

# Tight Focusing of LG beams with arbitrary polarisation states

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## A brief introduction to the theory of tightly focused paraxial beams

The electric field  $E(\rho, \phi, z)$  at the focal plane for an aplanatic lens with focal length  $f$  is given by the Richards and Wolf's Diffraction Integral (RWDI) [10.1098/rspa.1959.0200]:

$$\mathbf{E}(\rho, \phi, z) = -\frac{if}{\lambda} \int_0^\alpha \int_0^{2\pi} \sin\theta \mathbf{A}(\theta, \phi) \exp[i k \rho \sin\theta \cos(\phi - \phi)] \times \exp[i k z \cos\theta] d\theta d\phi,$$

where  $\alpha = \text{ArcSin}(NA/n)$  is the maximum angular aperture of the objective,  $n$  is the refractive index of the medium of the medium where the focusing occurs,  $NA$  is the lens numerical aperture and  $\mathbf{A}$  is the transformation of the input field  $\mathbf{U}(\mathbf{r}) = (U_x, U_y, 0)$  after the objective. Explicitly,  $\mathbf{A}$  is given by

$$\begin{aligned} A_x &= \sqrt{\cos\theta} \{ [\cos\theta + (1 - \cos\theta) \sin^2 \phi] U_x + (\cos\theta - 1) \cos\phi \sin\phi U_y \}, \\ A_y &= \sqrt{\cos\theta} \{ (\cos\theta - 1) \cos\phi \sin\phi U_x + [\cos\theta + (1 - \cos\theta) \cos^2 \phi] U_y \}, \\ A_z &= -\sqrt{\cos\theta} \sin\theta \{ \cos\phi U_x + \sin\phi U_y \}. \end{aligned}$$

It must be noticed that the 3D electric field at the focus satisfies Maxwell's equations. Since closed solutions for RWDI are hard or even impossible to obtain, the components of the electric field under tight focusing in this notes are found by numerical integration.

For the case of an arbitrary polarised Laguerre-Gauss beam, a simplified set of integrals can be obtained and you can find them in the PDF file.

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## Instructions

I designed the notebook to evaluate the *Definitions* section automatically, so you directly can start to play with the calculations. Here, I present a small introduction to the functions.

- `Efz0[xmax, dx, l, NA, n] = {EXx, EXy, EXz, EYx, EYy, EYz}`: Calculates the Cartesian contributions of a focused LG with topological charge  $l$  due to a lens with numerical aperture  $NA$  occurring in a linear and homogeneous medium with refractive index  $n$ . The calculation is performed in a numerical window

within the intervals  $-x_{\max} \leq x \leq x_{\max}$  and  $-y_{\max} \leq y \leq y_{\max}$ .

- **Eprop[xmax, dx, zmin, zmax, dz, /, NA, n]={EXx, EXy, EXz, EYx, EYy, EYz]**: Calculates the Cartesian contributions of a focused LG for different planes  $z_{\min} \leq z \leq z_{\max}$  with topological charge  $/$  due to a lens with numerical aperture **NA** occurring in a linear and homogeneous medium with refractive index **n**. The calculation is performed in a numerical window within the intervals  $-x_{\max} \leq x \leq x_{\max}$  and  $-y_{\max} \leq y \leq y_{\max}$  with resolution of  $dx$ .

- **Ep[ $\alpha, \beta$ ]={Ex, Ey, Ez}**: Given the array of the Cartesian contributions of the field it constructs the focused field given the parameters of the incident field. The incident beam is constructed  

$$\mathbf{E}_{\text{in}} = (\cos \alpha \hat{\mathbf{x}} + e^{i\beta} \sin \alpha \hat{\mathbf{y}}) U(\mathbf{r})$$

- **Int[ $\psi$ ]** Plots the intensity profile of the complex scalar amplitude  $\psi$ .

- **Ph[ $\psi$ ]** Plots the phase profile of the complex scalar amplitude  $\psi$ .

