

Matrix-valued finite elements

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In collaboration with

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Thanks:



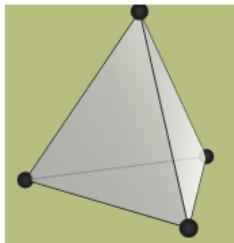
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Outline

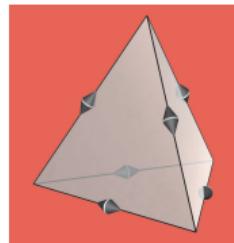
- 1 Introduction
- 2 Mass-Conserving Stress-yielding (MCS) method
- 3 Viscous stress elements
- 4 Other matrix finite elements
- 5 A unifying 2-complex of Sobolev spaces
- 6 Regular decompositions
- 7 A complementary diagram

Lowest order scalar (\mathbb{R}) and vector (\mathbb{V}) valued finite elements

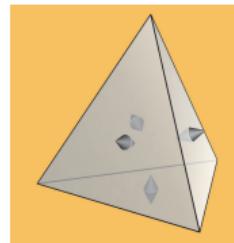
(\mathbb{R} -valued)



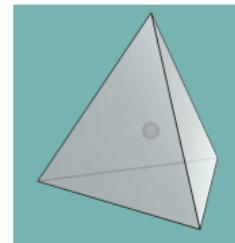
(\mathbb{V} -valued)



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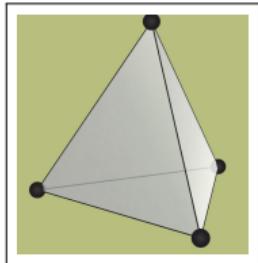
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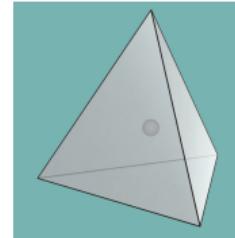
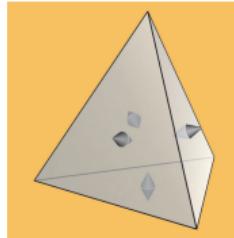
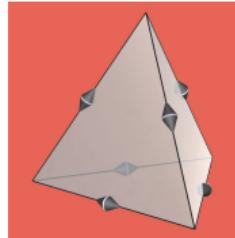
(\mathbb{V} -valued)

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(\mathbb{R} -valued)



grad, continuous



Lagrange finite element space

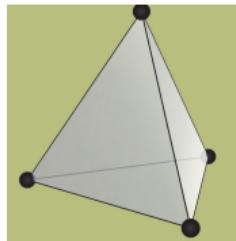
[Courant 1943]

Natural operator: grad

Continuous across elements

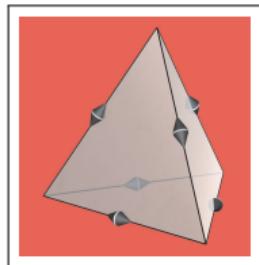
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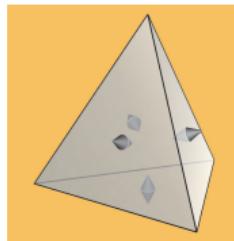
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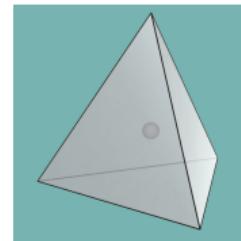


curl, “ t -continuous”

(\mathbb{V} -valued)



(\mathbb{R} -valued)



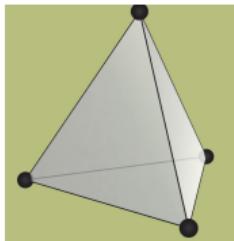
Nédélec finite element space [Nédélec 1980]

Natural operator: curl

Inter-element continuity: t -component

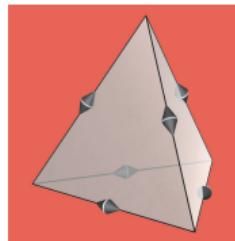
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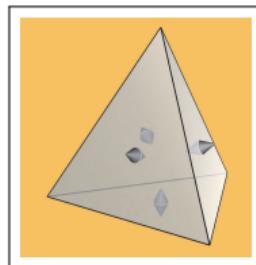
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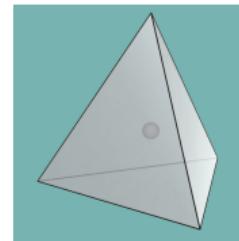
curl, " t -continuous"

(\mathbb{V} -valued)



div, " n -continuous"

(\mathbb{R} -valued)



Raviart-Thomas finite element space

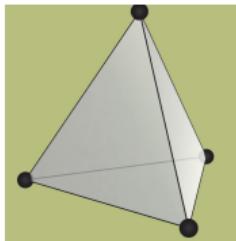
[Raviart+Thomas 1977]

Natural operator: div

Inter-element continuity: n -component

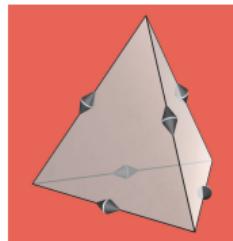
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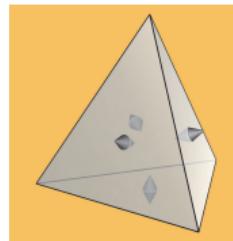
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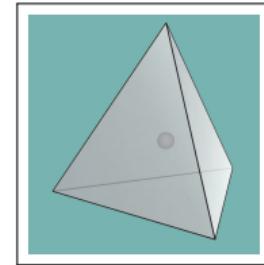
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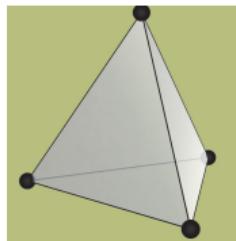


Discontinuous Galerkin finite element space

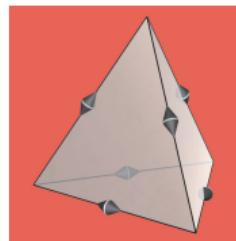
Inter-element continuity: None

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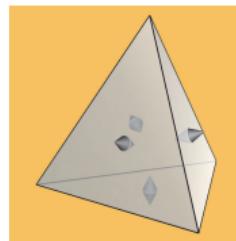
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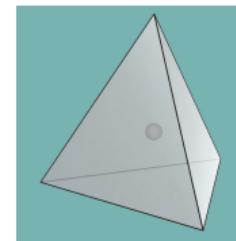
(\mathbb{V} -valued)



(\mathbb{V} -valued)



(\mathbb{R} -valued)



$\xrightarrow{\text{grad}}$

$\xrightarrow{\text{curl}}$

$\xrightarrow{\text{div}}$

grad, continuous

curl, “ t -continuous”

div, “ n -continuous”

These finite element spaces are now collectively understood using the de Rham complex.

Where do matrix-valued finite elements fit? And second-order differential operators?

Next example: curl div operator and “ nt -continuous” matrix elements for viscous stresses.

Stokes system

Textbook version:

$$\begin{aligned}-\nu \Delta u + \nabla p &= f && \text{in } \Omega \\ \operatorname{div} u &= 0 && \text{in } \Omega \\ u &= 0 && \text{on } \partial\Omega\end{aligned}$$

Physical stress-based version:

$$\begin{aligned}\frac{1}{2\nu}\sigma - \varepsilon(u) &= 0 && \text{in } \Omega \\ \operatorname{div} \sigma - \nabla p &= -f && \text{in } \Omega \\ \operatorname{div} u &= 0 && \text{in } \Omega \\ u &= 0 && \text{on } \partial\Omega\end{aligned}$$

Notation:

- $\nu(x)$ = kinematic viscosity
- $f(x)$ = body force
- $u(x)$ = fluid velocity
- $p(x)$ = kinematic pressure
- $\sigma(x)$ = viscous stress
- $\varepsilon(u) = \frac{1}{2}(\nabla u + (\nabla u)') = \operatorname{sym}(\nabla u)$.

What topology on σ pairs well with $H(\operatorname{div})$ -topology on u ?

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Kelvin-Helmholtz instability simulation

[Schroeder+John+Lederer+Lehrenfeld+Lube+Schöberl 2019]

Different results from different discretizations

Source: P. Lederer's thesis

Mass conservation

Mass is neither created nor destroyed.

- Mass = $\int \rho$, density integrated.
- Incompressibility \implies the material derivative $\frac{D\rho}{Dt} = 0$.

Equivalent constraint on fluid velocity u :

- Continuity equation (mass conservation) $\implies \frac{D\rho}{Dt} + \rho \operatorname{div} u = 0$.
- Hence, mass conservation in incompressible fluids $\iff \boxed{\operatorname{div} u = 0}$.

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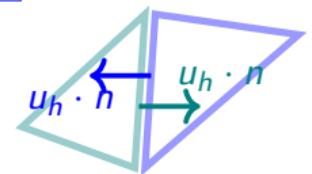
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A piecewise polynomial fluid velocity approximation u_h has $\boxed{\operatorname{div} u_h = 0}$ if

- its divergence is zero pointwise within elements,
- and $[\![u_h \cdot n]\!] = 0$, i.e., u_h has at least n -continuity.



How can methods gain exact mass conservation?

Discrete version of the incompressibility constraint in a space V_h with n -continuity:

$$(\operatorname{div} u_h, q_h) = 0 \quad \text{for all } q_h \text{ in the discrete pressure space } Q_h.$$

Hence u_h is exactly mass conserving if $\operatorname{div} V_h \subseteq Q_h$.

Two approaches to methods with mass conservation:

$H(\operatorname{div})$ -based

H^1 -based

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$H(\operatorname{div})$ -based

H^1 -based

$$H(\operatorname{div}) = \{u : u_i \in L_2, \sum_i \partial_i u_i \in L_2\}$$

$$H^1 \otimes \mathbb{V} = \{u : u_i \in L_2, \partial_j u_i \in L_2\}$$

(This is vector H^1 .)

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Two approaches to methods with mass conservation:

$H(\operatorname{div})$ -based

When $V_h \subset H(\operatorname{div})$:

- ⊕ Range of divergence is simple.
- ⊖ Discretizing viscous term is complex.

H^1 -based

When $V_h \subset H^1 \otimes \mathbb{V}$:

- ⊖ Range of divergence is complex.
- ⊕ Viscous term easily discretized.

How can methods gain exact mass conservation?

Discrete version of the incompressibility constraint in a space V_h with n -continuity:

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Two approaches to methods with mass conservation:

$H(\operatorname{div})$ -based

[Cockburn+Kanschat+Schötzau 2005]

⋮

[Cockburn+G+Nguyen+Peraire+Sayas 2011]

[Linke+Merdon 2016]

[G+Lederer+Schöberl 2019] (MCS method)

H^1 -based

[Scott+Vogelius 1985]

⋮

[Guzmán+Scott 2019]

[Neilan 2020]

[Ainsworth+Parker 2021]

Punchlines of Mass-Conserving Stress-yielding (MCS) method

- *Structure-preservation:*

Numerical solutions are exactly mass conserving.

Scheme is pressure robust.

- *Viscous stress σ approximation:* New matrix-valued finite elements with *continuous shear component* (or nt -component) used for approximating σ . ($\llbracket t \cdot \sigma n \rrbracket \equiv \llbracket \sigma_{nt} \rrbracket = 0$.)

- *Optimally convergent in* velocity (u), pressure (p), viscous stress (σ), and vorticity (ω).

- *More features:*

- ▶ No stabilization parameters.
- ▶ Easy to incorporate stress boundary conditions.
- ▶ No vertex unknowns, facet couplings only: permits easy hybridization.
- ▶ Element mappings (even curvilinear) are straightforward.
- ▶ Stable for ∇u -forms as well as the more physical $\varepsilon(u)$ -formulations.

