



Higher-Order DGFEM Transport Calculations on Polytope Meshes for Massively-Parallel Architectures

Michael W. Hackemack

Chair: Jean C. Ragusa

Committee Members: Marvin L. Adams, Jim E. Morel, Nancy M. Amato

External Advisor: Troy Becker

Department of Nuclear Engineering
Texas A&M University
College Station, TX, 77843, USA
mike_hack@tam.u.edu

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The Continuous-Energy Transport Equation

$$[\boldsymbol{\Omega} \cdot \nabla + \sigma_t(\mathbf{r}, E)] \psi(\mathbf{r}, E, \boldsymbol{\Omega}) = \int_{4\pi} \int_0^\infty \sigma_s(\mathbf{r}, E', E, \boldsymbol{\Omega}', \boldsymbol{\Omega}) \psi(\mathbf{r}, E', \boldsymbol{\Omega}') dE' d\boldsymbol{\Omega}' + Q(\mathbf{r}, E, \boldsymbol{\Omega})$$

\mathbf{r} - neutron position (cm)

E - neutron energy (eV)

$\boldsymbol{\Omega}$ - neutron solid angle (steradians)

$\psi(\mathbf{r}, E, \boldsymbol{\Omega})$ - angular flux

$Q(\mathbf{r}, E, \boldsymbol{\Omega})$ - distributed neutron source

$\sigma_t(\mathbf{r}, E)$ - total macroscopic cross section (1/cm)

$\sigma_s(\mathbf{r}, E', E, \boldsymbol{\Omega}', \boldsymbol{\Omega})$ - total

Polytope Grid Motivation

- Other physics communities are now employing polytope grids due to decreased cell/face counts (CFD in particular)
- They allow for transition elements between different domain regions
- Hanging nodes from non-conforming meshes are not necessary
- Independently-generated simplicial grids (*i.e.* created in parallel) can be stitched together with polytopes without communicating the whole mesh across processors

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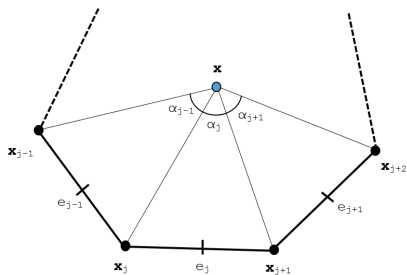
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Linear Basis Functions on 2D Polygons



Piecewise Linear Functions

Mean Value Coordinates

Maximum Entropy Coordinates

Quadratic Serendipity Basis Functions on 2D Polygons

Linear Basis Functions on 3D Polyhedra

The Diffusion Equation and Boundary Conditions

The Diffusion Equation

$$-\nabla \cdot D \nabla \Phi(\mathbf{r}) + \sigma \Phi(\mathbf{r}) = q(\mathbf{r}), \quad \mathbf{r} \in \mathcal{D}$$

Boundary Conditions

$$\begin{aligned} \Phi(\mathbf{r}) &= \Phi_0(\mathbf{r}), & \mathbf{r} \in \partial \mathcal{D}^d \\ -D \partial_n \Phi(\mathbf{r}) &= J_0(\mathbf{r}), & \mathbf{r} \in \partial \mathcal{D}^n \\ \frac{1}{4} \Phi(\mathbf{r}) + \frac{1}{2} D \partial_n \Phi(\mathbf{r}) &= J^{inc}(\mathbf{r}), & \mathbf{r} \in \partial \mathcal{D}^r \end{aligned}$$

Symmetric Interior Penalty (SIP) Form

Bilinear Form

$$\begin{aligned}
 a(\Phi, b) = & \left\langle D\vec{\nabla}\Phi, \vec{\nabla}b \right\rangle_{\mathcal{D}} + \left\langle \sigma\Phi, b \right\rangle_{\mathcal{D}} \\
 & + \left\{ \kappa_e^{SIP} \llbracket \Phi \rrbracket, \llbracket b \rrbracket \right\}_{E_h^i} - \left\{ \llbracket \Phi \rrbracket, \{ \{ D\partial_n b \} \} \right\}_{E_h^i} - \left\{ \{ \{ D\partial_n \Phi \} \}, \llbracket b \rrbracket \right\}_{E_h^i} \\
 & + \left\{ \kappa_e^{SIP} \Phi, b \right\}_{\partial\mathcal{D}^d} - \left\{ \Phi, D\partial_n b \right\}_{\partial\mathcal{D}^d} - \left\{ D\partial_n \Phi, b \right\}_{\partial\mathcal{D}^d} + \frac{1}{2} \left\{ \Phi, b \right\}_{\partial\mathcal{D}^r}
 \end{aligned}$$

Linear Form

$$\ell(b) = \left\langle q, b \right\rangle_{\mathcal{D}} - \left\{ J_0, b \right\}_{\partial\mathcal{D}^n} + 2 \left\{ J_{inc}, b \right\}_{\partial\mathcal{D}^r}$$

