

# CONTACT STRUCTURES & WEINSTEIN<sup>3</sup>

Toronto, April 7<sup>th</sup>, 2024

by Marullo Ghini Bettol (1006836676)

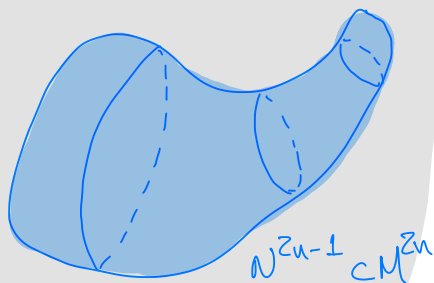
- REFERENCES:
- Thurston, 3-dim. Geometry & Topology
  - Hutchings, "Tambe's proof of the Weinstein conjecture in dimension three".

MOTIVATION FOR CONTACT STRUCTURES: Contact structures can be thought of as odd-dimensional counterparts to symplectic structures. More precisely, contact structures arose in the context of Hamiltonian systems with some sort of symmetry.

↖  $(M, \omega)$  symplectic  $\Rightarrow \dim M$  is even (b/c  $\omega$  non-degenerate)

The symmetry of the system (e.g. conservation of momentum, conservation of energy, etc.) makes the configuration space have less degrees of freedom. This amounts to the dimension of the configuration space to decrease.

Configuration space  $(M^{2n}, \omega)$



→ Since  $\dim M$  is even, in systems w/ symmetries, the motion becomes confined to odd-dimensional spaces.

Q: How to study this (since odd-dim. spaces do not admit symplectic structures)?

A: Contact structures.

History: Contact structures first studied by Hamilton / Jacobi / Huygens, later by Darboux, and others.

Because of some symmetry, configuration space reduces to a submfd.

Focus the discussion on 3-manifolds from now on  $\downarrow$

**Def: (PLANE FIELD)** A plane field  $\xi$  on a manifold  $M$  is a subbundle of  $TM$  such that  $\xi_p = T_p M \cap \xi$  is a 2-dimensional subspace of  $T_p M$  for each  $p \in M$ .

**Ex 1:** Let  $M^3 := \Sigma \times S^1$ ,  $\Sigma$  is a surface. Let  $(x, \theta)$  be local coordinates for  $\Sigma \times S^1$ . Then, for each  $p = (x, \theta)$ , set  $\xi_p := T_x \Sigma \subset T_p M$ .  $\leftarrow$  plane field on  $M$

**Ex 2:** Let  $\alpha \in \Omega^1(M)$ ; i.e., at each  $p \in M$ ,  $\alpha_p: T_p M \rightarrow \mathbb{R}$  linear. Then,  $\ker \alpha_p$  is either a plane or all of  $T_p M$ . Assuming  $\alpha$  never has all of  $T_p M$  as its kernel, then  $\xi := \ker \alpha$  is a plane field.  $\hookrightarrow$  Exm: In the previous example,  $\alpha = d\theta$  defines  $\xi$ .

Proposed by Weinstein (1974)

**Def: (CONTACT STRUCTURE)** A plane field  $\xi$  is a contact structure if for any 1-form  $\alpha$  with  $\xi = \ker \alpha$ , we have  $\alpha \wedge d\alpha \neq 0$ .  
"CONTACT FORM"  $\nearrow$   $\alpha$   $\nwarrow$  locally or globally defined

Exm: From calculus,  $\alpha \wedge d\alpha \neq 0 \iff d\alpha|_{\xi} \neq 0$ .

Obs: The plane field in the first example above is NOT a contact structure  $\leftarrow$  That plane field was defined by  $\alpha = d\theta$ .  
But  $d\alpha = d(d\theta) = 0$ .