Deriving the MCS formulation

Standard form:

$$-\nu \Delta u + \text{grad } p = 0$$

$$\text{div } u = 0$$

$$u = 0$$

With σ :

$$\frac{1}{2\nu} \sigma - \varepsilon(u) = 0 \quad \text{in } \Omega$$

$$\text{div } \sigma - \text{grad } p = -f \quad \text{in } \Omega$$

$$\text{div } u = 0 \quad \text{in } \Omega$$

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Notation:

$$\operatorname{dev} m := m - \frac{1}{3}(\operatorname{tr} m)\delta, \quad m \text{ in } \mathbb{R}^{3 \times 3}$$

$$\sigma = \operatorname{dev} \sigma, \text{ in } \mathbb{T} \quad \mathbb{T} := \operatorname{dev}(\mathbb{R}^{3 \times 3}) \equiv \{\text{traceless matrices}\}$$

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$$-\nu \Delta u + \operatorname{grad} p = 0$$

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With σ :

$$\frac{1}{2\nu} \operatorname{dev} \sigma - \nabla u + \omega = 0 \quad \text{in } \Omega$$

$$\operatorname{div} \sigma - \operatorname{grad} p = -f \quad \text{in } \Omega$$

$$\operatorname{div} u = 0 \quad \text{in } \Omega$$

$$\sigma - \sigma' = 0 \quad \text{in } \Omega$$

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$$\mathbb{T} := \operatorname{dev}(\mathbb{R}^{3 \times 3}) \equiv \{\text{traceless matrices}\}$$

$$\omega := \text{vorticity, in } \mathbb{K}$$

$$\mathbb{K} := \operatorname{skw}(\mathbb{R}^{3 \times 3}) \equiv \{\text{skew-symmetric matrices}\}$$

Deriving the MCS formulation

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$$\operatorname{div} u = 0$$

$$u = 0$$

With σ :

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$$u = 0 \quad \text{on } \partial\Omega$$

Test with $\tau : \Omega \rightarrow \mathbb{T}$, $v : \Omega \rightarrow \mathbb{V}$, $q : \Omega \rightarrow \mathbb{R}$, $\eta : \Omega \rightarrow \mathbb{K}$:

$$\left(\frac{1}{2\nu} \operatorname{dev} \sigma, \operatorname{dev} \tau \right) + (u, \operatorname{div} \tau) + (\omega, \tau) = 0$$

$$(\operatorname{div} \sigma, v) + (p, \operatorname{div} v) = -(f, v)$$

$$(\operatorname{div} u, q) = 0$$

$$(\sigma, \eta) = 0.$$

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$$(\operatorname{div} \sigma, v) + (p, \operatorname{div} v) = -(f, v)$$

$$(\operatorname{div} u, q) = 0$$

$$(\sigma, \eta) = 0.$$

Derivation is completed by replacing the integrals $(\operatorname{div} \sigma, v)$ by **functional actions** $\langle \operatorname{div} \sigma, v \rangle$.

The viscous stress space

MCS formulation: Find $\sigma \in \Sigma$, $u \in \mathring{H}(\text{div})$, $p \in L_{2,\mathbb{R}}$, $\omega \in L_2 \otimes \mathbb{K}$ satisfying

$$\begin{aligned} \left(\frac{1}{2\nu} \operatorname{dev} \sigma, \operatorname{dev} \tau \right) + \langle \operatorname{div} \tau, u \rangle + (\omega, \tau) &= 0, & \tau \in \Sigma, \\ \langle \operatorname{div} \sigma, v \rangle + (p, \operatorname{div} v) &= -\langle f, v \rangle, & v \in \mathring{H}(\text{div}), \\ (\operatorname{div} u, q) &= 0, & q \in L_{2,\mathbb{R}}, \\ (\sigma, \eta) &= 0, & \eta \in L_2 \otimes \mathbb{K}. \end{aligned}$$

Minimal requirement for any σ in Σ is $\operatorname{div} \sigma \in \mathring{H}(\text{div})^*$, the dual space of $\mathring{H}(\text{div})$.

Theorem: $\mathring{H}(\text{div})^* = H^{-1}(\text{curl})$.

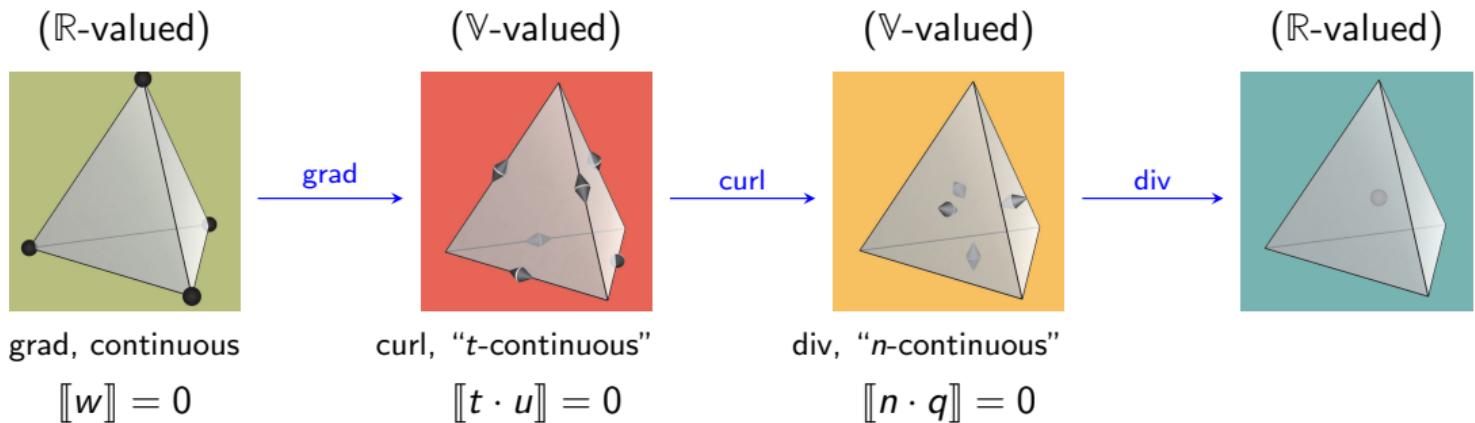
Space for viscous fluid stresses: $\Sigma = \{\tau \in L_2 \otimes \mathbb{T} : \operatorname{curl} \operatorname{div}(\tau) \in H^{-1}\}$.

Notation: ① $\mathbb{T} = \{\text{traceless matrices}\}$, ② $\mathbb{K} = \{\text{skew symmetric matrices}\}$, ③ $L_{2,\mathbb{R}} = \{p \in L_2 : (p, 1) = 0\}$,
④ $H^{-1}(\text{curl}) = \{\phi \in H^{-1} \otimes \mathbb{V} : \operatorname{curl} \phi \in H^{-1} \otimes \mathbb{V}\}$, ⑤ $\mathring{H}(\text{div}) = \{w \in H(\text{div}) : w \cdot n|_{\partial\Omega} = 0\}$.

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The *nt*-continuity



For the **curl div** operator, we make “*nt-continuous*” matrix elements for viscous stresses, i.e.,

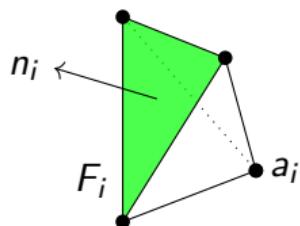
$$\llbracket t \cdot \sigma n \rrbracket = 0$$

for any tangent vector t and normal vector n on element boundaries.

Designing finite elements for viscous stresses

Lowest order case: *Is it possible to construct a constant matrix function whose nt -component vanishes on ∂T except on one face?*

Tetrahedron T

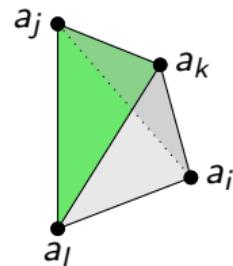


- Consider $\sigma = n_j \otimes (n_k \times n_l)$

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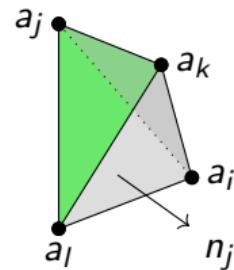


- Consider $\sigma = n_j \otimes (n_k \times n_l)$
- $\sigma n = 0$ on $F_k \cup F_l$

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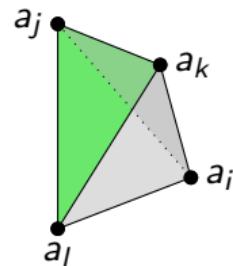


- Consider $\sigma = n_j \otimes (n_k \times n_l)$
- $\sigma n = 0$ on $F_k \cup F_l$ and σn_j is collinear with n_j , so $\sigma n_j \cdot t|_{F_j} = 0$.

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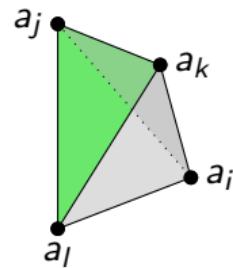


- Consider $\sigma = n_j \otimes (n_k \times n_l)$
- Thus $t \cdot \sigma n = 0$ on all faces except F_i .

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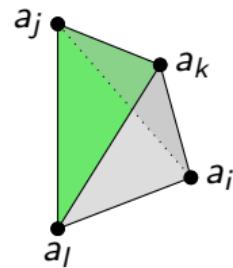


- Consider $\sigma = \text{dev}[n_j \otimes (n_k \times n_l)]$
- Thus $t \cdot \sigma n = 0$ on all faces except F_i .

Designing finite elements for viscous stresses

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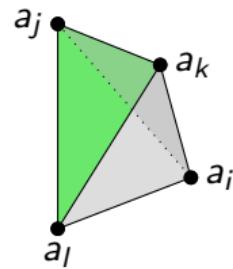


- Consider $\sigma = \text{dev}[n_j \otimes (n_k \times n_l)]$ and $\sigma = \text{dev}[n_l \otimes (n_j \times n_k)]$.
- Thus $t \cdot \sigma n = 0$ on all faces except F_i .

Designing finite elements for viscous stresses

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Tetrahedron T

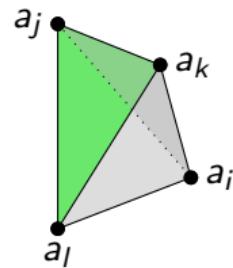


- Consider $\sigma = \text{dev}[n_j \otimes (n_k \times n_l)]$ and $\sigma = \text{dev}[n_l \otimes (n_j \times n_k)]$.
- Taking all faces, we get a basis for $\mathbb{T} = \{m \in \mathbb{R}^{3 \times 3} : \text{tr } m = 0\}$.

Designing finite elements for viscous stresses

Lowest order case: *Is it possible to construct a constant matrix function whose nt -component vanishes on ∂T except on one face?*

Tetrahedron T



- Consider $\sigma = \text{dev}[n_j \otimes (n_k \times n_l)]$ and $\sigma = \text{dev}[n_l \otimes (n_j \times n_k)]$.
- Taking all faces, we get a basis for $\mathbb{T} = \{m \in \mathbb{R}^{3 \times 3} : \text{tr } m = 0\}$.

More general shape functions: $\lambda_i^{\alpha_i} \lambda_j^{\alpha_j} \lambda_k^{\alpha_k} \text{dev}[\nabla \lambda_i \otimes (\nabla \lambda_j \times \nabla \lambda_k)]$.

Unisolvant finite element

Ciarlet-style finite element definition:

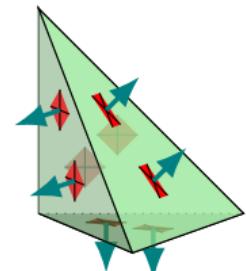
- *Geometry:* Tetrahedron T
- *Space:* $P_k(T) \otimes \mathbb{T} = \{m : T \rightarrow \mathbb{T} \mid m_{ij} \text{ is a polynomial of degree } \leq k\}$ for any $k \geq 0$.
- *Degrees of freedom:*

$$\int_F \sigma_{nt} \cdot r, \quad \text{tangential } r \text{ with } r_i \in P_k(F) \text{ on each face } F,$$
$$\int_T \sigma : \varsigma, \quad \varsigma \in P_{k-1}(T) \otimes \mathbb{T}, \quad \text{if } k \geq 1.$$

Theorem: These degrees of freedom are **unisolvant** for $P_k(T) \otimes \mathbb{T}$.

They generate global degrees of freedom suitable for enforcing
 nt -continuity.

[G+Lederer+Schöberl 2019]

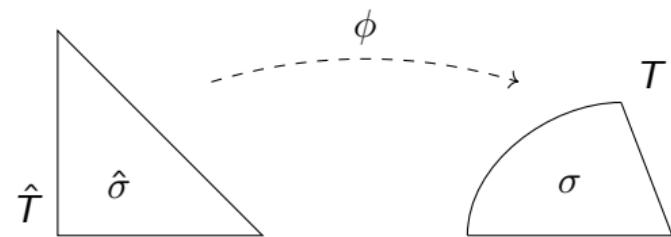


The $k = 0$ case

Mapping

- Let $T = \text{image of reference element } \hat{T}$ under (possibly curvilinear) map ϕ ,
- and define the **curl div pullback** by

$$\hat{\sigma} = \det(\nabla\phi) (\nabla\phi)^t (\sigma \circ \phi) (\nabla\phi)^{-t}.$$



Theorem: Let $F = \phi(\hat{F})$ for a facet \hat{F} of \hat{T} with normal \hat{n} and $\hat{t} \in \hat{n}^\perp$. Let n and t be the mapped normal and tangent on F . Then

- There is a nonzero c_F depending only on $\phi|_{\hat{F}}$ such that $\hat{t} \cdot \hat{\sigma} \hat{n} = c_F (t \cdot \sigma n) \circ \phi$,
- The mapped $\hat{\sigma}$ is traceless if and only if σ is.

