Layered-SQT usage

M. Grady Saunders mgs8033@rit.edu

1. Introduction

This document describes layered-sqt, a command-line tool for simulating the Bidirectional Scattering Distribution Function (BSDF) which emerges from a layered assembly. In this context, a *layered assembly* is a theoretical construction consisting of N layers separated by N+1 participating media for $N \geq 1$. A layer is an infinite plane which is offset along (and normal to) the z-axis, and which is associated with a constituent BSDF that describes how light is scattered upon intersection. A medium is thought to occupy the space between adjacent layers or, in boundary cases, the spaces above and below the top and bottom layers respectively. See figure 1 for clarification.

The *emergent BSDF* is the BSDF observed in the limit as one backs infinitely far away from the assembly or, identically, as the assembly shrinks to an infinitesmial point. To ensure that the emergent BSDF is well-defined, layered-sqt requires that participating media be homogeneous, i.e., that scattering properties be independent of spatial location. Furthermore, for simplicity/tractability, layered-sqt does not account for wavelength-dependence.

1.1. Basics of LSQT format

The structure of a layered assembly is easy enough to convey in plain-text with rudimentary syntax, so this is the format layered-sqt accepts as input. By convention, we refer to this format as "LSQT format", and we suffix associated filenames with the extension .lsqt (though this suffix is not strictly required for the program to run).

An LSQT file is therefore a line-by-line plain-text description of a layered assembly from top to bottom. So, the first line describes the top medium, the second line describes the top layer beneath the top medium, the third line describes the medium beneath the top layer, and so on until the bottom medium. That being the case, odd-numbered lines describe media and even-numbered lines describe layers. For media as well as layers, the general syntax is

Name key1=val1 key2=val2

where Name is described by keyword arguments key1 and key2—importantly, this syntax is whitespace-delimited, so keyword arguments of the form key=val must not contain

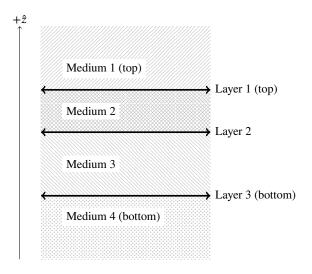


Figure 1: An example diagram of a layered assembly with 3 layers and 4 participating media.

whitespace. It is also worth mentioning that keyword arguments may appear in any order.

1.2. Basics of program usage

As layered-sqt is a command-line tool, it runs on the command line, whereby it scans command-line arguments for (optional) configuration flags and a (required) input filename appearing somewhere as a positional argument, i.e., an argument not consumed by a flag. It then parses the input file, simulates the emergent BSDF, and writes the results to a RAW-format file ready for conversion to SQT-format via raw2sqt. For instance,

simulates a layered assembly as described in example.lsqt with 50,000 paths, and writes the emergent BSDF in plaintext RAW-format to example.raw. To convert the BSDF to SQT-format, which is necessary for use in DIRSIG, run

\$./raw2sqt example.raw

which will write a new file example.sqt.

Two of the most common flags appear above, namely -p (or --path-count) to specify the number of paths used in the simulation and -o (or --output) to specify the output filename. To see a list of all acceptable flags with brief descriptions and default values, pass -h (or --help), or simply run layered-sqt with no filename. As an aside, layered-sqt verifies parameters specified in command-line flags as well as keyword arguments specified in the input LSQT file. In the event that something has an unreasonable value, the program issues an error and fails in a controlled manner.

Lastly, layered-sqt recognizes the single dash filename "-" as standard input. So, it is possible to pipe the (presumably LSQT format) output of a script into layered-sqt directly, if this happens to be convenient. As a trivial example,

```
$ cat example.lsqt | \
   ./layered-sqt -p 50000 -o example.raw -
```

is equivalent to just passing example.lsqt.

2. Tutorial

This section provides a series of progressively more interesting "learn-by-doing" tutorials, which should naturally introduce the concepts and features in layered-sqt. To get started, create a working directory lsqt-tutorial and link the relevant programs for easy access. From one of the Carlson CIS servers, type the following commands:

```
$ mkdir lsqt-tutorial & cd lsqt-tutorial
$ TMP_PATH=/cis/phd/mgs8033/layered-sqt/bin
$ ln -s $TMP_PATH/layered-sqt
$ ln -s $TMP_PATH/layered-sqt-lssview
$ ln -s /dirs/pkg/dirsig/bin/raw2sqt
```

2.1. Hello, world

Let's start with something mind-numbingly simple—a 60% reflective Lambertian surface. First, create and enter a sub-directory tutorial1.

```
$ mkdir tutorial1 & cd tutorial1
```

Using whatever text-editing program you prefer, create a plain-text LSQT file Lambertian.lsqt with three lines.

```
Medium
Layer z=0 Lambertian fR=0.6
Medium
```

As an aside, LSQT format input is syntax-highlighted in this document. However, syntax-highlighting files/add-ons are not available (yet) for any text editor.

Now, run layered-sqt on Lambertian.lsqt,

\$../layered-sqt Lambertian.lsqt

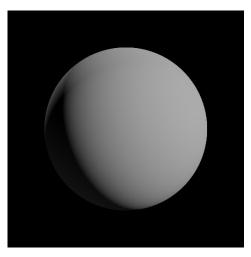


Figure 2: Image preview of a 60% reflective Lambertian BRDF output by layered-sqt-lssview, as in tutorial §2.1.

which should write two new files: Lambertian.lsqt.lss and Lambertian.lsqt.raw. As stated in the introduction, we may convert the plain-text RAW file to a binary SQT file for use with DIRSIG by running raw2sqt,

\$../raw2sqt Lambertian.lsqt.raw

which generates the final SQT file Lambertian.lsqt.sqt suitable for use in a DIRSIG scene.

