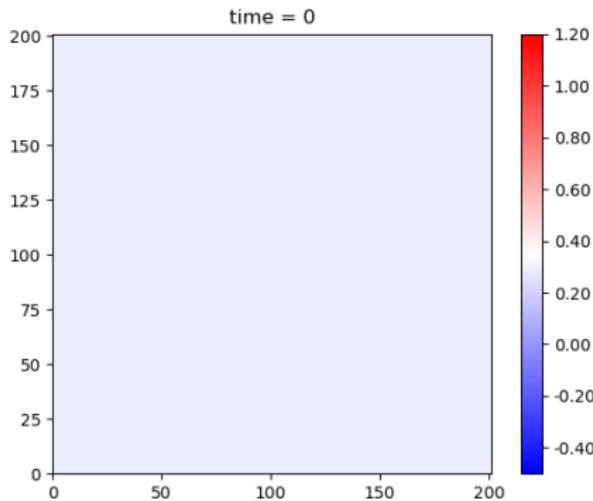


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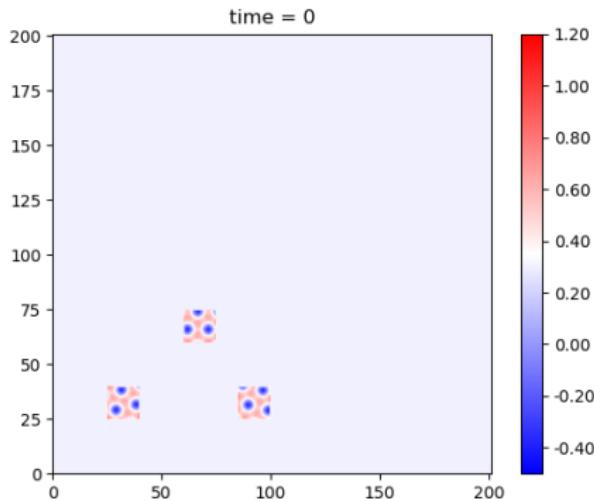


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- ◊ liquid phase characterized by a constant value of  $\varphi$
- ◊ solid phase characterized by a spatially varying periodic function  $\varphi$  that inherits the symmetry and periodicity of the crystal lattice



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Phase Field Crystal Equation (Elder et al. 2004)

$$\frac{\partial \varphi}{\partial t} = \nabla \cdot (\mathcal{M} \nabla (\varphi^3 + 2\Delta\varphi + (1-\varepsilon)\varphi + \Delta^2\varphi)) \quad \text{on } \Omega \times (0, T).$$

# PFC equation in the mixed form

$$\frac{\partial \varphi}{\partial t} - \nabla \cdot (\mathcal{M} \nabla \mu) = 0,$$
$$\varphi^3 + (1 - \varepsilon)\varphi + 2\Delta\varphi + \Delta^2\varphi - \mu = 0,$$

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with either periodic boundary conditions or natural boundary conditions

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Notation:

- $H^s(\Omega)$  denote the Sobolev spaces of order  $s \geq 1$ ,
- $Z := \{z \in H^2(\Omega) \mid \mathbf{n} \cdot \nabla z = 0 \text{ on } \partial\Omega\}.$

# Weak Formulation

Find  $(\varphi, \mu)$  such that

$$\begin{aligned}\varphi &\in L^\infty(0, T; Z) \cap L^2(0, T; H^3(\Omega)), \\ \partial_t \varphi &\in L^2(0, T; H_N^{-1}(\Omega)), \\ \mu &\in L^2(0, T; H^1(\Omega)),\end{aligned}$$

and for almost all  $t \in (0, T)$

$$\begin{aligned}\langle \partial_t \varphi, \nu \rangle + (\mathcal{M} \nabla \mu, \nabla \nu) &= 0 \quad \forall \nu \in H^1(\Omega) \\ ((\varphi)^3 + (1 - \epsilon)\varphi, \psi) - 2(\nabla \varphi, \nabla \psi) + a(\varphi, \psi) - (\mu, \psi) &= 0 \quad \forall \psi \in Z\end{aligned}$$

with  $a(u, v) := \int_{\Omega} (\nabla^2 u : \nabla^2 v) dx,$

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[Pawlow et al., 2013]

# Numerical Schemes: Challenges

## Phase Field Crystal Equation

$$\begin{aligned}\frac{\partial \varphi}{\partial t} - \nabla \cdot (\mathcal{M} \nabla \mu) &= 0, \\ \varphi^3 + (1 - \varepsilon)\varphi + 2\Delta\varphi + \Delta^2\varphi - \mu &= 0 \quad \text{on } \Omega \times (0, T).\end{aligned}$$

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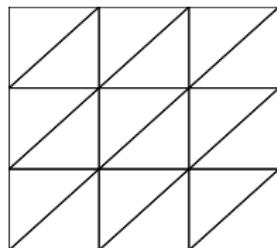
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- Fourier-spectral Method: Li and Shen, 2020 (Scalar Auxiliary Variable approach)  
Yang and Han, 2017 (Invariant Energy Quadratization)

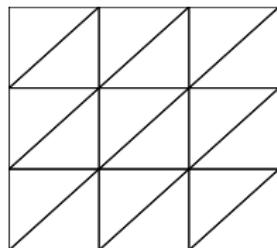
# Numerical Schemes: Our Approach

- Space discretization: Relax the  $C^1$ -continuity, use  $C^0$ -Interior Penalty Method
- Time discretization: Use Eyre's convex splitting scheme known to be uniquely solvable and unconditionally stable

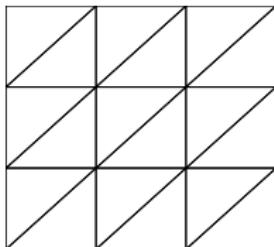
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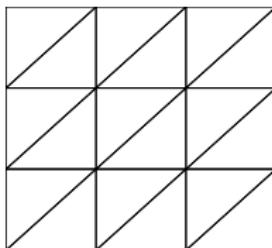


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**Assume this partition is geometrically-conforming and shape-regular.**

$\mathcal{T}_h$ : collection of all elements  $K$

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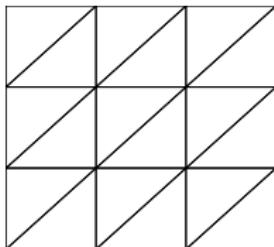
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- $h_K = \text{diameter of triangle } K$ ,  $h = \max_{K \in \mathcal{T}_h} h_K$

# Spatial Discretization



- $\Omega = \bigcup_{K \in \mathcal{T}_h} K$ ,

**Assume this partition is geometrically-conforming and shape-regular.**

$\mathcal{T}_h$ : collection of all elements  $K$

- $h_K = \text{diameter of triangle } K$ ,  $h = \max_{K \in \mathcal{T}_h} h_K$

- $\mathcal{E}_h$ : collection of all edges  $e$  wrt  $\mathcal{T}_h$

# Classical $C^1$ Finite Element Method

Find  $(\varphi, \mu) : [0, T] \rightarrow Z \times H^1(\Omega)$  s.t. for almost all  $t \in (0, T)$

$$\langle \partial_t \varphi, \nu \rangle + (\mathcal{M} \nabla \mu, \nabla \nu) = 0 \quad \forall \nu \in H^1(\Omega)$$

$$(\varphi^3 + (1 - \epsilon)\varphi, \psi) - 2(\nabla \varphi, \nabla \psi) + a(\varphi, \psi) - (\mu, \psi) = 0 \quad \forall \psi \in Z$$

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$$\langle \partial_t \varphi_h, \nu \rangle + (\mathcal{M} \nabla \mu_h, \nabla \nu) = 0,$$

$$((\varphi_h)^3 + (1 - \epsilon)\varphi_h, \psi) - 2(\nabla \varphi_h, \nabla \psi) + a_h^{IP}(\varphi_h, \psi) - (\mu_h, \psi) = 0$$

$\forall \nu \in V_h, \psi \in Z_h$  holds for almost all  $t \in (0, T)$ .

# Spatial Discretization using C<sup>0</sup>-IP Method

$$V_h := \{v \in C(\bar{\Omega}) \mid v|_K \in P_1(K) \ \forall K \in \mathcal{T}_h\}$$

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$R_h : H^1(\Omega) \rightarrow V_h$  is a Ritz projection operator such that

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$P_h : Z \rightarrow Z_h$  is a Ritz projection operator such that

$$a_h^{IP}(P_h\varphi - \varphi, \xi) + (1 - \epsilon)(P_h\varphi - \varphi, \xi) = 0 \quad \forall \xi \in Z_h, \quad (P_h\varphi - \varphi, 1) = 0.$$

# Spatial Discretization using C<sup>0</sup>-IP Method

$a_h^{IP} : Z_h \times Z_h \rightarrow \mathbb{R}$  according to

$$a_h^{IP}(\xi_h, \psi_h) := \sum_{K \in \mathcal{T}_h} \int_K \nabla^2 \xi_h : \nabla^2 \psi_h \, dx + J(\xi_h, \psi_h), \quad \xi_h, \psi_h \in Z_h$$

where

$$\begin{aligned} J(\xi_h, \psi_h) := & \sum_{e \in \mathcal{E}_h} \int_e \left( [n_e \cdot \nabla \xi_h]_e \{n_e \cdot \nabla^2 \psi_h n_e\}_e + \{n_e \cdot \nabla^2 \xi_h n_e\}_e [n_e \cdot \nabla \psi_h]_e \right) ds \\ & + \sum_{e \in \mathcal{E}_h} \int_e \frac{\alpha}{h_e} [n_e \cdot \nabla \xi_h]_e [n_e \cdot \nabla \psi_h]_e \, ds, \quad \xi_h, \psi_h \in Z_h \end{aligned}$$

and  $\alpha > 0$  is a penalty parameter.

