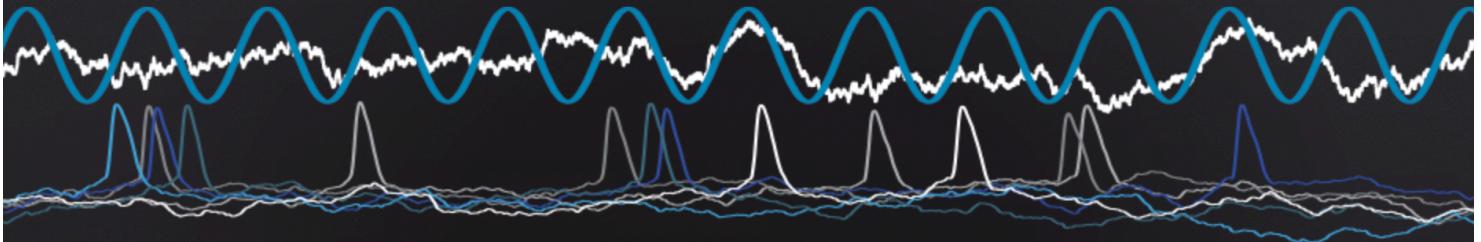


# ELECTROPHYSIOLOGICAL SIGNALS



GENERATION AND CHARACTERISATION

Michele GIUGLIANO  
**Neuroelectronics**

References for today's class content

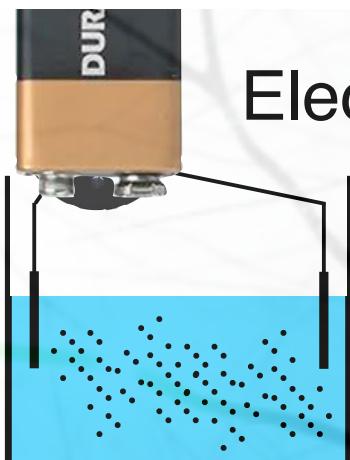
**supporting your own study and understanding**

## Chapters from

- Weiss TF (1996) “*Cellular Biophysics*” vol. 1, MIT Press.
- Johnston & Wu, 1995 “Foundations of Cellular Neurophysiology”
- Sterratt et al. (2011) “*Principles of Computational Modelling...*”
- Abbott LF, Dayan P (2001) “*Theoretical Neuroscience*”

# Origin of Bioelectricity

- we deal with **electrochemical systems** (in water): anions, cations
- we **measure** electrostatic potentials by **electrodes** (e.g. **AgCl** based)
- we **observe** -70mV at “rest” and swings up to +30mV in <1 ms
- starting from **the def.** of potential, we use **superposition of effects**
- *anchored charges?* ambiguous but intuition of **asymmetric concentrations**
- *moving charges?* **drift & diffuse**; there is a (cell) **membrane**
  - with **capacitive** properties as well as with ionic **permeability**
  - what is **permeability**? how **ions** flow through the membrane?



## Electrochemical systems

aqueous solution (**solvent**)  
+  
molecules (e.g. ions) (**solute**)  
(e.g.,  $\text{Na}^+$ ,  $\text{K}^+$ ,  $\text{Ca}^{++}$ ,  $\text{Mg}^{++}$ ,  $\text{Cl}^-$ )

Dissociation of salts by **solvation** (e.g.  $\text{NaCl} \leftrightarrow \text{Na}^+ + \text{Cl}^-$ ) into electrically charged particles. Globally, **electroneutrality** holds!

in **solutions**: charge carriers are *cations*(+) and *anions*(-)

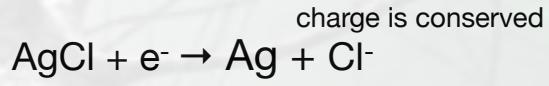
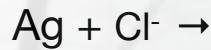
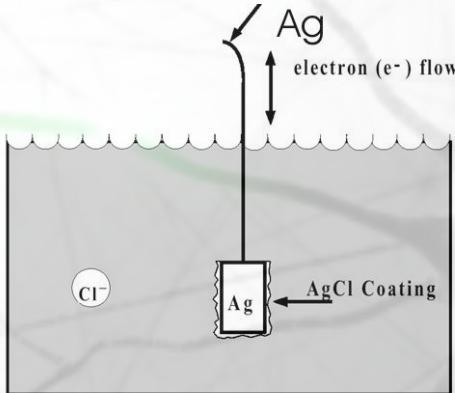
in **metals**: charge carriers are *electrons*(-)

(in semiconductors: charge carriers are *electrons*(-) and *holes*(+))

**electrodes**: *anodic* (+) and *catodic* electrodes (-)

# Silver-chloride junction

ions  $\neq$  electrons, distinct “currency” for charge exchanges



*this is a model!*

A simple electrical circuit model. A zigzag line representing a resistor is connected to a horizontal line representing ground. The symbol for a resistor, a zigzag line, is placed below the line.

$$\Delta V = R I$$



- ~ measurement of electrical potentials in solution
- ~ “injection” of external (ionic) currents

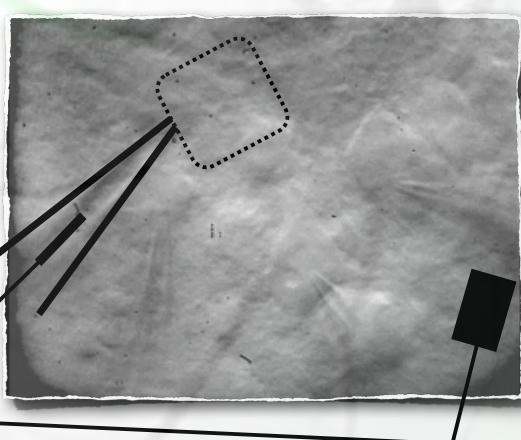
(see sodium hypochlorite, NaClO)

## Our task for today: understanding why

there is a membrane electrical potential (“at rest”)  
(in every cell, not just in nerve cells!)