Now, what is the LSS file? *LSS* is an initialism for *LSQT-Slice*, which is the name of the internal file format layered-sqt uses to store simulation data, and it is useful 1) for previewing the layered BSDF *without setting up and rendering a DIRSIG scene*, and 2) for simulating a layered BSDF progressively, with multiple runs of layered-sqt.

Run layered-sqt-lssview to preview the Lambertian BRDF.

\$../layered-sqt-lssview Lambertian.lsqt.lss

This should write a new file, Lambertian.lsqt.lss.png, which is a 512x512 rendering of the BSDF applied to a ball, shown in figure 2. When connected to the CIS server with X-window support (from a Linux machine, log-in with ssh-X), this image may be viewed by running eog.

\$ eog Lambertian.lsqt.lss.png

Importantly, layered-sqt-lssview is *not* a full-blown path-tracer, and may not perfectly represent how the BSDF will appear in DIRSIG. It only accounts for the direct (first bounce) contributions of a few directional light sources, and it further uses tone-mapping and sRGB correction, such that the output image is not suitable for any radiometric analysis. The intended use of this preview image is to determine if a simulated BSDF is suitably convergent/noise-free.

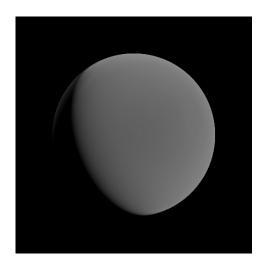


Figure 3: Image preview of a 60% reflective Lambertian BRDF with a rough microsurface output by layered-sqt-lssview, as in tutorial §2.2.

For more information on the Lambertian BSDF, refer to the documentation in §3.2.2.

2.2. Hello, world, with roughness

Let's expand on §2.1 by simulating a 60% reflective Lambertian BRDF with a rough microsurface—we assume that, microscopically, the surface geometry is not perfectly smooth, but is instead characterized by a distribution of normals about the primary surface normal. See the documentation in §3.2.3 for a more detailed explanation.

First, return to the lsqt-tutorial directory, and create and enter a new subdirectory tutorial2.

```
$ cd ..
$ mkdir tutorial2 & cd tutorial2
```

As before, create a new plain-text LSQT file Rough.lsqt with three lines.

```
Medium
Layer z=0 MicrosurfaceLambertian fR=0.6
    alpha=2.4
Medium
```

Above, fR=0.6 specifies that the Lambertian BRDF is 60% reflective, and alpha=2.4 sets the roughness parameter α to be 2.4—this is very rough. Note that setting alpha=0 is effectively equivalent to using Lambertian with fR=0.6 and fT=0.

Now, run layered-sqt on Rough.lsqt to generate files Rough.lsqt.lss and Rough.lsqt.raw. This should take longer to simulate than before, due to the computational complexity of the microsurface model.

```
$ ../layered-sqt Rough.lsqt
```

When this is complete, use layered-sqt-lssview to generate Rough.lsqt.lss.png, shown in figure 3.

\$../layered-sqt-lssview Rough.lsqt.lss

Notice that the microsurface roughness causes increased reflectance at grazing angles, and decreased reflectance elsewhere.

To produce the final SQT file suitable for use in a DIRSIG scene, run raw2sqt as in the first tutorial.

```
$ ../raw2sqt Rough.lsqt.raw
```

We'll omit this step in the next few tutorial sections, as it is always the same. We'll return to this in tutorial §2.6 however, which details the process for linking up an SQT file in a DIRSIG scene.

2.3. A simple substrate

Let's now consider a simple two-layer assembly to model a substrate. We imagine a 60% reflective diffuse Lambertian surface with a glossy dielectric coating. Before continuing, again return to the lsqt-tutorial directory, and create and enter a new subdirectory tutorial3.

```
$ cd ..
$ mkdir tutorial3 & cd tutorial3
```

Then create a new plain-text LSQT file Substr.lsqt with five lines as below.

```
Medium
Layer z=1 MicrosurfaceDielectric alpha=0.2
Medium eta=1.4
Layer z=0 Lambertian fR=0.6
Medium
```

This assembly has two layers, appearing on lines 2 and 4. The first layer is positioned at *z*-height 1 with a MicrosurfaceDielectric—this is similar to MicrosurfaceLambertian, but it applies the microsurface ideology to the dielectric Fresnel mirror BSDF instead of the Lambertian BRDF. This layer forms the surface of our glossy dielectric coating. The second layer is positioned at *z*-height 0 with a Lambertian as in tutorial §2.1.

The medium between these layers on line 3 is attributed a refractive index η (eta) of 1.4. Note that *every* medium in an LSQT file has a refractive index, which defaults to 1 if left unspecified. Note also that MicrosurfaceDielectric depends on the refractive indices of the media above and below the layer to which it is assigned, as it generalizes the dielectric Fresnel mirror BSDF.