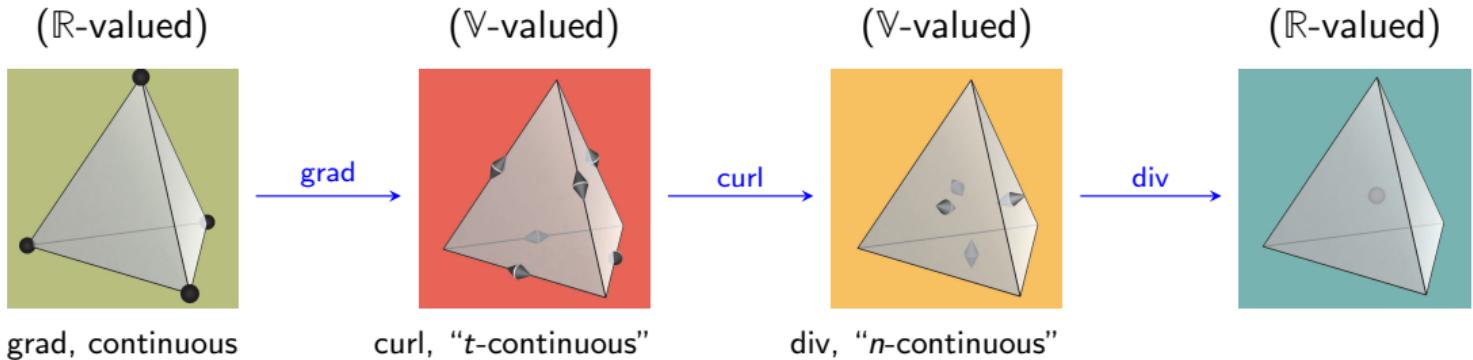
Other MCS variables have standard pullbacks (e.g., standard Piola for $H(\text{div})$ -velocity).

Thus all elements used in the MCS method have natural mappings.

Outline

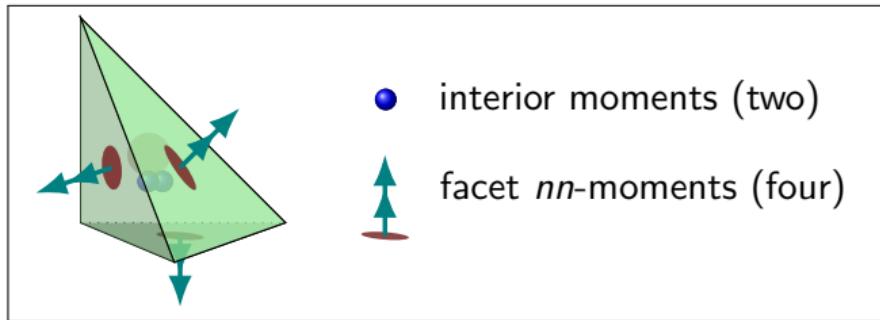
- 1 Introduction
- 2 Mass-Conserving Stress-yielding (MCS) method
- 3 Viscous stress elements
- 4 Other matrix finite elements
- 5 A unifying 2-complex of Sobolev spaces
- 6 Regular decompositions
- 7 A complementary diagram

Where do matrix-valued finite elements fit?



- Where does the nt -continuous \mathbb{T} -valued matrix element fit?
- How about matrix elements with tt -continuity?
- How about matrix elements with nn -continuity?

More matrix-valued elements



\mathbb{S} -valued HHJ element, aka TDNNS element

Natural operator: div div

Continuous nn -component

Coupling dofs:

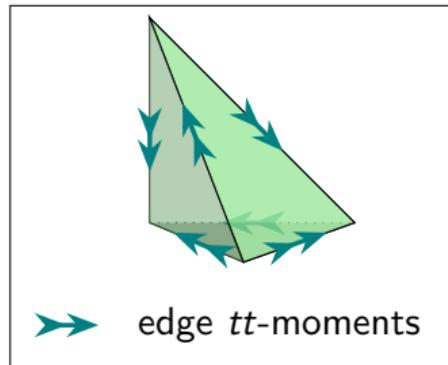
$$\sigma \mapsto \int_F n \cdot \sigma n$$

[Hellan 1967] [Herrmann 1967] [Johnson 1973]

[Comodi 1989] [Pechstein+Schöberl 2011]

Notation: ① $\text{sym}(\tau) = \frac{1}{2}(\tau + \tau^\top)$ for matrices $\tau \in \mathbb{R}^{3 \times 3}$. ② $\mathbb{S} = \text{sym}(\mathbb{R}^{3 \times 3})$, space of symmetric matrices.

More matrix-valued elements



Coupling dofs on each edge:

$$g \mapsto \int_E t \cdot gt$$

\mathbb{S} -valued Regge element

Natural operator: $\operatorname{curl}^\top \operatorname{curl}$ or “inc”

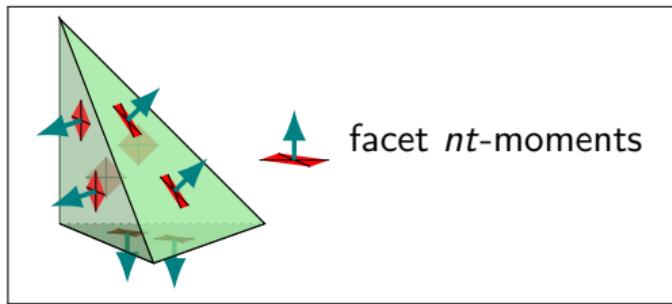
Continuous tt -component

[Regge 1961] [Christiansen 2011] [Li 2018]

[G+Neunteufel+Schöberl+Wardetsky 2023]

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More matrix-valued elements



Two coupling dofs on each face:

$$\tau \mapsto \int_F t \cdot \tau n$$

\mathbb{T} -valued viscous stress element

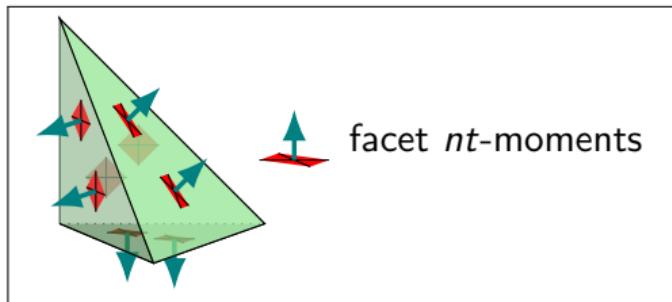
[G+Lederer+Schöberl 2019]

Natural operator: curl div

Continuous nt -component

Notation: ① $\text{sym}(\tau) = \frac{1}{2}(\tau + \tau^\top)$ for matrices $\tau \in \mathbb{R}^{3 \times 3}$. ② $\mathbb{S} = \text{sym}(\mathbb{R}^{3 \times 3})$, space of symmetric matrices.
③ $\text{dev } \tau = \tau - \frac{1}{3}(\text{tr } \tau)I$ for matrices τ , ④ $\mathbb{T} = \text{dev } \mathbb{R}^{3 \times 3}$ = space of traceless matrices.

More matrix-valued elements



Two coupling dofs on each face:

$$\tau \mapsto \int_F t \cdot \tau n$$

\mathbb{T} -valued viscous stress element

[G+Lederer+Schöberl 2019]

Natural operator: curl div

Continuous nt-component

Recent understanding: An algebraic structure with matrix-valued Sobolev space functions connected by curl div, curl \mathbb{T} curl, div div.

Notation: ① $\text{sym}(\tau) = \frac{1}{2}(\tau + \tau^\top)$ for matrices $\tau \in \mathbb{R}^{3 \times 3}$. ② $\mathbb{S} = \text{sym}(\mathbb{R}^{3 \times 3})$, space of symmetric matrices.
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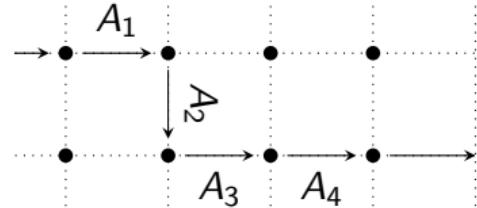
Outline

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Complexes from the 2-complex

We use “commutative diagrams”

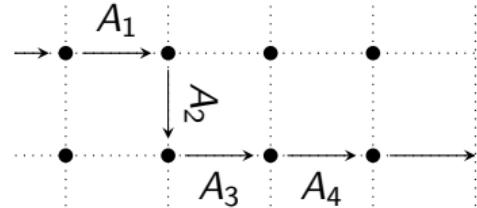
- whose “vertices” or “objects” are function spaces,
- whose “arrows” or “morphisms” are differential operators.



Complexes from the 2-complex

We use “commutative diagrams”

- whose “vertices” or “objects” are function spaces,
- whose “arrows” or “morphisms” are differential operators.



-
- Compositions of morphisms are called “paths.”
 - A path is a **complex** if compositions of 2 successive morphisms vanish:

$$A_{j+1} \circ A_j = 0.$$

- A complex is **exact** if $\text{range}(A_j) = \ker(A_{j+1})$.
- A path is a **2-complex** if the compositions of 3 successive morphisms vanish: [Olver 1982]

$$A_{j+2} \circ A_{j+1} \circ A_j = 0.$$

Scalar (\mathbb{R}) and vector (\mathbb{V}) valued fields in 3D

On a bounded contractible domain $\Omega \subset \mathbb{R}^3$ with Lipschitz boundary, let \mathcal{D} denote the space of infinitely differentiable functions compactly supported on Ω . The **de Rham complex**

$$0 \longrightarrow \mathcal{D} \xrightarrow{\text{grad}} \mathcal{D} \otimes \mathbb{V} \xrightarrow{\text{curl}} \mathcal{D} \otimes \mathbb{V} \xrightarrow{\text{div}} \mathcal{D} \longrightarrow \mathbb{R}$$

is exact.

Notation: $\mathbb{V} = \mathbb{R}^3$. Tensor product $X \otimes \mathbb{V}$ is identifiable with the Cartesian product $X \times X \times X$.

Scalar (\mathbb{R}) and vector (\mathbb{V}) valued fields in 3D

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is exact, and so is the Sobolev de Rham complex,

$$0 \longrightarrow \mathring{H}^1 \xrightarrow{\text{grad}} \mathring{H}(\text{curl}) \xrightarrow{\text{curl}} \mathring{H}(\text{div}) \xrightarrow{\text{div}} L_2 \longrightarrow \mathbb{R}$$

where

$$\begin{aligned}\mathring{H}(\text{curl}) &= \overline{\mathcal{D} \otimes \mathbb{V}}^{\|\cdot\|_{H(\text{curl})}}, & \mathring{H}(\text{div}) &= \overline{\mathcal{D} \otimes \mathbb{V}}^{\|\cdot\|_{H(\text{div})}} \\ \|u\|_{H(\text{curl})}^2 &= \|u\|_{L_2}^2 + \|\text{curl } u\|_{L_2}^2, & \|q\|_{H(\text{div})}^2 &= \|q\|_{L_2}^2 + \|\text{div } q\|_{L_2}^2.\end{aligned}$$

Notation: $\mathbb{V} = \mathbb{R}^3$. Tensor product $X \otimes \mathbb{V}$ is identifiable with the Cartesian product $X \times X \times X$.

\mathbb{S} and \mathbb{T} valued matrix fields

$$\mathcal{D} \xrightarrow{\text{grad}} \mathcal{D} \otimes \mathbb{V} \xrightarrow{\text{curl}} \mathcal{D} \otimes \mathbb{V} \xrightarrow{\text{div}} \mathcal{D}$$

$$\mathcal{D} \otimes \mathbb{V}$$

$$\mathcal{D} \otimes \mathbb{S}$$

$$\mathcal{D} \otimes \mathbb{T}$$

$$\mathcal{D} \otimes \mathbb{V}$$

Notation: ① $\text{sym } \tau = \frac{1}{2}(\tau + \tau^\top)$ for matrices τ . ② $\mathbb{S} = \text{sym } \mathbb{R}^{3 \times 3}$, space of **symmetric** matrices.

③ $\text{dev } \tau = \tau - \frac{1}{3}(\text{tr } \tau)I$ for matrices τ , ④ $\mathbb{T} = \text{dev } \mathbb{R}^{3 \times 3} =$ space of **traceless** matrices.

\mathbb{S} and \mathbb{T} valued matrix fields

$$\mathcal{D} \xrightarrow{\text{grad}} \mathcal{D} \otimes \mathbb{V} \xrightarrow{\text{curl}} \mathcal{D} \otimes \mathbb{V} \xrightarrow{\text{div}} \mathcal{D}$$

$$\mathcal{D} \otimes \mathbb{V} \xrightarrow{\text{def}} \mathcal{D} \otimes \mathbb{S} \xrightarrow{\text{curl}} \mathcal{D} \otimes \mathbb{T} \xrightarrow{\text{div}} \mathcal{D} \otimes \mathbb{V}$$

- **def v** $\equiv \varepsilon(v) = \text{sym grad } v$, $\text{grad } v$ = Jacobian matrix of vector field v ,

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- def $v \equiv \varepsilon(v) = \text{sym grad } v$, grad v = Jacobian matrix of vector field v ,
- curl and div on matrix-valued functions act row-wise.
- curl of a symmetric matrix is traceless.