$-70\text{ mV}$

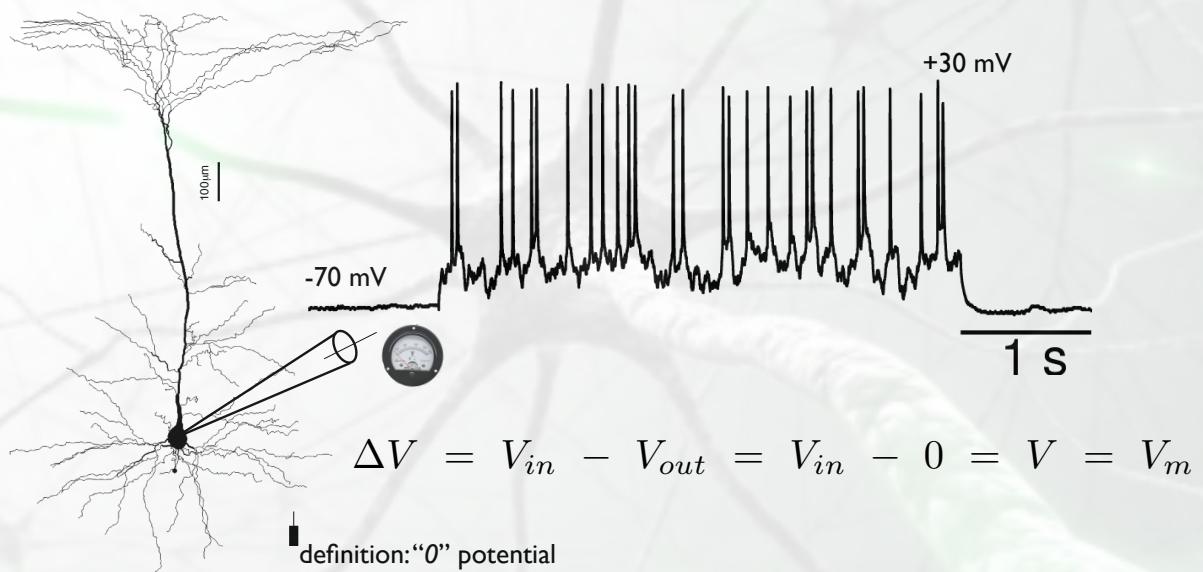


$$V_{in} - V_{out} = V_{in} - 0 = \\ = V = V_m$$

# Understanding, from “first (biophysical) principles”

- the existence of a *difference* of electrical potential - across the plasmatic membrane of a *living* cell...  $V_{in} - V_{out}$ 
  - electrical (electrostatic) potential?? measured in solutions? involving electrolytes??
  - the membrane? which physical (electrical) equivalent?
- ...in terms of its selective ionic permeability and of ionic flows.
- Towards a full (electrical) equivalent model of a cell membrane.

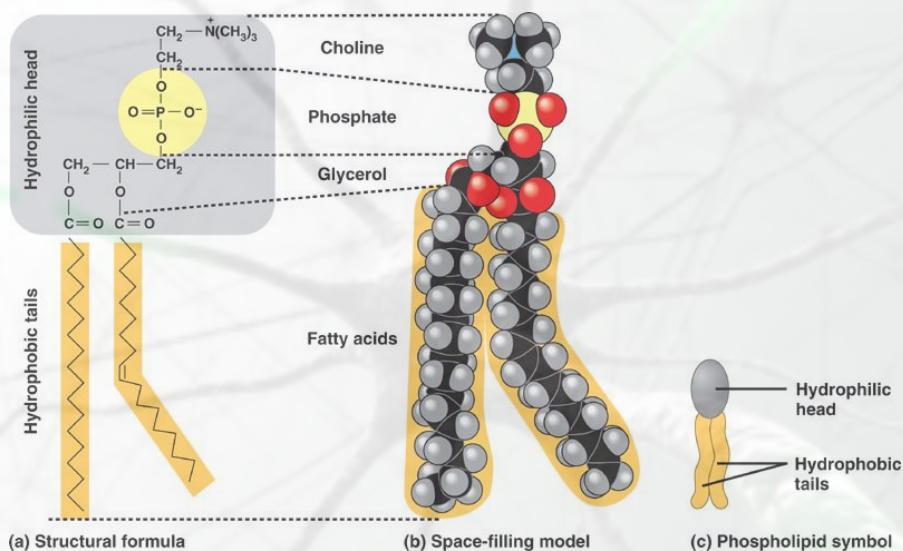
Such an understanding is **essential to explain** the **excitable electrical properties** of the cell membrane of neurons



# Origin of Bioelectricity

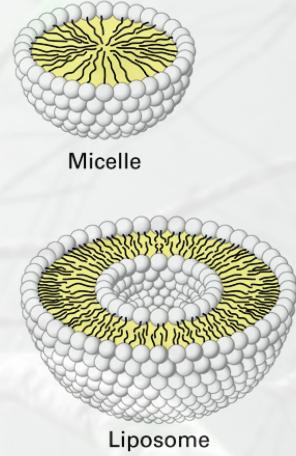
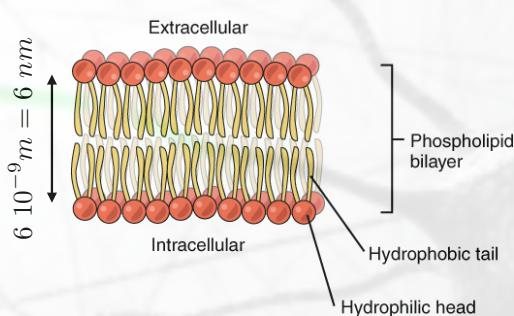
- we deal with **electrochemical systems** (in water): anions, cations
- we **measure** electrostatic potentials by **electrodes** (e.g. **AgCl** based)
- we **observe** -70mV at “rest” and swings up to +30mV in <1 ms
- we measure it across a thin membrane. What is a (cell) membrane?
- starting from **the def.** of potential, we use **superposition of effects**
- anchored charges? ambiguous but intuition of **asymmetric concentrations**
- moving charges? **drift & diffuse**; there is a (cell) **membrane**
  - with **capacitive** properties as well as with ionic **permeability**
  - what is **permeability**? how **ions flow through** the membrane?

The plasmatic membrane of biological cells is (mostly) made of...



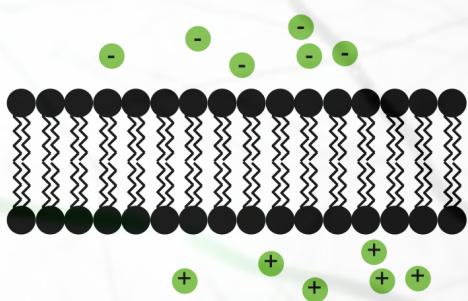
self-assembly of polymers (phospholipids)

The plasmatic membrane of biological cells is (mostly) made of...



- **hydrophilic** heads (phosphate groups, charged)
- **hydrophobic** tails (fatty acids)
- self-organise in water, into segregated compartments (emergence of life!!)

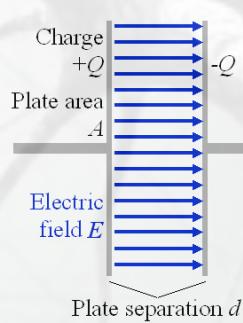
### Plasmatic membrane: a physical equivalent?



Two volumes, where *charges move freely*  
(as in an electrical **conductor**)  
...separated by a *barrier* that cannot be  
crossed (as in an **insulator** or dielectric)  
**Does this ring any bell?**

*this is a model!*

$C$



$$C = \frac{\Delta Q}{\Delta V}$$

$$C = \frac{\epsilon A}{d}$$



## Exercise

What is the membrane capacitance (per area unit) ?  
 What charge distribution “delta” leads to 70mV?

$$C = \frac{\epsilon A}{d} \quad c = 1 \mu F/cm^2$$

$$C = \frac{7 \cdot 8.85 \cdot 10^{-12} \text{ } F/m \text{ } A}{6 \cdot 10^{-9} m} = 0.010 \text{ } F/m^2 \text{ } A = 1 \mu F/cm^2 \text{ } A$$

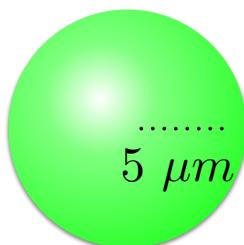
$$C = \frac{\Delta Q}{\Delta V}$$

$$\Delta Q = 1 \mu F/cm^2 \cdot 70 mV = 7 \cdot 10^{-8} C/cm^2$$

**Unique membrane properties and enhanced signal processing in human neocortical neurons**

Guy Eyal<sup>1</sup>, Matthijs B Verhoog<sup>2</sup>, Guilherme Testa-Silva<sup>2</sup>, Yair Deitcher<sup>3</sup>,  
 Johannes C Lodder<sup>2</sup>, Ruth Benavides-Piccione<sup>4,5</sup>, Juan Morales<sup>6</sup>,  
 Javier DeFelipe<sup>4,5</sup>, Christiaan PJ de Kock<sup>2</sup>, Huibert D Mansvelder<sup>2</sup>, Idan Segev<sup>1,3\*</sup>