Run layered-sqt as usual and preview the resulting LSS file with layered-sqt-lssview.

```
$ ../layered-sqt Substr.lsqt
$ ../layered-sqt-lssview Substr.lsqt.lss
```

Now, upon viewing Substr.lsqt.lss.png, we might suspect that something has gone wrong. As shown in figure 4, there is quite a bit of low frequency noise and artifacting. What is going on?

Behind the scenes, layered-sqt estimates values of the emergent BSDF at a finite number of sample directions using Monte Carlo integration, i.e., integration by averaging random evaluations. So, two central parameters controlling simulation fidelity are 1) the number of sample directions and 2) the number of random evaluations used to estimate the BSDF at each sample direction. By default, layered-sqt uses 80 sample directions and 10,000 random evaluations of potential light paths to estimate the emergent BSDF.

We can change the number of directions ω_i on the command line using the flag -wi (or --wi-count), and we can change the number of paths using the flag -p (or --path-count). In general, we should think to increase the number of directions to increase *resolution*, and we should think to increase the number of paths to reduce *noise*.

In the case of our simple substrate BRDF, let's increase the number of directions from 80 to 300, and the number of paths from 10,000 to 100,000. Run layered-sqt with the following flags:

```
$ ../layered-sqt -R -wi=300 \
   -p=100000 -rp=0.3 Substr.lsqt
```

This should take a minute or two to compute. Upon previewing the BRDF with layered-sqt-lssview, we see something much more sensible, as shown in figure 5.

The flag -R (or -- restart) tells layered-sqt to ignore the LSS file and restart the simulation from scratch. If the LSS file exists, then layered-sqt will ignore all flags except -p/--path-count by default, and add this number of paths to the existing simulation data (it will also display a message to remind us that is doing this). So, -R is necessary to regenerate sample directions.

We have specified the number of sample directions with -wi, but it remains a non-trivial problem to determine what these directions ought to be—layered-sqt implements a presumably novel algorithm to form a so-called *Redundancy Reduced Sample Set (RRSS)* of directions by initially sampling *more* directions than the user requests, then choosing the subset of those directions that minimizes a quantitative measure of "redundancy" with respect to a desirable density, being the normalization of the non-cosine-weighted BSDF.

The flag -rp (or --rrss-path-frac) thus specifies the fraction of paths (as specified by -p) used to initially estimate the emergent BSDF in forming the RRSS of directions ω_i . There is also the flag -rx (or --rrss-oversampling), to specify the RRSS oversampling multiplier. (We discuss this flag for completeness, but in practice it is never really necessary to specify -rx, as the default of 4 is generally sufficient.) For example,

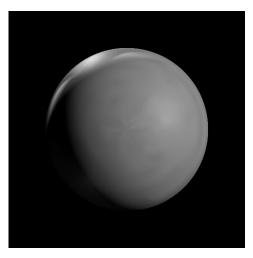


Figure 4: Image preview of the simple substrate BRDF in tutorial §2.3, as output from running layered-sqt with default settings. As is evident, the BRDF would benefit from more samples (higher resolution), and more simulated paths (less noise).

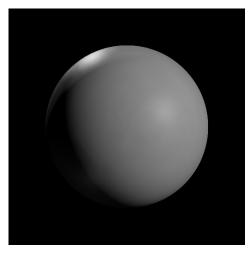


Figure 5: Image preview of the simple substrate BRDF in tutorial §2.3, as output from running layered-sqt with higher settings. In particular, this uses 300 incident directions (instead of the default 80) and 100,000 paths (instead of the default 10,000).

```
$ layered-sqt \
  -wi=100 -p=20000 -rp=0.25 -rx=5 File.lsqt
```

samples $100 \times 5 = 500$ directions initially, and uses $20000 \times 0.25 = 5000$ paths to estimate the emergent BSDF at each of these 500 directions, then forms an RRSS containing 100 of these 500 directions, and lastly uses 20000 - 5000 = 15000 paths to improve the estimates of the emergent BSDF at each of these 100 directions (such that the total number of paths is 20000).

In general, increasing -rp and -rx leads to better sample placement, at the (potentially extreme) cost of computation time. We discuss practical simulation strategies using these flags, including incremental simulation with multiple runs of layered-sqt, over the remaining tutorial sections.

2.4. Adding dust

Let's now add a layer of back-scattering dust to the surface of the substrate BRDF in tutorial §2.3. As usual, return to the lsqt-tutorial directory, and create and enter a new subdirectory tutorial4.

```
$ cd ..
$ mkdir tutorial4 & cd tutorial4
```

Next create another plain-text LSQT file Dusty.lsqt with three layers and thus seven lines.

```
Medium
Layer z=2 Null
Medium mus=0.2 HenyeyGreenstein g=-0.4
Layer z=1 MicrosurfaceDielectric alpha=0.2
Medium eta=1.4
Layer z=0 Lambertian fR=0.2
Medium
```

We attribute a Null to the topmost layer appearing on line 2, such that this layer only exists to separate media, and no scattering happens at the layer itself. The medium above on line 1 is vacuum as usual. The medium below on line 3 models our layer of dust. Lines 4 through 7 are just lines 2 through 5 from the previous section, except we have reduced f_R from 60% to 20% to exaggerate the dust effect.