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$$\begin{array}{ccccccc} \mathcal{D} & \xrightarrow{\text{grad}} & \mathcal{D} \otimes \mathbb{V} & \xrightarrow{\text{curl}} & \mathcal{D} \otimes \mathbb{V} & \xrightarrow{\text{div}} & \mathcal{D} \\ \downarrow \text{grad} & & \downarrow \text{def} & & \downarrow \frac{1}{2} \top \text{dev grad} & & \downarrow \frac{1}{3} \text{grad} \\ \mathcal{D} \otimes \mathbb{V} & \xrightarrow{\text{def}} & \mathcal{D} \otimes \mathbb{S} & \xrightarrow{\text{curl}} & \mathcal{D} \otimes \mathbb{T} & \xrightarrow{\text{div}} & \mathcal{D} \otimes \mathbb{V} \end{array}$$

- Connect vertically.

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\mathbb{S} and \mathbb{T} valued matrix fields

$$\begin{array}{ccccccc} \mathcal{D} & \xrightarrow{\text{grad}} & \mathcal{D} \otimes \mathbb{V} & \xrightarrow{\text{curl}} & \mathcal{D} \otimes \mathbb{V} & \xrightarrow{\text{div}} & \mathcal{D} \\ \downarrow \text{grad} & & \downarrow \text{def} & \text{C} & \downarrow \frac{1}{2} \top \text{dev grad} & & \downarrow \frac{1}{3} \text{grad} \\ \mathcal{D} \otimes \mathbb{V} & \xrightarrow{\text{def}} & \mathcal{D} \otimes \mathbb{S} & \xrightarrow{\text{curl}} & \mathcal{D} \otimes \mathbb{T} & \xrightarrow{\text{div}} & \mathcal{D} \otimes \mathbb{V} \end{array}$$

- The diagram commutes.

$$\text{curl def} = \frac{1}{2} \top \text{dev grad curl}$$

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A commutative diagram showing the relationships between spaces of vector fields (\mathcal{D}), symmetric matrices (\mathbb{S}), and traceless matrices (\mathbb{T}). The top row consists of \mathcal{D} , $\mathcal{D} \otimes \mathbb{V}$, $\mathcal{D} \otimes \mathbb{V}$, and \mathcal{D} . The bottom row consists of $\mathcal{D} \otimes \mathbb{V}$, $\mathcal{D} \otimes \mathbb{S}$, $\mathcal{D} \otimes \mathbb{T}$, and $\mathcal{D} \otimes \mathbb{V}$. Vertical arrows are labeled "grad", "def", and $\frac{1}{2} \top \text{dev grad}$. A diagonal arrow from $\mathcal{D} \otimes \mathbb{V}$ to $\mathcal{D} \otimes \mathbb{S}$ is labeled "def". A diagonal arrow from $\mathcal{D} \otimes \mathbb{V}$ to $\mathcal{D} \otimes \mathbb{T}$ is labeled "div". A vertical arrow from $\mathcal{D} \otimes \mathbb{V}$ to $\mathcal{D} \otimes \mathbb{V}$ is labeled $\frac{1}{3} \text{grad}$. A red circle highlights the term $\frac{1}{2} \top \text{dev grad}$.

- The diagram commutes.

$$\text{curl def} = \frac{1}{2} \top \text{dev grad curl}$$

$$\frac{1}{2} \text{div } \top \text{dev grad } u = \frac{1}{3} \text{grad div}$$

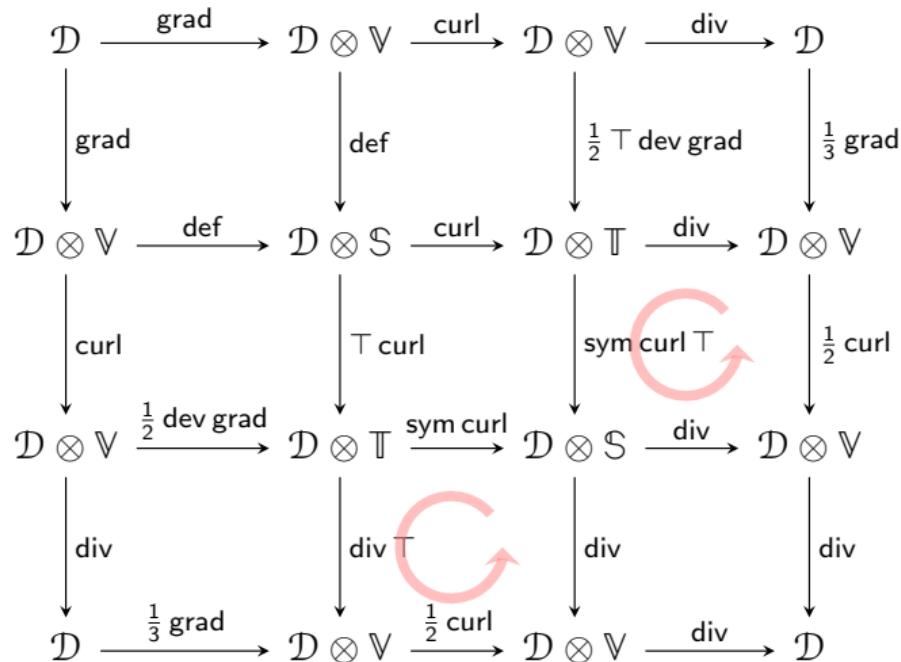
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Completing the picture

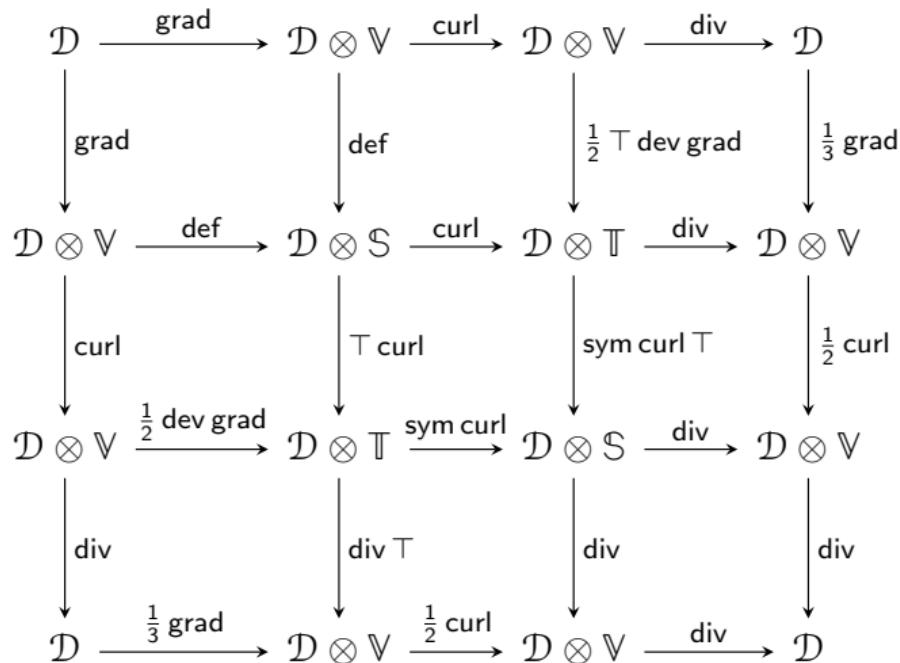
$$\begin{array}{ccccccc} \mathcal{D} & \xrightarrow{\text{grad}} & \mathcal{D} \otimes \mathbb{V} & \xrightarrow{\text{curl}} & \mathcal{D} \otimes \mathbb{V} & \xrightarrow{\text{div}} & \mathcal{D} \\ \downarrow \text{grad} & & \downarrow \text{def} & & \downarrow \frac{1}{2} \top \text{ dev grad} & & \downarrow \frac{1}{3} \text{ grad} \\ \mathcal{D} \otimes \mathbb{V} & \xrightarrow{\text{def}} & \mathcal{D} \otimes \mathbb{S} & \xrightarrow{\text{curl}} & \mathcal{D} \otimes \mathbb{T} & \xrightarrow{\text{div}} & \mathcal{D} \otimes \mathbb{V} \\ \downarrow \text{curl} & & \downarrow \top \text{ curl} & & \downarrow \text{sym curl } \top & & \downarrow \frac{1}{2} \text{ curl} \\ \mathcal{D} \otimes \mathbb{V} & \xrightarrow{\frac{1}{2} \text{ dev grad}} & \mathcal{D} \otimes \mathbb{T} & \xrightarrow{\text{sym curl}} & \mathcal{D} \otimes \mathbb{S} & \xrightarrow{\text{div}} & \mathcal{D} \otimes \mathbb{V} \\ \downarrow \text{div} & & \downarrow \text{div } \top & & \downarrow \text{div} & & \downarrow \text{div} \\ \mathcal{D} & \xrightarrow{\frac{1}{3} \text{ grad}} & \mathcal{D} \otimes \mathbb{V} & \xrightarrow{\frac{1}{2} \text{ curl}} & \mathcal{D} \otimes \mathbb{V} & \xrightarrow{\text{div}} & \mathcal{D} \end{array}$$

Completing the picture



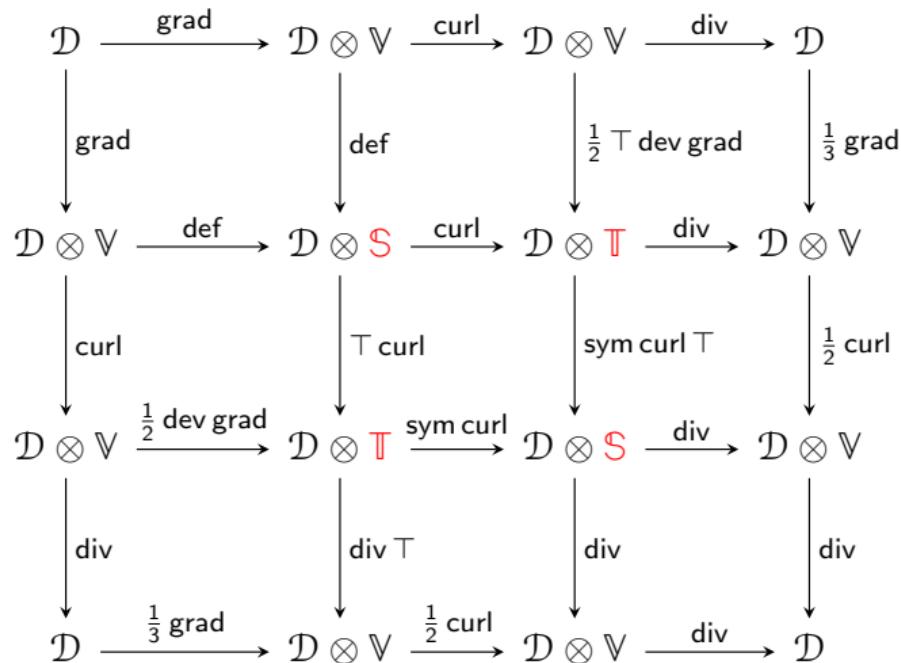
$$\text{div sym curl } \top = \frac{1}{2} \text{ curl div} .$$

Completing the picture



Note the symmetry in the diagram about the diagonal.

Completing the picture



The four middle spaces extend to new Sobolev spaces of matrix fields.

New and old Sobolev spaces

Let

$$H^{-1} = (\mathring{H}^1)^*,$$

$$\tilde{H}^{-1} = (H^1)^*.$$

Define new Sobolev spaces of matrix fields:

$$\tilde{H}_{\text{cc}} := \{g \in \tilde{H}^{-1} \otimes \mathbb{S} : \operatorname{curl} g \in \tilde{H}^{-1} \otimes \mathbb{T}, \operatorname{curl}^\top \operatorname{curl} g \in \tilde{H}^{-1} \otimes \mathbb{S}\}.$$

$$\tilde{H}_{\text{cd}} := \{\tau \in \tilde{H}^{-1} \otimes \mathbb{T} : \operatorname{div} \tau \in \tilde{H}^{-1} \otimes \mathbb{V}, \operatorname{sym} \operatorname{curl} \tau \in \tilde{H}^{-1} \otimes \mathbb{S}, \operatorname{curl} \operatorname{div} \tau \in \tilde{H}^{-1} \otimes \mathbb{V}\}.$$

$$\tilde{H}_{\text{cd}\top} := \top \tilde{H}_{\text{cd}}.$$

$$\tilde{H}_{\text{dd}} := \{\sigma \in \tilde{H}^{-1} \otimes \mathbb{S} : \operatorname{div} \sigma \in \tilde{H}^{-1} \otimes \mathbb{V}, \operatorname{div} \operatorname{div} \sigma \in \tilde{H}^{-1}\}.$$

Define $H_{\text{cc}}, H_{\text{cd}}, H_{\text{cd}\top}, H_{\text{dd}}$ similarly, replacing \tilde{H}^{-1} by H^{-1} .