**Therefore, there is a difference of charge across the membrane!**



523  $10^8$  ions

**50 billions**

$1.6 \cdot 10^8$  ions/ $\mu m^2$  (**shell**)

$2.56 \cdot 10^{-11} C/\mu m^2$

$1.6 \cdot 10^{-19} C/ion$

$$7 \cdot 10^{-8} C/cm^2 = 7 \cdot 10^{-8} C/10^8 \mu m^2 \approx 10^{-15} C/\mu m^2$$



70mV

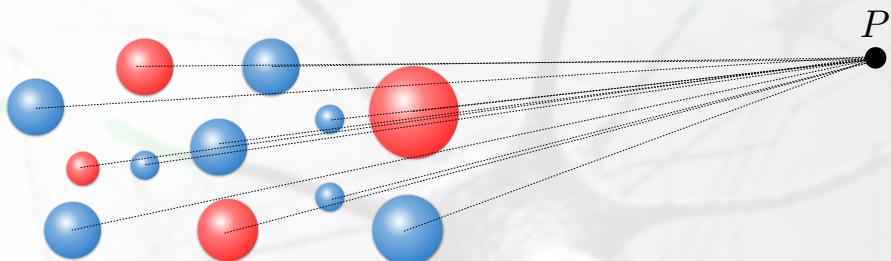
$$\frac{10^{-15}}{1.6 \cdot 10^{-19}} ions/\mu m^2 \approx 6 \cdot 10^3$$

**~1000 ions of difference, across the membrane**

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REMINDER: *superposition of the effects*



$$V_{total}(P) = \frac{1}{4\pi\epsilon_r\epsilon_0} \left( \frac{Q_1}{r_{P-Q_1}} + \frac{Q_2}{r_{P-Q_2}} + \dots + \frac{Q_M}{r_{P-Q_M}} \right)$$



It is like the *weighted sum* of the inverse of the distances...

**Exercise:** (discrete) distribution of charge  
anchored/fix/glued in free space  
(restrained from self-organising)

[https://phet.colorado.edu/sims/html/charges-and-fields/latest/charges-and-fields\\_en.html](https://phet.colorado.edu/sims/html/charges-and-fields/latest/charges-and-fields_en.html)

DEMO TIME!

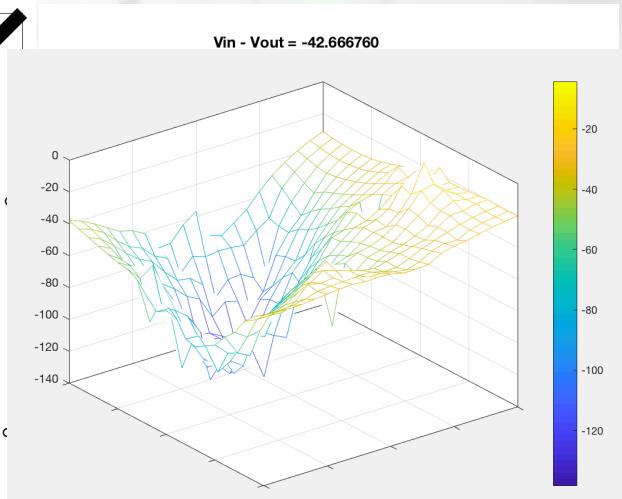
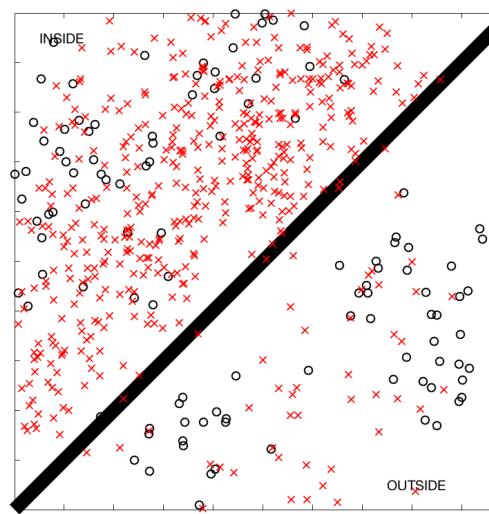
**Exercise:** (discrete) distribution of charge  
anchored/fix/glued in free space  
(restrained from self-organising)

Given the potential (measured),  
which is the corresponding distribution of charge??  
?????

$$\Delta V = V_{in} - V_{out} \quad -70 \text{ mV}$$

[https://phet.colorado.edu/sims/html/charges-and-fields/latest/charges-and-fields\\_en.html](https://phet.colorado.edu/sims/html/charges-and-fields/latest/charges-and-fields_en.html)

$$\Delta V = V_{in} - V_{out} \text{????? } -70 \text{ mV}$$

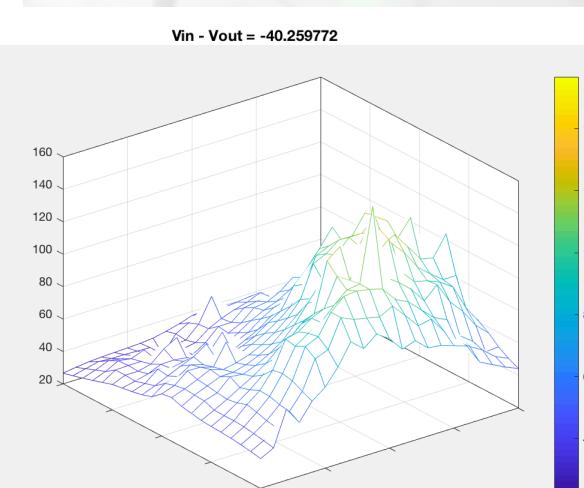
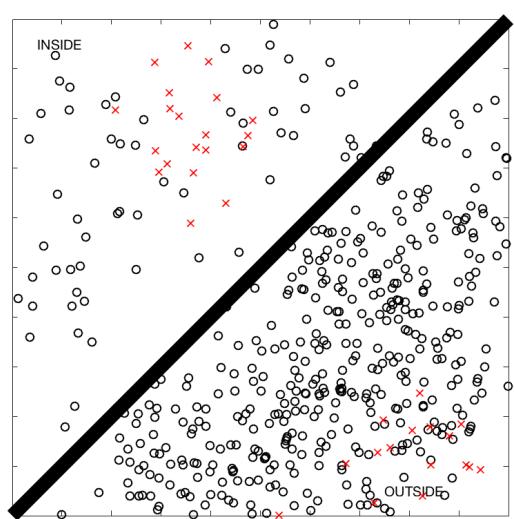


**o: positively charged ions  
x: negatively charged ions**

Negative ions inside..



$$\Delta V = V_{in} - V_{out} \text{????? } -70 \text{ mV}$$



**o: positively charged ions  
x: negatively charged ions**

Unequal distribution of positive ions..

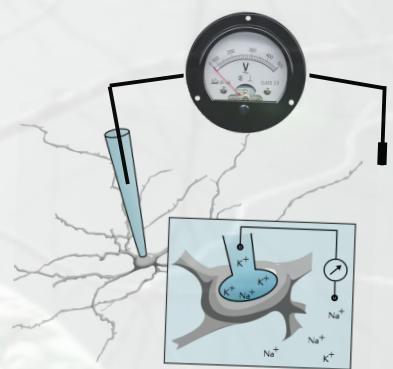


# Important points

- ions are free to **move**... (to diffuse and to experience *drifting forces*, due to their own electric fields)
- thin (insulating) **membrane** across which we measure a potential...

