

Building a Shell Model Code: The Pairing Model and the sd-Shell (and comparing results with NuShellX)

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Abstract

We have first implemented the pairing model which have a analytical solution (to benchmark the code). Then implemented the sd shell —> more general shell-model program that allows you to study general nuclear structure problems.

developing our own shell-model code that can perform shell-model studies of the oxygen isotopes using standard effective interactions (provided by us) using as example the 1s0d shell as model space.

We have also used the NushellX code in order to perform more advanced shell-model studies and compare the results obtained with your own shell-model code to those of NushellX

I. INTRODUCTION

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II. THEORY

i. Exact results: The Pairing Model

The pairing model is one of a few systems that has a simple analytic solution. It is therefore very useful to use this model to check that the numerical solution matches the analytic solution while developing the code.

The pairing problem consists of a system where the fermions combines together in pairs of two, one with spin up and one with spin down. This model does not allow so-called 'breaking of pairs' meaning that the pairs of particles always will be coupled together. An excitation must excite two particles at the same time.

Mathematically the pairing model can be described by a simplified Hamiltonian consisting of an unperturbed one-body operator H^0 and a perturbation so-called 'pairing interaction term' H^I . We will limit ourselves to at most two-body interactions so that we can write the operators as:

$$\hat{H}_0 = \zeta \sum_{p,\sigma} (p-1) \hat{a}_{p\sigma}^\dagger \hat{a}_{p\sigma}, \quad (1)$$

$$\hat{V} = -\frac{1}{2}g \sum_{p,q} \hat{a}_{p+}^\dagger \hat{a}_{p-}^\dagger \hat{a}_{q-} \hat{a}_{q+}, \quad (2)$$

so that the full Hamiltonian is given by the sum of the unperturbed term and the interacting part $\hat{H} = \hat{H}_0 + \hat{V}$. The fermion creation and annihilation operators are given by \hat{a}_p^\dagger and \hat{a}_q respectively and p,q,r,s represent all possible single-particle quantum numbers. To simplify the expressions we set the spacing between successive single-particle states given by $\zeta = 1$.

We will let the single-particle states $|p\rangle$ be eigenfunctions of the one-particle operator \hat{h}_0 . The above Hamiltonian acts in turn on various many-body Slater determinants constructed from the single-basis defined by the one-body operator \hat{h}_0 .

The two-body operator \hat{V} consists of one term:

$$\langle q_+q_- | \hat{V} | s_+s_- \rangle = -g \quad (3)$$

representing the pairing contribution with (for simplicity) constant strength g . The labeling requires that for a given matrix element $\langle pq | \hat{V} | rs \rangle$ the states p and q (or r and s) must have opposite spin ($\sigma = \pm 1$)

It can be shown that the unperturbed Hamiltonian \hat{H}_0 and \hat{V} commute with the spin projection \hat{S}_z and the total spin \hat{S}^2 . This allows us to block-diagonalize the full Hamiltonian. For the pairing case we have assumed that $S = 0$ giving the so-called 'no-broken pair' case.

Constructing the Hamiltonian matrix As an analytic solution we chose a system consisting of only four particles with a single-particle space consisting of only the four lowest levels $p = 1, 2, 3, 4$. In our system every level p contains two particles, one with spin up and one with spin down. The goal is to set up all possible Slater determinants and the Hamiltonian matrix using second quantization and find all eigenvalues by diagonalizing the Hamiltonian matrix.

We construct the basis:

$$|\Phi_0\rangle = \begin{pmatrix} 1 \\ 0 \\ 0 \\ 0 \\ 0 \\ 0 \end{pmatrix}, \quad |\Phi_1\rangle = \begin{pmatrix} 0 \\ 1 \\ 0 \\ 0 \\ 0 \\ 0 \end{pmatrix}, \quad \dots \quad |\Phi_5\rangle = \begin{pmatrix} 0 \\ 0 \\ 0 \\ 0 \\ 0 \\ 1 \end{pmatrix},$$

where

$$\begin{aligned}
 |\Phi_0\rangle &= \hat{a}_{2+}^\dagger \hat{a}_{2-}^\dagger \hat{a}_{1+}^\dagger \hat{a}_{1-}^\dagger |0\rangle = \hat{P}_2^+ \hat{P}_1^+ |0\rangle, \\
 |\Phi_1\rangle &= \hat{P}_3^+ \hat{P}_1^+ |0\rangle, \\
 |\Phi_2\rangle &= \hat{P}_4^+ \hat{P}_1^+ |0\rangle, \\
 |\Phi_3\rangle &= \hat{P}_3^+ \hat{P}_2^+ |0\rangle, \\
 |\Phi_4\rangle &= \hat{P}_4^+ \hat{P}_2^+ |0\rangle, \\
 |\Phi_5\rangle &= \hat{P}_4^+ \hat{P}_3^+ |0\rangle.
 \end{aligned}$$

(also include a figure with the states drawn?)