Two de Rham complexes

Recall the standard Sobolev de Rham complex:

$$0 \longrightarrow \mathring{H}^1 \xrightarrow{\text{grad}} \mathring{H}(\text{curl}) \xrightarrow{\text{curl}} \mathring{H}(\text{div}) \xrightarrow{\text{div}} L_{2,\mathbb{R}}.$$

A version with lower regularity:

$$L_{2,\mathbb{R}} \xrightarrow{\text{grad}} \tilde{H}_{\mathcal{RT}}^{-1}(\text{curl}) \xrightarrow{\text{curl}} \tilde{H}_{\mathcal{ND}}^{-1}(\text{div}) \xrightarrow{\text{div}} \tilde{H}_{\mathcal{P}_1}^{-1}.$$

Notation:

$$L_{2,\mathbb{R}} = \{u \in L_2 : \int_{\Omega} u = 0\},$$

$$\tilde{H}_{\mathcal{RT}}^{-1}(\text{curl}) = \{v \in \tilde{H}^{-1} : \text{curl } v \in \tilde{H}^{-1}, v(r) = 0 \text{ for all } r \in \mathcal{RT}\}, \quad \mathcal{RT} = \{a + bx : a \in \mathbb{V}, b \in \mathbb{R}\},$$

$$\tilde{H}_{\mathcal{ND}}^{-1}(\text{div}) = \{q \in \tilde{H}^{-1} : \text{div } q \in \tilde{H}^{-1}, q(r) = 0 \text{ for all } r \in \mathcal{ND}\}, \quad \mathcal{ND} = \{a + d \times x : a, d \in \mathbb{V}\},$$

$$\tilde{H}_{\mathcal{P}_1}^{-1} = \{w \in \tilde{H}^{-1} : w(p) = 0 \text{ for all } p \in \mathcal{P}_1\}, \quad \mathcal{P}_k = \text{polynomials of degree } \leq k.$$

Two de Rham complexes

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(Why remove \mathcal{RT} ? $u \in L_{2,\mathbb{R}} \implies (\text{grad } u)(a + bx) = -u(\text{div}(a + bx)) = -b \int_{\Omega} u = 0.$)

Notation:

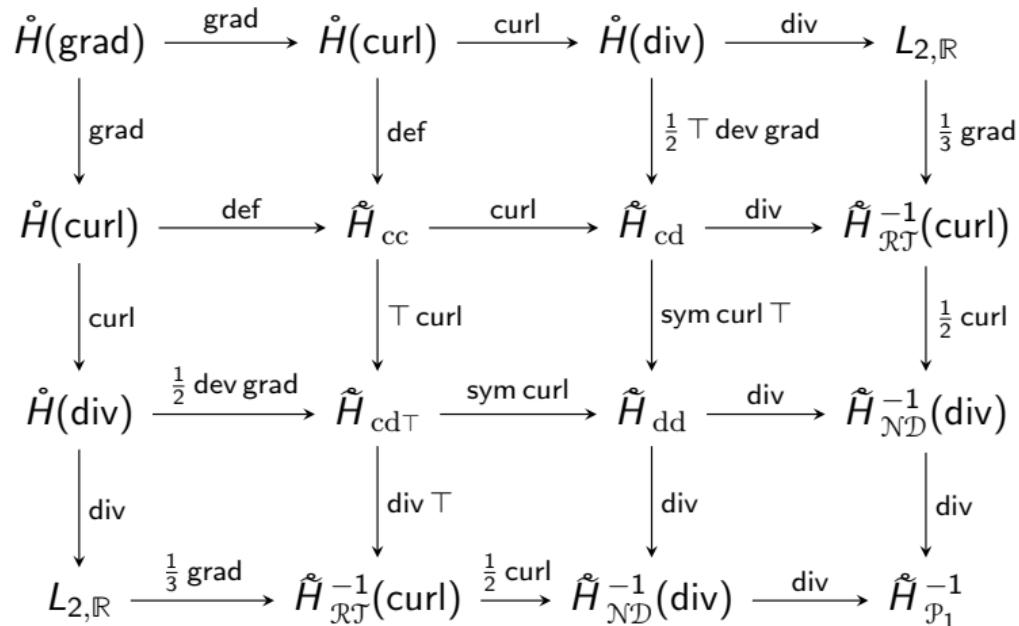
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$$\tilde{H}_{\mathcal{ND}}^{-1}(\text{div}) = \{q \in \tilde{H}^{-1} : \text{div } q \in \tilde{H}^{-1}, q(r) = 0 \text{ for all } r \in \mathcal{ND}\}, \quad \mathcal{ND} = \{a + d \times x : a, d \in \mathbb{V}\},$$

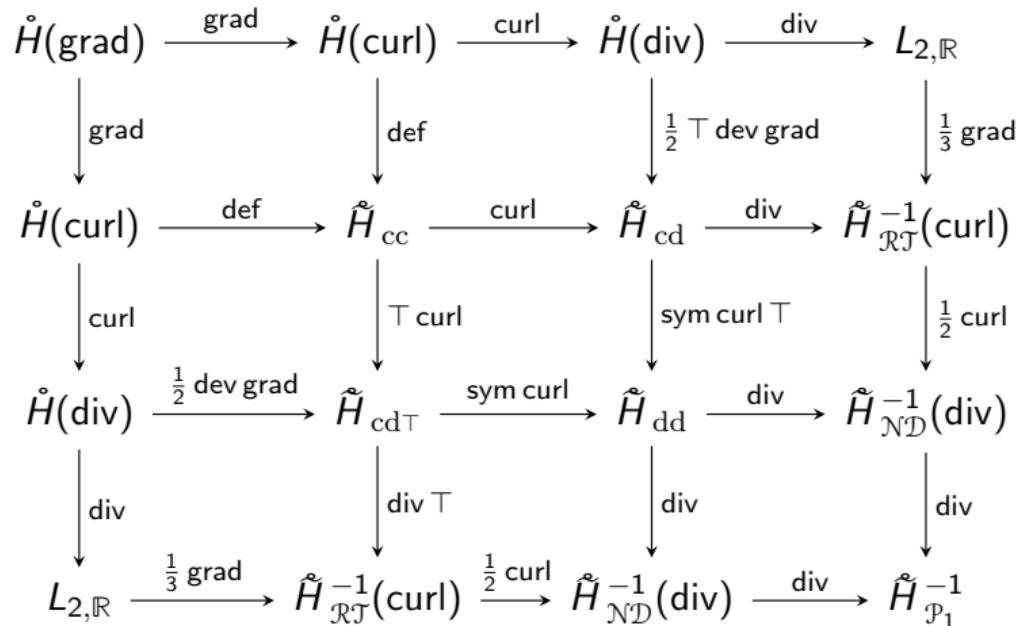
$$\tilde{H}_{\mathcal{P}_1}^{-1} = \{w \in \tilde{H}^{-1} : w(p) = 0 \text{ for all } p \in \mathcal{P}_1\}, \quad \mathcal{P}_k = \text{polynomials of degree } \leq k.$$

The diagram with Sobolev spaces



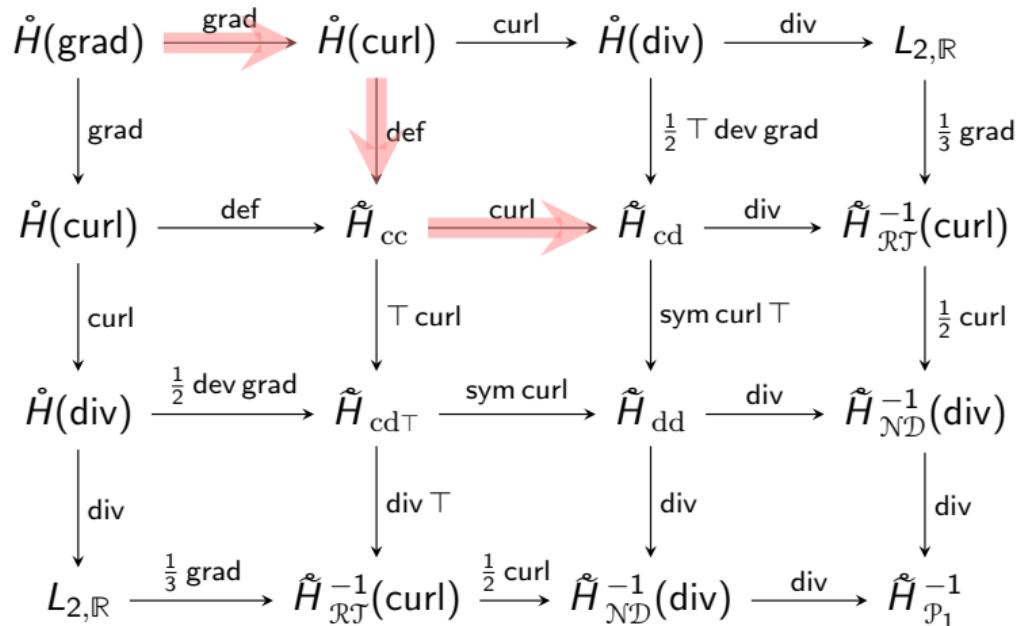
Theorem: The diagram **commutes** and every differential operator in it is **continuous** (with respect to the norms of the indicated domains and codomains) and has **closed range**.

2-complexes



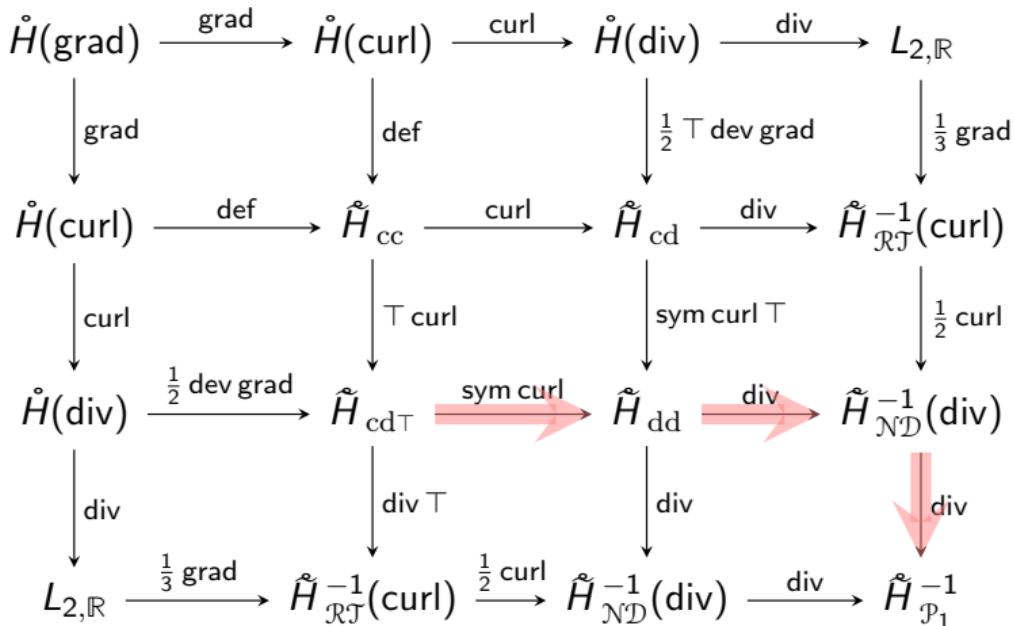
Theorem: All paths in the diagram are 2-complexes.

2-complexes



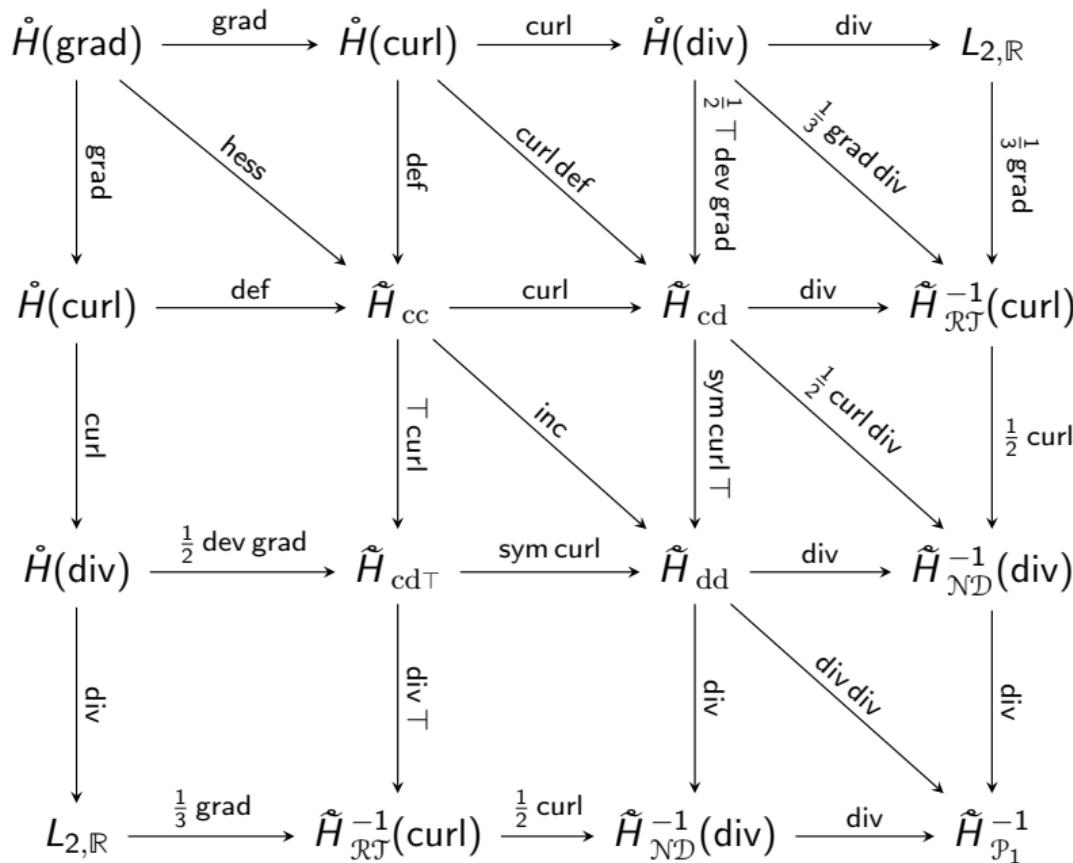
E.g., $\text{curl} \circ \text{def} \circ \text{grad } w = \frac{1}{2} \top \text{dev grad curl grad } w = 0$.

