# Higher-Order DGFEM Transport Calculations on Polytope Meshes for Massively-Parallel Architectures

# Michael W. Hackemack

Chair: Jean C. Ragusa

Committee Members: Marvin L. Adams, Jim E. Morel, Nancy M. Amato

External Advisor: Troy Becker

Department of Nuclear Engineering Texas A&M University College Station, TX, USA 77843 mike\_hack@tamu.edu



# Outline

- Overview
  - The DGFEM  $S_N$  Transport Equation
  - Polytope Grid Motivation
- Polytope Finite Element Basis Functions
  - Linear Basis Functions on 2D Polygons
  - Quadratic Serendipity Basis Functions on 2D Polygons
  - Linear Basis Functions on 3D Polyhedra
- 3 Diffusion Synthetic Acceleration on Polytopes
  - Theory
  - MIP Diffusion Form
- Proposed Work and Current Status
- Ongoing Work

# The Continuous-Energy Transport Equation

# Transport Equation

$$\left[\mathbf{\Omega}\cdot
abla+\sigma_t(\mathbf{r},E)
ight]\psi(\mathbf{r},E,\Omega)=\int_{A\pi}\int_0^\infty\,\sigma_s(\mathbf{r},E',E,\Omega',\Omega)\psi(\mathbf{r},E',\Omega')dE'd\Omega'+Q(\mathbf{r},E,\Omega)$$

#### **Boundary Conditions**

$$\psi(\mathbf{r}, E, \Omega) = \psi^{inc}(\mathbf{r}, E, \Omega) + \int_{4\pi} \int_0^\infty \beta(\mathbf{r}, E', E, \Omega', \Omega) \psi(\mathbf{r}, E', \Omega') dE' d\Omega'$$

#### Term Definitions

r - neutron position

E - neutron energy

 $\Omega$  - neutron solid angle

 $\psi(\mathbf{r}, E, \mathbf{\Omega})$  - angular flux

 $Q(\mathbf{r}, E, \Omega)$  - distributed neutron source

 $\sigma_t(\mathbf{r}, E)$  - total macroscopic cross section

 $\sigma_s(\mathbf{r}, E', E, \Omega', \Omega)$  - total macroscopic scattering cross section

 $\beta(\mathbf{r}, E', E, \Omega', \Omega)$  - boundary albedo

# The multigroup $S_N$ equations

Overview OOOO

$$\left(\mathbf{\Omega}_{\textit{m}}\cdot
abla+\sigma_{t,\textit{g}}
ight)\psi_{\textit{m},\textit{g}} = \sum_{\textit{g}'=1}^{\textit{G}}\sum_{p=0}^{N_{\textit{p}}}rac{2p+1}{4\pi}\sigma_{\textit{s},\textit{p}}^{\textit{g}'
ightarrow\textit{g}}\sum_{n=-p}^{p}\phi_{\textit{p},n,\textit{g}'}Y_{\textit{p},\textit{n}}(\mathbf{\Omega}_{\textit{m}})+Q_{\textit{m},\textit{g}}$$

# Multigroup Method

$$\psi_{\mathsf{g}} = \int_{\mathsf{F}}^{\mathsf{E}_{\mathsf{g}-1}} \psi(\mathsf{E}) \, \mathsf{d}\mathsf{E}$$

#### $S_N$ Discretization

$$egin{align} \phi_{p,n} &\equiv \int_{4\pi} d\Omega \, \psi(\mathbf{\Omega}) \, Y_{p,n}(\mathbf{\Omega}), \ \sigma_{\mathsf{s},p} &\equiv \int_{-1}^{1} \, d\mu \, \sigma_{\mathsf{s}}(\mu_0) P_p(\mu_0) \ \end{array}$$

$$\mu_{\rm 0} \equiv {\bf \Omega}' \cdot {\bf \Omega}$$

$$\sigma_s(\mathbf{\Omega}'\cdot\mathbf{\Omega})\equiv rac{1}{2\pi}\sigma_s(\mu)$$

