Higher-Order DGFEM Transport Calculations on Polytope Meshes for Massively-Parallel Architectures

Michael W. Hackemack

Chair: Jean C. Ragusa

Committee Members: Marvin L. Adams, Jim E. Morel, Nancy M. Amato

External Advisor: Troy Becker

Department of Nuclear Engineering Texas A&M University College Station, TX, USA 77843 mike_hack@tamu.edu



Outline

- Overview
 - The DGFEM S_N Transport Equation
 - Polytope Grid Motivation
- Polytope Finite Element Basis Functions
 - Linear Basis Functions on 2D Polygons
 - Quadratic Serendipity Basis Functions on 2D Polygons
 - Linear Basis Functions on 3D Polyhedra
- 3 Diffusion Synthetic Acceleration on Polytopes
 - Theory
 - MIP Diffusion Form
- Proposed Work and Current Status
- Ongoing Work

The Continuous-Energy Transport Equation

Transport Equation

$$\left[\mathbf{\Omega}\cdot
abla+\sigma_t(\mathbf{r},E)
ight]\psi(\mathbf{r},E,\Omega)=\int_{4\pi}\int_0^\infty\,\sigma_s(\mathbf{r},E',E,\Omega',\Omega)\psi(\mathbf{r},E',\Omega')dE'd\Omega'+Q(\mathbf{r},E,\Omega)$$

Boundary Conditions

$$\psi(\mathbf{r}, E, \Omega) = \psi^{inc}(\mathbf{r}, E, \Omega) + \int_{4\pi} \int_0^\infty \beta(\mathbf{r}, E', E, \Omega', \Omega) \psi(\mathbf{r}, E', \Omega') dE' d\Omega'$$

Term Definitions

r - neutron position

E - neutron energy

 Ω - neutron solid angle

 $\psi(\mathbf{r}, E, \mathbf{\Omega})$ - angular flux

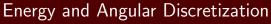
 $Q(\mathbf{r}, E, \Omega)$ - distributed neutron source

 $\sigma_t(\mathbf{r}, E)$ - total macroscopic cross section

 $\sigma_s(\mathbf{r}, E', E, \Omega', \Omega)$ - total macroscopic scattering cross section

 $\beta(\mathbf{r}, E', E, \Omega', \Omega)$ - boundary albedo

3 / 35



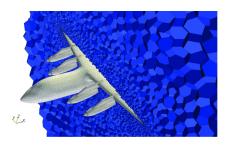


Overview

Spatial Discretization

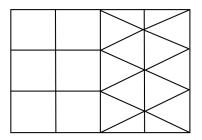


- Other physics communities are now employing polytope grids due to decreased cell/face counts (CFD in particular)



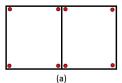


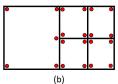
- Other physics communities are now employing polytope grids due to decreased cell/face counts (CFD in particular)
- They allow for transition elements between different domain regions





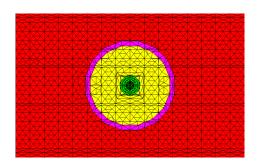
- Other physics communities are now employing polytope grids due to decreased cell/face counts (CFD in particular)
- They allow for transition elements between different domain regions
- Hanging nodes from non-conforming meshes are not necessary



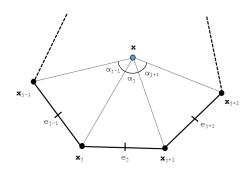




- Other physics communities are now employing polytope grids due to decreased cell/face counts (CFD in particular)
- They allow for transition elements between different domain regions
- Hanging nodes from non-conforming meshes are not necessary
- Independently-generated simplicial grids (i.e. created in parallel) can be stitched together with polytopes without communicating the whole mesh across processors







2D Linear Basis Function Properties - Barycentric Coordinates

- $\mathbf{0} \ \lambda_i > 0$
- $\sum_{i} \lambda_{i} = 1$
- $\lambda_i(\mathbf{x}_i) = \delta_{ii}$



Wachspress Rational Functions



Piecewise Linear Functions



Mean Value Coordinates





Quadratic Serendipity Basis Functions on 2D Polygons





POLYFEM 00000000



The Diffusion Equation and Boundary Conditions

The Diffusion Equation

$$-\nabla \cdot D \nabla \Phi(\mathbf{r}) + \sigma \Phi(\mathbf{r}) = q(\mathbf{r}), \qquad \mathbf{r} \in \mathcal{D}$$