Similarly to refractive index η (eta), every medium in an LSQT file has two volume scattering parameters—the scattering coefficient μ_s (mus) and the absorption coefficient μ_a (mua), which default to zero if left unspecified. Scattering and absorption events happen according to Beer's Law in a homogeneous medium. That is, the probability of scattering within a particular distance d is given by an exponential distribution $1 - \exp(-\mu_s d)$, where μ_s is the distribution parameter, and analogously for absorbing within a particular distance d and μ_a . We thus identify the units of μ_s and μ_a as "inverse distance", and we interpret the reciprocals $1/\mu_s$ and

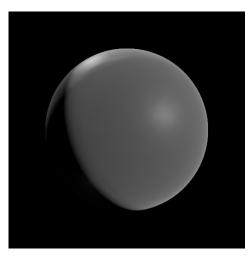


Figure 6: Image preview of the modified simple substrate BRDF with back-scattering dust in tutorial §2.4, where the bottom Lambertian BRDF layer is reduced to 20% reflectance to exaggerate the dust effect.

 $1/\mu_a$ as the mean distances between scattering and absorption events respectively. Note that in the limit $\mu_s = \mu_a = 0$ (the default values), the mean distances tend to infinity, such that scattering and absorption events cease to exist.

To simulate a thin layer of back-scattering dust, we attribute a scattering coefficient μ_s (mus) of 0.2 and a Henyey-Greenstein with shape parameter g of -0.4 (more on this in a moment). We choose μ_s to be 0.2 such that the mean distance between scattering events is $1/\mu_s = 5$ units. This is somewhat large in comparison to the thickness of the dust layer, which is 1 unit (the difference of the z-heights of layers on lines 2 and 4), so the dust appears suitably "thin". When plucking this number out of the sky, this is the line of reasoning to use.

The Henyey-Greenstein phase function is widely used in computer graphics due to its simplicity and intuitive parameterization (although it was originally intended to model scattering of interstellar dust clouds). The phase function is equipped with a single shape parameter $g \in (-1, +1)$, which is identically the mean scattering cosine. That is, it tends to forward scatter as $g \to +1$ and back scatter as $g \to -1$. See the documentation in §3.1.1 for more information.

Now run layered-sqt on Dusty.lsqt with the same flags as in the previous section.

```
$ ../layered-sqt -wi=300 \
  -p=100000 -rp=0.3 Dusty.lsqt
```

Notice that we do not need -R this time because no LSS file exists yet. As before, this should take a minute or two to compute. The image preview as output by layered-sqt-lssview is shown in figure 6.

2.5. Conductors

So far we have considered only dielectric, or electrically-insulating, materials. In the context of simulation, a dielectric material is thought to have a real-valued refractive index η , such that absorption is modeled by a simple multiplier (at a surface) or a Beer's Law coefficient μ_a (in a volume). On the other hand, a conductive material is thought to have a complex-valued index of refraction $\eta + i\kappa$.

At a boundary between between a dielectric material with refractive index η_1 and a conductive material with refractive index $\eta_2 + i\kappa$, the relative refractive index is thus a complex number

$$\eta_{\rm rel} = \frac{\eta_1}{\eta_2 + i\kappa}.$$

Analogously to the dielectric case, we may apply Fresnel's equations with such a complex refractive index. Contrary to the dielectric case however, the transmitted wave is now *evanescent*, and diminishes quickly to zero. The coefficient κ is thus termed the *absorption coefficient* of the conductive material. Note that this is a different "theory" of absorption than the dielectric case. Regardless, the idea of κ as a material property is spiritually equivalent to μ_a . For this reason, layered-sqt interprets μ_a as κ when the Microsurface-Conductive BRDF is used.

Once again, return to the lsqt-tutorial directory and create and enter a new subdirectory tutorial5.

```
$ cd ..
$ mkdir tutorial5 & cd tutorial5
```

Next create a new plain-text LSQT file Silver.lsqt with three lines.

```
Medium
Layer z=0 MicrosurfaceConductive alpha=0.2
Medium eta=0.051585 mua=3.9046
```

The top medium on line 1 is vacuum with $\eta=1$ by default as usual. The layer on line 2 uses a MicrosurfaceConductive BRDF with relatively smooth roughness $\alpha=0.2$. The bottom medium on line 3 models silver (Ag) at $0.5876\,\mu m$, with refractive index

$$\eta + i\kappa = 0.051585 + i3.9046$$
,

obtained from RefractiveIndex.INFO. Note that this configuration would not make proper sense if MicrosurfaceDielectric were used, in which case $\eta=0.051585$ would be interpreted as the real refractive index of a dielectric, which should not be less than 1, and μ_a would be interpreted as a Beer's Law coefficient of an infinite medium. Fortunately, layered-sqt ought to catch invalid configurations such as this, and report an error accordingly.

This BRDF has a reasonably sharp specular feature, so we want more incident directions than usual. Run layered-sqt as follows.

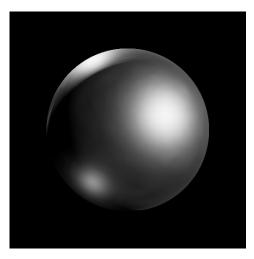


Figure 7: Image preview of the conductive BRDF in tutorial §2.5.

\$../layered-sqt -wi=600 -p=1 Silver.lsqt

This may seem odd—600 directions, but only 1 path? By default, the MicrosurfaceConductive BRDF is deterministic (i.e., multiple scattering is disabled, see sections §3.2.3–§3.2.5), such that it is resolvable with 1 path. Note that this is a very specific use case, being 1 layer with a deterministic BRDF, designed to illustrate why layered-sqt allows the path count to be 1 (99.9% of the time, it is unwise to do this).