$$P_{
ho}(\mathbf{\Omega}'\cdot\mathbf{\Omega})\equivrac{1}{2\pi}P_{
ho}(\mu)$$

ENGINEER

# Spatial Discretization



# Iterative Procedure

#### Classic Source Iteration

$$egin{aligned} \Psi^{(\ell+1)} &= \mathbf{L}^{-1} \left( \mathsf{M} \mathbf{\Sigma} \Phi^{(\ell)} + \mathbf{Q} 
ight) \ \Phi^{(\ell+1)} &= \mathbf{D} \mathbf{L}^{-1} \left( \mathsf{M} \mathbf{\Sigma} \Phi^{(\ell)} + \mathbf{Q} 
ight) \ \Phi &= \mathbf{D} \Psi \end{aligned}$$

#### Operator Terms

L - streaming + collision operator

M - moment-to-discrete operator

D - discrete-to-moment operator

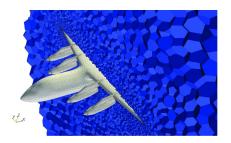
Σ - scattering operator

#### Transport Sweep

The operation  $\mathbf{L}^{-1}$  can be performed in myriad ways. For this work, we will use the matrix-free, full-domain transport sweep.

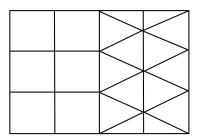


- Other physics communities are now employing polytope grids due to decreased cell/face counts (CFD in particular)



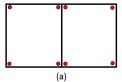


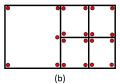
- Other physics communities are now employing polytope grids due to decreased cell/face counts (CFD in particular)
- They allow for transition elements between different domain regions
- Hanging nodes from non-conforming meshes are not necessary
- Independently-generated simplicial grids (i.e. created in parallel) can be stitched togethe
  with polytopes without communicating the whole mesh across processors





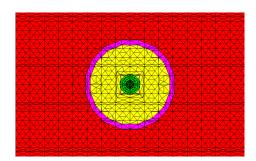
- Other physics communities are now employing polytope grids due to decreased cell/face counts (CFD in particular)
- They allow for transition elements between different domain regions
- Hanging nodes from non-conforming meshes are not necessary







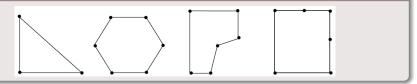
- Other physics communities are now employing polytope grids due to decreased cell/face counts (CFD in particular)
- They allow for transition elements between different domain regions
- Hanging nodes from non-conforming meshes are not necessary
- Independently-generated simplicial grids (i.e. created in parallel) can be stitched together with polytopes without communicating the whole mesh across processors



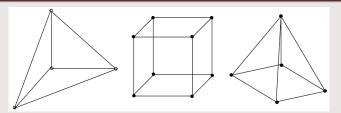


## Polytope Finite Elements

# 2D arbitrary convex/concave polygons

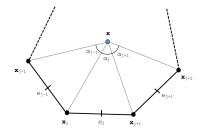


#### 3D convex polyhedra





# Linear Basis Functions on 2D Polygons

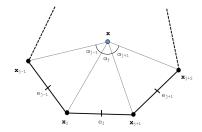


# Basis Function Properties - Barycentric Coordinates

- $\mathbf{0} \ \lambda_i \geq \mathbf{0}$
- $\sum_{i} \lambda_{i} = 1$
- $\lambda_i(\mathbf{x}_i) = \delta_{ii}$



# Wachspress Rational Functions

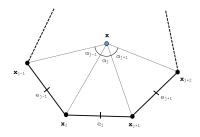


$$\lambda_i^W(\mathbf{x}) = \frac{w_i(\mathbf{x})}{\sum_i w_i(\mathbf{x})}, \qquad w_j(\mathbf{x}) = \frac{A(\mathbf{x}_{j-1}, \mathbf{x}_j, \mathbf{x}_{j+1})}{A(\mathbf{x}, \mathbf{x}_{i-1}, \mathbf{x}_j) A(\mathbf{x}, \mathbf{x}_i, \mathbf{x}_{j+1})}$$

$$A(\mathbf{x}_1, \mathbf{x}_2, \mathbf{x}_3) = \frac{1}{2} \begin{vmatrix} 1 & 1 & 1 \\ x_1 & x_2 & x_3 \\ y_1 & y_2 & y_3 \end{vmatrix}$$