Ex 3: Take  $\mathbb{R}^3$  with standard Cartesian coordinates  $(x, y, z)$  and the 1-form

$$\alpha = dz + x dy.$$

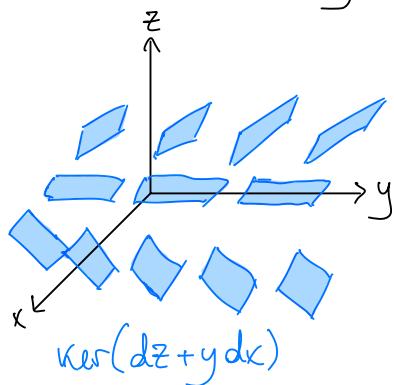
→ Sometimes called "STANDARD CONTACT STRUCTURE" on  $\mathbb{R}^3$

Note that  $d\alpha = dx \wedge dy \longrightarrow \alpha \wedge d\alpha = dz \wedge dx \wedge dy \neq 0$

i.e.,  $\alpha$  is a contact form in  $\mathbb{R}^3$  and  $\xi := \ker \alpha$  is a contact structure

Compute  $\xi$ : at  $(x, y, z)$ ,

$$\xi = \text{span} \left\{ \frac{\partial}{\partial x}, x \frac{\partial}{\partial z}, -\frac{\partial}{\partial y} \right\}$$



Ex 4: Again, take  $\mathbb{R}^3$  but w/ cylindrical coordinates  $(r, \theta, z)$  and

$$\alpha = dz + r^2 d\theta.$$

Then  $\alpha \wedge d\alpha = 2r dr \wedge d\theta \wedge dz \neq 0 \Rightarrow \xi := \ker \alpha$  is a contact structure

At  $(r, \theta, z)$ ,  $\xi = \text{span} \left\{ \frac{\partial}{\partial r}, r^2 \frac{\partial}{\partial z} - \frac{\partial}{\partial \theta} \right\}.$

→ if  $r=0$  (i.e., on the  $z$ -axis),  $\xi$  is horizontal.  
As we move out, the planes twist clockwise

Def: Two contact structures  $\xi_1$  and  $\xi_2$  on  $M$  are contactomorphic if there is a diffeomorphism  $f: M \rightarrow M$  st.  $f_* (\xi_1) = \xi_2$ .

← orientation preserving

e.g.: the structures from Exs 3 and 4 are contactomorphic

**Ex 5:** Take the unit 3-sphere  $S^3$  and

$$\beta = (x_1 dy_1 - y_1 dx_1 + x_2 dy_2 - y_2 dx_2) \Big|_{S^3}$$

Then  $\kappa := \ker \beta$  is a contact structure on  $S^3$ .

Note:  $(S^3 \setminus \{N\}, \kappa|_{S^3 \setminus \{N\}})$  is contactomorphic to  $(\mathbb{R}^3, \tilde{\kappa})$

(can see this using stereographic coordinates).

**REEB VECTOR FIELDS:** Much like symplectic forms defining Hamiltonian vector fields on symplectic manifolds, contact forms have a "special" vector field associated to them

Consider  $(M, \omega)$  and  $H \in C^\infty(M, \mathbb{R})$ .  
Then, Hamiltonian vec. field is defined as  $X_H := -(\omega^\#)^{-1}(dH)$ .  
Flow of  $X_H$  determines the motion.

French Mathematician Georges Reeb

**Def: (REEB VECTOR FIELD)** A contact form  $\alpha$  on  $M$  determines a vector field  $R_\alpha$  on  $M$ , called the Reeb vector field, characterized by

$$d\alpha(R_\alpha, \cdot) = 0, \quad \alpha(R_\alpha) = 1$$

$$\Leftrightarrow R_\alpha \in \ker d\alpha$$

**Prop:** If  $\alpha$  is a contact form for some contact structure, the Reeb vector field  $R$  is unique since  $\dim \ker d\alpha = 1$  and we need  $\alpha(R) = 1$ .

$d\alpha|_{T_p M}$  is a skew-symmetric form of maximal rank  $2n-1$ .

**Prop:** Different contact forms whose kernels produce the same contact  $\gamma$

structure will output different Reeb vector fields (i.e., the dynamical system will be different).