2-complexes

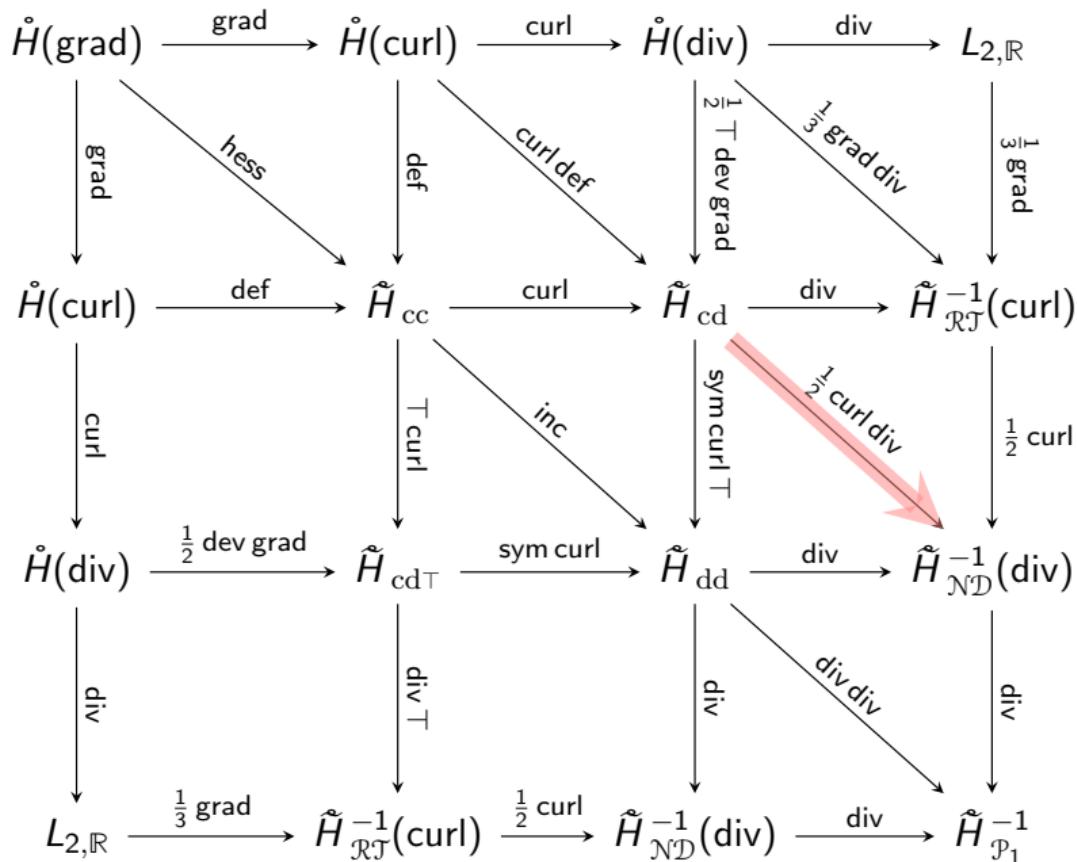


E.g., $\text{div} \circ \text{div} \circ \text{sym curl } \tau = \frac{1}{2} \text{ div curl div } \top \tau = 0$.

Diagonal second-order operators



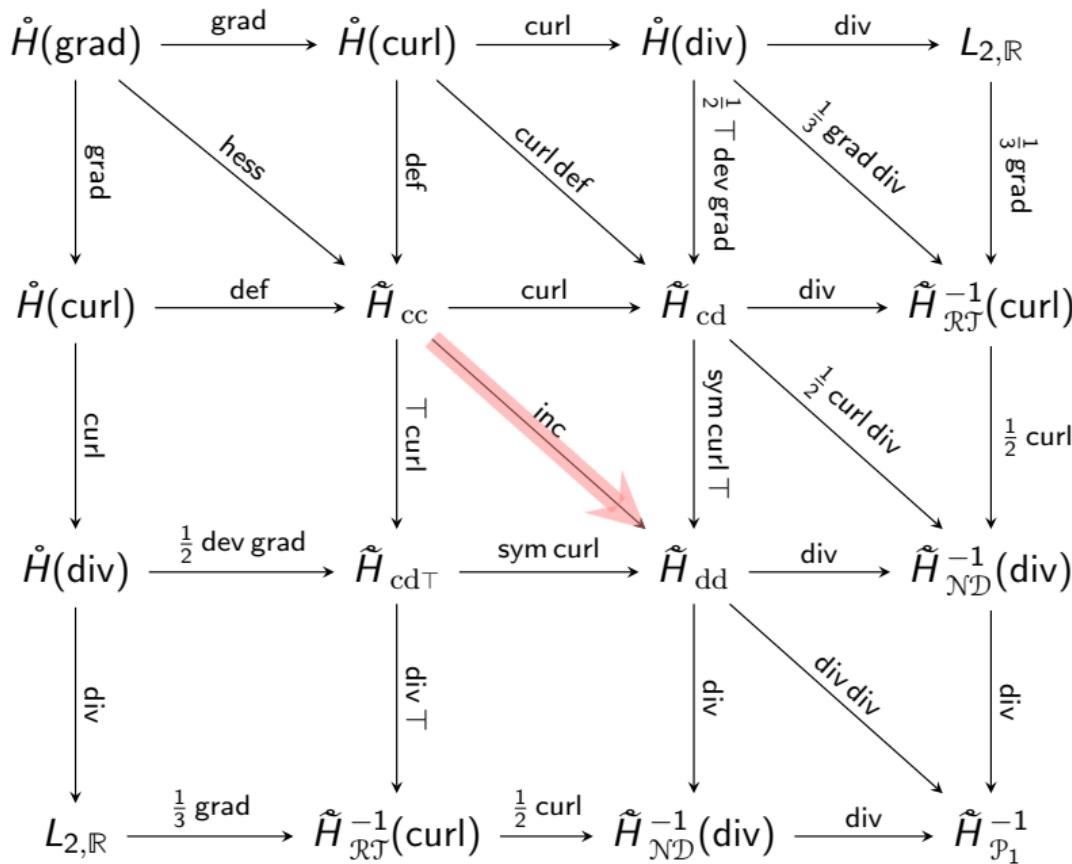
Diagonal second-order operators



Operators that appeared

- in MCS:
curl div

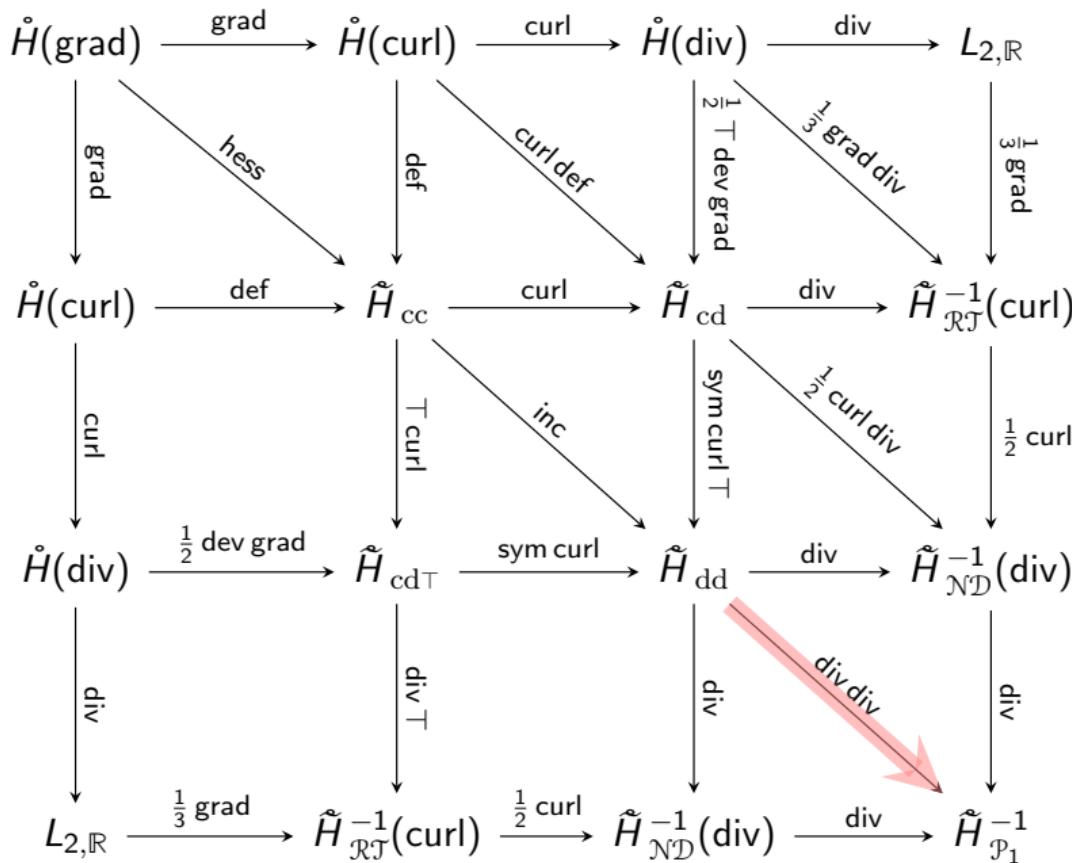
Diagonal second-order operators



Operators that appeared

- in MCS:
 curl div
- in curvature:
 $\text{inc} = \text{curl } \top \text{ curl}$

Diagonal second-order operators



Operators that appeared

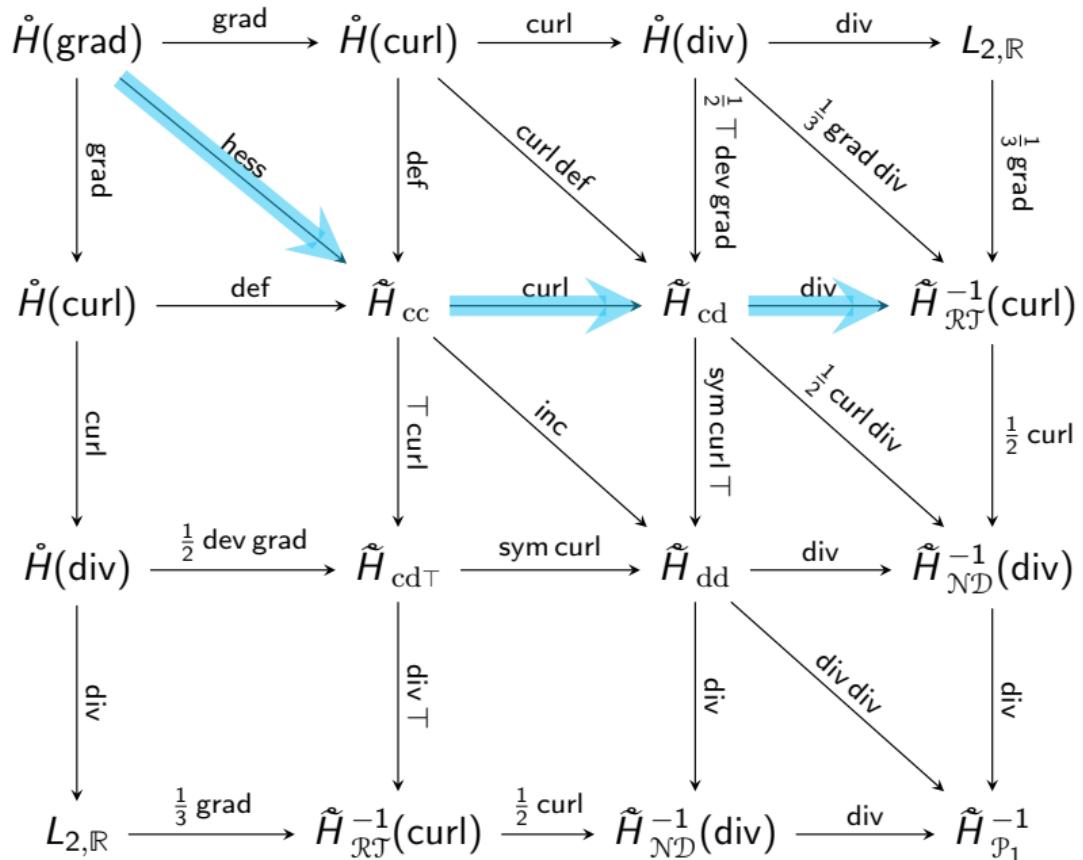
- in MCS:
 curl div
- in curvature:
 $\text{inc} = \text{curl } \top \text{ curl}$
- in TDNNS:
 div div