# Piecewise Linear (PWL) Functions



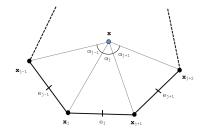
$$\lambda_i^{PWL}(\mathbf{x}) = t_i(\mathbf{x}) + \alpha_i t_c(\mathbf{x})$$

 $t_i$  - standard 2D linear function; 1 at vertex i that linearly decreases to 0 to the cell center and the adjoining vertices

 $t_c$  - 2D tent function; 1 at cell center and linearly decreases to 0 to each cell vertex  $\alpha_i = \frac{1}{N_{ij}}$  - weight parameter for vertex i



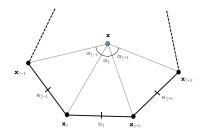
#### Mean Value Coordinates



$$\lambda_i^{MV}(\mathbf{x}) = \frac{w_i(\mathbf{x})}{\sum_i w_i(\mathbf{x})}, \qquad w_j(\mathbf{x}) = \frac{\tan(\alpha_{j-1}/2) + \tan(\alpha_j/2)}{|\mathbf{x}_i - \mathbf{x}|}$$



#### Maximum Entropy Coordinates



$$\lambda_i^{ME}(\mathbf{x}) = \frac{w_i(\mathbf{x})}{\sum_i w_i(\mathbf{x})}, \qquad w_j(\mathbf{x}) = m_j(\mathbf{x}) \exp(-\omega^* \cdot (\mathbf{x}_j - \mathbf{x}))$$

$$\omega^* = \operatorname{argmin} F(\omega, \mathbf{x}) \qquad F(\omega, \mathbf{x}) = \operatorname{In} \left( \sum_j w_j(\mathbf{x}) \right)$$



# Summary of the 2D Linear Basis Functions

POLYFEM

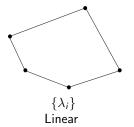
Basis Function	Dimension	Polytope Types	Analytical/Numerical	Direct/Iterative
Wachspress	2D/3D	Convex*	Numerical	Direct
PWL	1D/2D/3D	Convex/Concave	Analytical	Direct
Mean Value	2D**	Convex/Concave	Numerical	Direct
Max Entropy	1D/2D/3D	Convex/Concave	Numerical	Iterative***

- \* weak convexity for Wachspress coordinates does not cause blow up
- \*\* mean value 3D analogue only applicable to tetrahedron
- \*\*\* maximum entropy minimization solved via Newton's Method



# Quadratic Serendipity Basis Functions on 2D Polygons

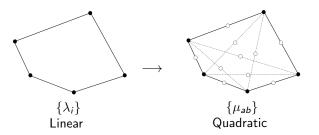
- lacktriangle Form the linear barycentrric functions  $\{\lambda_i\}$
- 2 Construct the pairwise products  $\{\mu_{ab}\}$
- **③** Eliminate the interior nodes to form a serendipity basis  $\{\xi_{ij}\}$





# Quadratic Serendipity Basis Functions on 2D Polygons

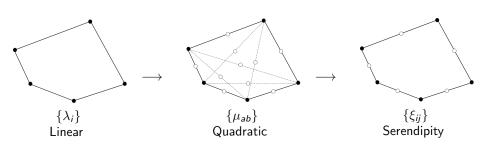
- Form the linear barycentrric functions  $\{\lambda_i\}$
- 2 Construct the pairwise products  $\{\mu_{ab}\}$





# Quadratic Serendipity Basis Functions on 2D Polygons

- Form the linear barycentrric functions  $\{\lambda_i\}$
- **②** Construct the pairwise products  $\{\mu_{ab}\}$
- **3** Eliminate the interior nodes to form a serendipity basis  $\{\xi_{ij}\}$