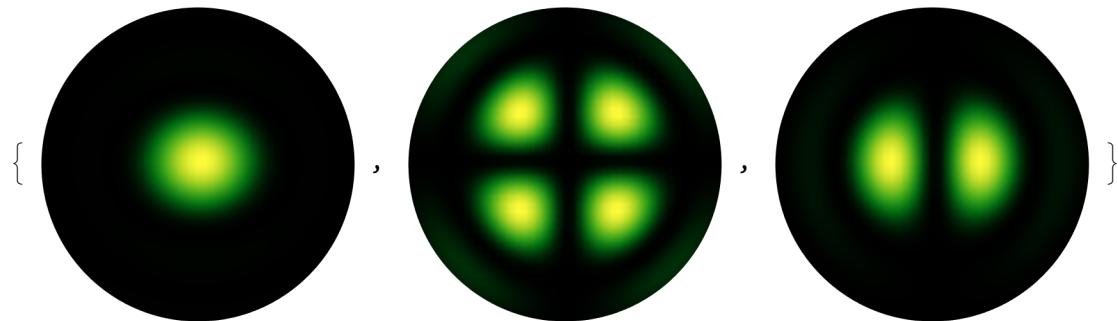
## Example 1 : LG00

```
In[203]:= AbsoluteTiming[E00 = Efz0[1.5, 0.07, 0, 0.95, 1]][1]
Out[203]= 39.222
```

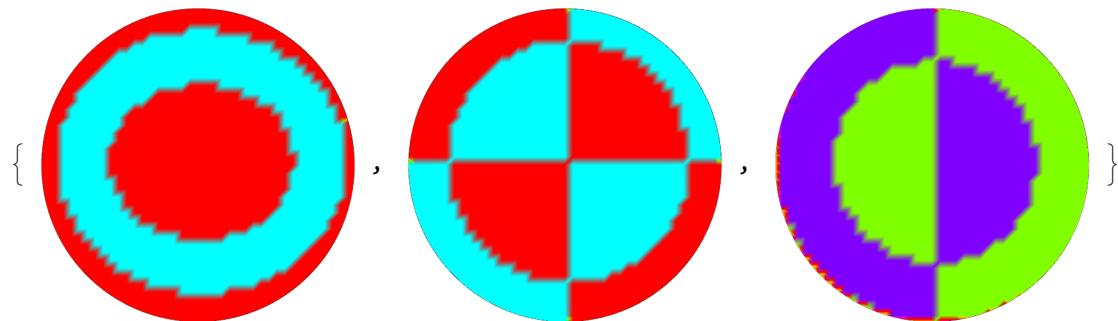
In[210]:=

```
(* Linearly polarised in x*)
Map[Int, Ep[E00, 0, 0]]
Map[Ph, Ep[E00, 0, 0]]
```