Run layered-sqt-lssview on Silver.lsqt.lss to obtain Silver.lsqt.lss.png, as in figure 7. Consider rerunning layered-sqt (don't forget -R) with fewer incident directions to see how artifacts manifest in the image preview. Such artifacts may be less noticeable once converted into SQT, which implements a smoother interpolation method than that used in layered-sqt-lssview.

2.6. Using SQT in a DIRSIG scene

We have now spent some time running basic simulations in layered-sqt. So, let's take a brief hiatus to review how to actually use an SQT file in a DIRSIG scene. Presenting a step-by-step guide to setting up a DIRSIG scene in general is out of the scope of this document—we'll instead describe the relevant parts of a pre-made scene, which is available at /cis/phd/mgs8033/lsqt-scene. The top-level directory structure is as follows.

```
lsqt-scene/
    geometry/
    materials/
    shots/
    lsqt-scene.scene
```

The geometry subdirectory contains a few GList (XML)

files that describe the geometry in the scene. In this example scene, there is only a ground plane and a row of 3 spheres positioned on top of it. Importantly, the GList description of an object specifies its material ID number (inside the <matid>tag). Here, the ground plane is assigned material ID 1, and the spheres are assigned material IDs 2, 3, and 4.

The materials subdirectory contains materials.mat, the main material description file for the scene, as well as the SQT files corresponding to Substr.lsqt, Dusty.lsqt, and Silver.lsqt. Each entry in materials.mat defines a material, assigns it a name and ID number, and associates relevant data files with it. To illustrate, the relevant portion of the substrate material entry is shown below.

```
MATERIAL_ENTRY {
   ID = 2
   NAME = Dusty
   ...
   SURFACE_PROPERTIES {
     REFLECTANCE_PROP_NAME = SphericalDataBRDF
     REFLECTANCE_PROP {
        DATA_FILENAME = Substr.lsqt.sqt
     }
   }
}
```

This defines a material with ID 2 and name "Substr" which uses an SQT BRDF. Specifically, the property name SphericalDataBRDF indicates to DIRSIG that we want to use an SQT file to describe the BRDF of the material, which is then specified as Substr.lsqt.sqt.

TODO

3. Documentation

This section documents the properties of participating media and scattering layers in more detail, and enumerates the available scattering models for media and layers, and how to specify them in LSQT format. Note that "phase function", "BRDF", and "BSDF" are specific types of scattering models. In particular,

- a *phase function* is a scattering model which accounts for volume-scattering within a participating medium,
- a BRDF, i.e., a Bidirectional Reflectance Distribution Function, is a scattering model which accounts for only reflection at a surface, and
- a BSDF, i.e., a Bidirectional Scattering Distribution Function, is a scattering model which accounts for reflection and transmission at a surface.

Note that phase functions are normalized by definition, as volume-absorbtion is accounted for by a separate parameter. BRDFs and BSDFs, however, account for absorption

as well as scattering, and so are only normalized if the model is non-absorbing (in other words, perfectly energy conserving). To state all of this more rigorously, a phase function p is normalized with respect to integration over incident directions ω_i , such that

$$\int_{\mathcal{S}^2} p(\omega_o \to \omega_i) \, \mathrm{d}\omega_i = 1.$$

However, a BRDF/BSDF f is only normalized as such if it is non-absorbing,

$$\int_{S^2} f(\omega_o \to \omega_i) d\omega_i = 1 \iff f \text{ is non-absorbing.}$$

As this ought to suggest—this document follows the convention that BRDFs and BSDFs contain an implicit cosine-weighting with respect to incident direction. This is often written explicitly elsewhere in the literature, such that the above normalization condition is written as

$$\int_{\mathcal{S}^2} f(\omega_o \to \omega_i) |\cos \theta_i| \, \mathrm{d}\omega_i = 1.$$

We effectively make the substitution $f \leftarrow f |\cos \theta_i|$.

3.1. Participating media

A medium is characterized by its absolute index of refraction $\eta > 0$, its absorption coefficient $\mu_a \ge 0$, its scattering coefficient $\mu_s \ge 0$, and a local phase function model.

To specify a medium, simply type the name Medium followed by (optional) keywords arguments eta, mua, and mus for η , μ_a , and μ_s respectively. Keyword arguments for a medium have default values eta=1, mua=0, mus=0, such that any may be omitted for brevity. If $\mu_s > 0$, the keyword arguments must be followed by the name of the desired phase function, followed by any keyword arguments for the phase function. If $\mu_s = 0$, it is not necessary to specify a phase function since no scattering occurs.

As the defaults may suggest, layered-sqt considers that vacuum is just a medium with refractive index 1 which neither absorbs nor scatters. So,

Medium

specifies vacuum.

3.1.1. Henyey-Greenstein phase

The Henyey-Greenstein phase function is given by

$$p(\omega_o \to \omega_i) = \frac{1}{4\pi} \frac{1 - g^2}{(1 + g^2 + 2g\omega_o \cdot \omega_i)^{3/2}}$$

with shape parameter $g \in (-1, 1)$. It may be helpful to note that g is the mean cosine of the scattered direction with respect to the incident direction, such that p becomes forward scattering as $g \to 1$, back scattering as $g \to -1$, and uniformly/isotropically scattering as $g \to 0$. See d'Eon's write-up [1, p. 19] for more information.

