

# Nuclear density from *ab initio* theory

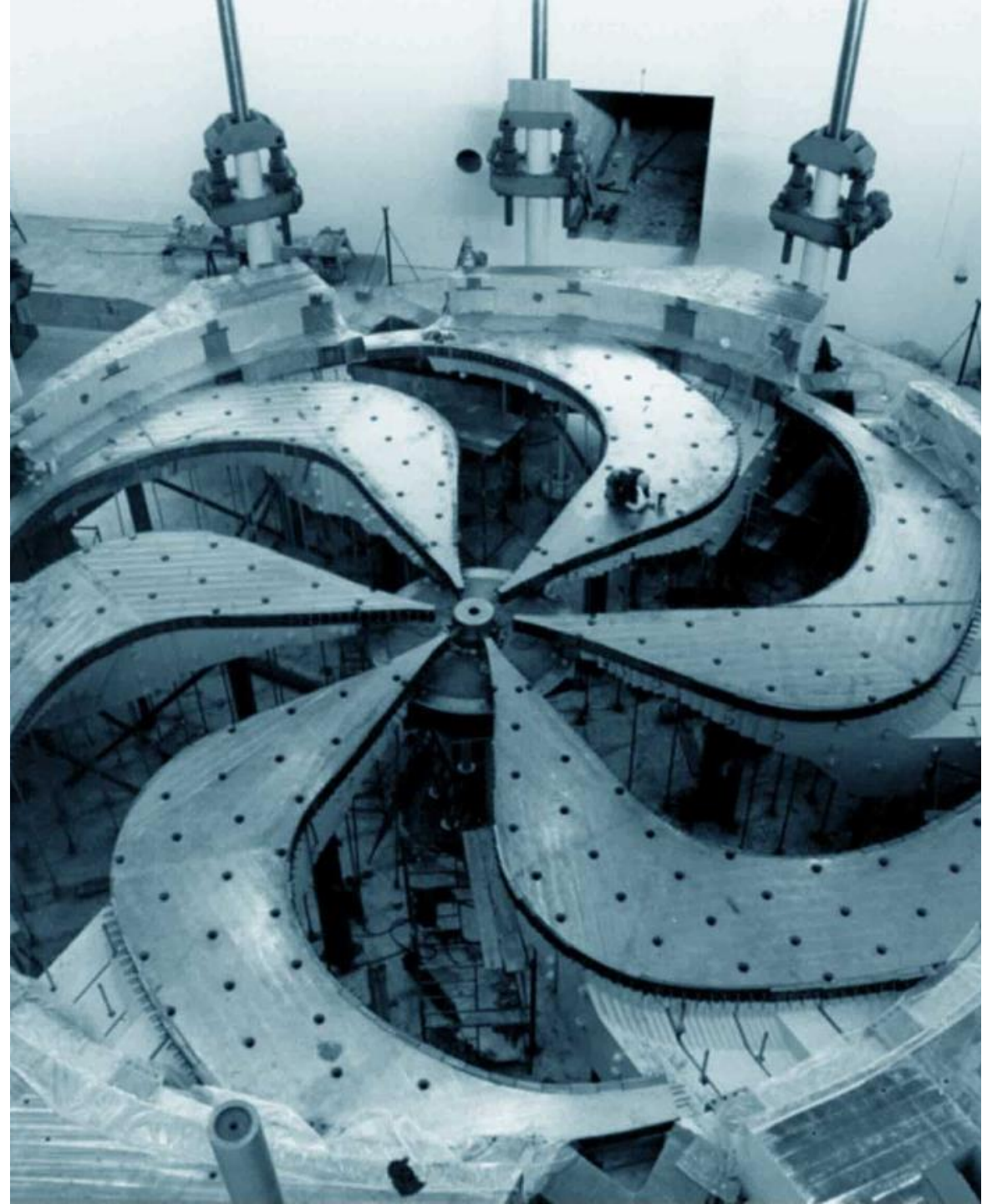
Michael Gennari

TRIUMF – Undergraduate researcher from  
University of Waterloo

In collaboration with

Petr Navrátil and Matteo Vorabbi

2018-08-11



# Outline

## 1. **Ab Initio No-Core Shell Model (NCSM)**

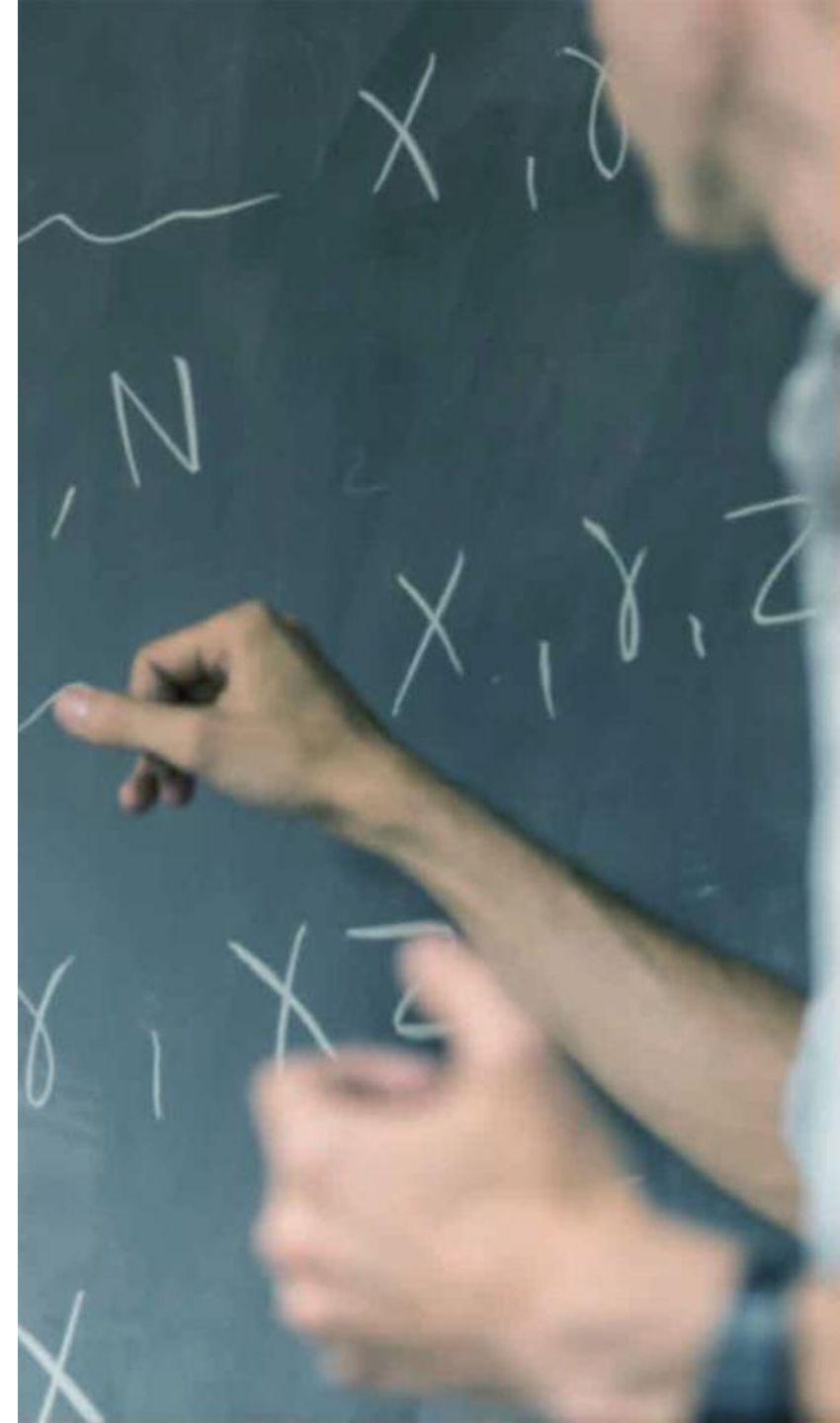
- a) Introduction to NCSM
- b) Factorization of spurious COM motion
- c) Coordinate space nonlocal density

## 2. **Chiral Effective Field Theory Interactions (EFT)**

- a) Chiral expansion in many-body interactions
- b) Convergence


## 3. **Results**

- a) Nonlocal density
- b) Applications to optical potentials
- c) Applications to density functional theory