# Boundedness of $a_h^{IP}(\cdot, \cdot)$

## Lemma (Boundedness of $a_h^{IP}(\cdot, \cdot)$ )

There exists positive constants  $C_{cont}$  and  $C_{coer}$  such that for choices of the penalty parameter  $\alpha$  large enough we have

$$\begin{aligned} a_h^{IP}(w_h, v_h) &\leq C_{cont} \|w_h\|_{2,h} \|v_h\|_{2,h} \quad \forall w_h, v_h \in Z_h, \\ C_{coer} \|w_h\|_{2,h}^2 &\leq a_h^{IP}(w_h, w_h) \quad \forall w_h \in Z_h, \end{aligned}$$

where the constants  $C_{cont}$  and  $C_{coer}$  depend only on the shape regularity of  $\mathcal{T}_h$ .

where the  $C^0$ -IP Norm is:

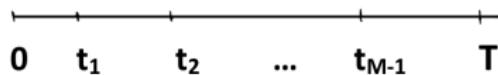
$$\|\xi_h\|_{2,h}^2 := \sum_{K \in \mathcal{T}_h} |\xi_h|_{H^2(K)}^2 + \sum_{e \in \mathcal{E}_h} \alpha \|h_e^{-\frac{1}{2}} [\mathbf{n}_e \cdot \nabla \xi_h]_e\|_{L^2(e)}^2.$$

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Notation:  $\varphi^m$  approximate  $\varphi$  at time  $t_m$ .

Numerical time derivative w.r.t.  $\tau$ :

$$\delta_\tau \varphi^m := \frac{\varphi^{m+1} - \varphi^m}{\tau}$$

# Convex Time Splitting Scheme

- Basic Idea:

$$\begin{aligned}\mu &= \delta_\varphi E(\varphi) = \delta_\varphi \left( \underbrace{E^+(\varphi)}_{\text{convex}} + \underbrace{E^-(\varphi)}_{\text{concave}} \right) \\ \implies \mu^m &= \underbrace{(\varphi^m)^3 + (1 - \epsilon)\varphi^m + \Delta^2\varphi^m}_{\delta_\varphi E^+(\varphi^m)} + \underbrace{2\Delta\varphi^{m-1}}_{\delta_\varphi E^-(\varphi^{m-1})}\end{aligned}$$

where  $E(\varphi) = \int_{\Omega} \left( \frac{\varphi^4}{4} + \frac{1-\epsilon}{2}\varphi^2 + \frac{1}{2}(\Delta\varphi)^2 - |\nabla\varphi|^2 \right) dx.$

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Given  $\varphi^0$ , find  $(\varphi^m, \mu^m)$  for  $1 \leq m \leq M$  by

$$\delta_\tau \varphi^m - \nabla \cdot (\mathcal{M} \nabla \mu^m) = 0,$$

$$(\varphi^m)^3 + (1 - \epsilon)\varphi^m + \Delta^2\varphi^m + 2\Delta\varphi^{m-1} - \mu^m = 0,$$

with boundary conditions

$$\partial_n \varphi^m = \partial_n \Delta \varphi^m = \partial_n \mu^m = 0.$$

# Fully Discrete $C^0$ -IP Method

Given  $\varphi_h^{m-1} \in Z_h$ , find  $\varphi_h^m, \mu_h^m \in Z_h \times V_h$  such that for all  $\nu_h \in V_h$ ,  $\psi_h \in Z_h$  it holds

$$(\delta_\tau \varphi_h^m, \nu_h) + (\mathcal{M} \nabla \mu_h^m, \nabla \nu_h) = 0$$

$$\left( (\varphi_h^m)^3 + (1 - \epsilon) \varphi_h^m, \psi_h \right) + a_h^{IP} (\varphi_h^m, \psi_h) - 2 (\nabla \varphi_h^{m-1}, \nabla \psi_h) - (\mu_h^m, \psi_h) = 0,$$

where  $\varphi_h^0 := P_h \varphi_0$  and  $\mu_h^0 \in V_h$  is defined as  $\mu_h^0 := R_h \mu_0$ .

## Remark

*The scheme satisfies the discrete conservation property*

$$(\varphi_h^m, 1) = (\varphi_h^0, 1) = (\varphi_0, 1) \text{ for any } 1 \leq m \leq M.$$

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$$F(\varphi_h^m) := \frac{1}{4} \|\varphi_h^m\|_{L^4}^4 + \frac{1-\epsilon}{2} \|\varphi_h^m\|_{L^2}^2 - \|\nabla \varphi_h^m\|_{L^2}^2 + \frac{1}{2} a_h^{IP} (\varphi_h^m, \varphi_h^m)$$

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## ③ Optimal error estimates: **Main Result!**

# Unconditional Unique Solvability

## Definition

Define the functional  $G_h : \mathring{Z}_h \rightarrow \mathbb{R}$

$$\begin{aligned} G_h(\varphi_h) := & \frac{\tau}{2} \left\| \frac{\varphi_h - \varphi_h^{m-1}}{\tau} \right\|_{-1,h}^2 + \frac{1}{2} a_h^{IP}(\varphi_h, \varphi_h) + \frac{1}{4} \|\varphi_h + \bar{\varphi}_0\|_{L^4(\Omega)}^4 \\ & + \frac{1-\epsilon}{2} \|\varphi_h + \bar{\varphi}_0\|_{L^2(\Omega)}^2 - 2 (\nabla \varphi_h^{m-1}, \nabla \varphi_h), \end{aligned}$$

where

$$\|v_h\|_{-1,h} = (\nabla T_h v_h, \nabla T_h v_h)^{1/2} = (v_h, T_h v_h)^{1/2} = (T_h v_h, v_h)^{1/2},$$

with  $T_h : \mathring{Z}_h \rightarrow \mathring{Z}_h$  defined as: given  $\zeta_h \in \mathring{Z}_h$ , find  $T_h \zeta_h \in \mathring{Z}_h$ :

$$(\nabla T_h \zeta_h, \nabla \chi_h) = (\zeta_h, \chi_h) \quad \forall \chi_h \in \mathring{Z}_h.$$

# Unconditional Unique Solvability

## Theorem

*The fully discrete  $C^0$ -IP scheme is uniquely solvable for any mesh parameters:  $\tau$  and  $h$  and for any  $\varepsilon < 1$ .*

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- ② Prove the one-to-one correspondence between solution of our scheme with the solution to the **zero mean formulation**
- ③ Existence of the unique solution to the **zero mean formulation** proved through existence of a minimizer for a functional  $G_h$ .



# Unconditional Stability

## Lemma (Discrete Energy Law)

Let  $(\varphi_h^m, \mu_h^m) \in Z_h \times V_h$  be a solution of the  $C^0$ -IP method. Then the following energy law holds for any  $h, \tau > 0$ :

$$\begin{aligned} F(\varphi_h^\ell) + \tau \sum_{m=1}^{\ell} \left\| \mathcal{M}^{1/2} \nabla \mu_h^m \right\|_{L^2(\Omega)}^2 \\ + \tau^2 \sum_{m=1}^{\ell} \left\{ \frac{(1-\epsilon)}{2} \|\delta_\tau \varphi_h^m\|_{L^2(\Omega)}^2 + \|\nabla \delta_\tau \varphi_h^m\|_{L^2(\Omega)}^2 \right. \\ \left. + \frac{1}{4} \|\delta_\tau (\varphi_h^m)^2\|_{L^2(\Omega)}^2 + \frac{1}{2} \|\varphi_h^m \delta_\tau \varphi_h^m\|_{L^2(\Omega)}^2 + \frac{1}{2} a_h^{IP}(\delta_\tau \varphi_h^m, \delta_\tau \varphi_h^m) \right\} \\ = F(\varphi_h^0), \quad 1 \leq \ell \leq M. \end{aligned}$$

# Uniform *a priori* estimates

## Lemma

Let  $(\varphi_h^m, \mu_h^m) \in Z_h \times V_h$  be the unique solution of  $C^0$ -IP scheme. Suppose that  $F(\varphi_h^0) \leq C$  independent of  $h$  and  $\epsilon < \frac{C_{coer} - 4}{C_{coer}} < 1$ . For any  $h, \tau > 0$ :

$$\max_{0 \leq m \leq M} \left[ \|\varphi_h^m\|_{L^4(\Omega)}^2 + \|\varphi_h^m\|_{L^2(\Omega)}^2 + \|\varphi_h^m\|_{2,h}^2 \right] \leq C$$

$$\max_{0 \leq m \leq M} \|\varphi_h^m\|_{H^1}^2 \leq C$$

$$\tau \sum_{m=1}^{\ell} \left\| \mathcal{M}^{1/2} \nabla \mu_h^m \right\|_{L^2(\Omega)}^2 \leq C$$

$$\tau^2 \sum_{m=1}^{\ell} \left\{ \|\nabla \delta_\tau \varphi_h^m\|_{L^2(\Omega)}^2 + \|(\varphi_h^m)^2 \delta_\tau (\varphi_h^m)^2\|_{L^2(\Omega)}^2 + \|\delta_\tau \varphi_h^m\|_{2,h}^2 \right\} \leq C$$

for some constant  $C$  that is independent of  $h, \tau$ , and  $T$ .

# Error Estimates

Assume additional regularities:

$$\begin{aligned}\varphi &\in L^\infty(0, T; H^3(\Omega)) \cap L^2(0, T; H^3(\Omega)), \\ \partial_t \varphi &\in L^2(0, T; H^3(\Omega)) \cap L^2(0, T; H_N^{-1}(\Omega)), \\ \partial_{tt} \varphi &\in L^2(0, T; L^2(\Omega)), \\ \mu &\in L^2(0, T; H^2(\Omega)), \\ \partial_t \mu &\in L^2(0, T; L^2(\Omega)).\end{aligned}$$

C<sup>0</sup>-IP Norm:

$$\|\xi_h\|_{2,h}^2 := \sum_{K \in \mathcal{T}_h} |\xi_h|_{H^2(K)}^2 + \sum_{e \in \mathcal{E}_h} \alpha \|h_e^{-\frac{1}{2}} [\mathbf{n}_e \cdot \nabla \xi_h]_e\|_{L^2(e)}^2.$$

Notation:

$$e^{\varphi,m} := \varphi^m - \varphi_h^m, \quad e^{\mu,m} := \mu^m - \mu_h^m.$$

Assumption:  $\mathcal{M} \equiv 1$ .