Complexes

Theorem:

The following paths are **exact** complexes:

The hessian complex



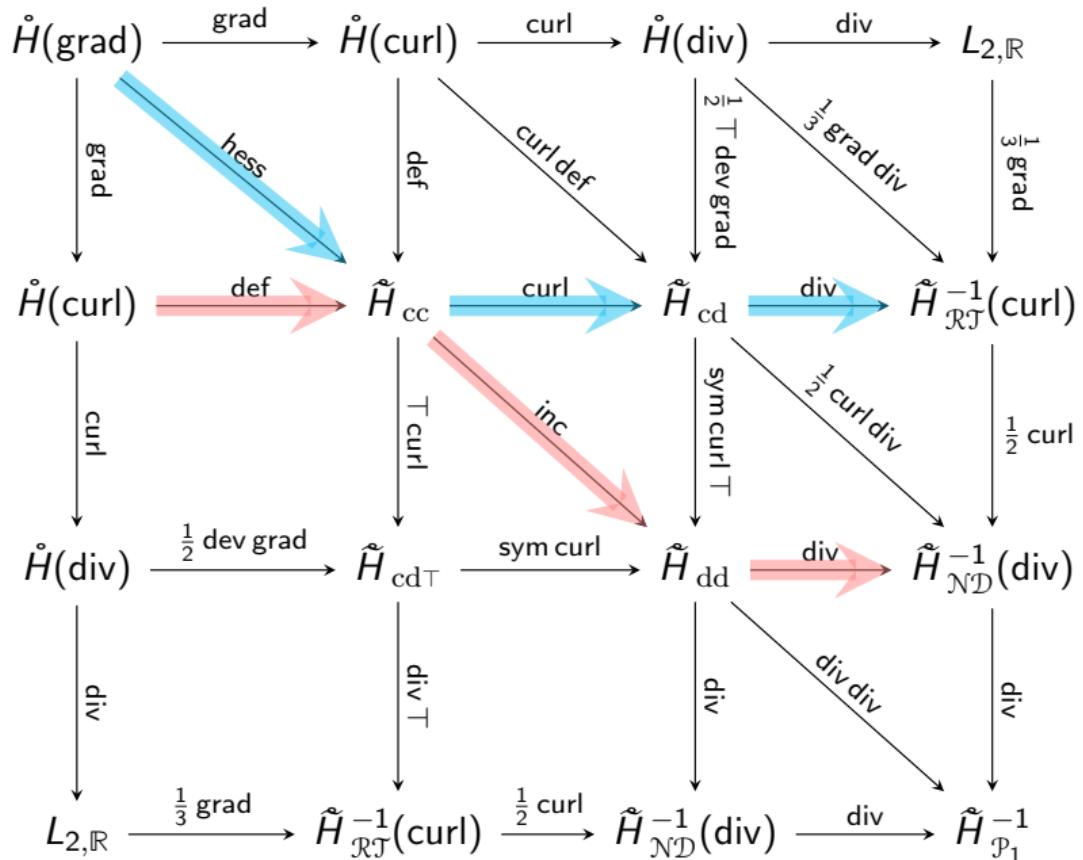
Complexes

Theorem:

The following paths are **exact** complexes:

The hessian complex

The elasticity complex



Complexes

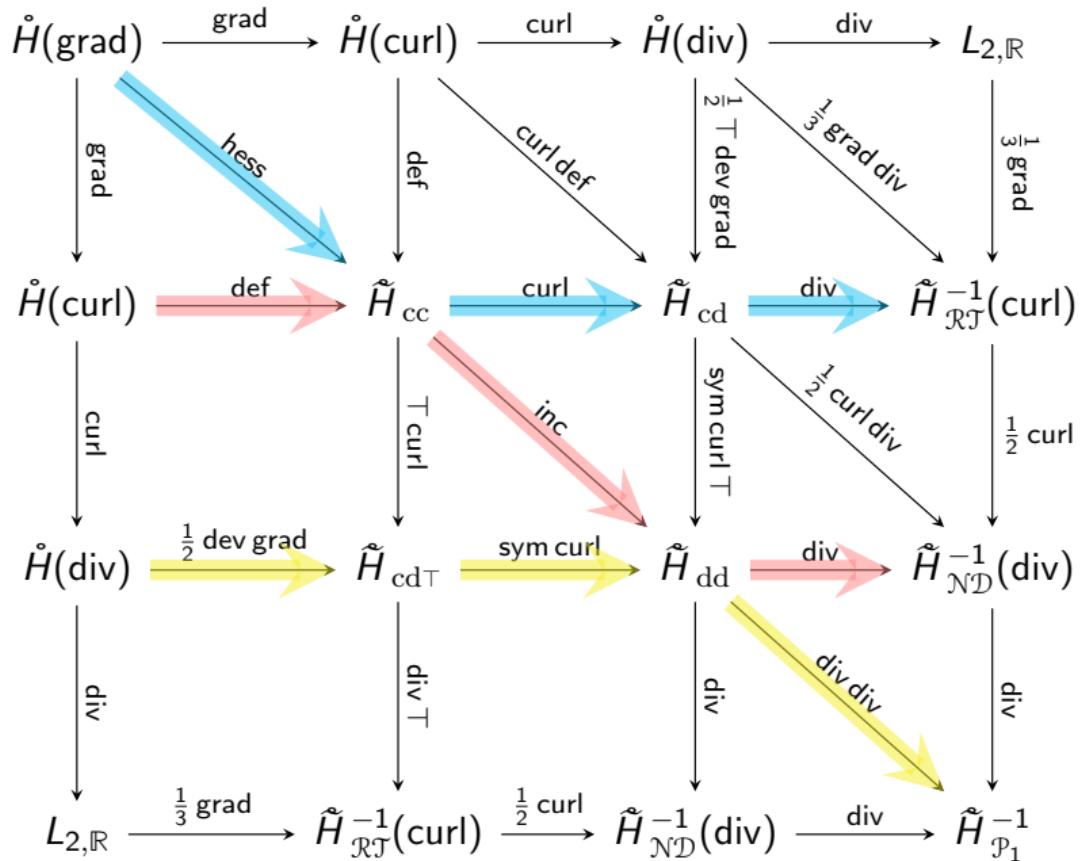
Theorem:

The following paths are **exact** complexes:

The hessian complex

The elasticity complex

The div div complex



Outline

- 1 Introduction
- 2 Mass-Conserving Stress-yielding (MCS) method
- 3 Viscous stress elements
- 4 Other matrix finite elements
- 5 A unifying 2-complex of Sobolev spaces
- 6 Regular decompositions
- 7 A complementary diagram

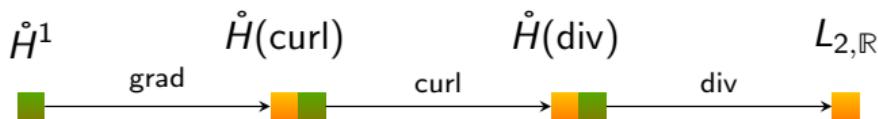
Generalized Bogovskii operators

- A classical regular right inverse of $\text{div}(\cdot)$ was developed by [Bogovskii 1979]. It was generalized to r -forms in n -dimensions by [Costabel+McIntosh 2010]. In 3D,

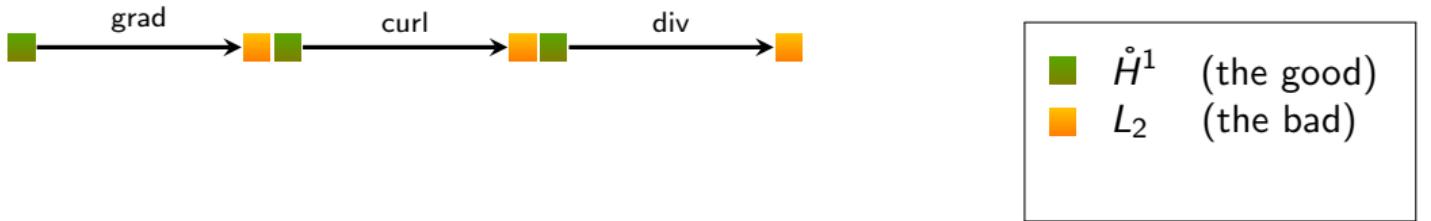
$$\begin{array}{ccccc} \mathring{H}^1 & \xrightarrow{\text{grad}} & \mathring{H}(\text{curl}) & \xleftarrow{\text{curl}} & \mathring{H}(\text{div}) & \xleftarrow{\text{div}} & L_{2,\mathbb{R}}, \\ & \xleftarrow{T_g} & & \xleftarrow{T_c} & & \xleftarrow{T_d} & \end{array}$$

where T_g , T_c , and T_d map into **more regular** subspaces with \mathring{H}^1 components.

- Consequently, the spaces can be split into **smooth** and **rough** parts:



Regular decompositions of the spaces



The standard case

There are continuous linear maps
 $S_c : \mathring{H}(\text{curl}) \rightarrow \mathring{H}^1$ and
 $S_d : \mathring{H}(\text{div}) \rightarrow \mathring{H}^1 \otimes \mathbb{V}$ such that

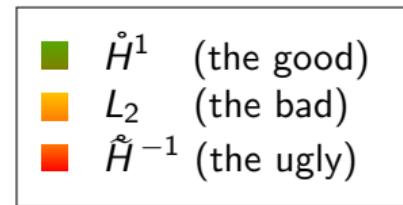
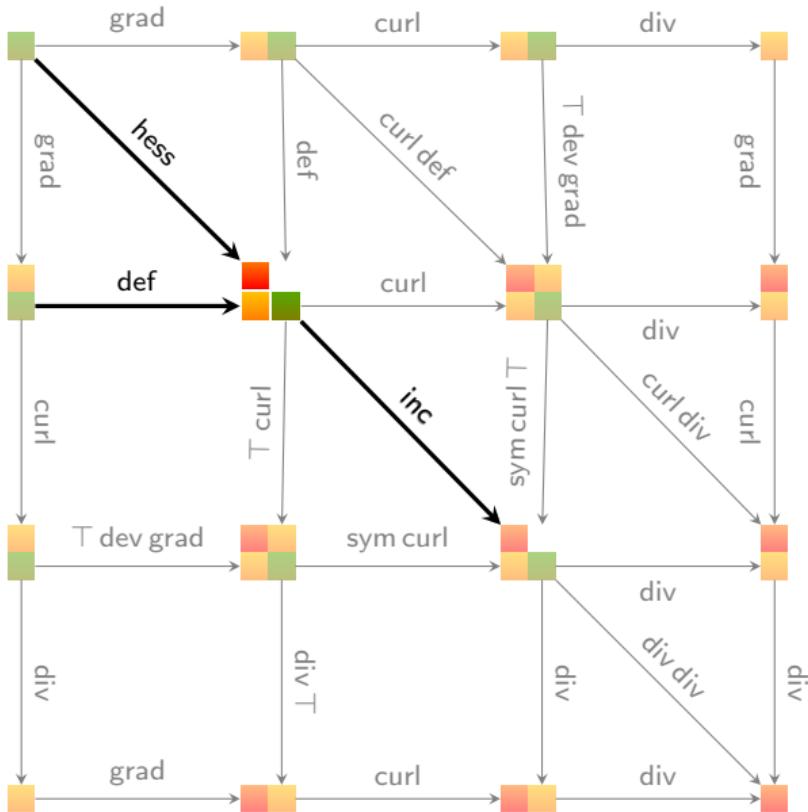
$$\forall u \in \mathring{H}(\text{curl}) :$$