# No-core shell model (NCSM)

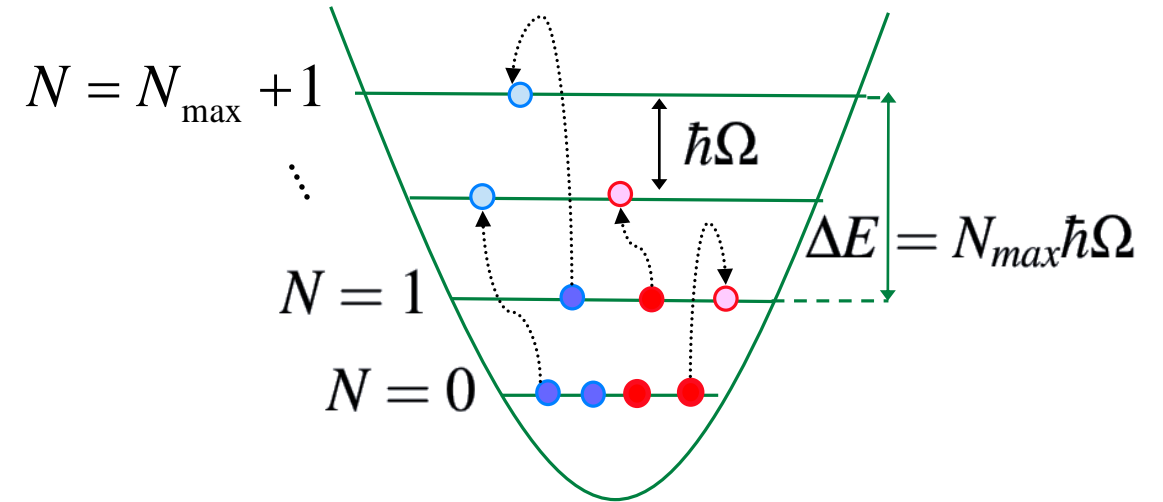
- NCSM is an *ab initio* approach to solve the many-body Schrödinger equation for bound states (narrow resonances) starting from [high-precision NN+NNN interactions](#)
- Uses large (but finite!) expansions in HO many-body basis states



$$\Psi^A = \sum_{N=0}^{N_{max}} \sum_i c_{Ni} \Phi_{Ni}^A$$


- Translational invariance of the internal wave function is preserved when single-particle Slater Determinant (SD) basis is used with  $N_{max}$  truncation

$$\langle \vec{r}_1 \cdots \vec{r}_A \vec{\sigma}_1 \cdots \vec{\sigma}_A \vec{\tau}_1 \cdots \vec{\tau}_A | A\lambda JM \rangle_{SD} = \langle \vec{\xi}_1 \cdots \vec{\xi}_{A-1} \vec{\sigma}_1 \cdots \vec{\sigma}_A \vec{\tau}_1 \cdots \vec{\tau}_A | A\lambda JM \rangle \varphi_{000}(\vec{\xi}_0)$$



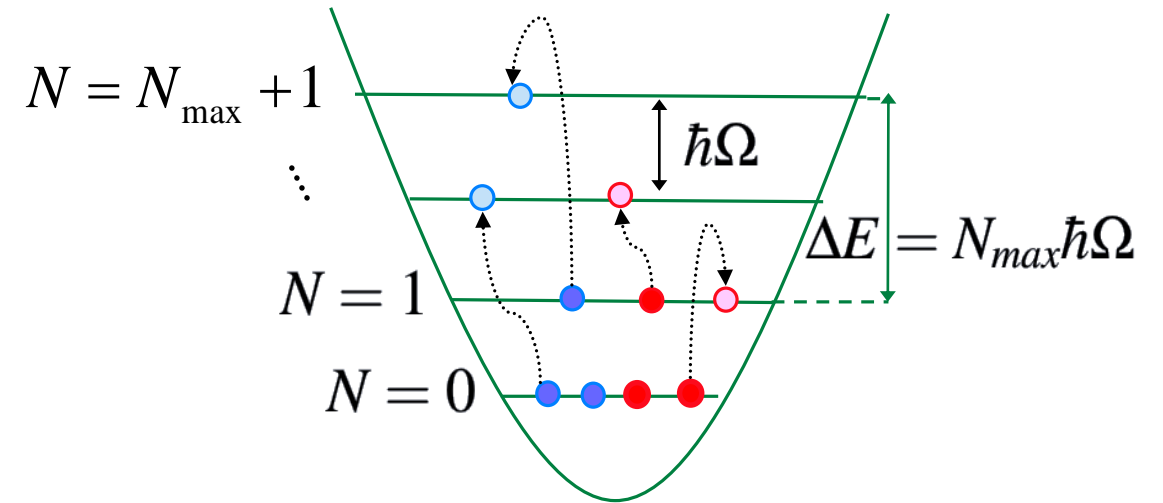
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Intrinsic  
wavefunction

COM  
wavefunction

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# Coordinate form of the density

$$\begin{aligned}
 & \langle A\lambda_f J_f M_f | \rho_{op}(\vec{r} - \vec{R}, \vec{r}' - \vec{R}) | A\lambda_i J_i M_i \rangle \\
 &= \left( \frac{A}{A-1} \right)^{\frac{3}{2}} \sum_{\hat{J}_f} \frac{1}{\hat{J}_f} (J_i M_i K k | J_f M_f) \left( Y_l^* \left( \widehat{\vec{r} - \vec{R}} \right) Y_{l'}^* \left( \widehat{\vec{r}' - \vec{R}} \right) \right)_k^{(K)} \\
 &\times R_{n,l} \left( \sqrt{\frac{A}{A-1}} |\vec{r} - \vec{R}| \right) R_{n',l'} \left( \sqrt{\frac{A}{A-1}} |\vec{r}' - \vec{R}| \right) \\
 &\times (M^K)_{n,l,n',l',n_1,l_1,n_2,l_2}^{-1} (-1)^{l_1+l_2+K+j_2-\frac{1}{2}} \hat{J}_1 \hat{J}_2 \hat{K} \begin{Bmatrix} j_1 & j_2 & K \\ l_2 & l_1 & 1/2 \end{Bmatrix} \\
 &\times \frac{(-1)}{\hat{K}} {}_{SD} \langle A\lambda_f J_f || (a_{n_1 l_1 j_1}^\dagger \tilde{a}_{n_2 l_2 j_2})^{(K)} || A\lambda_i J_i \rangle_{SD}
 \end{aligned}$$

PHYSICAL REVIEW C **97**, 034619 (2018)

## Microscopic optical potentials derived from *ab initio* translationally invariant nonlocal one-body densities

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Matteo Vorabbi,<sup>†</sup> Angelo Calci, and Petr Navrátil<sup>‡</sup>

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## Nonlocal translationally invariant density

- Translationally invariant nuclear density is obtained from intrinsic wavefunction
- Slater determinant description is advantageous for  $A > 4$
- When Slater determinant description is used, there is a spurious COM contribution
- It is possible to exactly remove this contamination

## Normalization

$$\int d\vec{x} \langle A\lambda JM | \rho_{op}^{phys}(\vec{x}) | A\lambda JM \rangle = A$$



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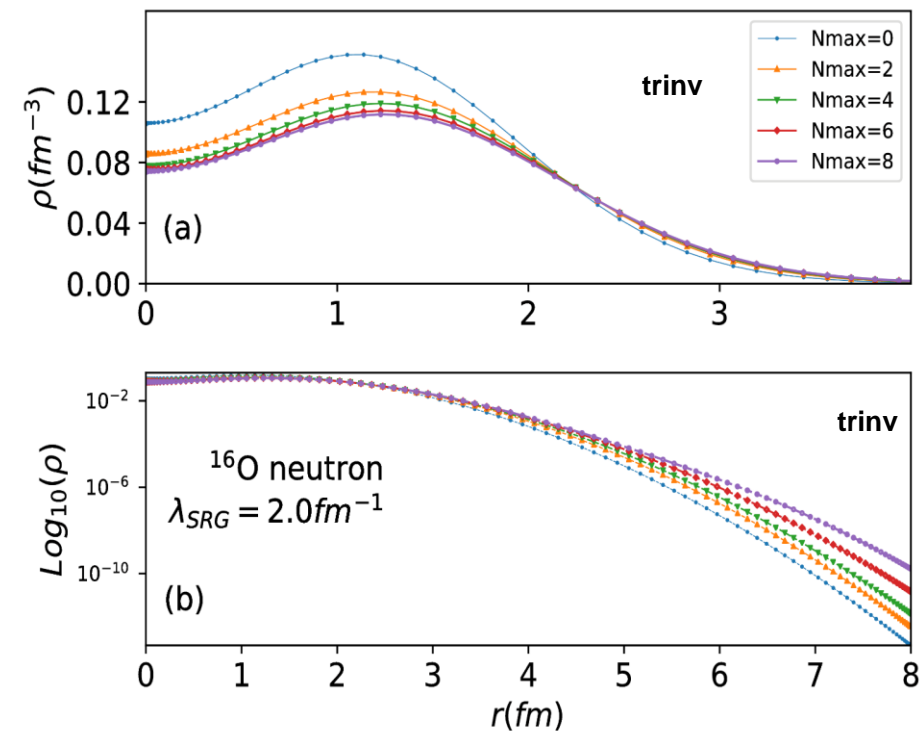
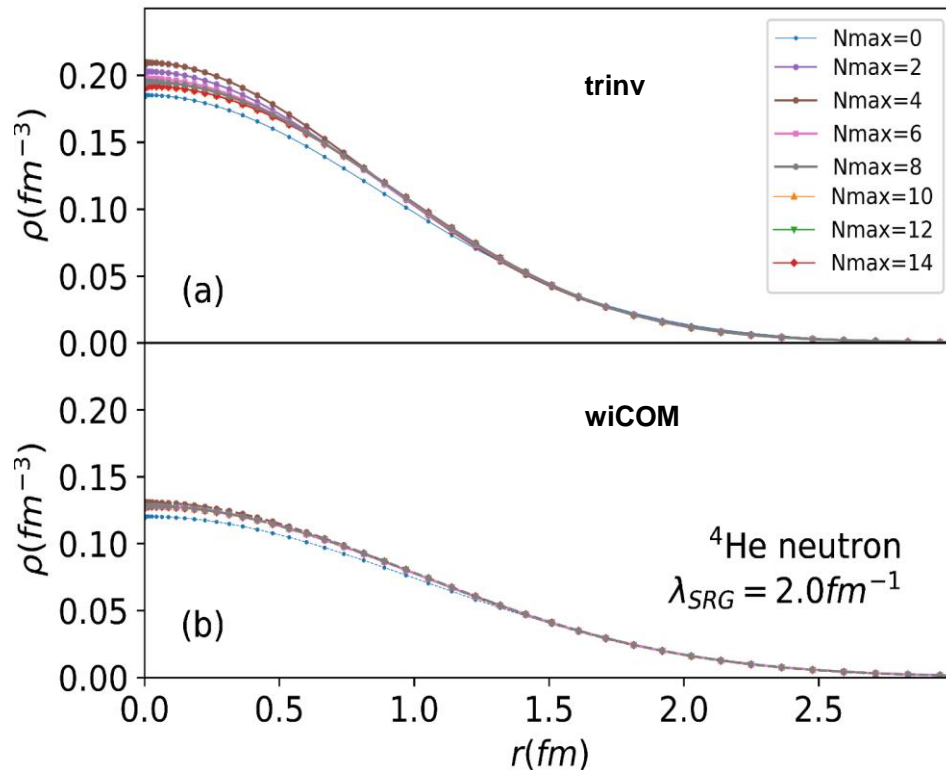
# NN and 3N interactions – $N^4\text{LO}(500)+3\text{N}^2\text{LO}$