# Error Analysis: Error Equation

**Weak Form:**

$$\langle \partial_t \varphi^m, \nu \rangle + (\nabla \mu^m, \nabla \nu) = 0 \quad \forall \nu \in H^1(\Omega) \quad (1)$$

$$\left( (\varphi^m)^3 + (1 - \epsilon)\varphi^m, \psi \right) - 2(\nabla \varphi^m, \nabla \psi) + a(\varphi^m, \psi) - (\mu^m, \psi) = 0 \quad \forall \psi \in Z. \quad (2)$$

**Fully Discrete C<sup>0</sup>-IP Form:**

$$(\delta_\tau \varphi_h^m, \nu) + (\nabla \mu_h^m, \nabla \nu) = 0 \quad \forall \nu \in V_h \quad (3)$$

$$\left( (\varphi_h^m)^3 + (1 - \epsilon)\varphi_h^m, \psi \right) - 2(\nabla \varphi_h^{m-1}, \nabla \psi) + a_h^{IP}(\varphi_h^m, \psi) - (\mu_h^m, \psi) = 0, \quad \forall \psi \in Z_h \quad (4)$$

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- (1) and (3)  $\implies (\delta_\tau e^{\varphi,m}, \nu_h) + (\nabla e^{\mu,m}, \nabla \nu_h) = (\delta_\tau \varphi^m - \partial_t \varphi^m, \nu_h),$   
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 $\nu_h \in V_h \subset H^1(\Omega).$
- Error equation based on (2) and (4) is not well-defined since  $Z_h \not\subset Z!$

# Error Analysis: Error Equation

**Weak Form:**

$$\left( (\varphi^m)^3 + (1 - \epsilon)\varphi^m, \psi \right) - 2(\nabla\varphi^m, \nabla\psi) + a(\varphi^m, \psi) - (\mu^m, \psi) = 0 \quad \forall \psi \in Z.$$

**Fully Discrete C<sup>0</sup>-IP Form:**

$$\left( (\varphi_h^m)^3 + (1 - \epsilon)\varphi_h^m, \psi \right) - 2(\nabla\varphi_h^{m-1}, \nabla\psi) + a_h^{IP}(\varphi_h^m, \psi) - (\mu_h^m, \psi) = 0, \quad \forall \psi \in Z_h$$

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$$\begin{aligned} & \left( (\varphi^m)^3 - (\varphi_h^m)^3 + (1 - \epsilon)e^{\varphi,m}, \psi \right) - 2(\nabla e^{\varphi,m}, \nabla\psi) + a(\varphi^m, \psi) - a_h^{IP}(\varphi_h^m, \psi) \\ & - (e^{\mu,m}, \psi) = 0 \end{aligned}$$

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Problem:  $\psi \in Z$  is not in  $H^{2+1/2}$  locally!  $\psi \in Z_h$  is not in  $H^2$  globally!

Remedy: Lift  $\psi \in Z_h$  into a finite dimensional subspace of  $Z$ .

# Error Analysis

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- **Weak Form with correction term:** find  $(\varphi^m, \mu^m) \in Z \times H^1(\Omega)$ :

$$(\partial_t \varphi^m, \nu_h) + (\nabla \mu^m, \nabla \nu_h) = 0 \quad \forall \nu_h \in V_h,$$

$$\begin{aligned} a_h^{IP}(\varphi^m, \psi_h) &+ \left( (\varphi^m)^3 + (1 - \epsilon)\varphi^m, \psi_h \right) - 2(\nabla \varphi^m, \nabla \psi_h) - (\mu^m, \psi_h) \\ &= \mathcal{L}(\varphi^m, \mu^m; \psi_h - E_h \psi_h) \quad \forall \psi_h \in Z_h \end{aligned}$$

where

$$\begin{aligned} \mathcal{L}(\varphi^m, \mu^m; \psi_h - E_h \psi_h) &:= a_h^{IP}(\varphi^m, \psi_h - E_h \psi_h) - (\mu^m, \psi_h - E_h \psi_h) \\ &+ ((\varphi^m)^3 + (1 - \epsilon)\varphi^m, \psi_h - E_h \psi_h) - 2(\nabla \varphi^m, \nabla \psi_h - \nabla E_h \psi_h). \end{aligned}$$

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- Solutions to weak form are consistent since  $a_h^{IP}(\varphi, E_h \psi) = a(\varphi, E_h \psi)$  for all  $\psi \in Z_h$ .

# Error Equation

Subtracting fully discrete form from the weak form with the correction term gives:

$$(\delta_\tau e^{\varphi,m}, \nu_h) + (\nabla e^{\mu,m}, \nabla \nu_h) = (\delta_\tau \varphi^m - \partial_t \varphi^m, \nu_h), \quad (5)$$

$$\begin{aligned} a_h^{IP}(e^{\varphi,m}, \psi_h) + ((1-\epsilon)e^{\varphi,m}, \psi_h) - 2(\nabla e^{\varphi,m-1}, \nabla \psi_h) - (e^{\mu,m}, \psi_h) = \\ - \left( (\varphi^m)^3 - (\varphi_h^m)^3, \psi_h \right) - 2(\nabla \varphi^{m-1} - \nabla \varphi^m, \nabla \psi_h) + \mathcal{L}(\varphi^m, \mu_h^m; \psi_h - E_h \psi_h). \end{aligned} \quad (6)$$

**Notation:**

$$\begin{aligned} e^{\varphi,m} &= e_P^{\varphi,m} + e_h^{\varphi,m}, & e_P^{\varphi,m} &:= \varphi^m - P_h \varphi^m, & e_h^{\varphi,m} &:= P_h \varphi^m - \varphi_h^m, \\ e^{\mu,m} &= e_R^{\mu,m} + e_h^{\mu,m}, & e_R^{\mu,m} &:= \mu^m - R_h \mu^m, & e_h^{\mu,m} &:= R_h \mu^m - \mu_h^m. \end{aligned}$$

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Set  $\nu_h = e_h^{\mu,m}$  in (5) and  $\psi_h = \delta_\tau e_h^{\varphi,m}$  in (6).

# Error Equation

$$\begin{aligned} & \|\nabla e_h^{\mu,m}\|_{L^2}^2 + a_h^{IP}(e_h^{\varphi,m}, \delta_\tau e_h^{\varphi,m}) + ((1-\epsilon)e_h^{\varphi,m}, \delta_\tau e_h^{\varphi,m}) - 2(\nabla e_h^{\varphi,m-1}, \nabla \delta_\tau e_h^{\varphi,m}) \\ &= (\delta_\tau \varphi^m - \partial_t \varphi^m, e_h^{\mu,m}) - (\delta_\tau e_P^{\varphi,m}, e_h^{\mu,m}) + (e_R^{\mu,m}, \delta_\tau e_h^{\varphi,m}) \\ &+ 2(\nabla \varphi^m - \nabla \varphi^{m-1}, \nabla \delta_\tau e_h^{\varphi,m}) - ((\varphi^m)^3 - (\varphi_h^m)^3, \delta_\tau e_h^{\varphi,m}) \\ &+ 2(\nabla e_P^{\varphi,m-1}, \nabla \delta_\tau e_h^{\varphi,m}) + \mathcal{L}(\varphi^m, \mu_h^m; \psi_h - E_h \psi_h) \end{aligned}$$

# Error Analysis

## Lemma

Let  $(\varphi^m, \mu^m)$  be a weak solution with the additional regularities. Then for any  $h, \tau > 0$  and any  $0 \leq m \leq M$ , we have

$$\|\delta_\tau e_h^{\varphi, m}\|_{-1, h}^2 \leq 4 \|\nabla e_h^{\mu, m}\|_{L^2}^2 + \frac{Ch^2}{\tau} \int_{t_{m-1}}^{t_m} \|\partial_s \varphi(s)\|_{H^2}^2 ds + C\tau \int_{t_{m-1}}^{t_m} \|\partial_{ss} \varphi(s)\|_{H^1}^2 ds$$

where the constant  $C$  may depend upon a Poincaré constant but does not depend on  $h$  or  $\tau$ .

where  $\|v_h\|_{-1, h} = (\nabla T_h v_h, \nabla T_h v_h)^{1/2} = (v_h, T_h v_h)^{1/2} = (T_h v_h, v_h)^{1/2}$  and  $T_h$  is the discrete inverse Laplacian.

# Error Analysis: LHS

$$\begin{aligned} & \|\nabla e_h^{\mu,m}\|_{L^2}^2 + a_h^{IP}(e_h^{\varphi,m}, \delta_\tau e_h^{\varphi,m}) + ((1-\epsilon)e_h^{\varphi,m}, \delta_\tau e_h^{\varphi,m}) - 2(\nabla e_h^{\varphi,m-1}, \nabla \delta_\tau e_h^{\varphi,m}) \\ &= (\delta_\tau \varphi^m - \partial_t \varphi^m, e_h^{\mu,m}) - (\delta_\tau e_P^{\varphi,m}, e_h^{\mu,m}) + (e_R^{\mu,m}, \delta_\tau e_h^{\varphi,m}) \\ &+ 2(\nabla \varphi^m - \nabla \varphi^{m-1}, \nabla \delta_\tau e_h^{\varphi,m}) - ((\varphi^m)^3 - (\varphi_h^m)^3, \delta_\tau e_h^{\varphi,m}) \\ &+ 2(\nabla e_P^{\varphi,m-1}, \nabla \delta_\tau e_h^{\varphi,m}) + \mathcal{L}(\varphi^m, \mu_h^m; \psi_h - E_h \psi_h) \end{aligned}$$