Out[210]=



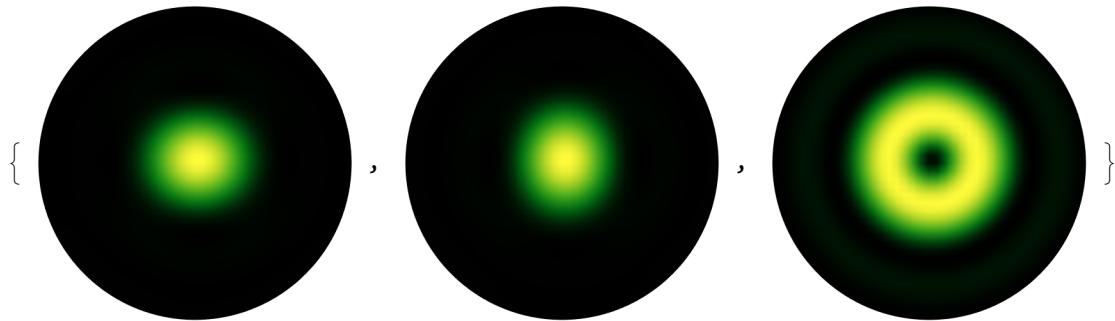
Out[211]=



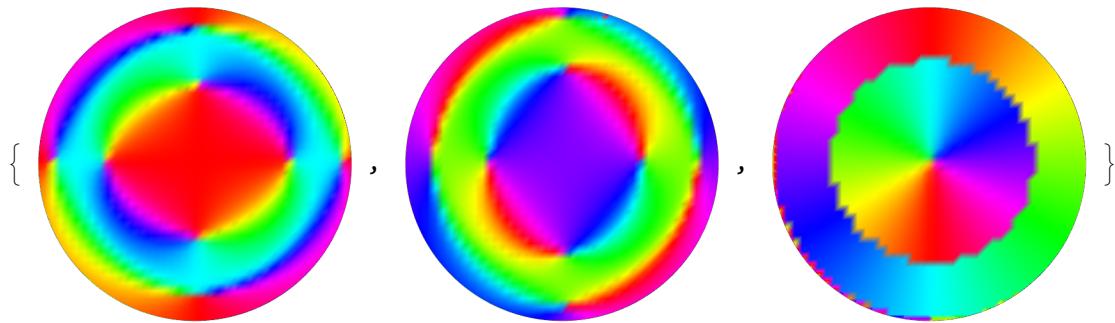
In[214]:=

```
(*Circular polarization in *)
Map[Int, Ep[E00, π / 4, -π / 2]]
Map[Ph, Ep[E00, π / 4, -π / 2]]
```

Out[214]=



Out[215]=



## Example 2 : LG01

In[216]:=

```
AbsoluteTiming[E01 = Efz0[1.5, 0.07, 1, 0.95, 1][[1]]]
```

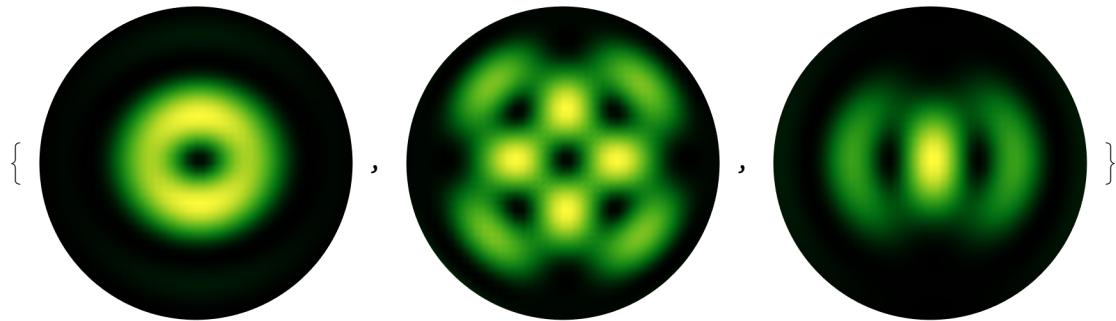
Out[216]=

```
29.8555
```

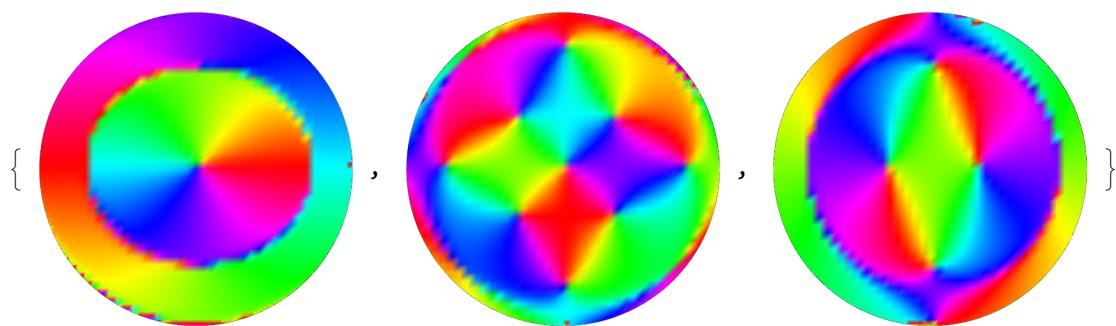
In[219]:=

```
(* Linearly polarised in x*)
Map[Int, Ep[E01, 0, 0]]
Map[Ph, Ep[E01, 0, 0]]
```