Boundary Conditions

$$\begin{aligned} \Phi(\mathbf{r}) &= \Phi_0(\mathbf{r}), & \mathbf{r} \in \partial \mathcal{D}^d \\ -D\partial_n \Phi(\mathbf{r}) &= J_0(\mathbf{r}), & \mathbf{r} \in \partial \mathcal{D}^n \\ \frac{1}{4} \Phi(\mathbf{r}) + \frac{1}{2} D\partial_n \Phi(\mathbf{r}) &= J^{inc}(\mathbf{r}), & \mathbf{r} \in \partial \mathcal{D}^r \end{aligned}$$



Symmetric Interior Penalty (SIP) Form

DSA on Polytopes 000000

Bilinear Form

$$\begin{split} \textbf{a}(\Phi,b) &= \left\langle D\nabla\Phi,\nabla b\right\rangle_{\mathcal{D}} + \left\langle \sigma\Phi,b\right\rangle_{\mathcal{D}} \\ &+ \left\{\kappa_e^{\textit{SIP}}\llbracket\Phi\rrbracket,\llbracket b\rrbracket\right\}_{E_h^i} - \left\{\llbracket\Phi\rrbracket,\{\{D\partial_nb\}\}\right\}_{E_h^i} - \left\{\{\{D\partial_n\Phi\}\},\llbracket b\rrbracket\right\}_{E_h^i} \\ &+ \left\{\kappa_e^{\textit{SIP}}\Phi,b\right\}_{\partial\mathcal{D}^d} - \left\{\Phi,D\partial_nb\right\}_{\partial\mathcal{D}^d} - \left\{D\partial_n\Phi,b\right\}_{\partial\mathcal{D}^d} + \frac{1}{2}\Big\{\Phi,b\Big\}_{\partial\mathcal{D}^d} \end{split}$$

Linear Form

$$\begin{split} \ell(\textit{b}) &= \left\langle \textit{q}, \textit{b} \right\rangle_{\mathcal{D}} - \left\{ \textit{J}_{0}, \textit{b} \right\}_{\partial \mathcal{D}^{n}} + 2 \Big\{ \textit{J}_{\textit{inc}}, \textit{b} \Big\}_{\partial \mathcal{D}^{l}} \\ &+ \Big\{ \kappa_{e}^{\textit{SIP}} \Phi_{0}, \textit{b} \Big\}_{\partial \mathcal{D}^{d}} - \Big\{ \Phi_{0}, \textit{D} \partial_{n} \textit{b} \Big\}_{\partial \mathcal{D}^{d}} \end{split}$$



SIP Penalty Coefficient

$$\kappa_e^{SIP} \equiv egin{cases} rac{C_B}{2} \left(rac{D^+}{h^+} + rac{D^-}{h^-}
ight) &, e \in E_h^i \ C_B rac{D^-}{h^-} &, e \in \partial \mathcal{D} \end{cases}$$

$$C_B = cp(p+1)$$

c - user defined constant $(c \geq 1)$

p - polynomial order of the finite element basis (1,2,3,...)

 $D^{(+/-)}$ - diffusion coefficient defined on the positive/negative side of a face $h^{(+/-)}$ - orthogonal projection defined on the positive/negative side of a face

$$u^{\pm} = \lim_{s \to 0^{\pm}} u(\mathbf{r} + s\mathbf{n})$$



Modified Interior Penalty (MIP) Form

DSA on Polytopes

00000

Diffusion Form

$$\begin{split} \left\langle D\nabla\Phi,\nabla b\right\rangle_{\mathcal{D}} + \left\langle \sigma\Phi,b\right\rangle_{\mathcal{D}} \\ + \left\{\kappa_{e}^{MIP}\llbracket\Phi\rrbracket,\llbracket b\rrbracket\right\}_{E_{h}^{i}} - \left\{\llbracket\Phi\rrbracket,\{\{D\partial_{n}b\}\}\right\}_{E_{h}^{i}} - \left\{\{\{D\partial_{n}\Phi\}\},\llbracket b\rrbracket\right\}_{E_{h}^{i}} \\ + \left\{\kappa_{e}^{MIP}\Phi,b\right\}_{\partial\mathcal{D}^{vac}} - \frac{1}{2}\Big\{\Phi,D\partial_{n}b\Big\}_{\partial\mathcal{D}^{vac}} - \frac{1}{2}\Big\{D\partial_{n}\Phi,b\Big\}_{\partial\mathcal{D}^{vac}} \\ = \left\langle q,b\right\rangle_{\mathcal{D}} \end{split}$$