Polarization Property:

$$\begin{aligned} a_h^{IP}(e_h^{\varphi,m}, \delta_\tau e_h^{\varphi,m}) &= \frac{1}{2} \delta_\tau a_h^{IP}(e_h^{\varphi,m}, e_h^{\varphi,m}) + \frac{\tau}{2} a_h^{IP}(\delta_\tau e_h^{\varphi,m}, \delta_\tau e_h^{\varphi,m}) \\ ((1-\epsilon)e_h^{\varphi,m}, \delta_\tau e_h^{\varphi,m}) &= \frac{(1-\epsilon)}{2} \delta_\tau \|e_h^{\varphi,m}\|_{L^2}^2 + \frac{(1-\epsilon)}{2} \|\delta_\tau e_h^{\varphi,m}\|_{L^2}^2 \\ -2(\nabla e_h^{\varphi,m-1}, \nabla \delta_\tau e_h^{\varphi,m}) &= \tau \|\nabla \delta_\tau e_h^{\varphi,m}\|_{L^2}^2 - \delta_\tau (\nabla e_h^{\varphi,m}, \nabla e_h^{\varphi,m}) \end{aligned}$$

## Error Analysis: First 3 RHS terms

$$(\delta_\tau \varphi^m - \partial_t \varphi^m, e_h^{\mu,m}) \leq C\tau \int_{t_{m-1}}^{t_m} \|\partial_{ss} \varphi(s)\|_{L^2}^2 ds + \frac{1}{12} \|\nabla e_h^{\mu,m}\|_{L^2}^2,$$

$$\begin{aligned} (\delta_\tau e_P^{\varphi,m}, e_h^{\mu,m}) &\leq C \|\delta_\tau e_P^{\varphi,m}\|_{L^2}^2 + \frac{1}{12} \|\nabla e_h^{\mu,m}\|_{L^2}^2 \\ &\leq \frac{C}{\tau} \int_{t_{m-1}}^{t_m} \|P_h \partial_s \varphi(s) - \partial_s \varphi(s)\|_{L^2}^2 ds + \frac{1}{12} \|\nabla e_h^{\mu,m}\|_{L^2}^2 \\ &\leq \frac{C}{\tau} \int_{t_{m-1}}^{t_m} \|\partial_s \varphi(s) - P_h \partial_s \varphi(s)\|_{2,h}^2 ds + \frac{1}{12} \|\nabla e_h^{\mu,m}\|_{L^2}^2, \end{aligned}$$

$$(e_R^{\mu,m}, \delta_\tau e_h^{\varphi,m}) \leq C \|\nabla e_R^{\mu,m}\|_{L^2}^2 + \frac{1}{36} \|\delta_\tau e_h^{\varphi,m}\|_{-1,h}^2,$$

## Error Analysis: next two RHS terms

$$\begin{aligned} 2(\nabla\varphi^m - \nabla\varphi^{m-1}, \nabla\delta_\tau e_h^{\varphi,m}) &= -2(\tau\Delta\delta_\tau\varphi^m, \delta_\tau e_h^{\varphi,m}) \\ &\leq 2\|\tau\nabla\Delta\delta_\tau\varphi^m\|_{L^2}\|\delta_\tau e_h^{\varphi,m}\|_{-1,h} \\ &\leq C\tau \int_{t_{m-1}}^{t_m} \|\partial_s\varphi(s)\|_{H^3}^2 ds + \frac{1}{36}\|\delta_\tau e_h^{\varphi,m}\|_{-1,h}^2. \end{aligned}$$

$$\begin{aligned} ((\varphi^m)^3 - (\varphi_h^m)^3, \delta_\tau e_h^{\varphi,m}) &\leq \left\| \nabla \left( (\varphi^m)^3 - (\varphi_h^m)^3 \right) \right\|_{L^2} \|\delta_\tau e_h^{\varphi,m}\|_{-1,h} \\ &= \left\| 3(\varphi^m)^2 \nabla\varphi^m - 3(\varphi_h^m)^2 \nabla\varphi_h^m \right\|_{L^2} \\ &\quad \times \|\delta_\tau e_h^{\varphi,m}\|_{-1,h} \\ &= 3 \left\| (\varphi^m + \varphi_h^m) \nabla\varphi^m e^{\varphi,m} + (\varphi_h^m)^2 \nabla e^{\varphi,m} \right\|_{L^2} \\ &\quad \times \|\delta_\tau e_h^{\varphi,m}\|_{-1,h} \end{aligned}$$

## Error Analysis: next two RHS terms

$$\begin{aligned} & \left( (\varphi^m)^3 - (\varphi_h^m)^3, \delta_\tau e_h^{\varphi,m} \right) \\ & \leq 3 \left( \|\varphi^m + \varphi_h^m\|_{L^6} \|\nabla \varphi^m\|_{L^6} \|e^{\varphi,m}\|_{L^6} + \|\varphi_h^m\|_{L^6}^2 \|\nabla e^{\varphi,m}\|_{L^6} \right) \\ & \quad \times \|\delta_\tau e_h^{\varphi,m}\|_{-1,h} \\ & \leq C \left( \|\nabla e_P^{\varphi,m}\|_{L^2} + \|\nabla e_h^{\varphi,m}\|_{L^2} + \|e_P^{\varphi,m}\|_{2,h} + \|e_h^{\varphi,m}\|_{2,h} \right) \\ & \quad \times \|\delta_\tau e_h^{\varphi,m}\|_{-1,h} \\ & \leq C \left( \|\nabla e_P^{\varphi,m}\|_{L^2} + \|\nabla e_h^{\varphi,m}\|_{L^2} + \|e_P^{\varphi,m}\|_{2,h} + \|e_h^{\varphi,m}\|_{2,h} \right) \\ & \quad \times \|\delta_\tau e_h^{\varphi,m}\|_{-1,h} \\ & \leq C \|e_P^{\varphi,m}\|_{2,h}^2 + C \|e_h^{\varphi,m}\|_{2,h}^2 + \frac{1}{36} \|\delta_\tau e_h^{\varphi,m}\|_{-1,h}^2. \end{aligned}$$

## Error Analysis: second last term

Using discrete product rule:

$$\begin{aligned} \left( a^{m-1}, \frac{b^m - b^{m-1}}{\tau} \right) &= \frac{1}{\tau} [(a^m, b^m) - (a^{m-1}, b^{m-1})] - \left( \frac{a^m - a^{m-1}}{\tau}, b^m \right) \\ &= \delta_\tau (a^m, b^m) - (\delta_\tau a^m, b^m), \end{aligned}$$

we have the following bound

$$\begin{aligned} 2 \left( \nabla e_P^{\varphi, m-1}, \nabla \delta_\tau e_h^{\varphi, m} \right) &= 2 \delta_\tau (\nabla e_P^{\varphi, m}, \nabla e_h^{\varphi, m}) - 2 (\nabla \delta_\tau e_P^{\varphi, m}, \nabla e_h^{\varphi, m}) \\ &\leq 2 \delta_\tau (\nabla e_P^{\varphi, m}, \nabla e_h^{\varphi, m}) + C \|\delta_\tau e_P^{\varphi, m}\|_{L^2}^2 + C \|e_h^{\varphi, m}\|_{2,h}^2 \\ &\leq 2 \delta_\tau (\nabla e_P^{\varphi, m}, \nabla e_h^{\varphi, m}) + \frac{C}{\tau} \int_{t_{m-1}}^{t_m} \|\partial_s \varphi(s) - P_h \partial_s \varphi(s)\|_{2,h}^2 ds \\ &\quad + C \|e_h^{\varphi, m}\|_{2,h}^2. \end{aligned}$$

# Error Analysis: last term

## Lemma

Suppose  $(\varphi^m, \mu^m)$  is a weak solution to the PFC equation, with the additional regularities. Then for any  $h, \tau > 0$  and any  $0 \leq m \leq M$  and any  $\beta > 0$ ,

$$\begin{aligned} & a_h^{IP} (\varphi^m, e_h^{\varphi,m} - E_h e_h^{\varphi,m}) + \left( (\varphi^m)^3 + (1 - \epsilon) \varphi^m, e_h^{\varphi,m} - E_h e_h^{\varphi,m} \right) \\ & - 2 (\nabla \varphi^m, \nabla (e_h^{\varphi,m} - E_h e_h^{\varphi,m})) - (\mu^m, e_h^{\varphi,m} - E_h e_h^{\varphi,m}) \leq C [Osc_j(\mu^m)]^2 + \\ & C \|e_P^{\varphi,m}\|_{2,h}^2 + \frac{C_{coer}}{4\beta} \|e_h^{\varphi,m}\|_{2,h}^2 \end{aligned}$$

and

$$\begin{aligned} & a_h^{IP} \left( \delta_\tau \varphi^m, e_h^{\varphi,m-1} - E_h e_h^{\varphi,m-1} \right) + \left( \delta_\tau \left( (\varphi^m)^3 + (1 - \epsilon) \varphi^m \right), e_h^{\varphi,m-1} - E_h e_h^{\varphi,m-1} \right) \\ & - 2 \left( \delta_\tau \nabla \varphi^m, \nabla (e_h^{\varphi,m-1} - E_h e_h^{\varphi,m-1}) \right) - \left( \delta_\tau \mu^m, e_h^{\varphi,m-1} - E_h e_h^{\varphi,m-1} \right) \\ & \leq C [Osc_j(\mu_t(t^*))]^2 + C \|e_P^{\varphi,m}\|_{2,h}^2 \end{aligned}$$

$$\begin{aligned}
& \|\nabla e_h^{\mu,m}\|_{L^2}^2 + \frac{1}{2} \delta_\tau a_h^{IP}(e_h^{\varphi,m}, e_h^{\varphi,m}) + \frac{\tau}{2} a_h^{IP}(\delta_\tau e_h^{\varphi,m}, \delta_\tau e_h^{\varphi,m}) \\
& + \frac{(1-\epsilon)}{2} \delta_\tau \|e_h^{\varphi,m}\|_{L^2}^2 + \frac{(1-\epsilon)\tau}{2} \|\delta_\tau e_h^{\varphi,m}\|_{L^2}^2 + \tau \|\nabla \delta_\tau e_h^{\varphi,m}\|_{L^2}^2 \\
& \leq \delta_\tau (\nabla e_h^{\varphi,m}, \nabla e_h^{\varphi,m}) + 2\delta_\tau (\nabla e_P^{\varphi,m}, \nabla e_h^{\varphi,m}) + \frac{6}{12} \|\nabla e_h^{\mu,m}\|_{L^2}^2 + C \|e_h^{\varphi,m}\|_{2,h}^2 \\
& + C \|e_h^{\varphi,m-1}\|_{2,h}^2 + C \|\nabla e_R^{\mu,m}\|_{L^2}^2 + C \|e_P^{\varphi,m}\|_{2,h}^2 + C [\text{Osc}_j(\mu_t(t^*))]^2 \\
& + C\tau \int_{t_{m-1}}^{t_m} \left[ \|\partial_s \varphi(s)\|_{H^3}^2 + \|\partial_{ss} \varphi(s)\|_{L^2}^2 \right] ds + \frac{C}{\tau} \int_{t_{m-1}}^{t_m} \|\partial_s \varphi(s) - P_h \partial_s \varphi(s)\|_{2,h}^2 ds \\
& + \delta_\tau a_h^{IP}(\varphi^m, e_h^{\varphi,m} - E_h e_h^{\varphi,m}) + \delta_\tau \left( (\varphi^m)^3 + (1-\epsilon)\varphi^m, e_h^{\varphi,m} - E_h e_h^{\varphi,m} \right) \\
& - 2\delta_\tau (\nabla \varphi^m, \nabla (e_h^{\varphi,m} - E_h e_h^{\varphi,m})) - \delta_\tau (\mu^m, e_h^{\varphi,m} - E_h e_h^{\varphi,m}).
\end{aligned}$$