## NN systematic from LO to $N^4\text{LO}$

- D. R. Entem, N. Kaiser, R. Machleidt, and Y. Nosyk, Phys. Rev. C 91, 014002 (2015)
- D. R. Entem, R. Machleidt, and Y. Nosyk, Phys. Rev. C 96, 024004 (2017)

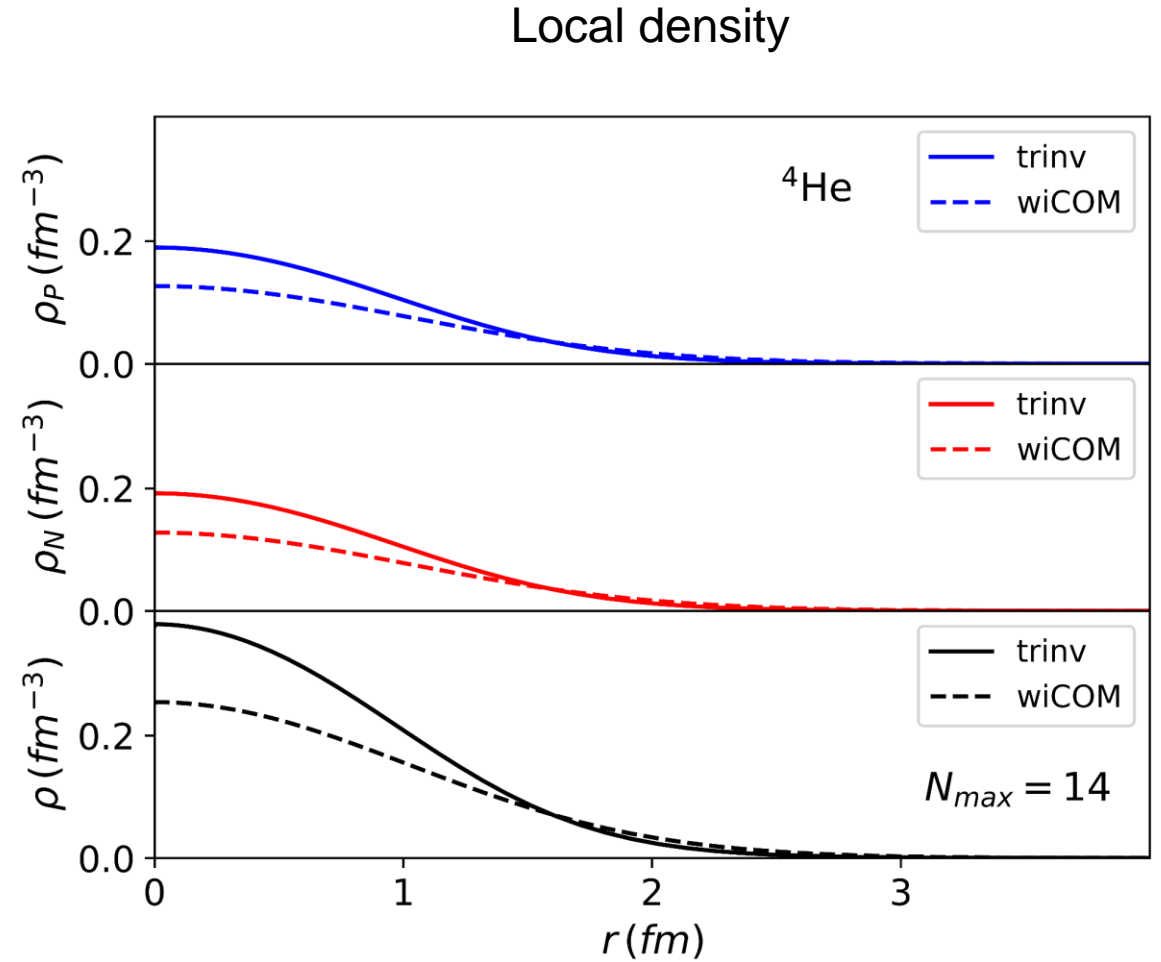
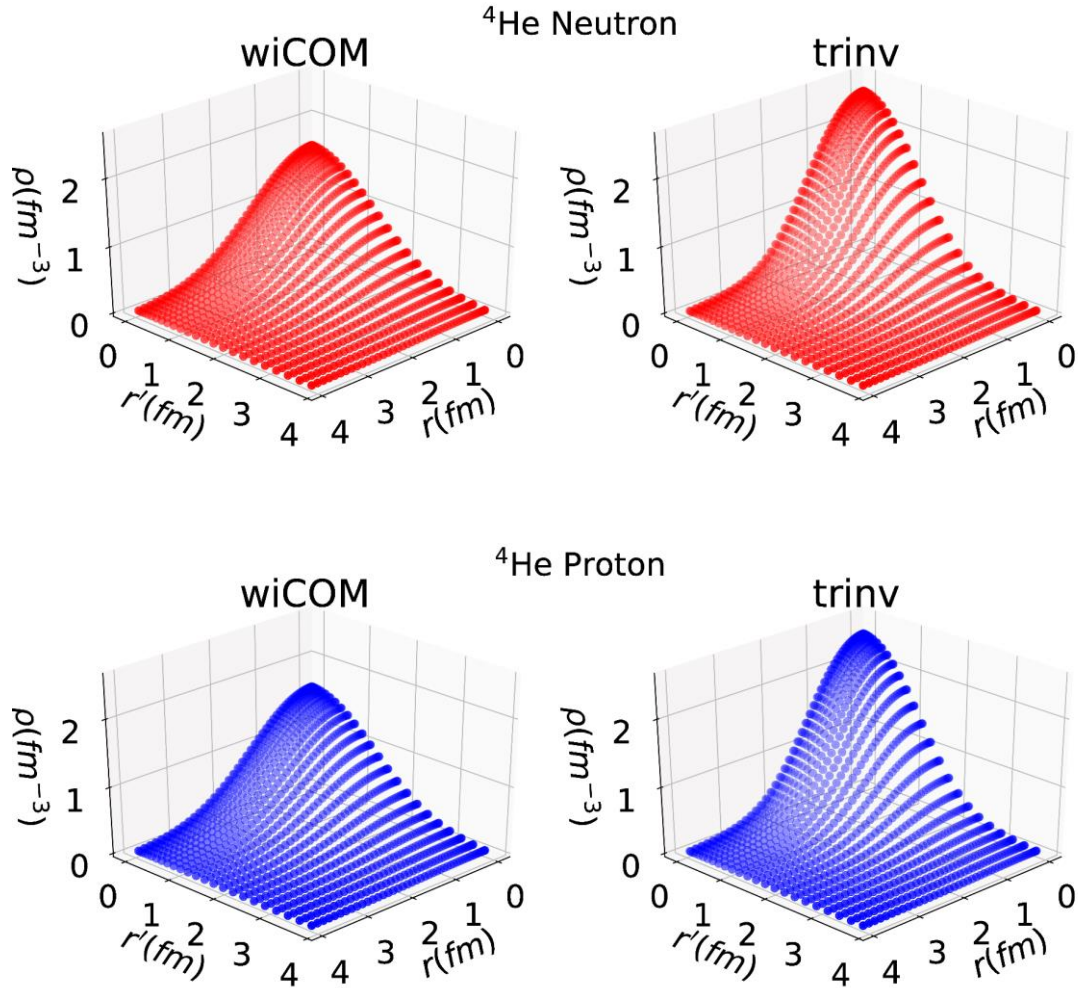
## 3N at $N^2\text{LO}$

- Navrátil, 650 MeV local cut-off and 500 MeV nonlocal cut-off



# Density of ground state ${}^4,{}^6,{}^8\text{He}$ , ${}^{12}\text{C}$ , ${}^{16}\text{O}$ with NN- $\text{N}^4\text{LO}(500)+3\text{NInI}$

