$$\rho_{,t} + (\rho v^i)_{;i} = 0 \tag{1}$$

$$(\rho v_j)_{,t} + (\rho v^i v_j + P \delta_j^i)_{;i} = -\rho \phi_{,j} + c^{-1} \sum_s \int_0^\infty (\kappa_{s\varepsilon} + \sigma_{s\varepsilon}^{\text{tr}}) F_{s\varepsilon j} d\varepsilon$$
(2)

$$\left[\rho\left(e + \frac{1}{2}\|v\|^{2}\right)\right]_{t} + \left[\rho v^{i}\left(e + \frac{1}{2}\|v\|^{2} + \frac{P}{\rho}\right)\right]_{i} = -\rho v^{i}\phi_{,i}$$
(3)

$$-\sum_{s} \int_{0}^{\infty} \left(j_{s\varepsilon} - c\kappa_{s\varepsilon} E_{s\varepsilon} - \frac{v^{i}}{c} (\kappa_{s\varepsilon} + \sigma_{s\varepsilon}^{\text{tr}}) F_{s\varepsilon i} \right) d\varepsilon$$
(4)

$$(\rho Y_e)_{,t} + (\rho Y_e v^i)_{;i} = \sum_{s} \int_0^\infty \xi_{s\varepsilon} (j_{s\varepsilon} - c\kappa_{s\varepsilon} E_{s\varepsilon}) d\varepsilon$$
 (5)

$$E_{s\varepsilon,t} + (F_{s\varepsilon}^i + v^i E_{s\varepsilon})_{;i} - v_{;j}^i \frac{\partial}{\partial \ln \varepsilon} P_{s\varepsilon i}^j = j_{s\varepsilon} - c\kappa_{s\varepsilon} E_{s\varepsilon}$$

$$\tag{6}$$

$$F_{s\varepsilon j,t} + (c^2 P_{s\varepsilon j}^i + v^i F_{s\varepsilon j})_{;i} + v_{;j}^i F_{s\varepsilon i} - v_{;k}^i Q_{s\varepsilon ji}^k - v_{;k}^i \frac{\partial}{\partial \ln \varepsilon} Q_{s\varepsilon ji}^k = -c(\kappa_{s\varepsilon} + \sigma_{s\varepsilon}^{tr}) F_{s\varepsilon j}$$

$$\tag{7}$$

where e is the specific internal energy, $s \in \{\nu_e, \bar{\nu}_e, \nu_x\}$,

$$\xi_{s\varepsilon} = \begin{cases} -(N_A \varepsilon)^{-1} & s = \nu_e, \\ (N_A \varepsilon)^{-1} & s = \bar{\nu}_e, \\ 0 & s = \nu_x, \end{cases}$$
 (8)

 $P = P(\rho, e, Y_e)$, and we use the M1 closure to truncate the radiation moment hierarchy by specifying the second and third moments in terms of the zeroth and first.

1 Algorithm Steps

The following describes an Euler push (i.e. 1^{st} -order in time). The generalization to 2^{nd} -order Runge-Kutta schemes is straightforward.

1.1 Step 1

Advance this subsystem by Δt :

$$\rho_{,t} + (\rho v^i)_{;i} = 0 \tag{9}$$

$$(\rho v_j)_{,t} + (\rho v^i v_j + P\delta_j^i)_{;i} = -\rho \phi_{,j}$$

$$\tag{10}$$

$$\left[\rho\left(e + \frac{1}{2}\|v\|^{2}\right)\right]_{t} + \left[\rho v^{i}\left(e + \frac{1}{2}\|v\|^{2} + \frac{P}{\rho}\right)\right]_{:i} = -\rho v^{i}\phi_{,i}$$
(11)

$$(\rho Y_e)_{,t} + (\rho Y_e v^i)_{;i} = 0 \tag{12}$$

$$E_{s\varepsilon,t} + (F_{s\varepsilon}^i + v^i E_{s\varepsilon})_{;i} = 0 (13)$$

$$F_{s\varepsilon j,t} + (c^2 P_{s\varepsilon j}^i + v^i F_{s\varepsilon j})_{;i} + v_{;i}^i F_{s\varepsilon i} - v_{;k}^i Q_{s\varepsilon ji}^k = 0$$
(14)

1.2 Step 2

Advance this subsystem by Δt :

$$E_{s\varepsilon,t} - v_{,j}^{i} \frac{\partial}{\partial \ln \varepsilon} P_{s\varepsilon i}^{j} = 0 \tag{15}$$

$$F_{s\varepsilon j,t} - v_{,k}^{i} \frac{\partial}{\partial \ln \varepsilon} Q_{s\varepsilon ji}^{k} = 0 \tag{16}$$

1.3 Step 3

Advance this subsystem by Δt :

$$(\rho v_j)_{,t} = c^{-1} \sum_{s} \int_0^\infty (\kappa_{s\varepsilon} + \sigma_{s\varepsilon}^{tr}) F_{s\varepsilon j} d\varepsilon$$
 (17)

$$\left[\rho\left(e + \frac{1}{2}\|v\|^2\right)\right]_{t} = \sum_{s} \int_0^\infty \left(\frac{v^i}{c}(\kappa_{s\varepsilon} + \sigma_{s\varepsilon}^{\text{tr}})F_{s\varepsilon i}\right) d\varepsilon \tag{18}$$

$$F_{s\varepsilon j,t} = -c(\kappa_{s\varepsilon} + \sigma_{s\varepsilon}^{\text{tr}}) F_{s\varepsilon j}$$
(19)

1.4 Step 4

Advance this subsystem by Δt :

$$(\rho e)_{,t} = -\sum_{s} \int_{0}^{\infty} (j_{s\varepsilon} - c\kappa_{s\varepsilon} E_{s\varepsilon}) d\varepsilon$$
 (20)

$$(\rho Y_e)_{,t} = \sum_{s} \int_0^\infty \xi_{s\varepsilon} (j_{s\varepsilon} - c\kappa_{s\varepsilon} E_{s\varepsilon}) d\varepsilon$$
 (21)

$$E_{s\varepsilon,t} = j_{s\varepsilon} - c\kappa_{s\varepsilon} E_{s\varepsilon} \tag{22}$$