Applying  $2\tau \sum_{m=1}^{\ell}$ , using the fact that  $e_h^{\varphi,0} = 0$  we obtain

$$\begin{aligned}
 & a_h^{IP} \left( e_h^{\varphi,\ell}, e_h^{\varphi,\ell} \right) + (1-\epsilon) \left\| e_h^{\varphi,\ell} \right\|_{L^2}^2 + \tau \sum_{m=1}^{\ell} \left\| \nabla e_h^{\mu,m} \right\|_{L^2}^2 \\
 & + \tau^2 \sum_{m=1}^{\ell} \left[ a_h^{IP} (\delta_\tau e_h^{\varphi,m}, \delta_\tau e_h^{\varphi,m}) + (1-\epsilon) \left\| \delta_\tau e_h^{\varphi,m} \right\|_{L^2}^2 + 2 \left\| \nabla \delta_\tau e_h^{\varphi,m} \right\|_{L^2}^2 \right] \\
 & \leq \frac{C_{coer}}{2\beta} \left\| e_h^{\varphi,\ell} \right\|_{2,h}^2 + \frac{8\beta}{C_{coer}} \left\| e_h^{\varphi,\ell} \right\|_{L^2}^2 + C \left\| e_P^{\varphi,\ell} \right\|_{2,h}^2 + C\tau \sum_{m=1}^{\ell} \left\| e_h^{\varphi,m} \right\|_{2,h}^2 \\
 & + C\tau \sum_{m=1}^{\ell} \left[ \left\| \nabla e_R^{\mu,m} \right\|_{L^2}^2 + \left\| e_P^{\varphi,m} \right\|_{2,h}^2 + [\text{Osc}_j(\mu_t(t^*))]^2 \right] \\
 & + C\tau^2 \int_{t_0}^{t_\ell} \left[ \left\| \partial_s \varphi(s) \right\|_{H^3}^2 + \left\| \partial_{ss} \varphi(s) \right\|_{L^2}^2 \right] ds + C \int_{t_0}^{t_\ell} \left\| \partial_s \varphi(s) - P_h \partial_s \varphi(s) \right\|_{2,h}^2 ds \\
 & + 2 \left[ C [\text{Osc}_j(\mu^\ell)]^2 + C \left\| e_P^{\varphi,\ell} \right\|_{2,h}^2 + \frac{C_{coer}}{4\beta} \left\| e_h^{\varphi,\ell} \right\|_{2,h}^2 \right],
 \end{aligned}$$

# Error Estimates: Main Result

## Theorem

Suppose  $(\varphi^m, \mu^m)$  is a weak solution to the weak form of the PFC equation, with the additional regularities. Then for any  $\tau, h > 0, \epsilon < \frac{C_{coer} - 16}{C_{coer}} < 1$  and any  $0 \leq \tau \leq M$ ,

$$\begin{aligned} & \|e_h^{\varphi,\ell}\|_{2,h}^2 + C \|e_h^{\varphi,\ell}\|_{L^2}^2 + C\tau \sum_{m=1}^{\ell} \|\nabla e_h^{\mu,m}\|_{L^2}^2 + \\ & C\tau^2 \sum_{m=1}^{\ell} \left[ \|\delta_{\tau} e_h^{\varphi,\ell}\|_{2,h}^2 + (1-\epsilon) \|\delta_{\tau} e_h^{\varphi,m}\|_{L^2}^2 + \|\nabla \delta_{\tau} e_h^{\varphi,m}\|_{L^2}^2 \right] \leq C^*(h^2 + \tau^2) \end{aligned}$$

where  $C^*$  may depend on the oscillations of  $\mu$  and  $\partial_t \mu$  and the final stopping time  $T$  but does not depend upon the spacial step size  $h$  or the time step size  $\tau$ .

# Our Numerical Scheme

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$$(\delta_\tau \varphi_h^m, \nu_h) + (\mathcal{M} \nabla \mu_h^m, \nabla \nu_h) = 0$$

$$\left( (\varphi_h^m)^3 + (1 - \epsilon) \varphi_h^m, \psi_h \right) + a_h^{IP} (\varphi_h^m, \psi_h) - 2 (\nabla \varphi_h^{m-1}, \nabla \psi_h) - (\mu_h^m, \psi_h) = 0,$$

where  $\varphi_h^0 := P_h \varphi_0$  and  $\mu_h^0 \in V_h$  is defined as  $\mu_h^0 := R_h \mu_0$ .

model parameters:  $\epsilon = 0.025$  and  $\mathcal{M} = 1$ , penalty parameter:  $\alpha = 20$ .

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- We use Newton iteration method to solve the discrete nonlinear system with tolerance of the Newton iteration is set to  $10^{-6}$ .  
The initial guess at each time step is taken as the numerical solution at the previous time level.  
One to three Newton's iterative steps are involved at each time step.

# Numerical Experiment I: Accuracy Test

Example (Hu, Wise, Wang, Lowengrub, 2009 )

$$\begin{aligned}\varphi_0(x, y) = & 0.07 - 0.02 \cos\left(\frac{2\pi(x - 12)}{32}\right) \sin\left(\frac{2\pi(y - 1)}{32}\right) \\ & + 0.02 \cos^2\left(\frac{\pi(x + 10)}{32}\right) \cos^2\left(\frac{\pi(y + 3)}{32}\right) \\ & - 0.01 \sin^2\left(\frac{4\pi x}{32}\right) \sin^2\left(\frac{4\pi(y - 6)}{32}\right)\end{aligned}$$

$$\Omega = (0, 32) \times (0, 32), \quad T = 10.$$

$$\mathcal{M} \equiv 1, \varepsilon = 0.025, \text{ and the penalty parameter } \alpha = 20.$$

# Numerical Experiment I: Accuracy Test

$$\|\xi_h\|_{2,h}^2 := \sum_{K \in \mathcal{T}_h} |\xi_h|_{H^2(K)}^2 + \sum_{e \in \mathcal{E}_h} \alpha \|h_e^{-\frac{1}{2}} [\mathbf{n}_e \cdot \nabla \xi_h]_e\|_{L^2(e)}^2.$$

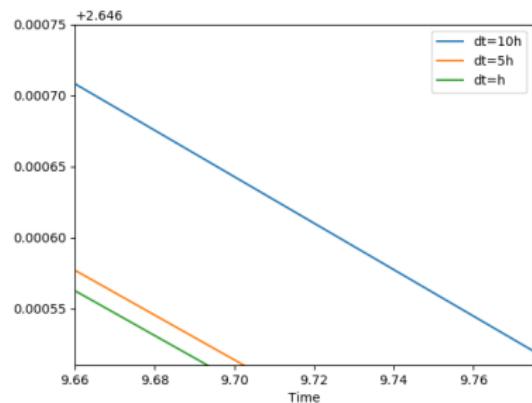
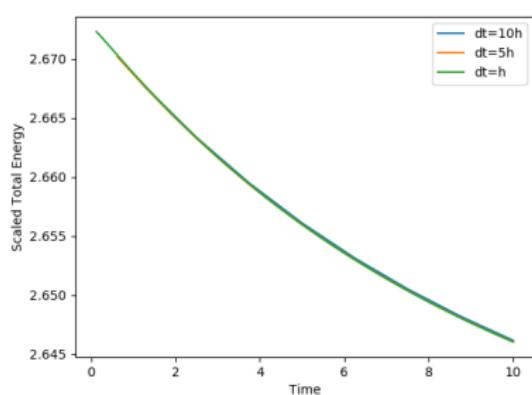
Mesh  $h = 32/512$  with  $\tau$  with  $\tau = 0.05h$  and  $T = 10$  as the ‘exact’ solution,  $\varphi_{exact}$ .

$$error_\varphi := \varphi_h - \varphi_{exact}$$

where  $\varphi_h$  indicates the solution on the mesh size  $h$ .

$h$	$\ error_\varphi\ _{2,h}$	rate	$\ error_\mu\ _{H^1}$	rate
$32/8$	0.08412	N/A	0.00522	N/A
$32/16$	0.05896	0.71329	0.00242	1.07627
$32/32$	0.03466	0.85058	0.00157	0.76970
$32/64$	0.01568	1.10514	0.00103	0.76082
$32/128$	0.00601	1.30482	0.00041	1.25840
$32/256$	0.00255	1.17707	0.00016	1.27362

# Numerical Experiment I: Unconditional Stability



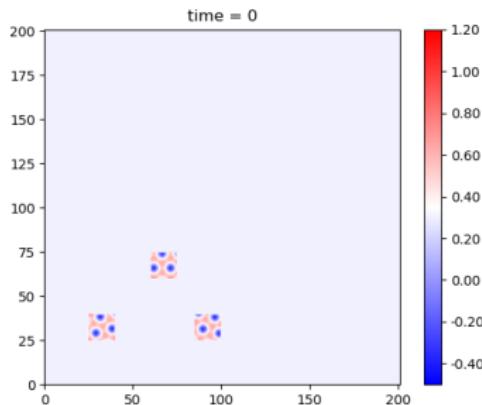
**Figure:** Unconditional stability demonstrated through the time evolution of the scaled total energy  $F/32^2$  for time step sizes  $dt = 10h, 5h, h$  with the spacial step size  $h = 32/256$ .