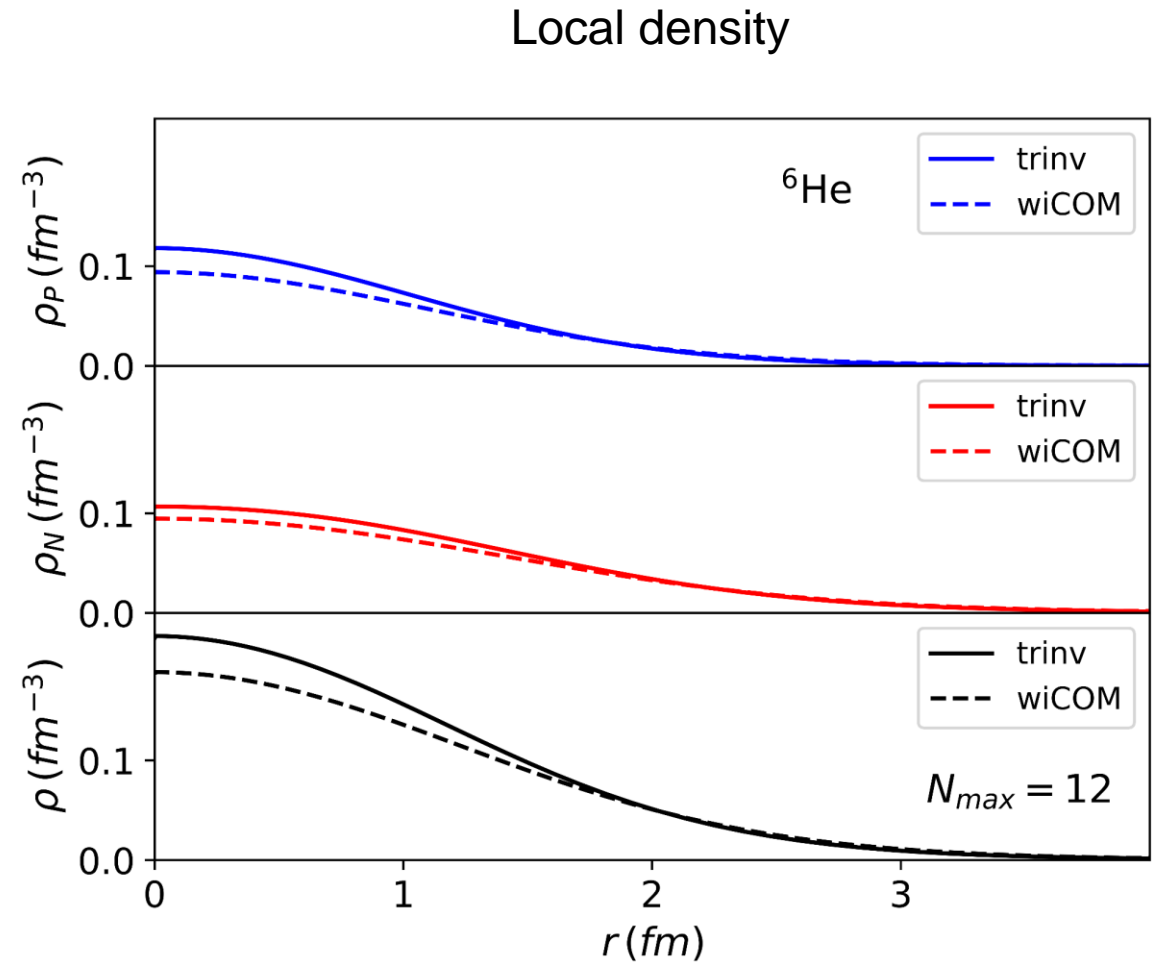
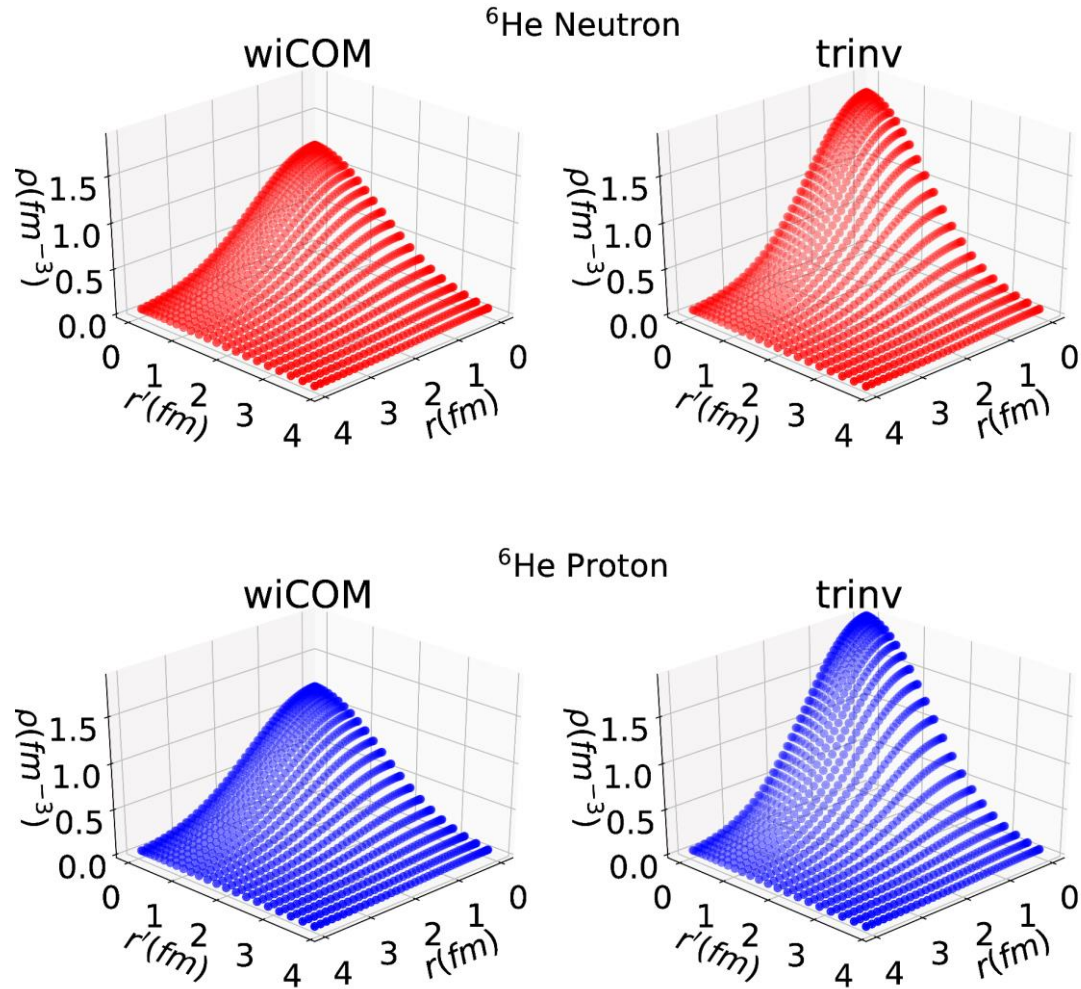
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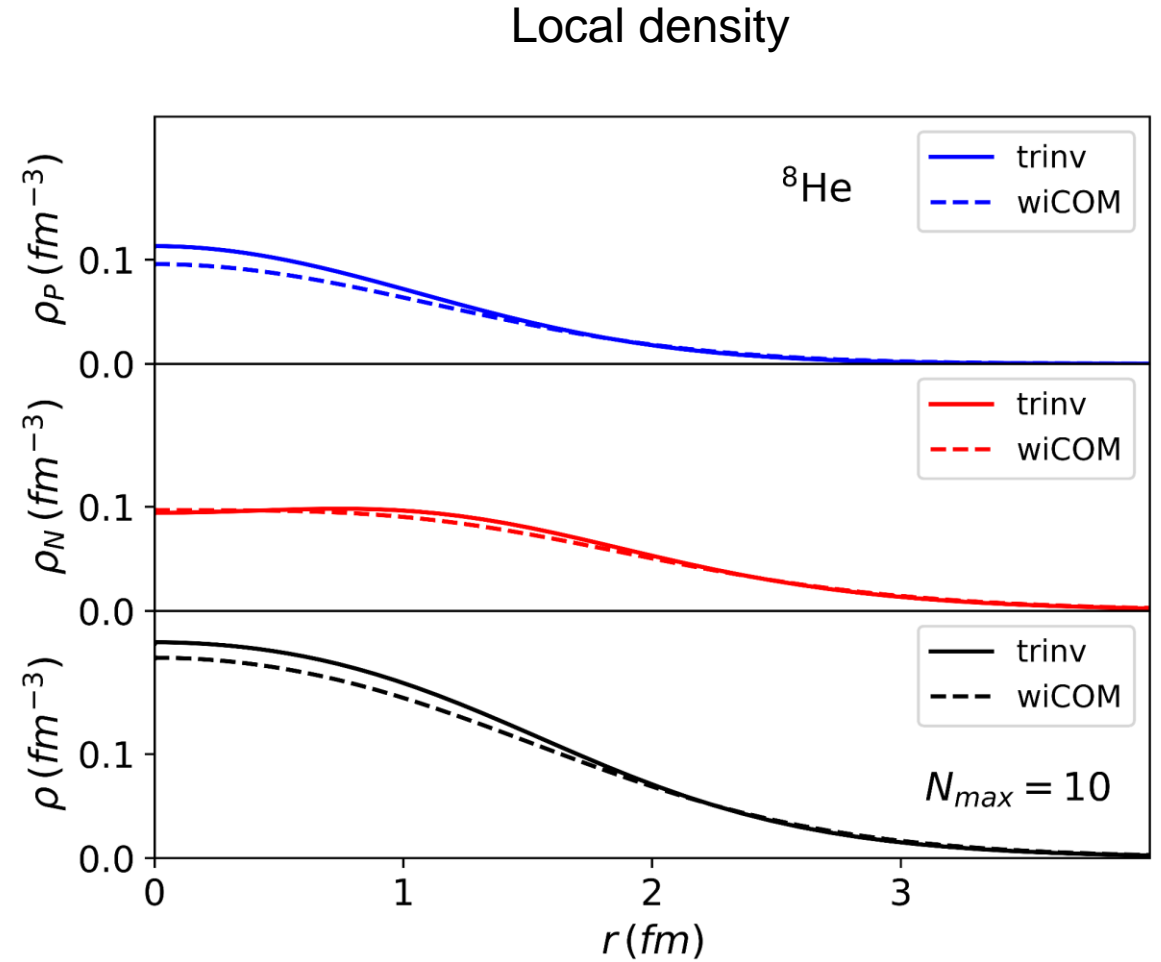