To specify the Henyey-Greenstein phase in LSQT format, use the name HenyeyGreenstein followed by the (optional) keyword argument g for g, which is zero by default. For example,

Medium mus=1 HenyeyGreenstein g=-0.22

specifies a moderately back-scattering medium.

3.1.2. Henyey-Greenstein 2-lobe phase

For convenience, layered-sqt provides a 2-lobe variant of the Henyey-Greenstein phase as described in §3.1.1. This is expressed as

$$p(\omega_o \to \omega_i) = (1 - b)p_0(\omega_o \to \omega_i) + bp_1(\omega_o \to \omega_i)$$

where $b \in [0, 1]$ is the "blend factor" which interpolates between Henyey-Greenstein phase functions p_0 and p_1 , which have shape parameters g_0 and g_1 respectively. This is useful to characterize asymmetric scattering.

To specify the Henyey-Greenstein 2-lobe phase in LSQT format, use the name HenyeyGreenstein2 followed by keyword arguments g0 for g_0 , g1 for g_1 , and g0 for g1, and g1 for g2 for g1, and g1 for g2 for example,

specifies a medium with a sharply forward-scattering component, which is otherwise mostly (65%) isotropic.

3.1.3. Rayleigh phase

The general form of the Rayleigh phase function is given by

$$p(\omega_o \to \omega_i) = \frac{3}{16\pi} \left[\frac{1+3\gamma}{1+2\gamma} + \frac{1-\gamma}{1+2\gamma} (\omega_o \cdot \omega_i)^2 \right]$$

where

$$\gamma = \frac{\rho}{2 - \rho}$$

where, in turn, $\rho \in [-1, 1]$ is a shaping parameter known as the *depolarization factor*. The simpler, more common form of the Rayleigh phase function is given by

$$p(\omega_o \to \omega_i) = \frac{3}{16\pi} (1 + (\omega_o \cdot \omega_i)^2),$$

which is obtained by setting $\rho = 0$. See d'Eon's write-up [1, p. 19] for more information.

The Rayleigh phase function prefers parallel scattering to perpendicular scattering. That is to say, scattering in directions perpendicular to the direction of propagation is diminished, and scattering is symmetric in terms of forward-versus back-scattering. This effect is maximized as $\rho \to -1$ and minimized as $\rho \to +1$. Note that p becomes uniformly/isotropically scattering in the limiting case $\rho = +1$.

To specify the Rayleigh phase function in LSQT format, use the name Rayleigh followed by the (optional) keyword argument rho for ρ , which is zero by default. For example,

Medium mus=1 Rayleigh rho=-0.5

specifies a medium with an exaggerated "Rayleigh effect" relative to the common form where $\rho = 0$.

3.1.4. SGGX phase

The Symmetric GGX (SGGX) phase function, given by Heitz et al. [4], is based on *microflake theory*, wherein we assume volume scattering happens due to surface scattering with infinitely many microscopic disjoint flakes. The general form of the phase function is

$$p(\omega_o \to \omega_i) = \int_{\mathcal{S}^2} p_m(\omega_m, \omega_o \to \omega_i) D_{\omega_o}(\omega_m) \, d\omega_m$$

where p_m is the microscopic phase function assigned to each flake, and D_{ω_o} is the distribution of *visible* flake normals for a particular viewing/outgoing direction ω_o .

The distribution of flake normals D is defined by its projected areas with respect to three orthogonal directions. In particular,

$$D(\omega_m) = \frac{1}{\pi \sqrt{\det \mathbf{S}} (\omega_m \cdot \mathbf{S}^{-1} \omega_m)^2}$$

where S is a 3 × 3 real symmetric matrix. The eigenvalues of S are squared projected areas, and the eigenvectors of S and the orthogonal coordinate directions. The distribution of *visible* normals D_{ω_n} is obtained from D as

$$D_{\omega_o}(\omega_m) = \frac{D(\omega_m)}{\sigma(\omega_o)} \langle \omega_o, \omega_m \rangle$$

where σ yields projected area along ω_o and

$$\langle \omega_o, \omega_m \rangle = \begin{cases} \omega_o \cdot \omega_m & (\omega_o, \omega_m) \text{ in same hemisphere,} \\ 0 & \text{otherwise.} \end{cases}$$

So, unlike the phase functions in the previous sections, the SGGX phase function has a global orientation, meaning that p depends on the specific values of directions ω_o and ω_i , and not just on the angle between them.

Due to the constraint in layered-sqt that the emergent BSDF must be isotropic, we cannot allow all configurations of S. If S is written as the eigendecomposition $S = Q^T \Lambda Q$, then 1) we fix Q = I, the identity matrix, and we 2) require that the eigenvalues of Λ describing projected area in the XY plane be identical. We thus obtain

$$\mathbf{S} = \begin{bmatrix} A_{\parallel}^2 & 0 & 0 \\ 0 & A_{\parallel}^2 & 0 \\ 0 & 0 & A_{\perp}^2 \end{bmatrix}$$

with free parameters $A_{\parallel} > 0$ and $A_{\perp} > 0$.