Out[219]=



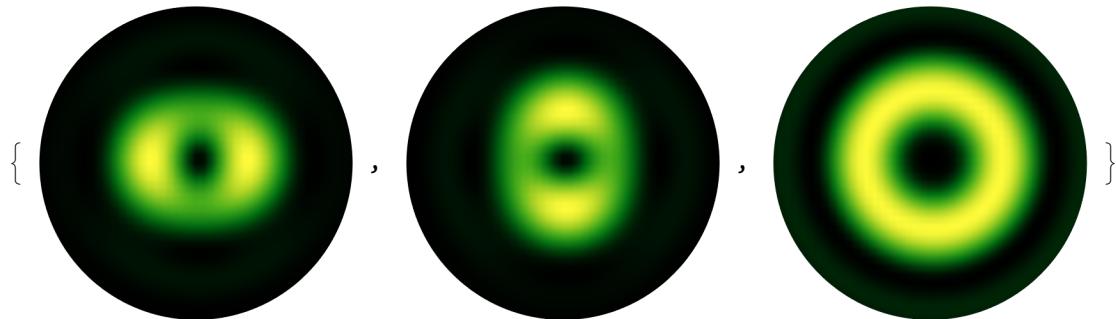
Out[220]=



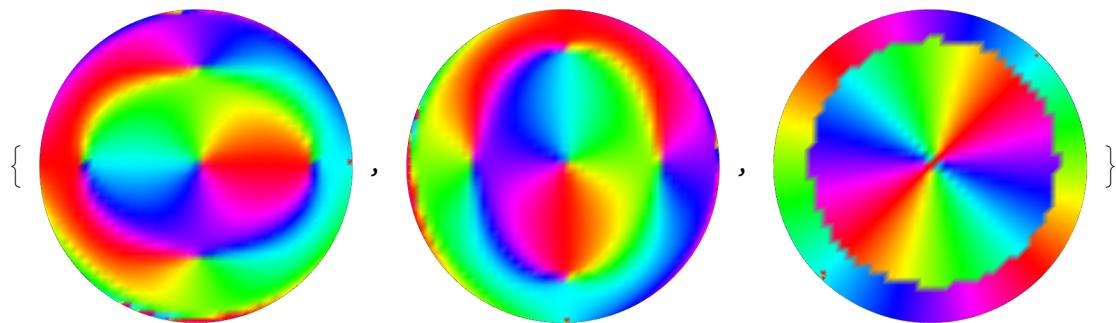
In[221]:=

```
(*Circular polarization in *)
Map[Int, Ep[E01, π / 4, +π / 2]]
Map[Ph, Ep[E01, π / 4, +π / 2]]
```

Out[221]=



Out[222]=



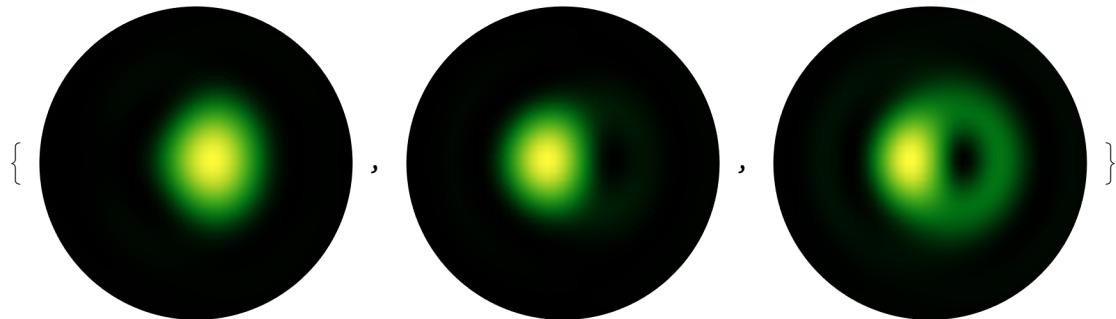
### Example 3 : Vector beams

```
In[*]:= AbsoluteTiming[E01 = Efz0[1.5, 0.07, 1, 0.95, 1]][[1]]
Out[*]= 40.7425
```