## The existence of a (resting) membrane potential: **facts**

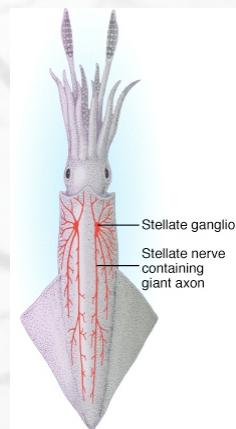
- global electro-neutrality holds => **isopotential** inside & outside the cell
- ions species ( $\text{Na}^+$ ,  $\text{K}^+$ ,  $\text{Cl}^-$ , ...) are **distributed differently** inside & outside
- the membrane bilayer is impermeable to water (and ions) (i.e. double layer of phospholipids - hydrophobic!!);
- conventionally, we measure electric *potential*  $V$ , inside with reference to outside the membrane
- if a (generic) cell is not *dead* (NO thermodynamical equilibrium = identity with its surrounding),  $V \sim -70 \text{ mV}$
- in excitable cells,  $V$  **may change abruptly** in time (e.g., neurons, myocytes, pancreatic  $\beta$ -cells)



from Sterratt et al., 2011

**Why ? How?**

## Ionic concentrations for the squid's giant-axon ("prep" used by Hodgkin & Huxley)



Ion	$K^+$	$Na^+$	$Cl^-$	$Ca^{2+}$
Concentration inside (mM)	400	50	40	$10^{-4}$
Concentration outside (mM)	20	440	560	10

from Sterratt et al., 2011



Why is there an electric potential across the cell membrane?



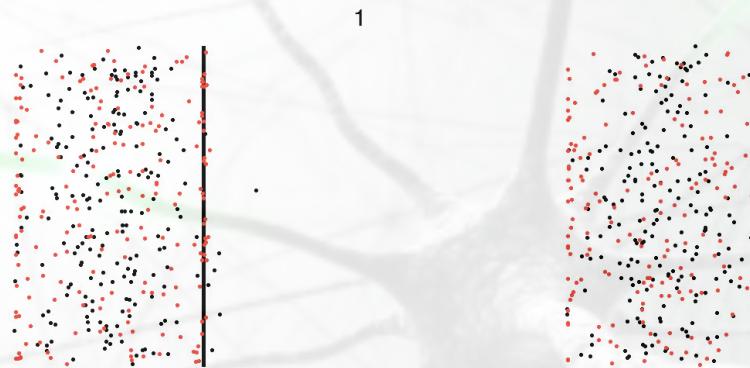
Because of a **heterogeneous distribution** of charges, across the membrane,...

**...due to its semi-permeability!**

multiple ion-species, NOT at the equilibrium  
distinct concentrations



# Semi-permeability... really? YES



permeable (non-selective) membrane

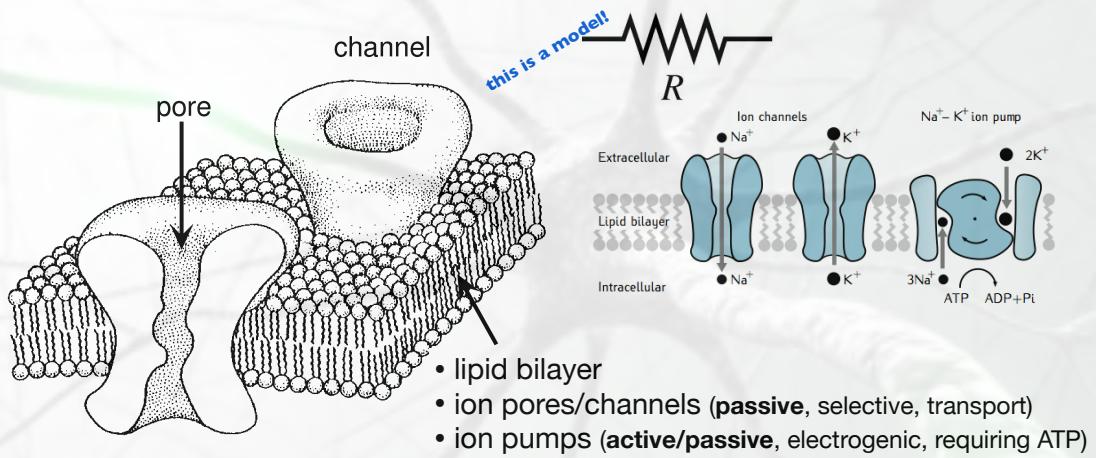
semi-permeable (selective, *black-only*) membrane

Monte Carlo simulation = [Coulomb's attractive forces](#) + friction + random collisions...

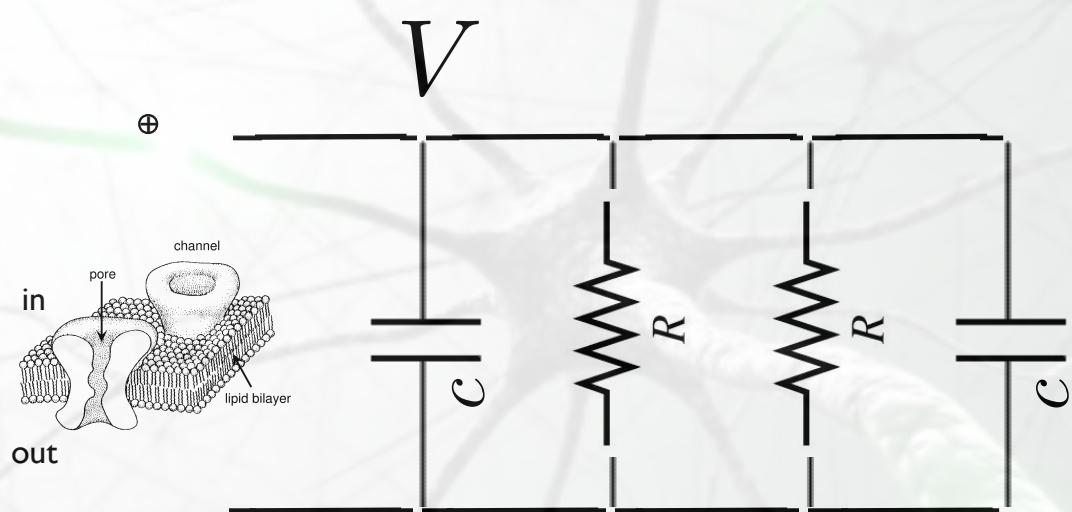
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## The plasmatic cell membrane is *selectively* permeable to specific ions