To specify the SGGX phase function in LSQT format, use the name Sggx followed by (optional) keyword arguments Apara for A_{\parallel} , Aperp for A_{\perp} , and type for the microphase function type (either Specular or Diffuse). For example,

```
Medium mus=1 Sggx Apara=0.5 type=Diffuse
```

By default, Apara=1, Aperp=1, and type=Specular. Note that only the ratio of A_{\parallel} to A_{\perp} matters, so it always acceptable to fix $A_{\perp}=1$ and just specify A_{\parallel} (or vice versa).

3.2. Layers

A layer is characterized by its z-height and a local BSDF model. As different BSDF models require different parameters, there is a different "type" of layer for each BSDF model implemented in layered-sqt.

Every layer accepts *z*-height as a parameter, so every line describing a layer must begin with the name Layer followed by the (required) keyword argument z. This is followed in turn by the name of the desired BSDF, followed by any keyword arguments for the BSDF. For example,

```
Layer z=1 Null
```

specifies a layer at *z*-height 1 with a null BSDF. As explained in the next sub-section, a null BSDF is a special case which accepts no additional keyword arguments.

The implementation requires that layers appear in top-tobottom order, such that the z-heights of subsequent layers in an LSQT file are strictly decreasing. If this condition is violated, the implementation issues an error and exits.

3.2.1. Null BSDF

At times, it is desirable to separate participating media without scattering at a layer. This is possible in layered-sqt by assigning a null BSDF with name Null, which may be interpreted as a 100% transmissive directional delta function in the direction a ray is already traveling. This is useful, for example, to model a layer of dust on top of a surface. To do so, we might specify the following LSQT file

```
Medium
Layer z=1 Null
Medium mus=1 HenyeyGreenstein g=0.2
Layer z=0 Lambertian fR=1
Medium
```

which describes a layer of forward-scattering "dust" on top of a 100% reflective Lambertian surface.

3.2.2. Lambertian BSDF

The Lambertian BSDF models a uniformly scattering surface which may reflect, transmit, or both. That is, the (cosine-weighted) scattering function is

$$f(\omega_o \to \omega_i) = \frac{|\cos \theta_i|}{\pi} \begin{cases} f_R & (\omega_o, \omega_i) \text{ in same hemisphere} \\ f_T & \text{otherwise} \end{cases}$$

where f_R and f_T specify the amount of incident energy reflected and transmitted respectively. Both f_R and f_T should be non-negative numbers such that $f_R + f_T \leq 1$ for this to be physically plausible. In the event that $f_R + f_T = 1$, this is perfectly energy-conserving.

To specify the Lambertian BSDF in LSQT format, use the name Lambertian followed by (optional) keyword arguments fR and fT for f_R and f_T respectively. By default, fR=1 and fT=0. For example,

specifies a Lambertian surface at *z*-height 2.2 which is 80% reflective, 10% transmissive, and 10% absorptive.

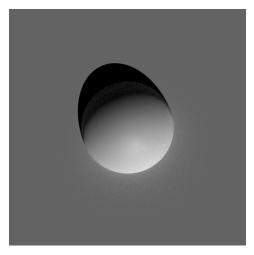
3.2.3. Microsurface Lambertian BSDF

A microsurface is thought to be an infinitesimally thin cloud of microfacets, where each facet is thought to scatter light according to another, simpler BSDF. The cloud is characterized geometrically by 1) a distribution of the slopes of the facets and 2) a distribution of the heights of the facets, which are assumed to uncorrelated, such that the facets are discontinuous. The distribution of slopes is parameterized by its so-called roughness α (more generally, anisotropic roughness α_x , α_y). As $\alpha \to \infty$, the distribution of slopes widens and the emergent BSDF appears rougher. As $\alpha \to 0$, the distribution of slopes collapses, recovering the initial, simpler BSDF in the limiting case $\alpha = 0$.

As given in [3], the microsurface Lambertian BSDF is a multiple scattering (stochastic) microfacet model wherein the facets are assumed to be Lambertian scatterers. As one might expect, this is particularly useful for modeling rough diffuse surfaces. The implementation in layered-sqt is parameterized by roughness $\alpha > 0$ and Lambertian BSDF coefficients $f_R \in [0,1]$ and $f_T \in [0,1]$, equivalent to those in section §3.2.2. For reference, $\alpha < 0.1$ is very smooth and $\alpha > 0.8$ is very rough.

To specify the microsurface Lambertian BSDF in LSQT format, use the name MicrosurfaceLambertian followed by keyword arguments alpha for α , fR for f_R , and fT for f_T . By default, alpha=0.5, fR=1, and fT=0. Furthermore, multiple scattering interactions may be disabled by the keyword argument

in which case the implementation computes single scattering interactions with the microsurface only—this may be desirable to speed up calculations. It is however important to note that $f_R + f_T = 1$ does *not* properly guarantee that the



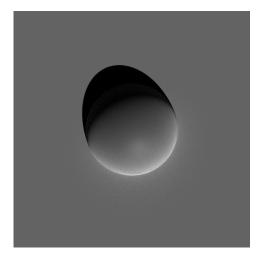


Figure 8: A sphere with an LSQT BRDF on a Lambertian ground plane rendered from above with DIRSIG5. The LSQT BRDF models a dusty dielectric substrate—from top to bottom, we have a vacuum medium, a null BSDF layer, a slightly back-scattering medium with g = -0.2, a decently smooth microsurface dielectric BSDF layer, a transparent medium with $\eta = 1.5$, and a Lambertian BRDF layer. On the left, the Lambertian layer has $f_R = 0.8$. On the right, the Lambertian layer has $f_R = 0.4$, such that the backscattering due to "dust" is more obvious. All other parameters are the same in both images.