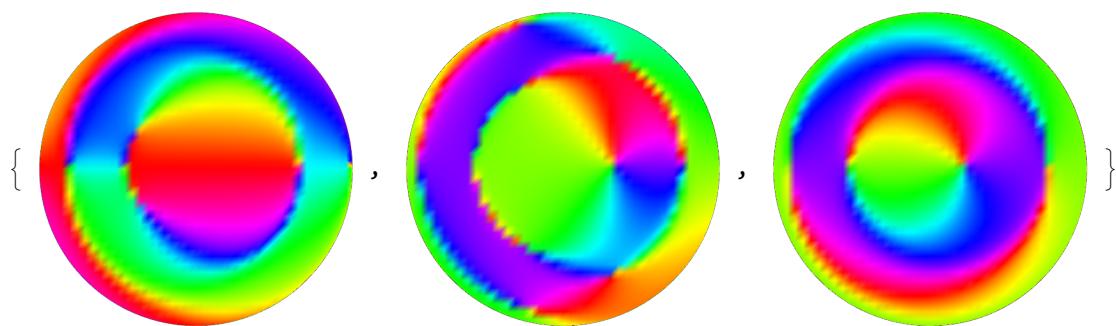
In[223]:=

```
(* Lemon Topology*)
Map[Int, Ep[E00, \[Pi]/4, \[Pi]/2] + Ep[E01, \[Pi]/4, -\[Pi]/2]]
Map[Ph, Ep[E00, \[Pi]/4, \[Pi]/2] + Ep[E01, \[Pi]/4, -\[Pi]/2]]
```

Out[223]=



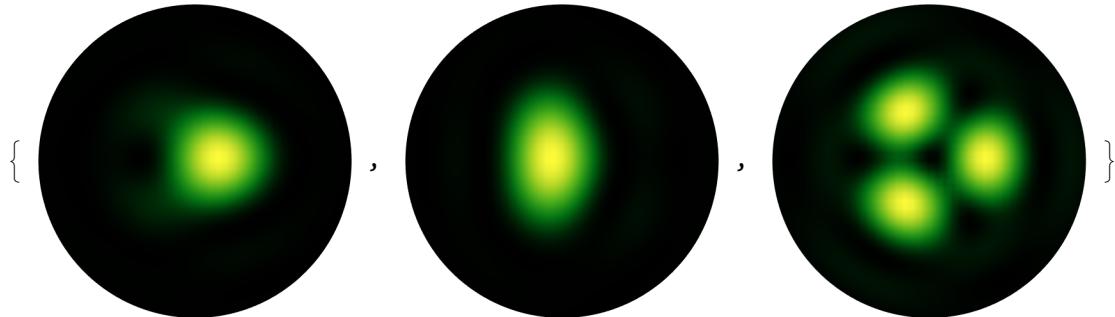
Out[224]=



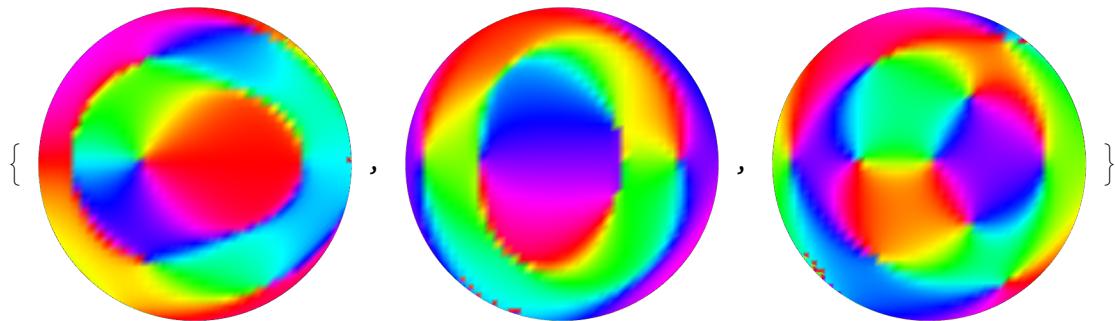
In[225]:=

```
(* Star Topology*)
Map[Int, Ep[E00, π/4, -π/2] + Ep[E01, π/4, π/2]]
Map[Ph, Ep[E00, π/4, -π/2] + Ep[E01, π/4, π/2]]
```

Out[225]=



Out[226]=



## Definitions

In[206]:=

```
Ph[ψ_] := ListDensityPlot[Arg[ψ], ColorFunction → (Hue[Rescale[#, {0, 2π}, {0, 1}]] &),
  PlotRange → All, ColorFunctionScaling → False, Frame → None,
  BoundaryStyle → Directive[Transparent, Thick], DataRange → {{-1.5, 1.5}, {-1.5, 1.5}},
  RegionFunction → Function[{x, y, f}, (x)^2 + (y)^2 < 2]];
Int[ψ_] := ListDensityPlot[Abs[ψ]^2,
  ColorFunction → ColorData["AvocadoColors"], PlotRange → All, Frame → None,
  BoundaryStyle → Directive[Transparent, Thick], DataRange → {{-1.5, 1.5}, {-1.5, 1.5}},
  RegionFunction → Function[{x, y, f}, (x)^2 + (y)^2 < 2]];
k = 2π;
lg[ρ_, ℓ_] := ρ^Abs[ℓ] e^{-ρ^2};
```