#### Pairwise products of the barycentric basis functions

# **Necessary Precision Properties**

$$\begin{split} \sum_{\textit{aa} \in V} \mu_{\textit{aa}} + \sum_{\textit{ab} \in E \cup D} 2\mu_{\textit{ab}} &= 1 \\ \sum_{\textit{aa} \in V} \mathbf{x}_{\textit{aa}} \mu_{\textit{aa}} + \sum_{\textit{ab} \in E \cup D} 2\mathbf{x}_{\textit{ab}} \mu_{\textit{ab}} &= \mathbf{x} \\ \sum_{\textit{aa} \in V} \mathbf{x}_{\textit{a}} \mathbf{x}_{\textit{a}}^{\mathsf{T}} \mu_{\textit{aa}} + \sum_{\textit{ab} \in E \cup D} \left( \mathbf{x}_{\textit{a}} \mathbf{x}_{\textit{b}}^{\mathsf{T}} + \mathbf{x}_{\textit{b}} \mathbf{x}_{\textit{a}}^{\mathsf{T}} \right) \mu_{\textit{ab}} &= \mathbf{x} \mathbf{x}^{\mathsf{T}} \end{split}$$

#### Further Notation/Notes

$$\mathbf{x}_{ab} = \frac{\mathbf{x}_a + \mathbf{x}_b}{2}, \qquad \mu_{ab} = \lambda_a \lambda_b$$

$$\mu_{ab}^{K}(\mathbf{r}) = 0, \qquad \{ab \in D, \, \mathbf{r} \in \partial K\}$$



#### Eliminate interior nodes to form serendipity basis

# Reduction Problem - $[\xi] := \mathbb{A}[\mu]$

$$\mathbb{A} = \left[ \begin{array}{cccc} c_{11}^{11} & \dots & c_{ab}^{11} & \dots & c_{(n-2)n}^{11} \\ \dots & \ddots & \vdots & \ddots & \vdots \\ c_{11}^{ij} & \dots & c_{ab}^{ij} & \dots & c_{(n-2)n}^{ij} \\ \dots & \ddots & \vdots & \ddots & \vdots \\ c_{11}^{n(n+1)} & \dots & c_{ab}^{n(n+1)} & \dots & c_{(n-2)n}^{n(n+1)} \end{array} \right]$$

#### Serendipity Precision Properties

$$\begin{split} \sum_{ii \in V} \xi_{ii} + \sum_{i(i+1) \in E} 2\xi_{i(i+1)} &= 1 \\ \sum_{ii \in V} \mathbf{x}_{ii} \xi_{ii} + \sum_{i(i+1) \in E} 2\mathbf{x}_{i(i+1)} \xi_{i(i+1)} &= \mathbf{x} \\ \sum_{ii \in V} \mathbf{x}_{i} \mathbf{x}_{i}^{T} \xi_{ii} + \sum_{i(i+1) \in E} \left( \mathbf{x}_{i} \mathbf{x}_{i+1}^{T} + \mathbf{x}_{i+1} \mathbf{x}_{i}^{T} \right) \xi_{i(i+1)} &= \mathbf{x} \mathbf{x}^{T} \end{split}$$

November 24, 2015

#### Special case - bilinear coordinates on the unit square

## Bilinear coordinates and quadratic extension

$$\begin{array}{lll} \lambda_1 = (1-x)(1-y) & \mu_{11} = (1-x)^2(1-y)^2 & \mu_{12} = (1-x)x(1-y)^2 \\ \lambda_2 = x(1-y) & \mu_{22} = x^2(1-y)^2 & \mu_{23} = x^2y(1-y) \\ \lambda_3 = xy & \mu_{33} = x^2y^2 & \mu_{34} = (1-x)xy^2 \\ \lambda_4 = (1-x)y & \mu_{44} = (1-x)^2y^2 & \mu_{41} = (1-x)^2y(1-y) \\ \mu_{13} = (1-x)x(1-y)y & \mu_{24} = (1-x)x(1-y)y \end{array}$$