MIP Penalty Term

$$\kappa_{\rm e}^{\it MIP} = \max(\frac{1}{4},\kappa_{\rm e}^{\it SIP})$$

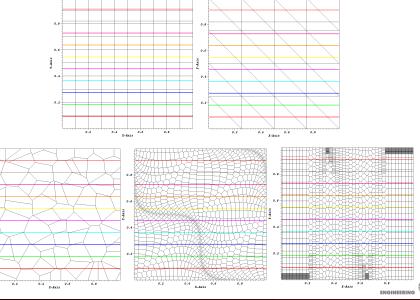








SIP exactly linear solutions on 3D polyhedral meshes using the PWL basis functions





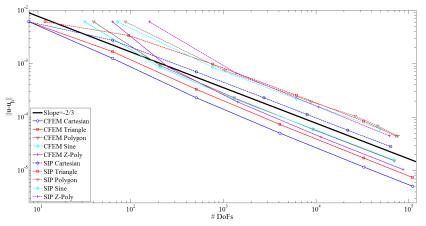
0.8

SIP convergence study - quadratic solution on 3D cube using the PWL basis functions

000000000000000

$$\Phi(x, y, z) = xyz(L_x - x)(L_y - y)(L_z - z)$$

$$L_x = L_x = L_x = 1.0$$





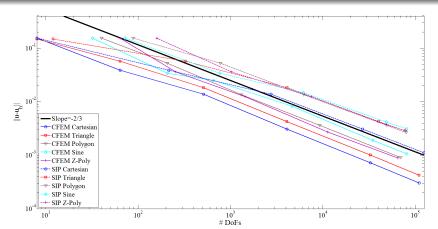


SIP convergence study - gaussian solution on 3D cube using the PWL basis functions

$$\Phi(x, y, z) = xyz(L_x - x)(L_y - y)(L_z - z) \exp(-(\mathbf{r} - \mathbf{r}_0) \cdot (\mathbf{r} - \mathbf{r}_0))$$

$$L_x = L_x = L_x = 1.0, \qquad \mathbf{r}_0 = (3/4, 3/4, 3/4)$$

Proposed Work and Current Status 000000000000000

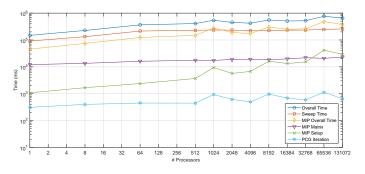




Proposed Work and Current Status 000000000000000

Proposed Work and Current Status 00000000000000

MIP DSA Timing Data with PDT on Vulcan using HYPRE



Problem Description

- Modified Zerr problem used optimal sweep aggregation parameters
 - homogeneous cube c=0.9999
 - 58 level-symmetric quadrature
- pointwise convergence tolerance of 1e-8
- precondition with MIP DSA using HYPRE PCG and AMG



Questions?

A special acknowledgment to the Department of Energy Rickover Fellowship Program in Nuclear Engineering, which provides strong support to its fellows and their professional development.





A stretch goal is to compare my method to Monte Carlo

I claim the following is the best way to show our method has practical importance, because continuous-energy Monte Carlo codes do exact particle tracking / kinematics and use very accurate cross sections. Such codes may attain higher fidelity in all respects than DRAGON.

Start with a 0-D problem to isolate energy discretization effects

- Come up with a reactor-themed problem
- Solve the same problem in PDT and MCNP or OpenMC
- **3** Choose QOI, such as k-eigenvalue, radial power profile, absorption/fission rates per nuclide, etc.
- Quantify how errors in PDT's QOI change as energy resolution is increased



November 24, 2015

A stretch goal is to compare my method to Monte Carlo

I claim the following is the best way to show our method has practical importance, because continuous-energy Monte Carlo codes do exact particle tracking / kinematics and use very accurate cross sections. Such codes may attain higher fidelity in all respects than DRAGON.

Start with a 0-D problem to isolate energy discretization effects

- Come up with a reactor-themed problem
- Solve the same problem in PDT and MCNP or OpenMC
- **3** Choose QOI, such as k-eigenvalue, radial power profile, absorption/fission rates per nuclide, etc.
- Quantify how errors in PDT's QOI change as energy resolution is increased

Build up problem complexity slowly: cylindricized pin cell with white boundary conditions, infinite lattice of pin cells, heterogeneous lattice of pin cells, etc.

- Quantify how errors in PDT's QOI change as spatial / angular / scattering moment resolution is increased
- Quantify how errors in PDT's QOI change as energy resolution is increased
- **3** ...
- Profit