# Numerical Experiment II: Crystal growth

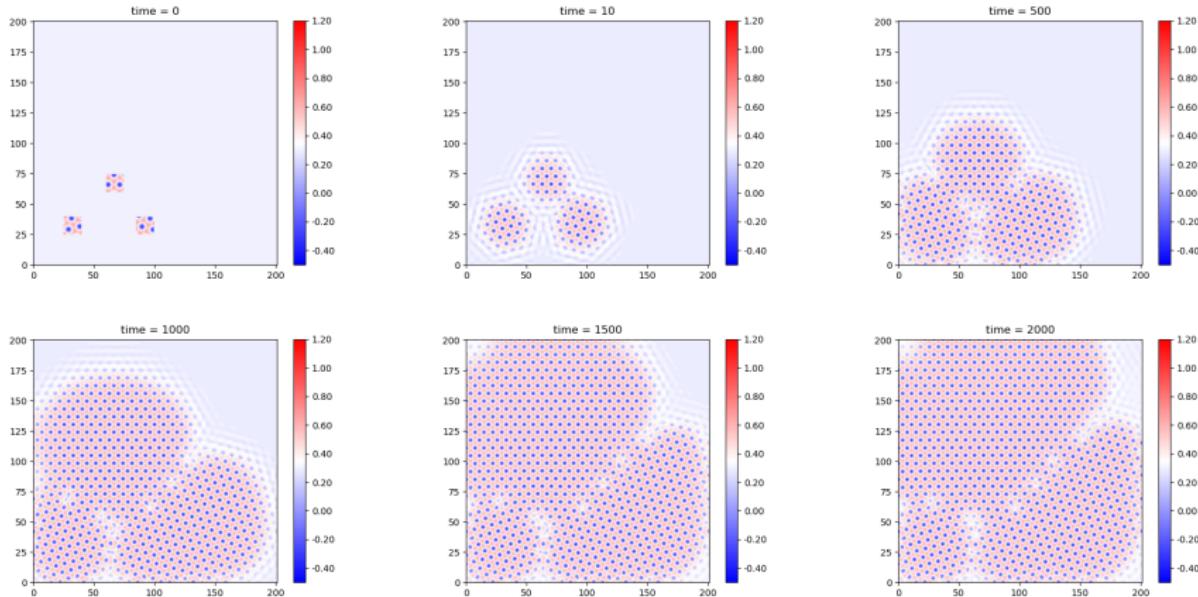
Example (Gomez, Nogueira, 2012)

$$\varphi_0(x, y) = \bar{\varphi} + C \left[ \cos\left(\frac{q}{\sqrt{3}}y\right) \cos(qx) - 0.5 \cos\left(\frac{2q}{\sqrt{3}}y\right) \right]$$

where  $\bar{\varphi} = 0.285$ ,  $C = 0.466$ ,  $q = 0.66$ ,  $\Omega = (0, 201) \times (0, 201)$ ,  $h = 201/402$ .

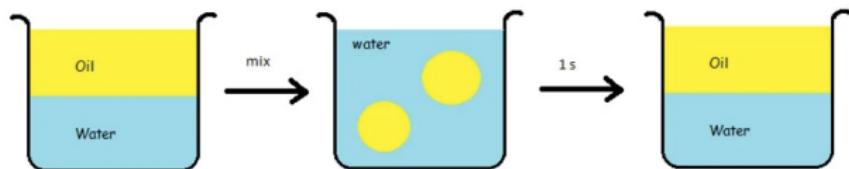


# Motion of liquid-crystal interfaces and grain boundaries

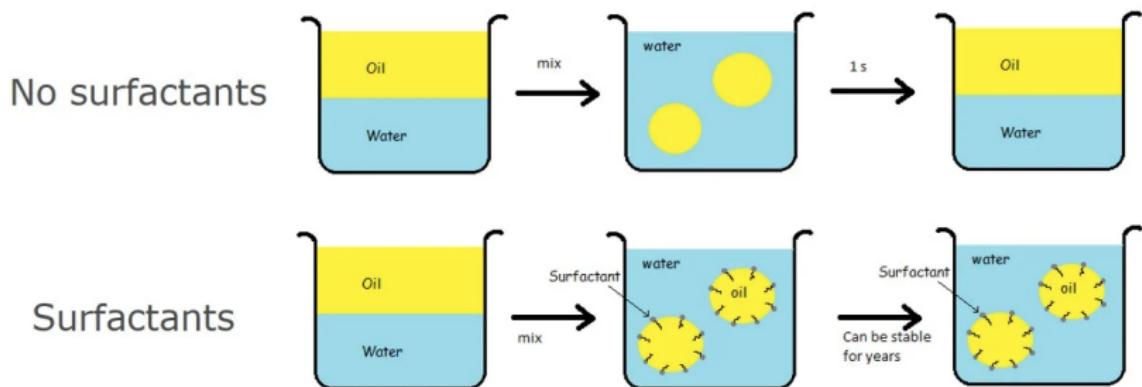


# Phase Field Model for Microemulsions: Motivation

No surfactants



# Phase Field Model for Microemulsions: Motivation



[Picture Courtesy: Biolin Scientific]

# Motivation: Microemulsion systems as skin drug delivery systems



[Source: Walgreens]

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- **Issue:** Topical cream formulations contain petrochemical ingredients (such as petrolatum, silicones).
- **Goal:** Develop skin drug delivery systems that contain natural and renewable sourced alternatives to these ingredients without compromising on the functionality.
- **Approach:** Provide computational tools to predict and assess the properties of novel and more sustainable alternatives to the toxic ingredients.

[Source: The Nabi Laboratory of Bioengineered Therapeutics, UTEP]

[Source: Walgreens]

# Mathematical Model for Microemulsions

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- $\varphi$  : the scalar order parameter indicating the local difference between oil and water concentrations.
- $\varphi = -1$  (water phase),  $\varphi = 1$  (oil phase) and  $\varphi = 0$  (microemulsions phase)
- $E(\varphi)$  : (Ginzburg-Landau free energy)

$$\underbrace{\int_{\Omega} \left\{ \frac{(\varphi^2 - a_0)}{2} |\nabla \varphi|^2 + \frac{\lambda}{2} (\Delta \varphi)^2 \right\} dx}_{\text{tendency to mix}} + \underbrace{\frac{\beta}{2} \int_{\Omega} (\varphi + 1)^2 (\varphi^2 + 0.5) (\varphi - 1)^2 dx}_{\text{tendency to separate}}$$

[Gompper et. al. 90]

# Mathematical Model for Microemulsions

Conservation Law:

$$\partial_t \varphi + \nabla \cdot \mathbf{j} = 0$$

- $\varphi$  : the scalar order parameter
- $\mathbf{j} = -\mathcal{M}\nabla\mu$ : mass flux
- $\mathcal{M}$ : mobility coefficient,  $\mu = \delta_\varphi E$ : chemical potential

$$E(\varphi) = \int_{\Omega} \left\{ \frac{(\varphi^2 - a_0)}{2} |\nabla \varphi|^2 + \frac{\lambda}{2} (\Delta \varphi)^2 + \frac{\beta}{2} (\varphi + 1)^2 (\varphi^2 + 0.5)(\varphi - 1)^2 \right\} dx$$

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$$\partial_t \varphi - \mathcal{M} \Delta \left( 3\beta(\varphi^5 - \varphi^3) + \varphi |\nabla \varphi|^2 - \nabla \cdot ((\varphi^2 - 4) \nabla \varphi) + \lambda \Delta^2 \varphi \right) = 0.$$

# Mathematical Model for Microemulsions

Compliment

$$\frac{\partial \varphi}{\partial t} - \nabla \cdot (\mathcal{M} \nabla \mu) = 0, \quad \text{in } \Omega^T := \Omega \times (0, T),$$

$$3\beta(\varphi^5 - \varphi^3) + \varphi |\nabla \varphi|^2 - \nabla \cdot ((\varphi^2 - a_0) \nabla \varphi) + \lambda \Delta^2 \varphi - \mu = 0, \quad \text{in } \Omega^T$$

with natural boundary conditions

$$\partial_n \varphi = \lambda \partial_n \Delta \varphi = \partial_n \mu = 0, \quad \text{on } \partial \Omega^T$$

and the initial value:

$$\varphi(0) = \varphi_0.$$

Notation:

- $H^s(\Omega)$  denote the Sobolev spaces of order  $s \geq 1$ ,
- $Z := \{z \in H^2(\Omega) \mid \partial_n z = 0 \text{ on } \partial \Omega\}$ . [Pawlow et al., 2011]

# Numerical Schemes: Existing Literature

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quasi-optimal error estimates without any discrete energy law

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- Hoppe/Linsemann 2019: Fully implicit backward Euler and  $C^0$ -IP Method quasi-optimal error estimates without any discrete energy law
- Diegel/Sharma 2022: Closely related literature is the  $C^0$ -IP framework developed for the Phase Field Crystal Equation based on Eyre's convex splitting scheme.

# Fully Discrete Scheme

Given  $\varphi_h^{m-1} \in Z_h$ , find  $(\varphi_h^m, \mu_h^m) \in Z_h \times V_h$  which satisfies

$$(\delta_\tau \varphi_h^m, \nu_h) + (M \nabla \mu_h^m, \nabla \nu_h) = 0, \quad \forall \nu_h \in V_h$$

$$\begin{aligned} & 3\beta ((\varphi_h^m)^5 - (\varphi_h^{m-1})^3, \psi_h) + ((\varphi_h^m)^2 \nabla \varphi_h^m, \nabla \psi_h) + (\varphi_h^m |\nabla \varphi_h^{m-1}|^2, \psi_h) \\ & - a_0 (\nabla \varphi_h^{m-1}, \nabla \psi_h) + \lambda a_h^{IP} (\varphi_h^m, \psi_h) - (\mu_h^m, \psi_h) = 0 \quad \forall \psi_h \in Z_h \end{aligned}$$

with initial data taken to be  $\varphi_h^0 := P_h \varphi_0$ .