#### Reduction matrix

#### Serendipity coordinates

$$\xi_{11} = (1 - x)(1 - y)(1 - x - y)$$

$$\xi_{22} = x(1 - y)(x - y)$$

$$\xi_{33} = xy(-1 + x + y)$$

$$\xi_{44} = (1 - x)y(y - x)$$

$$\xi_{12} = (1 - x)x(1 - y)$$

$$\xi_{23} = xy(1 - y)$$

$$\xi_{34} = (1 - x)xy$$

$$\xi_{41} = (1 - x)y(1 - y)$$



November 24, 2015

# Linear Basis Functions on 3D Polyhedra

# Linear basis functions and convex polyhedra only for 3D

- The 2D quadratic serendipity formulation is more arduous in 3D
- Intercell coupling is not straightforward for concave polyhedra
- Focus on 3D PWL functions MAXENT only other function for arbitrary polyhedra
- Focus on 3D parallelepipeds and extruded convex polygons (convex prisms)

#### 3D PWL basis functions

$$b_i(\mathbf{x}) = t_i(\mathbf{x}) + \sum_{f=1}^{F_i} \beta_f^i t_f(\mathbf{x}) + \alpha_i t_c(\mathbf{x})$$

 $t_i$  - standard 3D linear function; 1 at vertex i, linearly decreases to 0 to the cell center, each adjoining face center, and each adjoining vertex

 $t_c$  - 3D tent function; 1 at cell center, linearly decreases to 0

 $t_f$  - face tent function; 1 at face center, linearly decreases to 0 at each face vertex and cell center  $\alpha_i = \frac{1}{N_{t'}}$  - weight parameter for vertex i

 $\beta_f^i = \frac{N_V}{N_c}$  - weight parameter for face f touching vertex i



Theory goes here...



#### The diffusion equation is used as our low-order operator

000000

## The Diffusion Equation

Overview

$$-\nabla \cdot D\nabla \Phi(\mathbf{r}) + \sigma \Phi(\mathbf{r}) = q(\mathbf{r}), \qquad \mathbf{r} \in \mathcal{D}$$

#### General Boundary Conditions

$$\begin{split} \Phi(\mathbf{r}) &= \Phi_0(\mathbf{r}), \qquad \mathbf{r} \in \partial \mathcal{D}^d \\ - D \partial_n \Phi(\mathbf{r}) &= J_0(\mathbf{r}), \qquad \mathbf{r} \in \partial \mathcal{D}^n \\ \frac{1}{4} \Phi(\mathbf{r}) + \frac{1}{2} D \partial_n \Phi(\mathbf{r}) &= J^{inc}(\mathbf{r}), \qquad \mathbf{r} \in \partial \mathcal{D}^r \end{split}$$

#### Desirable diffusion form properties

- Can handle concave and degenerate polytope cells
- Symmetric Positive-Definite (SPD)
- Availability of suitable preconditioners
- Agnostic of directionality of interior faces

#### Symmetric Interior Penalty (SIP) Form

#### Bilinear Form

$$\begin{split} a(\Phi,b) &= \left\langle D \nabla \Phi, \nabla b \right\rangle_{\mathcal{D}} + \left\langle \sigma \Phi, b \right\rangle_{\mathcal{D}} \\ &+ \left\{ \kappa_e^{\textit{SIP}} \llbracket \Phi \rrbracket, \llbracket b \rrbracket \right\}_{E_h^i} - \left\{ \llbracket \Phi \rrbracket, \{\{D \partial_n b\}\} \right\}_{E_h^i} - \left\{ \{\{D \partial_n \Phi\}\}, \llbracket b \rrbracket \right\}_{E_h^i} \\ &+ \left\{ \kappa_e^{\textit{SIP}} \Phi, b \right\}_{\partial \mathcal{D}^d} - \left\{ \Phi, D \partial_n b \right\}_{\partial \mathcal{D}^d} - \left\{ D \partial_n \Phi, b \right\}_{\partial \mathcal{D}^d} + \frac{1}{2} \left\{ \Phi, b \right\}_{\partial \mathcal{D}^r} \end{split}$$