Given the above Slater determinants we can now compute the matrix elements $\langle \Phi_i | \hat{H} | \Phi_j \rangle$ using the Hamiltonian of Eq. (1). The one-body operator acts only on the diagonal and results in terms proportional with $(p1)$. The interaction will excite or deexcite a pair of particles from level q to level p . Using this it is easy to see that the Hamiltonian matrix becomes: (REWRITE THIS PARAGRAPH, ALL COPY)

$$\hat{H} = \begin{pmatrix} 2-g & -g/2 & -g/2 & -g/2 & -g/2 & 0 \\ -g/2 & 4-g & -g/2 & -g/2 & 0 & -g/2 \\ -g/2 & -g/2 & 6-g & 0 & -g/2 & -g/2 \\ -g/2 & -g/2 & 0 & 6-g & -g/2 & -g/2 \\ -g/2 & 0 & -g/2 & -g/2 & 8-g & -g/2 \\ 0 & -g/2 & -g/2 & -g/2 & -g/2 & 10-g \end{pmatrix}.$$

For a given g this matrix can be used as the analytic result to compare with the output from the shell model code for the pairing case.

Vary your results for values of $g \in [1, 1]$ (diagonalize, find eigvec \rightarrow plot figure) Comment the behavior of the ground state as function of g ?

III. BUILDING THE SHELL MODEL CODE

develop a code which sets up the above Hamiltonian matrices for two and four particles in 2 and 4 single-particles states (the same as what you did in exercises b) and c) and obtain the eigenvalues.

What did we choose?

- Decide whether you want to read from file the single-particle data and the matrix elements in m-scheme, or set them up internally in your code. The latter is the simplest possibility for the pairing model, whereas the first option gives you a more general code which can be extended to the more realistic cases discussed in the second part.
- Based on the single-particle basis, write a function which sets up all possible Slater determinants which have total $M = 0$. Test that this function reproduces the cases in b) and c). If you make this function more general, it can then be reused for say a shell-model calculation of sd-shell nuclei in the second part.
- Use the Slater determinant basis from the previous step to set up the Hamiltonian matrix.
- With the Hamiltonian matrix, you can finally diagonalize the matrix and obtain the final eigenvalues and test against the results of b) and c).

Include some of this? Or too much? :

The lecture slides contain a rather detailed recipe on how to construct a Slater determinant basis and how to set up the Hamiltonian matrix to diagonalize.

Table 1: *Example table*

Name		
First name	Last Name	Grade
John	Doe	7.5
Richard	Miles	2

i. first step....

ii. Unit tests and benchmarks

(have not done this?)

One obvious case is that of removing the two-body interaction. Then we have only the single-particle energies. For the case of degenerate single-particle orbits, that is one value of total single-particle angular momentum only j , with degeneracy $\Omega = 2j + 1$, one can show that the ground state energy E_0 is with n particles

$$E_0 = -\frac{g}{4}n(\Omega - n + 2) \quad (4)$$

Enlarge now your system to six and eight fermions and to $p = 6$ and $p = 8$ single-particle states, respectively. Run your program for a degenerate single-particle state with degeneracy Ω and test against the exact result for the ground state. Introduce thereafter a finite single-particle spacing and study the results as you vary g , as done in b) and c). Comment your results.

(describe the unit test that we have implemented and how)

IV. NuSHELLX

(describe what is, then put results in results section?)

V. RESULTS

VI. DISCUSSION

REFERENCES

[Figueredo and Wolf, 2009] Figueredo, A. J. and Wolf, P. S. A. (2009). Assortative pairing and life history strategy - a cross-cultural study. *Human Nature*, 20:317–330.

[Github of the TALENT School] HJ and Alex Brown *link goes here*,