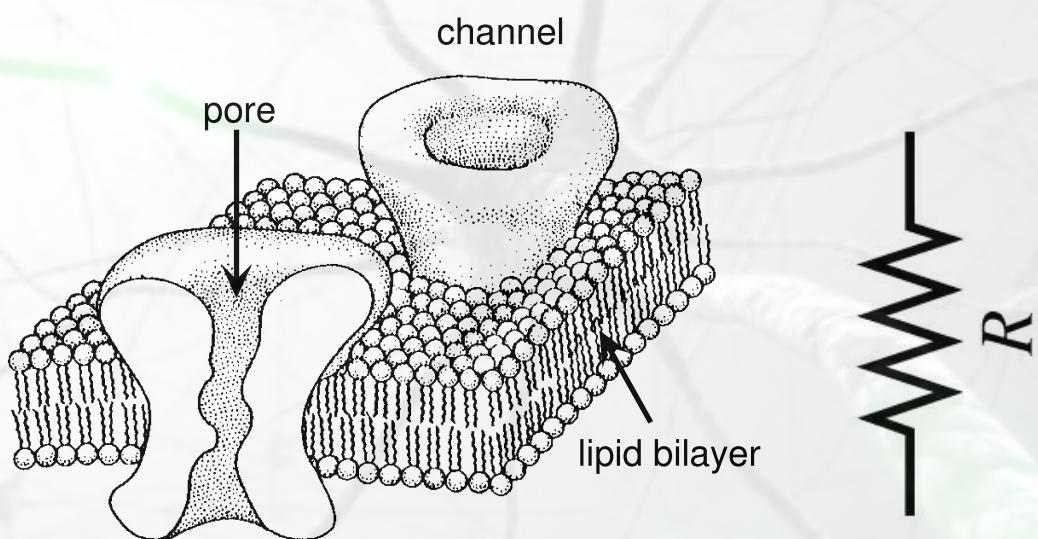


Towards a physical equivalent (electrical) circuit of the plasmatic cell membrane



Is the “resistor” a good model for an ion channel...?

We first need to understand how ions flow  
across a permeable membrane !!!



**Back to the key problem for today**

- heterogenous charge concentrations (inside & outside) and ..-70mV

**Tackle the self-consistence or self-organizing  
character of the resting potential.**

# Origin of Bioelectricity

- intuition for **ionic semi-permeability** (NOT identical permeabilities)
- single species, **equilibrium**: **Nernst** Potential
  - Na+, K+, Ca++, Cl-, ....?
- single species, non-equilibrium: **ion fluxes (ionic current densities)**
  - **Ohmic** approximation and non-Ohmic (**Goldmann eq.**)
- **multiple** species, **non**-equilibrium: **steady-state** hypothesis
  - **Goldman Hodgkin Huxley** equation(s) and **resting potential**

What kind of fluxes can occur in solution?

Diffusive and Drift fluxes

$$J = u c \left( -R T \frac{d}{dx} \ln[c(x)] \right) \quad J = u c \left( -z F \frac{dV}{dx} \right)$$

What does their knowledge predict?

The formation of a (Nernst) electrical potential in space, given an unequal concentration of ions (e.g. across a semi-permeable membrane)

The phenomenon of (electro-)diffusion





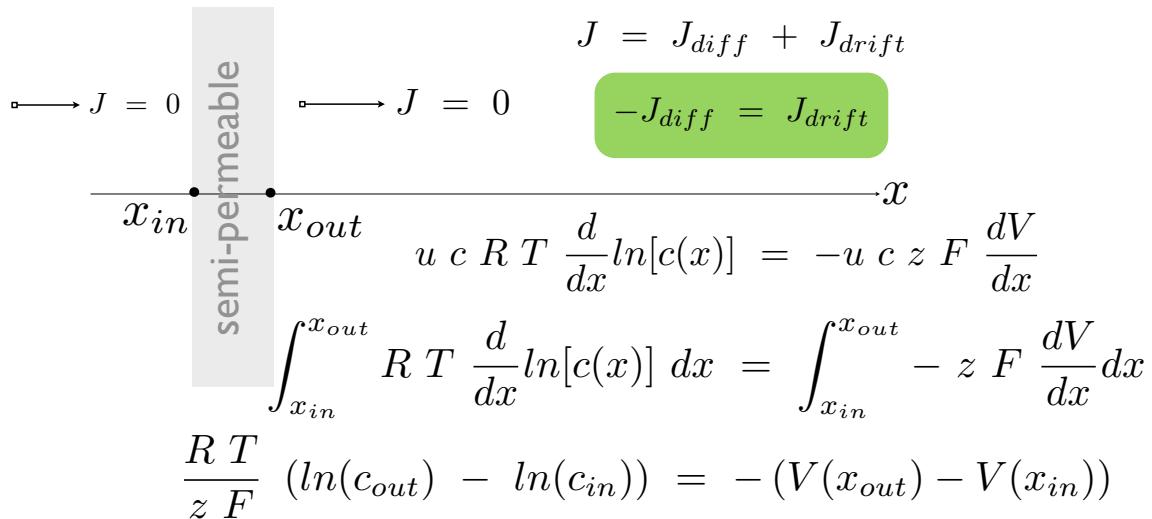
### Concept of the **definite integral** (i.e. fundamental theorem)

$$\int_a^b \frac{1}{x} dx = \ln(x)|_a^b = \ln(b) - \ln(a) = \ln(b/a)$$

### Concept of (Taylor's) expansion into a series of polynomials

$$f(x_0 + h) \approx f(x_0) + \frac{df(x)}{dx}|_{x_0} h$$

Two compartments and an existing concentration gradient  
 [hp: **single ion-specie, at equilibrium (i.e.  $J = 0$ )**]



$$V_{in} - V_{out} = \frac{R T}{z F} \ln \left( \frac{c_{out}}{c_{in}} \right) \quad \frac{R T}{z F} = \frac{K T}{z q}$$

Two compartments and an existing concentration gradient  
 [hp: **single** ion-specie, at **equilibrium (i.e.  $J = 0$ )**]



Walther Nernst  
 (1864-1941)

### Nernst equation

$$E_{Nernst} = V_{in} - V_{out} = \frac{R T}{z F} \ln \left( \frac{c_{out}}{c_{in}} \right)$$



*this is a modell*

Nernst equilibrium potential for that ion  
 or also known as “reversal potential” (more on it later)

It does NOT depend on the ion’s mobility!

$$\frac{R T}{z F} = \frac{K T}{z q} \xleftarrow{T = 300^\circ, z = 1} \approx 26mV$$



**Exercise**  
 what is the Nernst’s potential for individual ionic species, at room temperature ( $\sim 20^\circ\text{C} = 300\text{K}$ )?



Walther Nernst  
 (1864-1941)

Ion	K <sup>+</sup>	Na <sup>+</sup>	Cl <sup>-</sup>	Ca <sup>2+</sup>
Concentration inside (mM)	400	50	40	10 <sup>-4</sup>
Concentration outside (mM)	20	440	560	10

K<sup>+</sup>

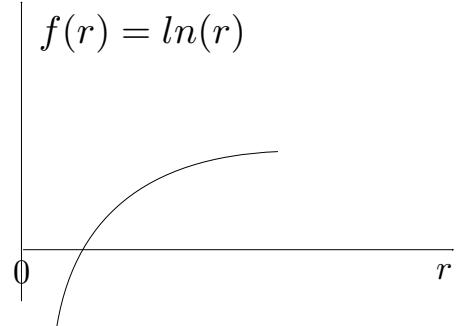
from Sterratt et al., 2011

$$f(r) = \ln(r)$$

Na<sup>+</sup>

Cl<sup>-</sup>

Ca<sup>2+</sup>



Ion	K <sup>+</sup>	Na <sup>+</sup>	Cl <sup>-</sup>	Ca <sup>2+</sup>
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from Sterrett et al., 2011

$$\frac{R T}{z F} = \frac{K T}{z q}$$

$$T = 300^\circ, z = 1$$

$$\approx 26mV$$



Walther Nernst  
(1864-1941)

$$E_{Nernst} = V_{in} - V_{out} = \frac{R T}{z F} \ln \left( \frac{c_{out}}{c_{in}} \right)$$

$$K^+ \quad 26mV \ln\left(\frac{20}{400}\right) = -77.88 mV$$

$$Na^+ \quad 26mV \ln\left(\frac{440}{50}\right) = +56.54 mV$$

$$Cl^- \quad -26mV \ln\left(\frac{560}{40}\right) = -68.6 mV$$

$$Ca^{2+} \quad 13mV \ln\left(\frac{10}{0.0001}\right) = +149.6 mV$$



How to get -70mV ??