## Field produced by the X-polarised component a LG beam

In[195]:=

```
(* Resulting X-
component of the polarization due to a initially X-polarised LG01 mode *)
LGxx[ρ_, φ_, z_, ℓ_, a_, b_] := e^(iℓφ) NIntegrate[ √Cos[θ] Sin[θ] lg[ (Sin[θ]/Sin[α]), ℓ]
  (1/2 (1 + Cos[θ]) BesselJ[ℓ, k ρ Sin[θ]] + 1/4 (1 - Cos[θ]) (BesselJ[ℓ + 2, k ρ Sin[θ]] e^(i2φ) +
   BesselJ[ℓ - 2, k ρ Sin[θ]] e^(-i2φ)) ) * Exp[I k z Cos[θ]], {θ, a, b}];

(* Resulting Y-component of the polarization due to a initially X-
polarised LG01 mode *) LGxy[ρ_, φ_, z_, ℓ_, a_, b_] := e^(iℓφ) (-1/(2 i)) NIntegrate[
  √Cos[θ] Sin[θ] lg[ (Sin[θ]/Sin[α]), ℓ] ((Cos[θ] - 1) (BesselJ[ℓ + 2, k ρ Sin[θ]] e^(2 iφ) -
   BesselJ[ℓ - 2, k ρ Sin[θ]] e^(-2 iφ)) ) * Exp[I k z Cos[θ]], {θ, a, b}];

(* Resulting Z-component of the polarization due to a initially X-
polarised LG01 mode *) LGxz[ρ_, φ_, z_, ℓ_, a_, b_] :=
  e^(iℓφ) (-i/(2)) NIntegrate[ √Cos[θ] Sin[θ]^2 lg[ (Sin[θ]/Sin[α]), ℓ] (BesselJ[ℓ + 1, k ρ Sin[θ]] e^(iφ) -
   BesselJ[ℓ - 1, k ρ Sin[θ]] e^(-iφ)) * Exp[I k z Cos[θ]], {θ, a, b}];
```

## Field produced by the Y-polarised component a LG beam

In[197]:=

```

(* Resulting X-
component of the polarization due to a initially Y-polarised LG01 mode *)
LGyx[ρ_, φ_, z_, ℓ_, a_, b_] := e^(iℓφ) (-1/(2 π)) NIntegrate[ Sqrt[Cos[θ]] Sin[θ] lg[(Sin[θ]/Sin[α]), ℓ]
  ((Cos[θ] - 1) (BesselJ[ℓ + 2, k ρ Sin[θ]] e^(2 i φ) - BesselJ[ℓ - 2, k ρ Sin[θ]] e^(-2 i φ))) *
  Exp[I k z Cos[θ]], {θ, a, b}];

(* Resulting X-
component of the polarization due to a initially Y-polarised LG01 mode *)
LGyy[ρ_, φ_, z_, ℓ_, a_, b_] :=
  e^(iℓφ)
  NIntegrate[ Sqrt[Cos[θ]] Sin[θ] lg[(Sin[θ]/Sin[α]), ℓ] (1/2 (1 + Cos[θ]) BesselJ[ℓ, k ρ Sin[θ]] -
  1/4 (1 - Cos[θ]) (BesselJ[ℓ + 2, k ρ Sin[θ]] e^(i 2 φ) + BesselJ[ℓ - 2, k ρ Sin[θ]] e^(-i 2 φ))) *
  Exp[I k z Cos[θ]], {θ, a, b}];

(* Resulting X-component of the polarization due to a initially Y-
polarised LG01 mode *) LGyz[ρ_, φ_, z_, ℓ_, a_, b_] :=
  e^(iℓφ) (-1/(2 π)) NIntegrate[ Sqrt[Cos[θ]] Sin[θ]^2 lg[(Sin[θ]/Sin[α]), ℓ] (BesselJ[ℓ + 1, k ρ Sin[θ]] e^(i φ) +
  BesselJ[ℓ - 1, k ρ Sin[θ]] e^(-i φ)) * Exp[I k z Cos[θ]], {θ, a, b}];