#### Linear Form

$$\begin{split} \ell(\textit{b}) &= \left\langle \textit{q}, \textit{b} \right\rangle_{\mathcal{D}} - \left\{ \textit{J}_{0}, \textit{b} \right\}_{\partial \mathcal{D}^{n}} + 2 \Big\{ \textit{J}_{\textit{inc}}, \textit{b} \Big\}_{\partial \mathcal{D}^{l}} \\ &+ \Big\{ \kappa_{e}^{\textit{SIP}} \Phi_{0}, \textit{b} \Big\}_{\partial \mathcal{D}^{d}} - \Big\{ \Phi_{0}, \textit{D} \partial_{n} \textit{b} \Big\}_{\partial \mathcal{D}^{d}} \end{split}$$



November 24, 2015

#### SIP Penalty Coefficient

Overview

$$\kappa_e^{SIP} \equiv \begin{cases} \frac{C_B}{2} \left( \frac{D^+}{h^+} + \frac{D^-}{h^-} \right) &, e \in E_h^i \\ C_B \frac{D^-}{h^-} &, e \in \partial \mathcal{D} \end{cases}$$

 $C_B = cp(p+1)$ 

$$c$$
 - user defined constant ( $c \geq 1$ )

p - polynomial order of the finite element basis (1, 2, 3, ...)

 $D^{(+/-)}$  - diffusion coefficient defined on the positive/negative side of a face  $h^{(+/-)}$  - orthogonal projection defined on the positive/negative side of a face

$$u^{\pm} = \lim_{s \to 0^{+}} u(\mathbf{r} + s\mathbf{n})$$



#### Modified Interior Penalty (MIP) Form

#### Diffusion Form

$$\begin{split} \left\langle D\nabla\Phi,\nabla b\right\rangle_{\mathcal{D}} + \left\langle \sigma\Phi,b\right\rangle_{\mathcal{D}} \\ + \left\{\kappa_{e}^{MIP}\llbracket\Phi\rrbracket,\llbracket b\rrbracket\right\}_{E_{h}^{i}} - \left\{\llbracket\Phi\rrbracket,\{\{D\partial_{n}b\}\}\right\}_{E_{h}^{i}} - \left\{\{\{D\partial_{n}\Phi\}\},\llbracket b\rrbracket\right\}_{E_{h}^{i}} \\ + \left\{\kappa_{e}^{MIP}\Phi,b\right\}_{\partial\mathcal{D}^{vac}} - \frac{1}{2}\left\{\Phi,D\partial_{n}b\right\}_{\partial\mathcal{D}^{vac}} - \frac{1}{2}\left\{D\partial_{n}\Phi,b\right\}_{\partial\mathcal{D}^{vac}} \\ = \left\langle q,b\right\rangle_{\mathcal{D}} \end{split}$$

#### MIP Penalty Term

$$\kappa_e^{MIP} = \max(rac{1}{4}, \kappa_e^{SIP})$$



#### Proposed Work

#### **POLYFEM**

- Analyze the different 2D linear polygonal basis functions for use in DGFEM transport calculations
- Perform the same analysis with the quadratic serendipity basis functions

#### MIP DSA

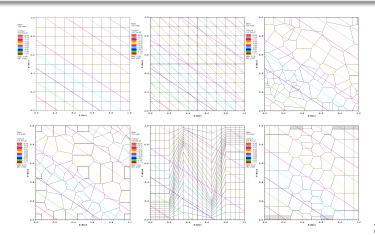
- Analyze the 2D linear/quadratic basis functions with DSA preconditioning through Fourier/numerical analysis
- Extend the analysis of MIP DSA to arbitrary convex 3D polyhedra
- Implement MIP DSA in PDT using HYPRE
  - Analyze the scalability of the method to high process counts
  - Perform parametric studies on aggregation/partitioning factors to generate a performance model of MIP DSA with HYPRE
  - Implement and perform analysis of two-grid acceleration at scale
  - Run real-world numerical experiments IM1