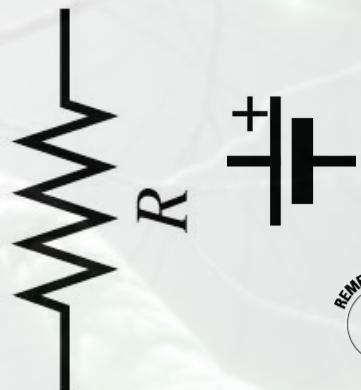
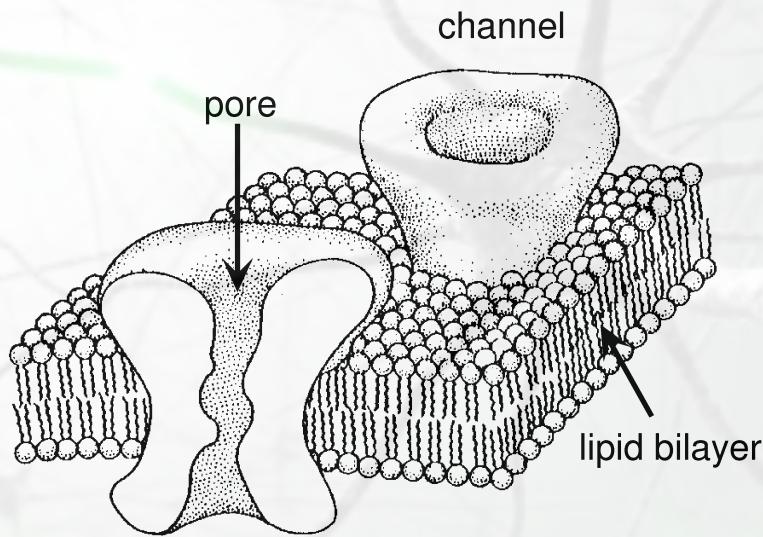
It must be from a combination of those potentials.

But.. how to combine them together??



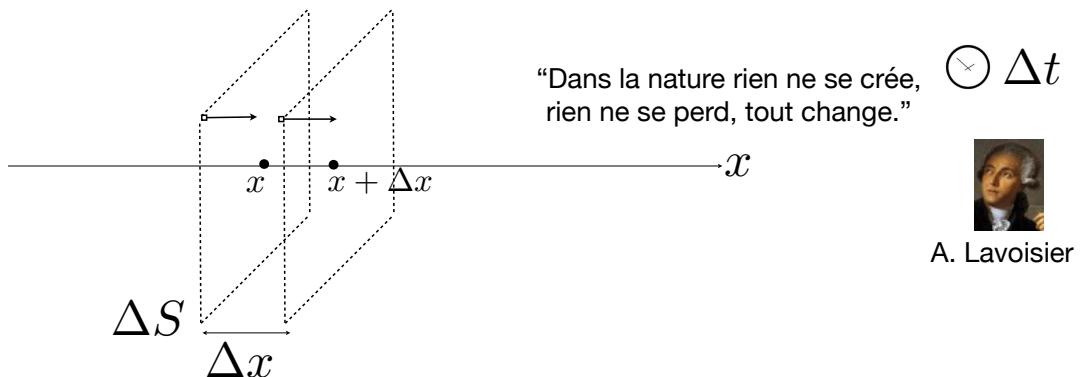
Is the “resistor” a good model for an ion channel...?  
**NOT ONLY A RESISTOR**

We first need to understand how ions flow across a permeable membrane !!!



## Electro-diffusion equation

invoking conservation of mass for charged particles in aq. solution



$$c(x, t + \Delta t) (\Delta S \Delta x) = \\ c(x, t) (\Delta S \Delta x) + J(x, t) \Delta t \Delta S - J(x + \Delta x, t) \Delta t \Delta S$$

## Electro-diffusion equation

invoking conservation of mass for charged particles in aq. solution

$$c(x, t + \Delta t) (\Delta S \Delta x) = c(x, t) (\Delta S \Delta x) + J(x, t) \Delta t \cancel{\Delta S} - J(x + \Delta x, t) \Delta t \cancel{\Delta S}$$

$$\frac{c(x, t + \Delta t) - c(x, t)}{\Delta t} = - \frac{J(x + \Delta x, t) - J(x, t)}{\Delta x}$$

$\downarrow$   
 $\Delta x \rightarrow 0$   
 $\Delta t \rightarrow 0$

$$\frac{\partial c(x, t)}{\partial t} = - \frac{\partial J(x, t)}{\partial x}$$

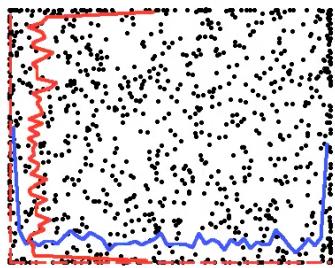
$$J = J_{diff} + J_{drift} \quad J = -D \frac{dc}{dx} - u c z F \frac{dV}{dx}$$

## Non-charged particles in aqueous solution Diffusion equation

$$\frac{\partial c(x, t)}{\partial t} = - \frac{\partial J(x, t)}{\partial x} \rightarrow \frac{\partial c}{\partial t} = -D \frac{\partial^2 c}{\partial x^2}$$

at the steady-state...  
 $c(x, t) = c(x)$

$$0 = -D \frac{\partial^2 c}{\partial x^2}$$



$$c(x) = k_1 + k_2 x$$

$$J(L, t) = -D \frac{\partial c}{\partial x} |_L = 0$$

$$c(x) = k_1$$

from the Nernst equation...  $V_{in} - V_{out} = 0$  this is a model!

single ion-specie, semi-permeable mem., NOT @equilibrium

$$J = - u R T c \frac{d}{dx} \ln(c) - u z F c \frac{d}{dx} V$$

$$J = - u c z F \left[ \frac{R T}{z F} \frac{d}{dx} \ln(c) + \frac{d}{dx} V \right]$$

**F** (Faraday constant)  
 $9.6 \cdot 10^4 \text{ C mol}^{-1}$

$$\frac{J z F}{u c z^2 F^2} = - \left[ \frac{R T}{z F} \frac{d}{dx} \ln(c) + \frac{d}{dx} V \right]$$

single ion-specie, semi-permeable mem., NOT @equilibrium

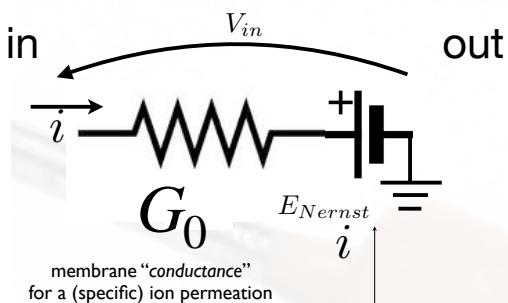
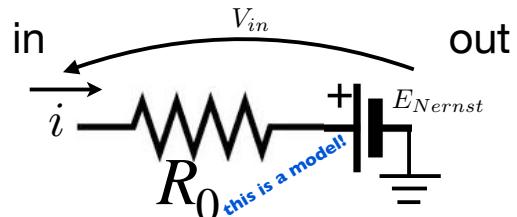
$$\int_{x_{in}}^{x_{out}} \frac{J z F}{u c z^2 F^2} dx = - \int_{x_{in}}^{x_{out}} \left[ \frac{R T}{z F} \frac{d}{dx} \ln(c) + \frac{d}{dx} V \right] dx$$