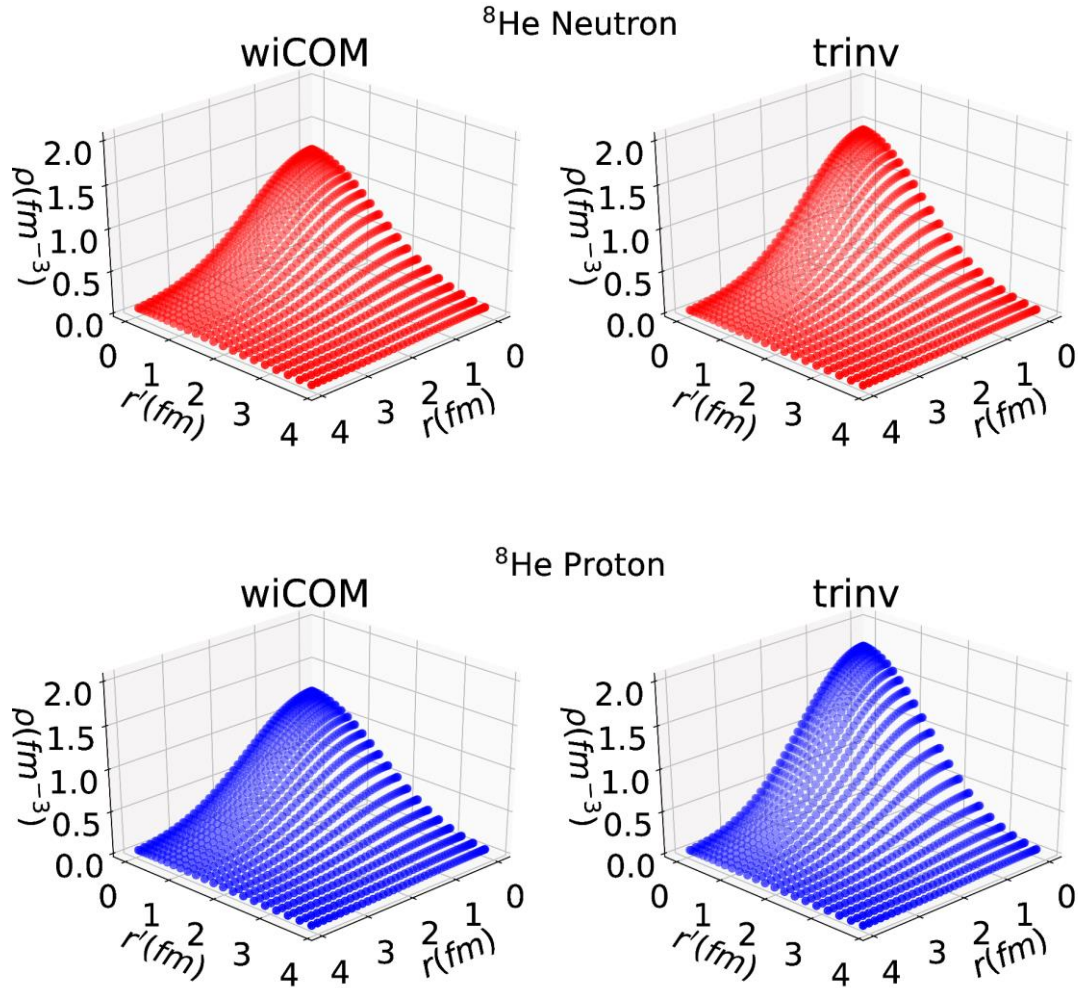
# Density of ground state ${}^4, {}^6, {}^8\text{He}$ , ${}^{12}\text{C}$ , ${}^{16}\text{O}$ with NN- $\text{N}^4\text{LO}(500)+3\text{NInl}$