#### 2D Exactly-Linear Transport Solutions - mean value coordinates

$$\mu \frac{\partial \psi}{\partial x} + \eta \frac{\partial \psi}{\partial y} + \sigma_t \psi = Q(x, y, \mu, \eta)$$

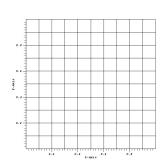
$$\psi(x, y, \mu, \eta) = ax + by + c\mu + d\eta + e, \qquad \phi(x, y) = 2\pi (ax + by + e)$$

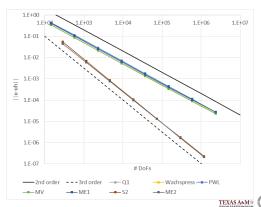




#### Convergence rates using MMS for the 2D polygonal basis functions

$$\psi(x,y) = \sin(\nu \frac{\pi x}{L_x}) \sin(\nu \frac{\pi y}{L_y})$$
$$\phi(x,y) = 2\pi \sin(\nu \frac{\pi x}{L_x}) \sin(\nu \frac{\pi y}{L_y})$$

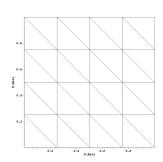


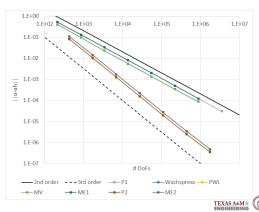




#### Convergence rates using MMS for the 2D polygonal basis functions

$$\psi(x,y) = \sin(\nu \frac{\pi x}{L_x}) \sin(\nu \frac{\pi y}{L_y})$$
$$\phi(x,y) = 2\pi \sin(\nu \frac{\pi x}{L_x}) \sin(\nu \frac{\pi y}{L_y})$$

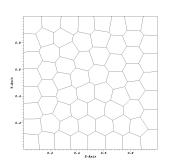


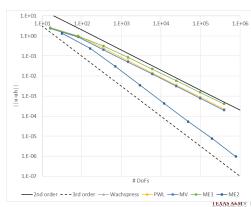




#### Convergence rates using MMS for the 2D polygonal basis functions

$$\psi(x,y) = \sin(\nu \frac{\pi x}{L_x}) \sin(\nu \frac{\pi y}{L_y})$$
$$\phi(x,y) = 2\pi \sin(\nu \frac{\pi x}{L_x}) \sin(\nu \frac{\pi y}{L_y})$$



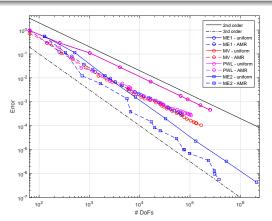


November 24, 2015

## Convergence rates using MMS and AMR for the 2D polygonal basis functions

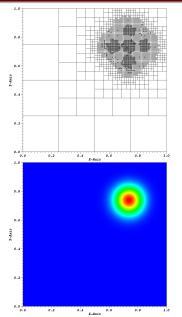
$$\psi(x,y) = x(L_x - x)y(L_y - y) \exp(-\frac{(x - x_0)^2 + (y - y_0)^2}{\gamma}),$$
  
$$\phi(x,y) = 2\pi x(L_x - x)y(L_y - y) \exp(-\frac{(x - x_0)^2 + (y - y_0)^2}{\gamma})$$