Ohmic approx  
non-Ohmic approx

$$i \stackrel{\text{current density}}{\leftarrow} J(zF) \frac{1}{uz^2 F^2} \int_{x_{in}}^{x_{out}} \frac{1}{c} dx \approx -\frac{RT}{zF} \ln \frac{c_{out}}{c_{in}} + (V_{in} - V_{out})$$

$i R_0 \approx (V_{in} - 0) - E_{Nernst}$

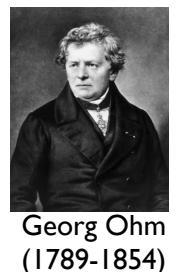
Ohmic approximation



membrane "conductance" for a (specific) ion permeation

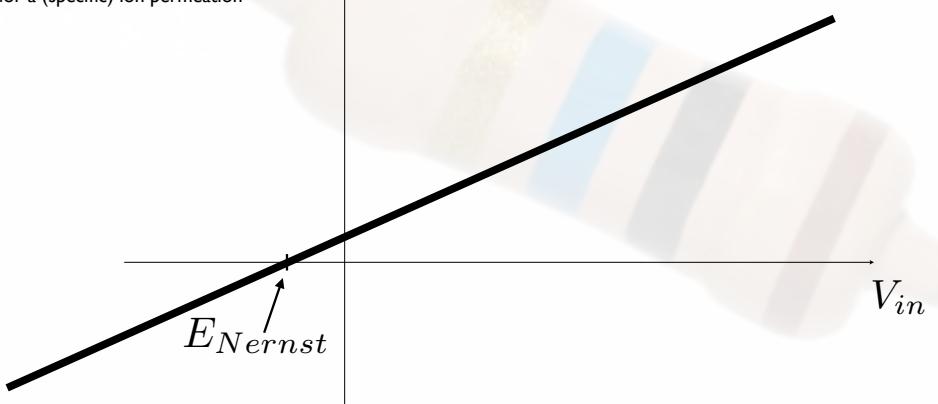
$i R_0 \approx (V_{in} - 0) - E_{Nernst}$

$i \approx G_0 (V_{in} - E_{Nernst})$



Georg Ohm  
(1789-1854)

$$G_0 = 1/R_0$$



This value is the Nernst equilibrium potential for that ionic species,  
or also known as "reversal potential"!



## Alternatives to the Ohmic-approximation?? **Goldman equation!**

**David E. Goldman**  
**(1910–1998)**

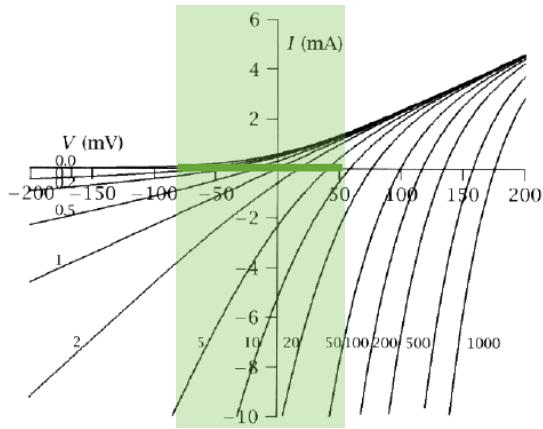


Figure 2.5 Current-voltage relations given by equation 2.7.17 (GHK current equation) for various values of  $[Cl_{out}]/[C_{in}]$  (indicated by small numbers near each curve).

- It accounts for strong (**non-linear**) **rectification** (inward or outward, depending on  $z$ , etc...)
- In the range of neuronal activity voltages and concentrations, often it can be **approximated by a straight line** (i.e. Ohmic approximation is its Taylor expansion!!)...  
e.g. For  $[Ca^{++}]$  very unbalanced concentrations, the Ohmic approximation is poor!

single ion-specie, semi-permeable mem., NOT @equilibrium

### Goldman equation: derivation!

$$J = u c F_{ext}$$

$$J = -u R T c \frac{d}{dx} \ln(c) - u z F c \frac{d}{dx} V$$

$$D = u R T$$

$$a = \frac{z F}{R T}$$

$$h(x) = u R T c(x) e^{a V(x)}$$

$$\frac{d}{dx} h(x) = u R T \left( e^{a V(x)} \frac{d}{dx} c(x) + c(x) a e^{a V(x)} \frac{d}{dx} V(x) \right)$$

$$\frac{d}{dx} h(x) = u R T c(x) e^{a V(x)} \left( \frac{1}{c(x)} \frac{d}{dx} c(x) + a \frac{d}{dx} V(x) \right)$$

$$\frac{d}{dx} h(x) = -J e^{a V(x)}$$

single ion-specie, semi-permeable mem., NOT @equilibrium

### Goldman equation: derivation!

$$D = u R T$$

$$\frac{d}{dx} h(x) = -J e^a V(x) \quad a = \frac{z F}{R T}$$

$$\int_{x_{in}}^{x_{out}} \frac{d}{dx} h(x) dx = - \int_{x_{in}}^{x_{out}} J e^a V(x) dx$$

$$h(x_{out}) - h(x_{in}) = - \int_{x_{in}}^{x_{out}} J e^a V(x) dx$$

- hp:  $\mathbf{J}$  does **NOT** depend on  $x$  (inside the membrane).
- hp: the electric field within the membrane is **uniform**; thus  $V(x)$  changes **linearly** (inside the membrane) - say  $\mathbf{V}(x) = \mathbf{m} x + \mathbf{p}$  (from  $V_{in}$  to  $V_{out}$ )

single ion-specie, semi-permeable mem., NOT @equilibrium

### Goldman equation: derivation!

$$D = u R T$$

$$a = \frac{z F}{R T}$$

$$h(x_{out}) - h(x_{in}) = -J \int_{x_{in}}^{x_{out}} e^{a(mx+p)} dx$$

$$h(x_{out}) - h(x_{in}) = -J \frac{1}{a m} \left( e^{a(m x_{out} + p)} - e^{a(m x_{in} + p)} \right)$$

$V(x) = \frac{x - x_{in}}{x_{out} - x_{in}} (V_{out} - V_{in}) + V_{in}$	$m = \frac{V_{out} - V_{in}}{x_{out} - x_{in}}$
--	---

$$h(x_{out}) - h(x_{in}) = -J \frac{1}{a m} \left( e^{a V_{out}} - e^{a V_{in}} \right)$$

$$h(x_{out}) - h(x_{in}) = -J \frac{1}{a} \frac{x_{out} - x_{in}}{V_{out} - V_{in}} \left( e^{a V_{out}} - e^{a V_{in}} \right)$$

single ion-specie, semi-permeable mem., NOT @equilibrium

### Goldman equation: derivation!

$$h(x) = u R T c(x) e^{a V(x)}$$

$$D = u R T$$

$$a = \frac{z F}{R T}$$

$$h(x_{out}) - h(x_{in}) = -J \frac{1}{a} \frac{x_{out} - x_{in}}{V_{out} - V_{in}} (e^{a V_{out}} - e^{a V_{in}})$$

$$u R T (c_{out} e^{a V_{out}} - c_{in} e^{a V_{in}}) = -J \frac{1}{a} \frac{x_{out} - x_{in}}{V_{out} - V_{in}} (e^{a V_{out}} - e^{a V_{in}})$$