# Existence of a solution

## Theorem

Let  $\lambda \geq \frac{3\beta|\varphi_0|^4 C_{P,1}}{2C_{coer}}$ , where  $C_{P,1}$  depends upon a Poincarè constant but does not depend upon  $h$  or  $\tau$ . Then, there exists a solution  $(\varphi_h^m, \mu_h^m) \in Z_h \times V_h$  to the scheme.

# Unconditional Energy Stability

$$F(\varphi) := \frac{\beta}{2} \|\varphi\|_{L^6}^6 - \frac{3\beta}{4} \|\varphi\|_{L^4}^4 + \frac{\beta|\Omega|}{4} + \frac{1}{2} \|\varphi \nabla \varphi\|_{L^2}^2 - \frac{a_0}{2} \|\nabla \varphi\|_{L^2}^2 + \frac{\lambda}{2} a_h^{IP}(\varphi, \varphi).$$

## Theorem (Discrete Energy Law)

Let  $(\varphi_h^m, \mu_h^m) \in Z_h \times V_h$  be a solution. Then the following energy law holds for any  $h, \tau > 0$ :

$$F(\varphi_h^\ell) + \tau \sum_{m=1}^{\ell} \left\| \sqrt{\mathcal{M}} \nabla \mu_h^m \right\|_{L^2}^2 \leq F(\varphi_h^0),$$

for all  $1 \leq \ell \leq M$ .

# Uniform *A Priori* Estimates

## Theorem

Let  $(\varphi_h^m, \mu_h^m) \in Z_h \times V_h$  be the  $C^0$  IP approximation. Suppose that  $F(\varphi_h^0) \leq C$  independent of  $h$  and that  $\lambda > \max \left\{ \frac{3\beta|\varphi_0|^4 C_{P,1}}{2C_{coer}}, \frac{a_0 C_{P,2}}{C_{coer}} \right\} > 0$  where  $C_{P,1}, C_{P,2}$  are Poincarè constants and do not depend on  $h$  or  $\tau$ . Then the following estimates hold for any  $\tau, h > 0$ :

$$\begin{aligned} \max_{0 \leq m \leq M} \|\varphi_h^m\|_{2,h}^2 &\leq C \\ \max_{0 \leq m \leq M} \left[ \|\varphi_h^m\|_{L^2}^2 + \|\nabla \varphi_h^m\|_{L^2}^2 + \|\varphi_h^m \nabla \varphi_h^m\|_{L^2}^2 + \|\varphi_h^m\|_{L^\infty}^2 \right] &\leq C^* \\ \tau \sum_{m=1}^{\ell} \left\| \sqrt{\mathcal{M}} \nabla \mu_h^m \right\|_{L^2}^2 &\leq C \end{aligned}$$

for some constants  $C^*, C$  that is independent of  $h, \tau$ , and  $T$ .

# Unique Solvability

## Theorem

Let  $\varphi_h^{m-1} \in Z_h$  be given and

$$\lambda > \max \left\{ \frac{3\beta |\bar{\varphi}_0|^4 C_{P,1}}{2C_{coer}}, \frac{a_0 C_{P,2}}{C_{coer}}, \frac{C^* C_{P,3}}{2C_{coer}} \right\} > 0,$$

where  $C^*$  is the constant from uniform a priori bounds and  $C_{P,1}, C_{P,2}, C_{P,3}$  are all Poincarè constants and do not depend on  $h$  or  $\tau$ .

The solution to the fully discrete scheme is unique for all  $h, \tau > 0$ .

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# Numerical Experiment I: Accuracy Test

## Example

$$\varphi_0(x, y) = 0.3 \cos(3x) + 0.5 \cos(y)$$

$$\Omega = [0, 2\pi]^2, T = 0.4.$$

$$\mathcal{M} = 10^{-3}, \lambda = 1$$

# Numerical Experiment I: First Order Convergence

- $\varphi_{256}$ : “exact” solution

$N$	$\ \varphi_{256} - \varphi_N\ _{2,h}$	rate	$\ \varphi_{256} - \varphi_N\ _{L^2}$	rate
8	6.1911	—	0.2625	—
16	1.9293	1.6045	0.0624	2.1039
32	0.5601	1.7221	0.0151	2.0685
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**Table:** Errors and convergence rates of the C<sup>0</sup>-IP method with  $\mathcal{M} = 10^{-3}$ ,  $\lambda = 1$ ,  $h = 2\sqrt{2}\pi/N$ ,  $\tau = 0.05/N$ .

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# Numerical Experiment II: Energy Dissipation

## Example

$$\varphi_0(x, y) = 0.3 \cos(3x) + 0.5 \cos(y)$$

$$\Omega := [0, 10]^2, \quad T = 5.$$

$\mathcal{M} = 10^{-3}$ ,  $\lambda = 1$ , penalty parameter  $\alpha = 8$

$$F(\varphi) = \frac{\beta}{2} \|\varphi\|_{L^6}^6 - \frac{3\beta}{4} \|\varphi\|_{L^4}^4 + \frac{\beta|\Omega|}{4} + \frac{1}{2} \|\varphi \nabla \varphi\|_{L^2}^2 - \frac{a_0}{2} \|\nabla \varphi\|_{L^2}^2 + \frac{\lambda}{2} a_h^{IP}(\varphi, \varphi).$$

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Track the scaled energy  $F(\varphi) - \frac{\beta|\Omega|}{4}$  for time step sizes  $\tau = 0.5, 0.25, 0.0125$ , and  $0.0625$ .

# Numerical Experiment II: Energy Dissipation

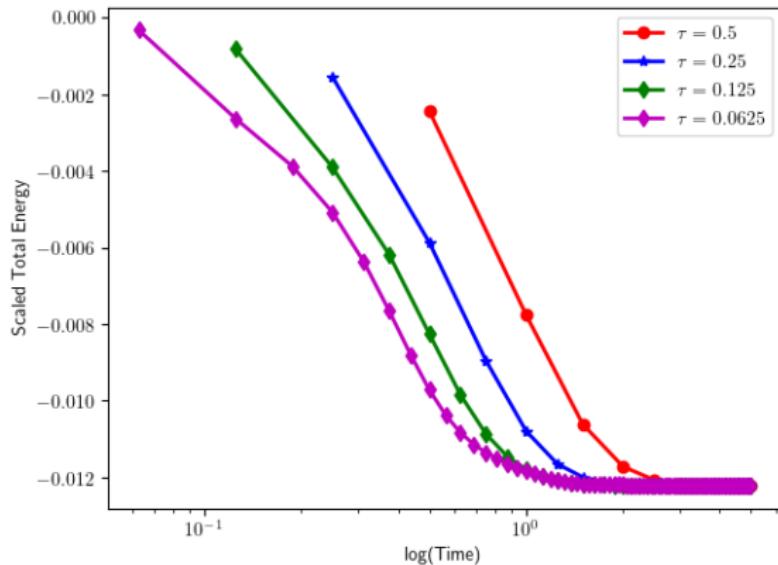
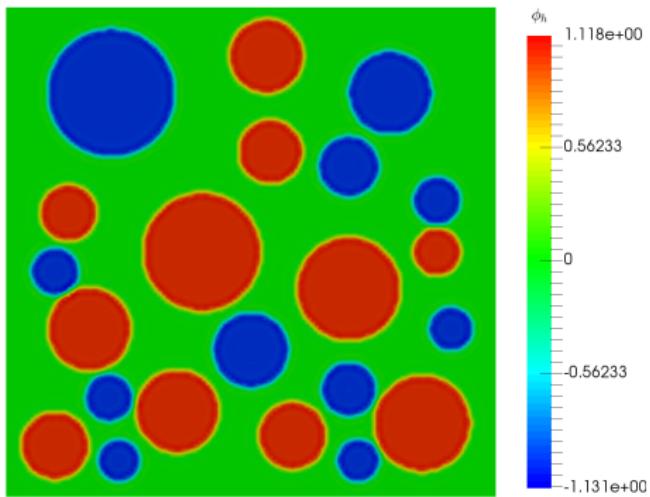


Figure: The time evolution of the scaled total energy  $F(\varphi) - \frac{\beta|\Omega|}{4}$ ,  $h = 10\sqrt{2}/128$

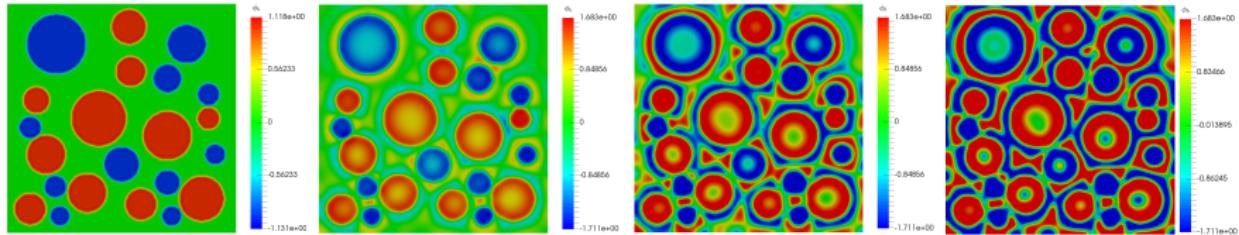
# Numerical Experiment III: Microemulsions Simulation

## Example

$$\Omega = (-5, 5)^2, T = 0.1, \tau = 1.1 \times 10^{-4}, h = 10\sqrt{2}/128, \mathcal{M} = 10, \lambda = 10^{-2}.$$