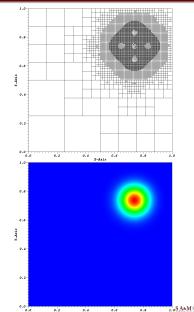
Proposed Work and Current Status



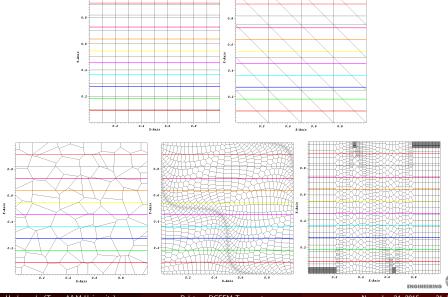


#### Linear ME cycle 15 (left) and quadratic ME cycle 08 (right)



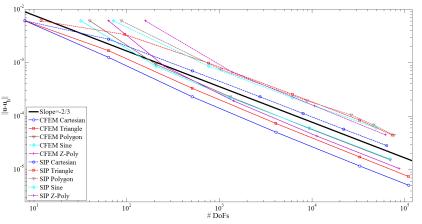


## SIP exactly linear solutions on 3D polyhedral meshes using the PWL basis functions



$$\Phi(x, y, z) = xyz(L_x - x)(L_y - y)(L_z - z)$$
  
 
$$L_x = L_x = L_x = 1.0$$

Proposed Work and Current Status 00000000000



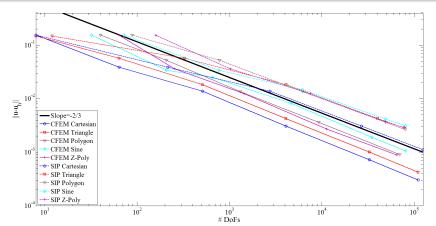




## SIP convergence study - gaussian solution on 3D cube using the PWL basis functions

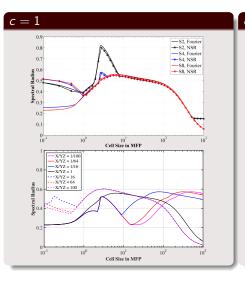
$$\Phi(x, y, z) = xyz(L_x - x)(L_y - y)(L_z - z) \exp(-(\mathbf{r} - \mathbf{r}_0) \cdot (\mathbf{r} - \mathbf{r}_0))$$

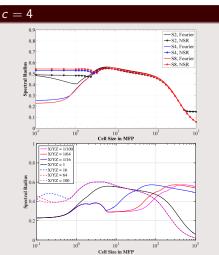
$$L_x = L_x = L_x = 1.0, \qquad \mathbf{r}_0 = (3/4, 3/4, 3/4)$$



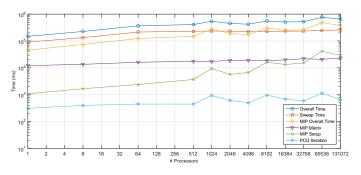


#### Fourier analysis - 3D PWL basis functions









## Problem Description

- Modified Zerr problem used optimal sweep aggregation parameters
  - homogeneous cube about 500 mfp and c=0.9999
  - 58 level-symmetric quadrature
- pointwise convergence tolerance of 1e-8
- SI precondition with MIP DSA using HYPRE PCG and AMG





Overview

#### Two-grid acceleration implementation in PDT

- Successfully implemented and deugged
  - Includes non-orthogonal mesh configurations
  - Includes multi-material configurations
- Have tested the two-grid methodology on a homogeneous graphite block as well as a block with an air duct
- Iteration counts for a very large configuration (very optically thick) are similar to simple infinite medium calculations

Materials	Unaccelerated Iterations	Accelerated Iterations	
Graphite Only	2027	21	
Graphite + Air Duct	2138	23	



#### POLYFEM Ongoing Work

- Finish quadratic serendipity extension for polygonal finite elements to Wachspress, PWL, and mean value.
- Determine the effects of numerical integration strategies on highly-distorted polygonal elements
- Perform analysis on benchmark cases using AMR
  - Searchlight problem (partially completed)
  - IAEA-EIR-2 problem

## MIP DSA Ongoing Work

- Finish the MIP DSA Fourier analysis for all the 2D polygonal basis functions (including the quadratic serendipity extension).
- Analyze the effects of AMR with polygonal basis functions on the MIP DSA PCG iteration counts (with and without bootstrapping)
- Remaining PDT work:
  - Complete parametric studies (a lot of the raw results have been compiled) of MIP DSA with HYPRE and generate a working performance model
  - Perform additional two-grid acceleration studies on simple geometries
  - Run IM1 problem with two-grid acceleration
  - Commit all acceleration additions to master for use in the CERT project



# Questions?

A special acknowledgment to the Department of Energy Rickover Fellowship Program in Nuclear Engineering, which provides strong support to its fellows and their professional development.