$$J = u R T a \frac{V_{in} - V_{out}}{x_{out} - x_{in}} \frac{(c_{out} e^{a V_{out}} - c_{in} e^{a V_{in}})}{e^{a V_{out}} - e^{a V_{in}}}$$

single ion-specie, semi-permeable mem., NOT @equilibrium

### Goldman equation: derivation!

$$D = u R T$$

$$a = \frac{z F}{R T}$$

$$J = u R T a \frac{V_{in} - V_{out}}{x_{out} - x_{in}} \frac{(c_{out} e^{a V_{out}} - c_{in} e^{a V_{in}})}{e^{a V_{out}} - e^{a V_{in}}}$$

$$J = u R T a \frac{V_{in} - V_{out}}{x_{out} - x_{in}} \frac{(c_{out} - c_{in} e^{a (V_{in} - V_{out})})}{1 - e^{a (V_{in} - V_{out})}}$$

$$J = P a V_{in} \frac{c_{out} - c_{in} e^{a V_{in}}}{1 - e^{a V_{in}}} \quad P = \frac{u R T}{x_{out} - x_{in}}$$

single ion-specie, semi-permeable mem., NOT @equilibrium

### Goldman equation: Taylor series expansion?

$$P = \frac{u R T}{x_{out} - x_{in}}$$

$$a = \frac{z F}{R T}$$

$$J(V_{in}) = P a V_0 \frac{c_{out} - c_{in} e^{a V_{in}}}{1 - e^{a V_{in}}}$$

$$V_0 = \frac{RT}{zF} \ln \frac{c_{out}}{c_{in}}$$

$$J(V) \approx J(V_0) + \frac{d}{dV} J(V)|_{V_0} (V - V_0)$$

**Note:**  
 $c_{out} - c_{in} e^{\ln c_{out}/c_{in}} = 0$

$$J(V_0) = P a V_0 \frac{c_{out} - c_{in} e^{a V_0}}{1 - e^{a V_0}} = 0$$

$$\frac{d}{dV} J(V)|_{V_0} = P a \frac{c_{in} c_{out}}{c_{in} - c_{out}} \ln \frac{c_{out}}{c_{in}} = \frac{uzF}{x_{out} - x_{in}} \frac{c_{in} c_{out}}{c_{in} - c_{out}} \ln \frac{c_{out}}{c_{in}}$$

**always < 0**

What kind of fluxes can occur in solution?

Diffusive and Drift fluxes

$$J = -D \frac{dc}{dx}$$

$$J = u c \left( -z F \frac{dV}{dx} \right)$$

What does their knowledge predict?

Conduction across an ion channel (non-equilibrium): ohmic approximation

The Goldman equation (i.e., non-ohmic approximation)

Semi-permeable membranes and multiple ionic species

The existence of the resting membrane potential. Goldman Hodgkin Huxley eq.

## The membrane potential multiple ion-species, NOT at the equilibrium

Many ionic species (e.g. K<sup>+</sup>, Na<sup>+</sup>, Cl<sup>-</sup>, Ca<sup>++</sup>, Mg<sup>++</sup>,...)

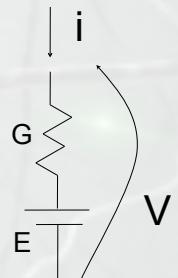
Biological membranes have distinct conductances (i.e. distinct  $G_h$ , h = Na, K, Cl...)

Ions have distinct reversal potentials (i.e. distinct  $E_h$ , h = Na, K, Cl...)

$$i_h = G_h(V - E_h) \quad G_h = 1/R_h \quad h = 1,2,3,\dots$$

Which is the total ionic current density across the membrane?

$$i_{tot} = i_1 + i_2 + i_3 + \dots$$



$$i_{tot} = G_1(V - E_1) + G_2(V - E_2) + G_3(V - E_3) + \dots + G_N(V - E_N)$$

$$= (G_1 + G_2 + G_3 + \dots + G_N)V - (G_1E_1 + G_2E_2 + G_3E_3 + \dots + G_NE_N)$$



## The membrane potential multiple ion-species, NOT at the equilibrium

At “rest” (steady-state), the **total current density vanishes...**

This is NOT (thermodynamical) equilibrium for each specie (i.e. death)

$$i_{tot} = 0$$

$$= (G_1 + G_2 + G_3 + \dots + G_N)V_{rest} - (G_1E_1 + G_2E_2 + G_3E_3 + \dots + G_NE_N)$$

$$V_{rest} = \frac{(G_1E_1 + G_2E_2 + G_3E_3 + \dots)}{(G_1 + G_2 + G_3 + \dots)}$$

It is **NOT** the algebraic sum of the individual Nernst potentials (one for each ion specie)! **Conductances** define the membrane potential!



## Goldman Hodgkin Katz equation

$$J_1 = P_1 a V_{in} \frac{c_{out\ 1} - c_{in\ 1} e^{a V_{in}}}{1 - e^{a V_{in}}}$$

$$J_2 = P_2 a V_{in} \frac{c_{out\ 2} - c_{in\ 2} e^{a V_{in}}}{1 - e^{a V_{in}}}$$

:

$$J_k = P_k a V_{in} \frac{c_{out\ k} - c_{in\ k} e^{a V_{in}}}{1 - e^{a V_{in}}}$$

$$J_{tot} = a \frac{V_{in}}{1 - e^{a V_{in}}} \left( \sum_k P_k c_{out\ k} - e^{a V_{in}} \sum_k P_k c_{in\ k} \right)$$

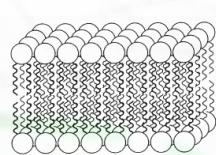
$$J_{tot} = 0 \xrightarrow{\text{"Weighted sum"}} \sum_k P_k c_{out\ k} = e^{a V_{in}} \sum_k P_k c_{in\ k}$$

$$V_{rest} = \frac{RT}{zF} \ln \left( \frac{\sum_k P_k c_{out\ k}}{\sum_k P_k c_{in\ k}} \right)$$

$$E_{Nernst} = V_{in} - V_{out} = \frac{R T}{z F} \ln \left( \frac{c_{out}}{c_{in}} \right)$$

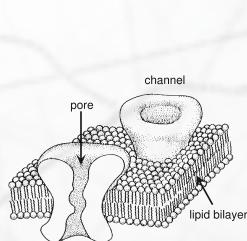
$$a = \frac{z F}{R T}$$

**hp:**  $z$  is exactly the same for all ionic species  
( $a$  does not depend on  $k$ )



$$C = \frac{\Delta Q}{\Delta V} \quad i = C \frac{dV}{dt}$$

$$c = 1 \mu F/cm^2$$



$$i = G (V_{in} - E) \quad E = \frac{RT}{zF} \ln \left( \frac{c_{out}}{c_{in}} \right)$$



## **Next stop: excitability**

### Concepts and Discoveries

**1872-1905:** Hermann proposed that propagation is an electrical self-stimulation of the axon by **inward currents** spreading passively from excited region to neighbouring unexcited regions;

**1902-1912:** Bernstein **proposed** that potentials might arise across a membrane that is selectively permeable and separates solutions of different ionic concentrations and that excitation involves an **increase in permeability**;

**1938:** Cole & Curtis, **experimentally did find** changes in ionic permeability

**1949:** Hodgkin & Katz, showed **inward currents**, by **selective permeability** to  $\text{Na}^+$

**1952:** Hodgkin & Huxley, described **how ionic permeability changes in time.**