# Evolution of the profile



**Figure:** Profiles at  $m = 0, 1, 2$  and  $3$ .

# Evolution of the profile

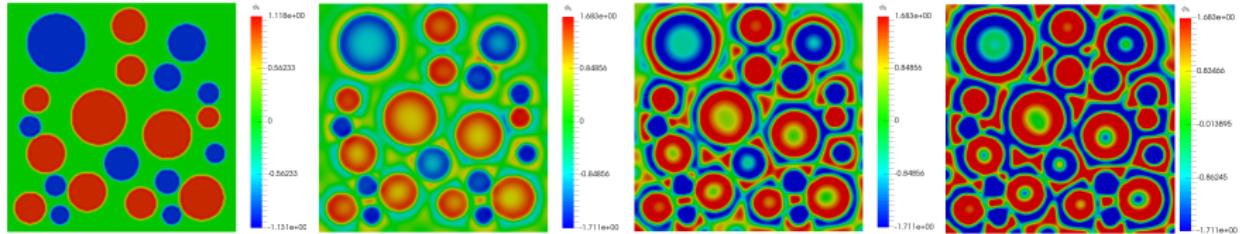


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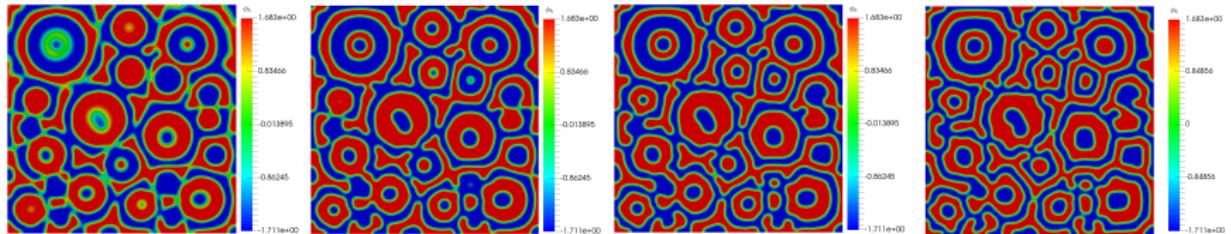


Figure: Profiles at  $m = 4, 11, 20$  and  $25$ .

# Evolution of the profile

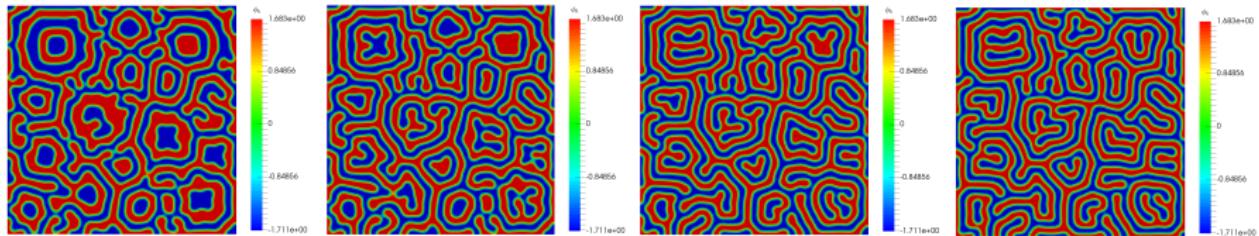


Figure: Profiles at  $m = 37, 74, 158$  and  $511$ .

# Evolution of the profile

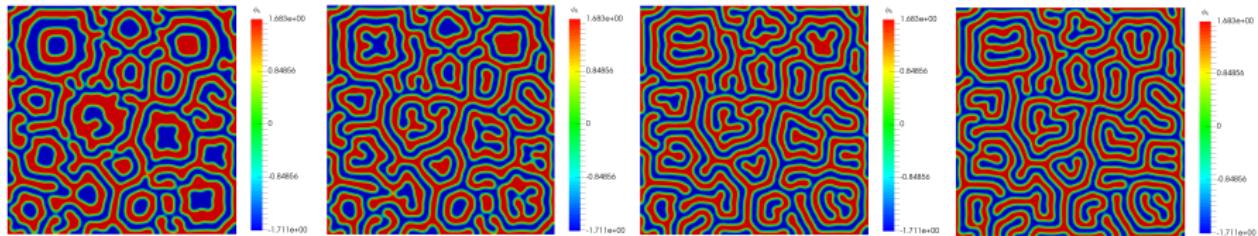


Figure: Profiles at  $m = 37, 74, 158$  and  $511$ .

- Different temporal scales capture different stages of phase-field evolution.

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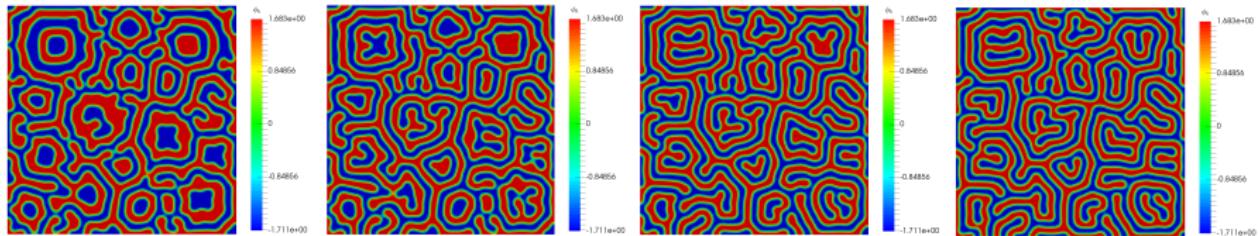


Figure: Profiles at  $m = 37, 74, 158$  and  $511$ .

- Different temporal scales capture different stages of phase-field evolution.
- Optimal time steps for each stage can differ by several orders of magnitude.

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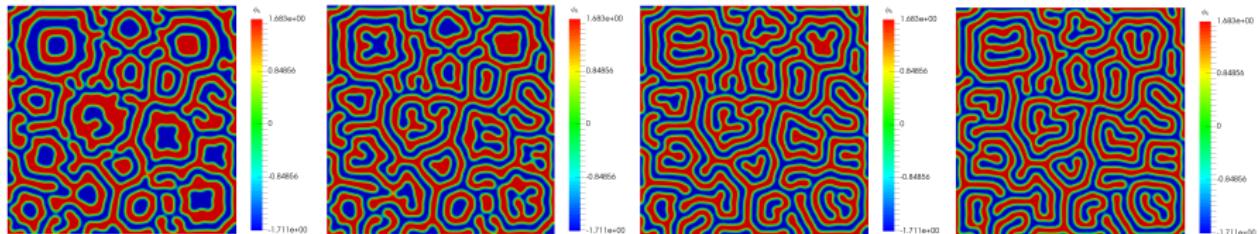


Figure: Profiles at  $m = 37, 74, 158$  and  $511$ .

- Different temporal scales capture different stages of phase-field evolution.
- Optimal time steps for each stage can differ by several orders of magnitude.
- Time-step adaptivity is crucial for accuracy and efficiency of the numerical scheme.

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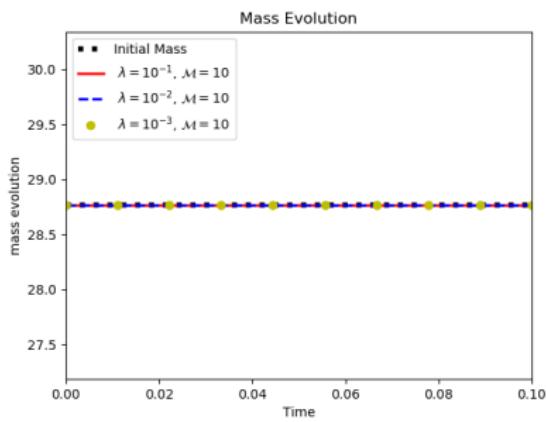
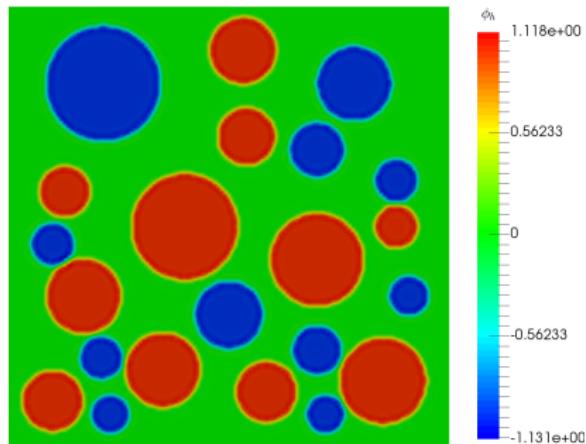
Thank you for your attention!

# Numerical Experiment: Discrete Mass Conservation

Initial Conditions:

## Example

$$\Omega = (-5, 5)^2, T = 0.1, \tau = 1.1 \times 10^{-4}, h = 10\sqrt{2}/128, \mathcal{M} = 10, \lambda = 10^{-1}, 10^{-2}, 10^{-3}.$$



# Effect of the decreasing $\lambda$

Theory suggests that increasing  $\lambda$  guarantees the existence and stability of the solution.

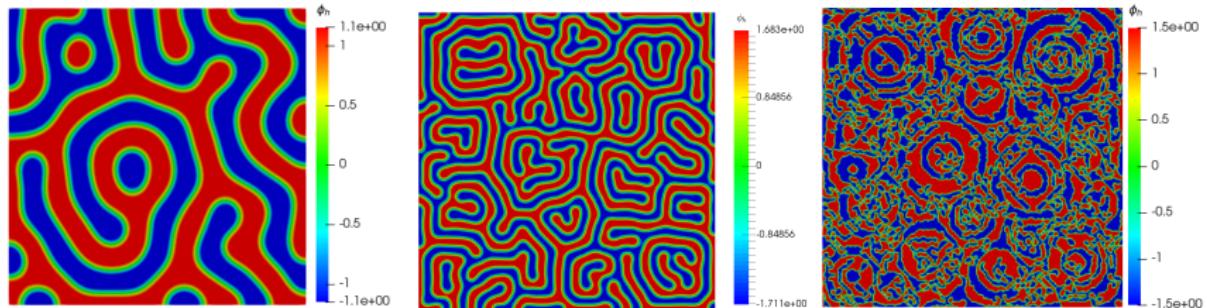


Figure:  $\lambda = 10^{-1}, 10^{-2}, 10^{-3}$

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