```

## Construction of the beams

In[200]:=

```

Efz0[xmax_, dx_, ℓ_, NA_, n_] := Block[{EXx, EXy, EXz, EYx, EYy, EYz, α = ArcSin[NA / n]},

EXx = ParallelTable[LGxx[ $\sqrt{x^2 + y^2}$ , ArcTan[x, y], 0, ℓ, 0, α],
{y, -xmax, xmax, dx}, {x, -xmax, xmax, dx}];

EXy = ParallelTable[LGxy[ $\sqrt{x^2 + y^2}$ , ArcTan[x, y], 0, ℓ, 0, α],
{y, -xmax, xmax, dx}, {x, -xmax, xmax, dx}];

EXz = ParallelTable[LGxz[ $\sqrt{x^2 + y^2}$ , ArcTan[x, y], 0, ℓ, 0, α],
{y, -xmax, xmax, dx}, {x, -xmax, xmax, dx}];

EYx = ParallelTable[LGyx[ $\sqrt{x^2 + y^2}$ , ArcTan[x, y], 0, ℓ, 0, α],
{y, -xmax, xmax, dx}, {x, -xmax, xmax, dx}];

EYy = ParallelTable[LGyy[ $\sqrt{x^2 + y^2}$ , ArcTan[x, y], 0, ℓ, 0, α],
{y, -xmax, xmax, dx}, {x, -xmax, xmax, dx}];

EYz = ParallelTable[LGyz[ $\sqrt{x^2 + y^2}$ , ArcTan[x, y], 0, ℓ, 0, α],
{y, -xmax, xmax, dx}, {x, -xmax, xmax, dx}];

Return[{EXx, EXy, EXz, EYx, EYy, EYz}]
];

```

In[201]:=

```
Efprop[xmax_, dx_, zmin_, zmax_, dz_, ℓ_, NA_, n_] :=
  Block[{EXx, EXy, EXz, EYx, EYy, EYz, α = ArcSin[NA / n]},
    EXx = ParallelTable[LGxx[ $\sqrt{x^2 + y^2}$ , ArcTan[x, y], z, ℓ, θ, α],
      {y, -xmax, xmax, dx}, {x, -xmax, xmax, dx}, {z, zmin, zmax, dz}];
    EXy = ParallelTable[LGxy[ $\sqrt{x^2 + y^2}$ , ArcTan[x, y], z, ℓ, θ, α],
      {y, -xmax, xmax, dx}, {x, -xmax, xmax, dx}, {z, zmin, zmax, dz}];
    EXz = ParallelTable[LGxz[ $\sqrt{x^2 + y^2}$ , ArcTan[x, y], z, ℓ, θ, α],
      {y, -xmax, xmax, dx}, {x, -xmax, xmax, dx}, {z, zmin, zmax, dz}];
    EYx = ParallelTable[LGyx[ $\sqrt{x^2 + y^2}$ , ArcTan[x, y], z, ℓ, θ, α],
      {y, -xmax, xmax, dx}, {x, -xmax, xmax, dx}, {z, zmin, zmax, dz}];
    EYy = ParallelTable[LGyy[ $\sqrt{x^2 + y^2}$ , ArcTan[x, y], z, ℓ, θ, α],
      {y, -xmax, xmax, dx}, {x, -xmax, xmax, dx}, {z, zmin, zmax, dz}];
    EYz = ParallelTable[LGyz[ $\sqrt{x^2 + y^2}$ , ArcTan[x, y], z, ℓ, θ, α],
      {y, -xmax, xmax, dx}, {x, -xmax, xmax, dx}, {z, zmin, zmax, dz}];
    Return[{EXx, EXy, EXz, EYx, EYy, EYz}]
  ];
```

In[202]:=

```
(*Constructing the arbitrary polarised*)
Ep[E_, α_, β_] := {Cos[α] * E[1] + Sin[α] Exp[I β] * E[4],
  Cos[α] * E[2] + Sin[α] Exp[I β] * E[5], Cos[α] * E[3] + Sin[α] Exp[I β] * E[6]};
```