9



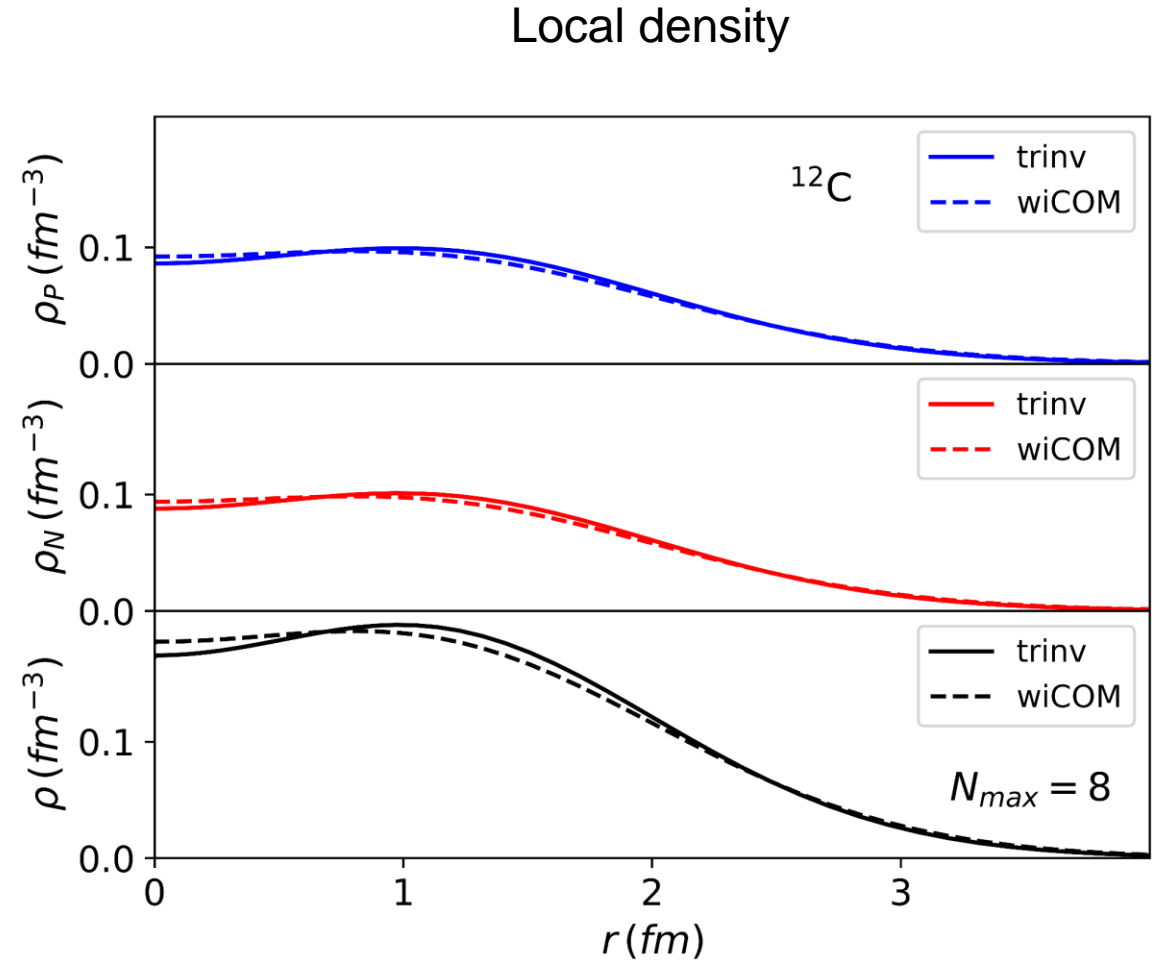
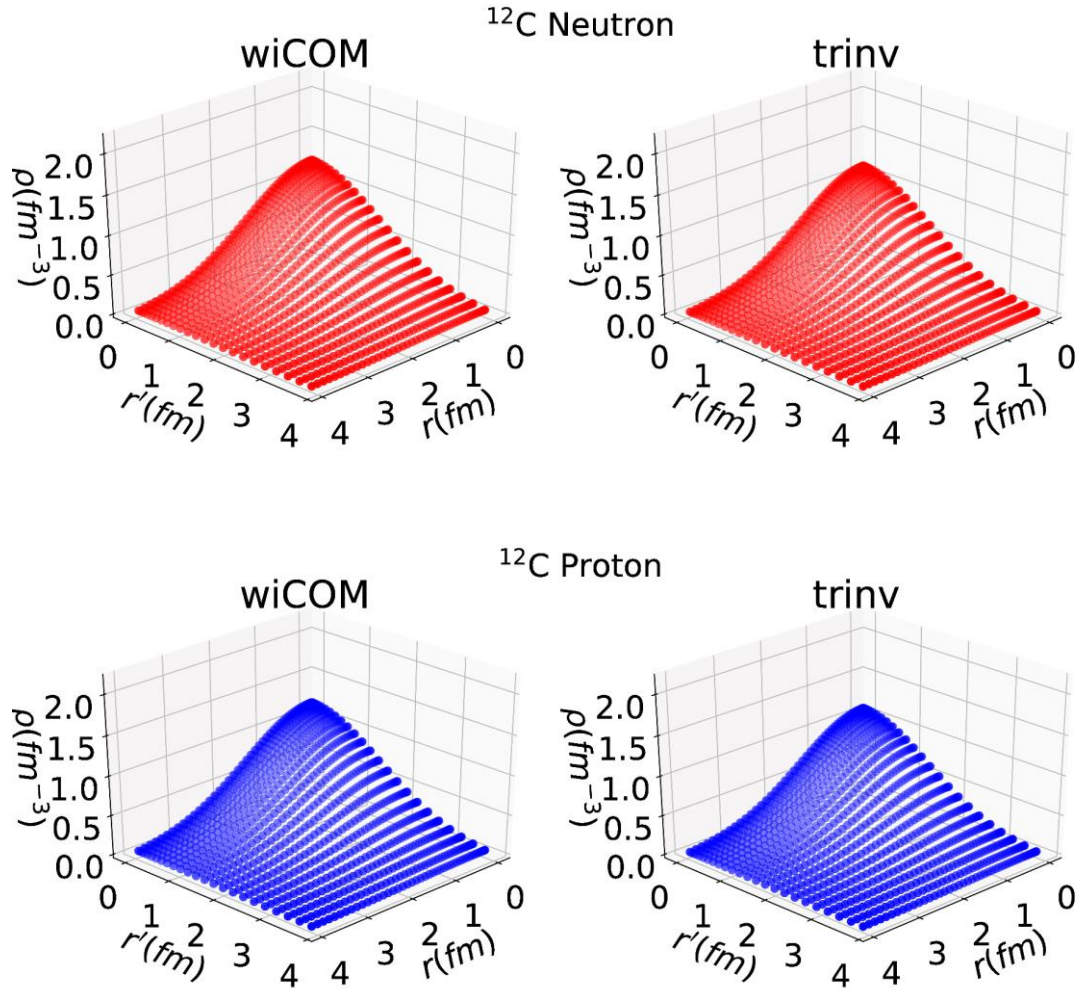
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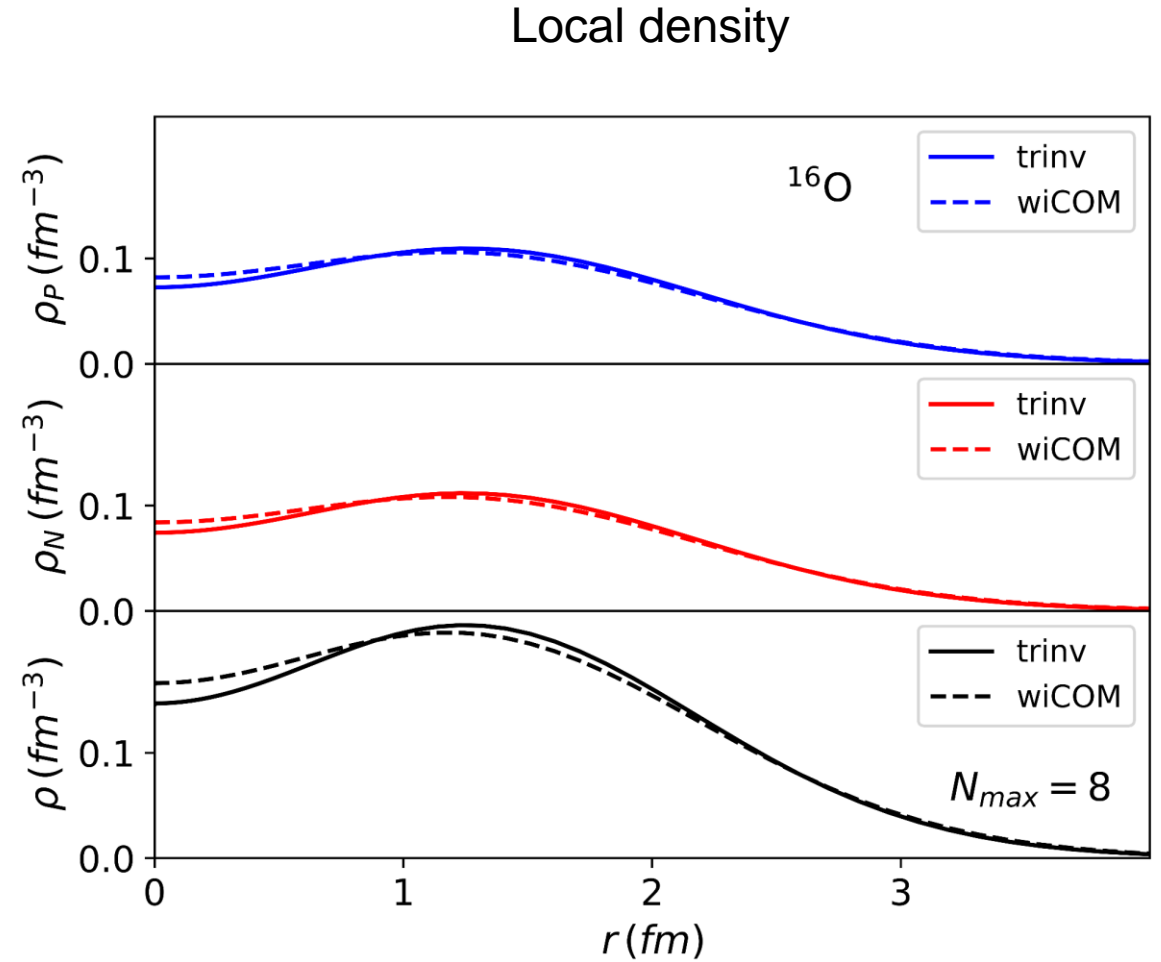
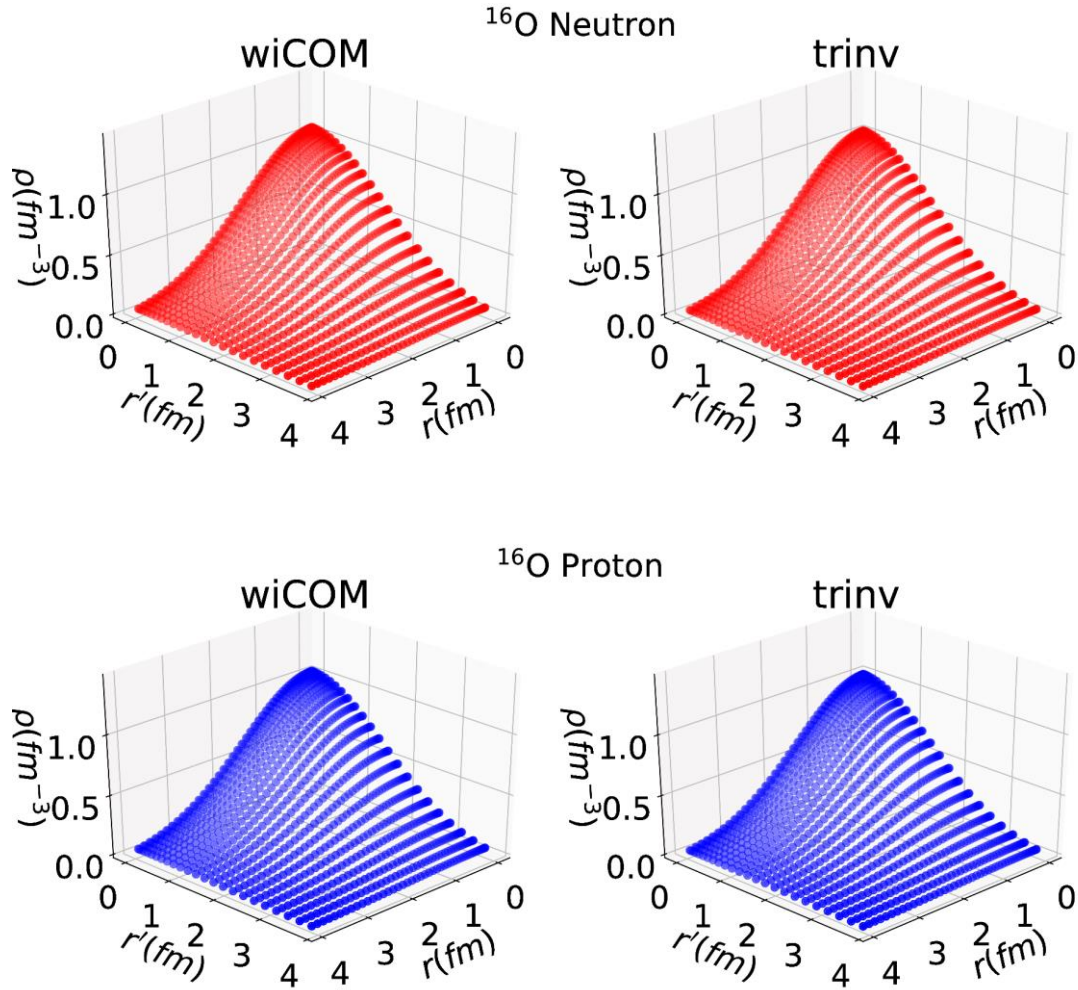
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11