**Ex 6:** Consider the following family of contact forms on the 3-sphere  $S^3 \subset \mathbb{R}^4$ :

$$\alpha_t := (x_1 dy_1 - y_1 dx_1) + (1+t)(x_2 dy_2 - y_2 dx_2),$$

where  $t \geq 0$  is a real parameter. The Reeb vector field of  $\alpha_t$  is

$$R_{\alpha_t} = \left( x_1 \frac{\partial}{\partial y_1} - y_1 \frac{\partial}{\partial x_1} \right) + \frac{1}{1+t} \left( x_2 \frac{\partial}{\partial y_2} - y_2 \frac{\partial}{\partial x_2} \right).$$

Note that the flow of  $R_{\alpha_0}$  defines the Hopf fibration. In particular, all orbits of  $R_{\alpha_0}$  are closed.

(If  $t \in \mathbb{R}_{\geq 0} \setminus \mathbb{Q}$ , however,  $R_{\alpha_t}$  has only  $\mathbb{Z}$  periodic orbits....)

---

**WEINSTEIN CONJECTURE:** Before moving on, we need a technical definition:

**Def: (CLOSED ORBIT)** Let  $M$  be a closed manifold and let  $X \in \mathcal{X}(M)$ . A closed orbit of  $X$  is a map

$$\gamma: \mathbb{R}/c\mathbb{Z} \longrightarrow M,$$

for some  $c > 0$ , satisfying

$$\frac{d}{dt} \gamma(t) = X(\gamma(t)).$$

GOAL: Understand when vector fields have closed orbits. Are there cases in which all vector fields have closed orbits?

Prmk: In some special cases such as the 3-torus, it is fairly easy to construct vector fields with no closed orbits.

However, things get complicated very soon: do all vector fields in  $S^3$  have closed orbits? As it turns out, the answer is NO. Examples of vector fields of increasing regularity on  $S^3$  with no closed orbits were produced by

Schweizer (1974)  $\rightarrow C^1$

Harrison (1998)  $\rightarrow C^2$

Kupurbeg (1994)  $\rightarrow C^\infty$

WEINSTEIN CONJECTURE (1978): Let  $M$  be a closed oriented odd-dimensional manifold with a contact form  $\alpha$ . Then, the associated Reeb vector field  $R_\alpha$  has a closed orbit.

$\rightarrow$  Remains open in several cases but it was proven for all closed 3-manifolds by Clifford Taubes in 2007.

Thm: (TAUBES, 2007) If  $M$  is a closed oriented 3-manifold with a contact form  $\alpha$ , then the Reeb vector field  $R_\alpha$  has a closed orbit.

Taubes proof used some fancy techniques like Seiberg-Witten theory.

**Ex 7:** Let  $Q$  be a smooth manifold. From a classical construction, there is a canonical 1-form  $\omega$  on the cotangent bundle  $T^*Q$ . Let  $\pi: T^*Q \rightarrow Q$  be the standard projection. If  $q \in Q$  and  $p \in T_q^*Q$ , then  $\omega: T_{(q,p)}T^*Q \rightarrow \mathbb{R}$  is given by

$$T_{(q,p)}T^*Q \xrightarrow{\pi_*} T_qQ \xrightarrow{p} \mathbb{R}.$$

Explicitly, if  $q^1, \dots, q^n, p_1, \dots, p_n$  are local coordinates on  $T^*Q$ , we can write

$$\omega = \sum_{i=1}^n p_i dq^i.$$

Important object when studying classical mechanics

It follows that  $d\omega$  is a symplectic form on  $T^*Q$ . Now, suppose we equip  $Q$  with a Riem. metric  $g$ . This induces a metric on  $T^*Q$  and we can consider the unit cotangent bundle

$$ST^*Q := \{p \in T^*Q : \|p\| = 1\}.$$

The restriction of  $\omega$  to  $ST^*Q$  gives a contact form. It turns out that the Reeb vector field  $R_\omega|_{ST^*Q}$  agrees with the geodesic flow under the identification  $T^*Q \simeq TQ$ .

isom. b/c  $g$  is Riemannian (in part, non-degenerate)

Thus, closed orbits in  $ST^*Q$  from  $R_\omega|_{ST^*Q}$  are equivalent to closed geodesics in  $Q$ !

If  $Q$  is compact, then  $ST^*Q$  is also compact, and we

can apply Weinstein's conjecture.



In this case, this is equivalent to Ljusternik-Fet  
affirming that every compact Riemannian manifold  
has at least one closed geodesic.