$$u = \text{grad}(S_c u) + T_c \text{curl } u,$$

$$\forall q \in \mathring{H}(\text{div}) :$$

$$q = \text{curl}(S_d q) + T_d \text{div } q.$$

Regular decompositions of the spaces

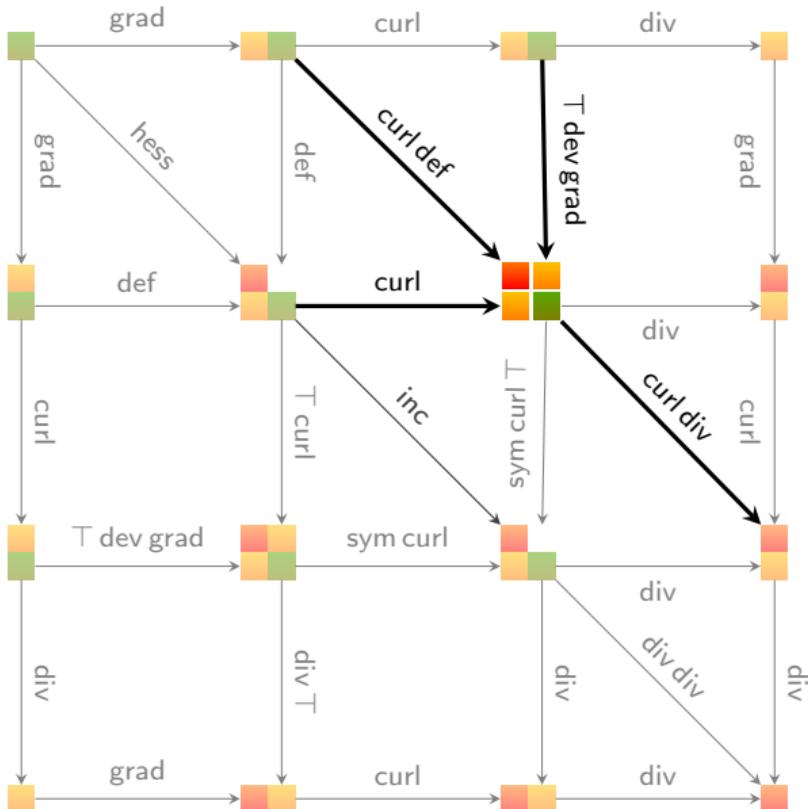


The \tilde{H}_{cc} case

Theorem: Any g in \tilde{H}_{cc} admits the decomposition

$$\begin{aligned} g = & \text{hess}(S_{cc}^{(1)} g) \\ & + \text{def}(S_{cc}^{(2)} g) \\ & + D_{cc} \text{inc } g. \end{aligned}$$

Regular decompositions of the spaces



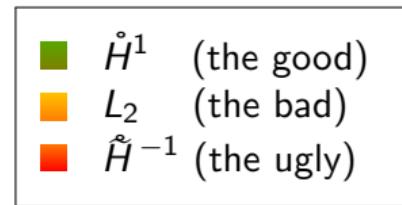
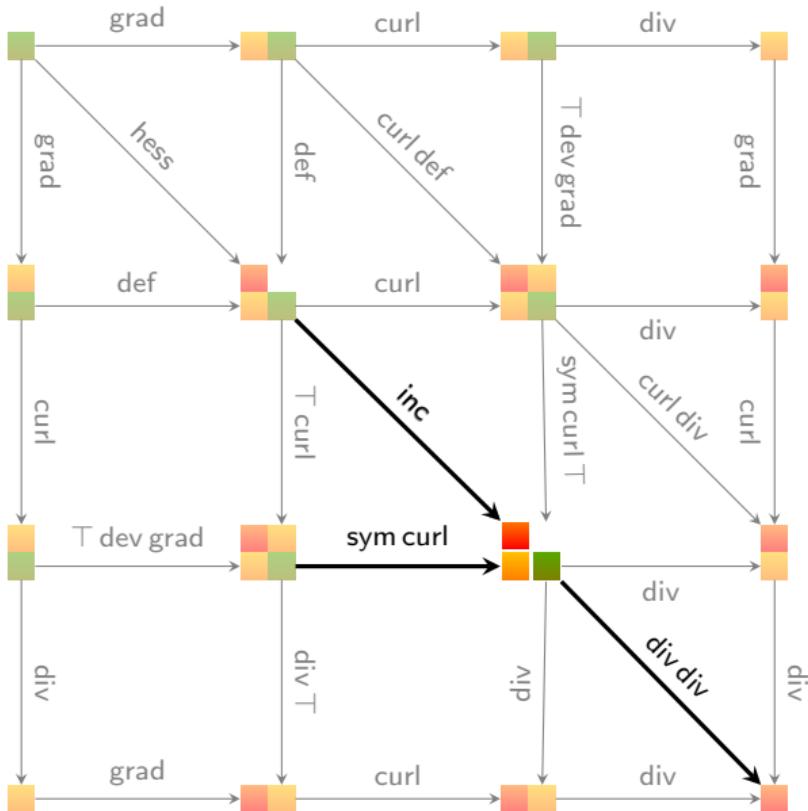
	H^1	(the good)
	L_2	(the bad)
	\tilde{H}^{-1}	(the ugly)

The \tilde{H}_{cd} case

Theorem: Any τ in \tilde{H}_{cd} admits the decomposition

$$\begin{aligned}\tau = & \text{curl def}(S_{cd}^{(1)} \tau) \\ & + \text{curl}(S_{cd}^{(2)} \tau) \\ & + \text{T dev grad}(S_{cd}^{(3)} \tau) \\ & + D_{cd} \text{curl div } \tau.\end{aligned}$$

Regular decompositions of the spaces



The \mathring{H}_{dd} case

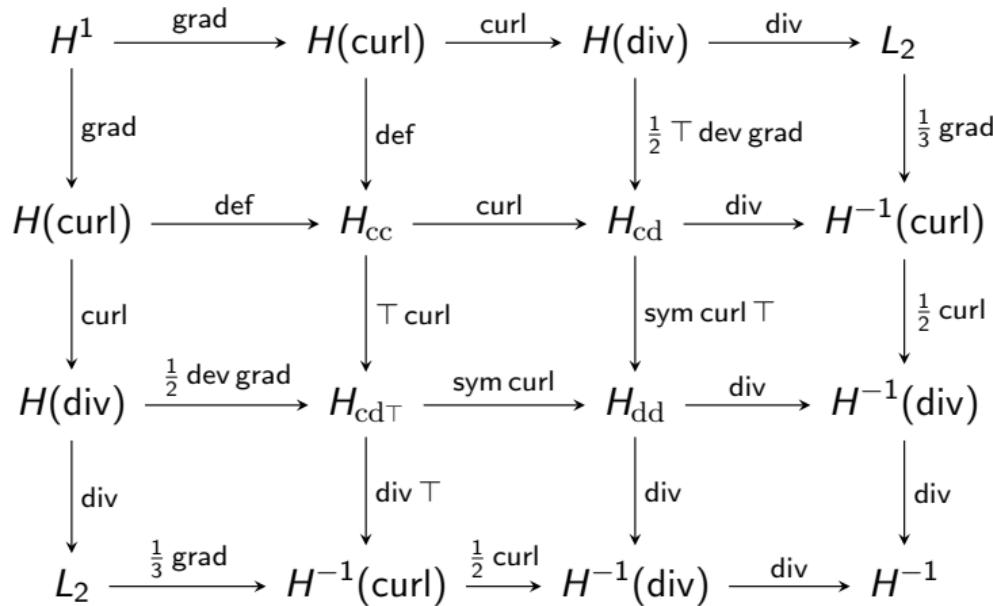
Theorem: Any σ in \mathring{H}_{dd} admits the decomposition

$$\begin{aligned}\sigma = & \text{inc}(S_{dd}^{(1)} \sigma) \\ & + \text{sym curl}(S_{dd}^{(2)} \sigma) \\ & + D_{dd} \text{div div } \sigma.\end{aligned}$$

Outline

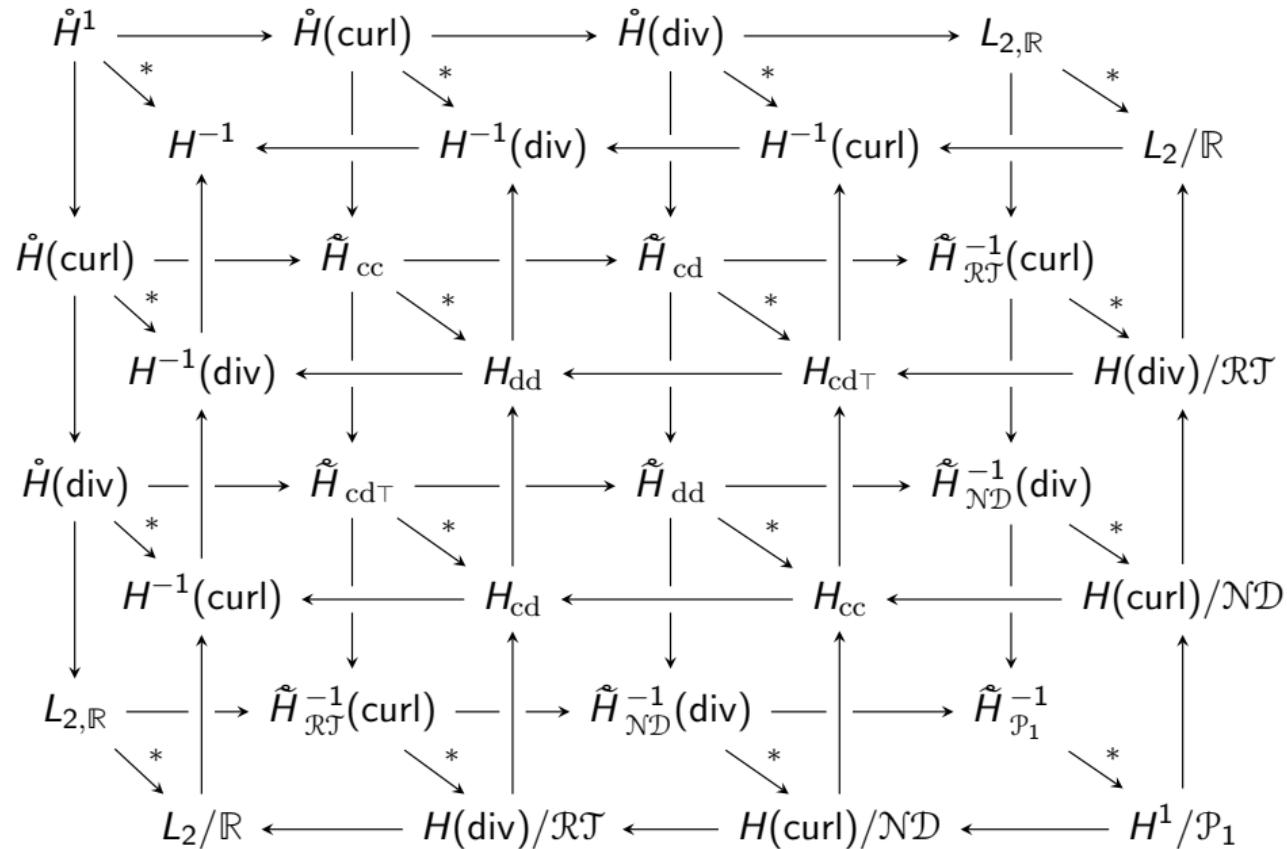
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Diagram for the case without boundary conditions



Prior results extend to this diagram also. (We can also start with H^1/\mathcal{P}_1 instead of H^1 and mod out appropriate subspaces as in the next slide.)

Duality



Duality identities

$$① (\mathring{H}^1)^* = H^{-1}$$

$$② (H^1)^* = \tilde{H}^{-1}$$

$$③ \mathring{H}(\text{div})^* = H^{-1}(\text{curl})$$

$$④ \mathring{H}(\text{curl})^* = H^{-1}(\text{div})$$

$$⑤ H(\text{div})^* = \tilde{H}^{-1}(\text{curl})$$

$$⑥ H(\text{curl})^* = \tilde{H}^{-1}(\text{div})$$

$$⑦ \tilde{H}_{cc}^* = H_{dd}$$

$$⑧ \tilde{H}_{cd}^* = H_{cd\top}$$

$$⑨ \tilde{H}_{dd}^* = H_{cc}$$

Conclusion

- Operator $\operatorname{curl} \operatorname{div}$ is natural for **viscous stresses** of Stokes flow with $H(\operatorname{div})$ velocity.
 - Simple **nt -continuous** matrix finite elements can be designed: [[G+Lederer+Schöberl 2019](#)].
 - The **MCS method** yield optimal rates for fluid stresses and is Reynolds robust.
 - MCS convergence analysis: [[G+Lederer+Schöberl 2020](#)], [[G+Kogler+Lederer+Schöberl 2023](#)].
-
- Other **nn -** and **tt -continuous** elements motivate study of accompanying Sobolev spaces.
 - They feature 2nd order operators **$\operatorname{curl} \operatorname{div}$, $\operatorname{div} \operatorname{div}$, and inc** on matrix fields.
-
- Two commuting diagrams that give **2-complexes** were seen. One is **dual** to the other.
 - They contain new Sobolev spaces **H_{cc} , H_{cd} , H_{dd}** of matrix fields with low-regularity.
 - Through **regular decompositions**, we understand these spaces better.
 - Recent ArXiV preprint: [[G+Hu+Schöberl 2025](#)].