BSDF is non-absorbing if multiple-scattering interactions are disabled. For small roughness values (say, α < 0.2), the energy carried by multiple scattering iterations is insignificant, so this may not be a concern. For greater roughness values however, multiple scattering is important to prevent energy loss.

3.2.4. Microsurface dielectric BSDF

The microsurface dielectric BSDF follows the same logical construction as the microsurface Lambertian BSDF, except the constituent BSDF assigned to the facets is the delta Fresnel mirror BSDF. The implementation in layered-sqt is parameterized by roughness $\alpha>0$, a scaling factor for the Fresnel mirror BRDF $k_R\in[0,1]$, and a scaling factor for the Fresnel mirror BTDF $k_T\in[0,1]$.

To specify the microsurface dielectric BSDF in LSQT format, use the name MicrosurfaceDielectric followed by keyword arguments alpha, kR, and kT for α , k_R , and k_T respectively. By default, alpha=0.5, kR=1, and kT=1.

Unlike the Lambertian case, the single-scattering term is directly computable due to the delta function in the dielectric Fresnel BSDF. This leads to noticeably faster simulations, so much so that multiple-scattering is disabled by default. Energy loss is less noticeable than the Lambertian case for moderately rough surfaces, but is still an issue for very rough surfaces (say $\alpha > 0.6$). As such, multiple-scattering may be enabled at the user's discretion by the keyword argument

to ensure proper energy conservation.

3.2.5. Microsurface conductive BRDF TODO

3.2.6. Oren-Nayar diffuse BRDF

As given in [5], the Oren-Nayar BRDF is a closed-form generalization of the Lambertian BRDF based on principles of microsurface theory. The original paper deals with quite a bit of geometry and includes defined angles of interest in the resulting expression. Some manipulation leads to a simpler expression,

$$f_s(\omega_o \to \omega_i) = f_R \frac{\left|\cos \theta_i\right|}{\pi} \left[A + B \frac{\max\left(\omega_o' \cdot \omega_i', 0\right)}{\max\left(\left|\cos \theta_o\right|, \left|\cos \theta_i\right|\right)} \right]$$

where the primed vectors ω'_o , ω'_i denote the tangential components of ω_o , ω_i with respect to the normal, such that

$$\begin{aligned} \omega_o' \cdot \omega_i' &= \omega_{o_x} \omega_{i_x} + \omega_{o_y} \omega_{i_y} \\ &= \sin(\theta_o) \sin(\theta_i) \cos(\phi_o - \phi_i) \end{aligned}$$

and constants A and B are given by

$$A = 1 - 0.5 \frac{\sigma^2}{\sigma^2 + 0.33}, \qquad B = 0.45 \frac{\sigma^2}{\sigma^2 + 0.09},$$

where $\sigma \in [0,0.8]$ denotes the standard deviation of the microfacet angle in radians, which is a parameter alongside the reflection coefficient $f_R \in [0,1]$. In the limit $\sigma = 0$, f_s is identically the Lambertian BRDF. As σ increases, f_s appears rougher. We place an arbitrary upper bound of 0.8 rad on σ , as this is sufficiently "rough", and the accuracy of the model diminishes as $\sigma \to \infty$.

The Oren-Nayar BRDF is essentially an approximation of the rigorous stochastic microsurface Lambertian BRDF given 20 years later in [3]. This model is still convenient to have around however, since it is characteristically similar for reasonably small roughnesses, and is evaluable in closed form.

To specify the Oren-Nayar BRDF in LSQT format, use the name OrenNayarDiffuse followed by (optional) keyword arguments fR for f_R and sigma for σ . By default, fR=1 and sigma=0.1.

References

- [1] E. d'Eon. A Hitchhiker's Guide to Multiple Scattering. An Incomplete Collection of Mostly Steady-State

 Monoenergetic Neutral-Particle Problem Solutions. 2016.

 URL: http://www.eugenedeon.com/project/a-hitchhikers-guide-to-multiple-scattering/.
- [2] A. A. Goodenough and S. D. Brown. "DIRSIG5: Next-Generation Remote Sensing Data and Image Simulation Framework". In: *IEEE Journal of Selected Topics in Applied Earth Observations and Remote Sensing* 10.11 (Nov. 2017), pp. 4818–4833. ISSN: 2151-1535. DOI: 10.1109/JSTARS.2017.2758964.
- [3] E. Heitz et al. "Multiple-scattering microfacet BSDFs with the Smith model". In: *ACM Transactions on Graphics* 35.4 (July 2016), pp. 1–58. ISSN: 0730-0301. DOI: 10.1145/2897824.2925943. URL: https://doi.org/10.1145/2897824.2925943.
- [4] E. Heitz et al. "The SGGX Microflake Distribution". In: ACM Trans. Graph. 34.4 (July 2015). ISSN: 0730-0301. DOI: 10.1145/2766988. URL: https://doi.org/10.1145/2766988.
- [5] M. Oren and S. K. Nayar. "Generalization of Lambert's reflectance model". In: Proceedings of the 21st Annual Conference on Computer Graphics and Interactive Techniques. (New York, NY, USA). SIGGRAPH'94. ACM, 1994, pp. 239–246. ISBN: 0-89791-667-0. DOI: 10.1145/192161.192213.