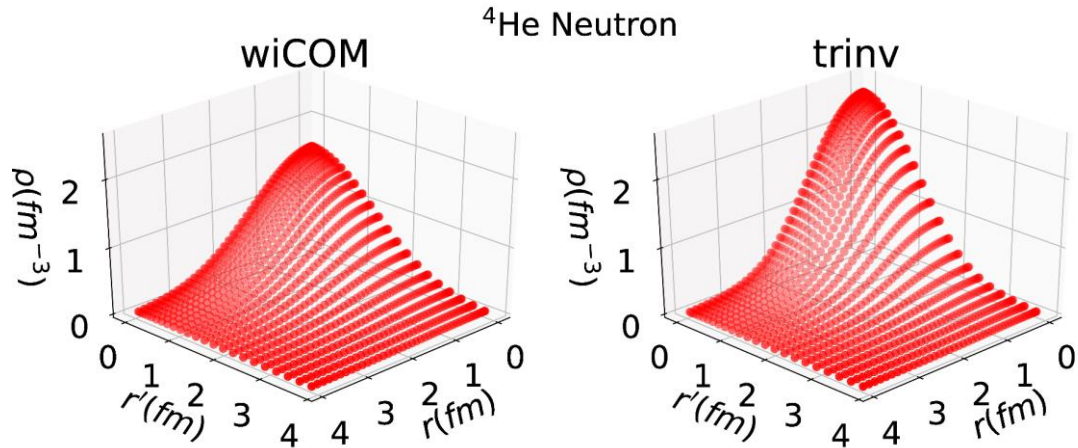
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12





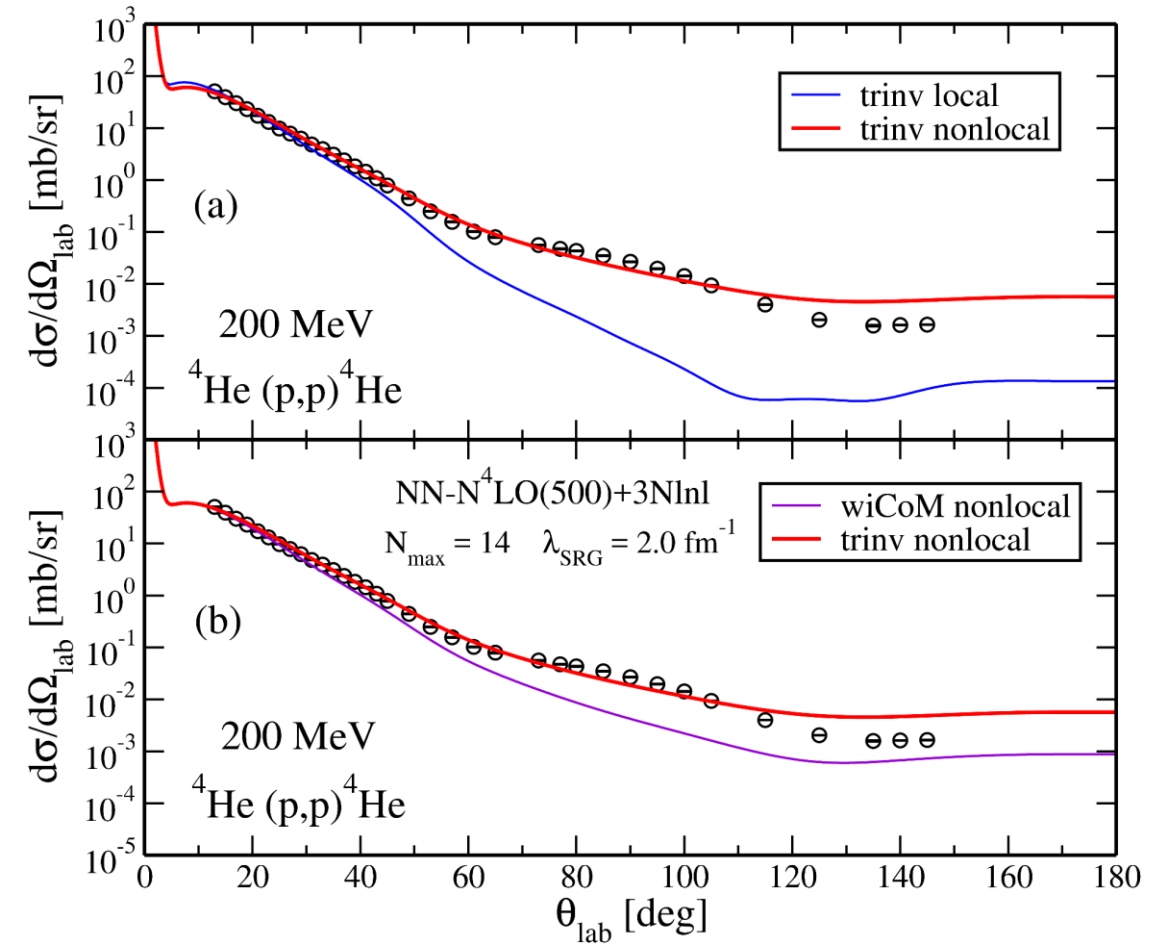
# Applications to optical potentials



- We can achieve an additional step towards consistent optical potential calculations by using the nonlocal density

$$U(\vec{q}, \vec{K}) = \sum_{N=n,p} \int d\vec{P} \quad \eta(\vec{q}, \vec{K}, \vec{P}) \quad t_{pN}(\vec{q}, \vec{K}, \vec{P}) \quad \rho_N(\vec{q}, \vec{P})$$

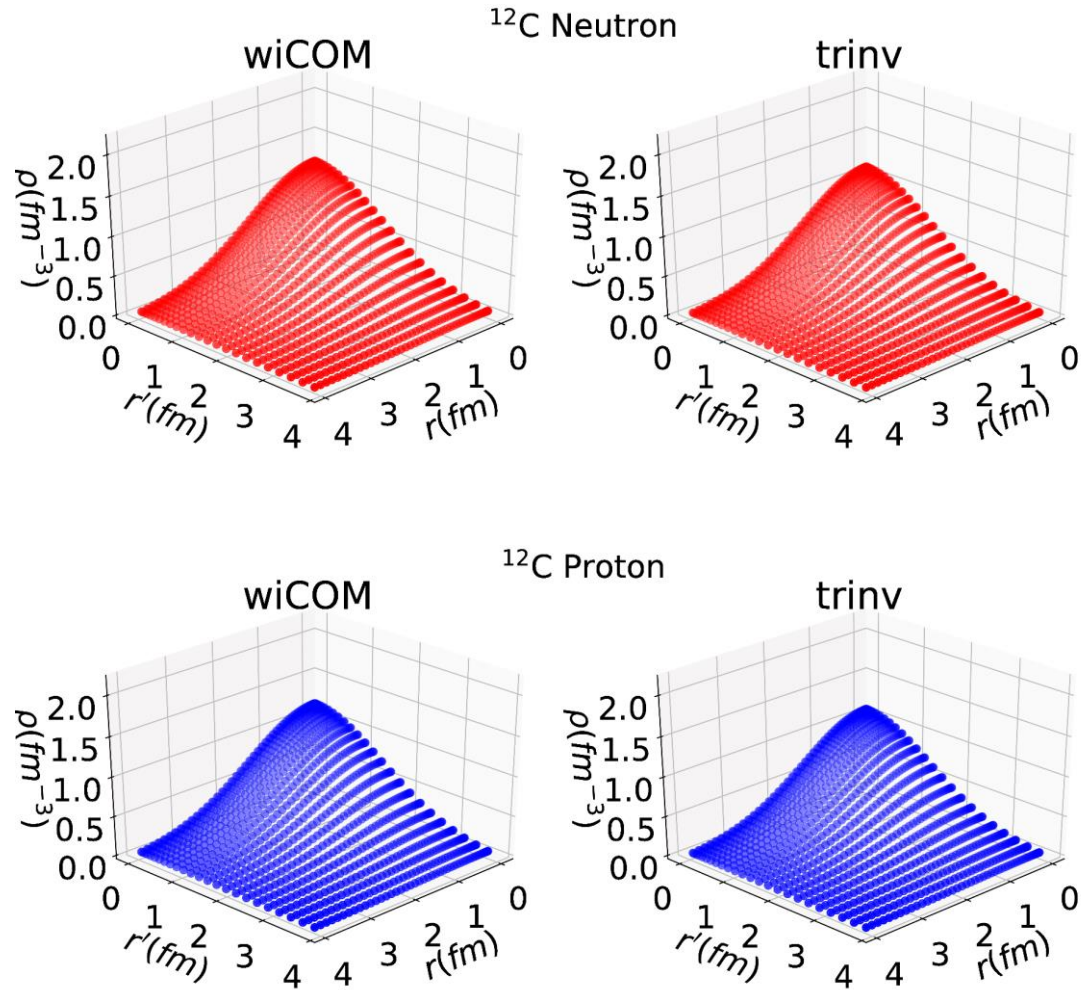
## Differential cross sections



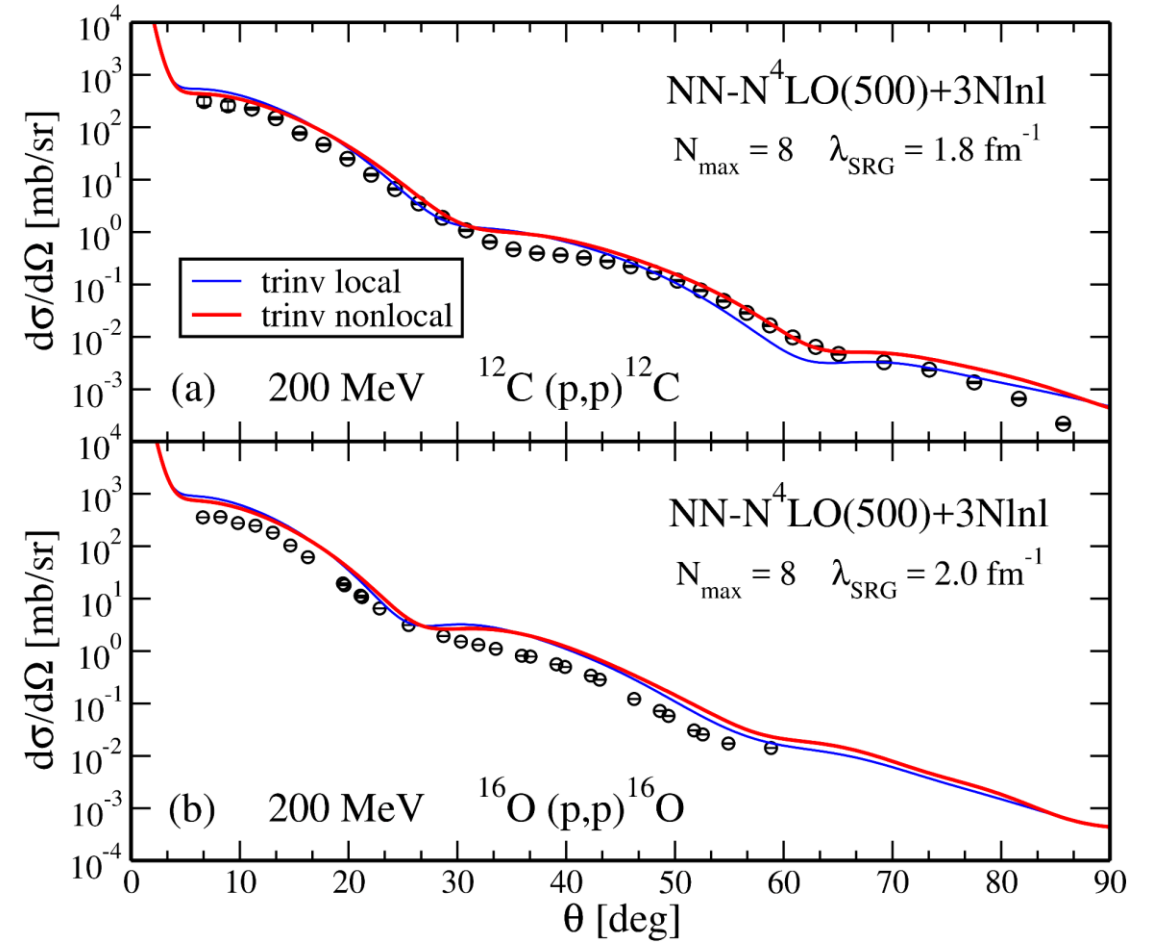


# Applications to optical potentials

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## Differential cross sections



# Applications to density functional theory (DFT)

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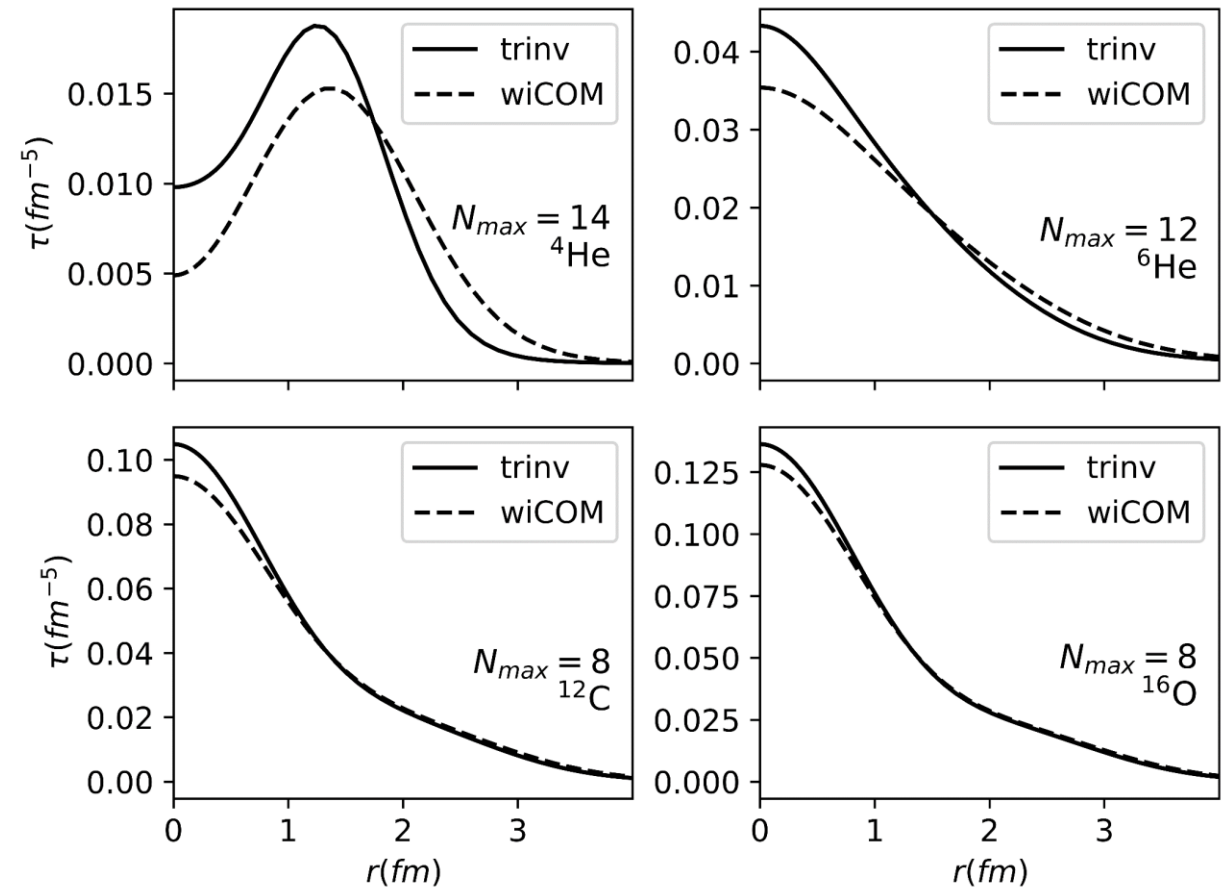
- The kinetic density, a DFT quantity used in energy density functionals, is calculable from *ab initio* wavefunctions by making use of the relations,

$$\tau_T(\vec{r}) = \left( \vec{\nabla} \cdot \vec{\nabla}' \rho_T(\vec{r}, \vec{r}') \right) |_{\vec{r}=\vec{r}'}$$

$$\frac{d}{dr} R_{n,l}(r) = \frac{l}{r} R_{n,l}(r) - \frac{1}{b} \left[ \sqrt{n+l+\frac{3}{2}} \cdot R_{n,l+1}(r) + \sqrt{n} \cdot R_{n-1,l+1}(r) \right]$$

- Effects of COM removal should be amplified in DFT quantities such as the kinetic density, due to the application of gradients on the nonlocal density

Kinetic density



# Conclusions and outlook

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- Conclusions
  - When the spurious COM contamination was removed, we observed significant differences in the nuclear structure of light systems
  - We can now use the more general, nonlocal density for optical potentials of nuclear reactions and *ab initio* calculations in DFT
  - More details on some of these results can be found in Phys. Rev. C 97, 034619 (2018)
- Outlook
  - We are now pursuing the use of these densities in calculations with natural orbitals, reducing necessary basis sizes and improving convergence
  - We will attempt to extend this to the translationally invariant one-body nuclear density matrix and further cut down basis sizes

Thank you  
**Merci**

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