

## PEU 416 Assignment 3

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# 1 Problem 1

(a)

$$\rho(t) = \rho_0 \left( \frac{a_0}{a(t)} \right)^3$$

$$\dot{\rho}(t) = -3\rho_0 \frac{a_0^3}{a^4} \dot{a} = -3 \frac{\dot{a}}{a} \rho_0 \left( \frac{a_0}{a} \right)^3 = -3H\rho(t)$$

(b)

$$g = \ddot{r} = -\frac{GM(r_i)}{r^2} = -\frac{4\pi}{3} \frac{Gr_i^3 \rho_i}{r^2}$$

$$a \equiv \frac{r}{r_i} \rightarrow r = ar_i, \quad \ddot{r} = \ddot{a}r_i$$

$$\ddot{a}r_i = -\frac{4\pi}{3} \frac{Gr_i \rho_i}{a^2}$$

$$\ddot{a} = -\frac{4\pi}{3} \frac{G\rho_i}{a^2} = -\frac{4\pi}{3} \frac{G(\rho(t)a^3)}{a^2} = -\frac{4\pi}{3} G\rho a \quad \text{for } a_0 = 1$$

$$\ddot{a} + \frac{4\pi}{3} G\rho a = 0$$

$$\dot{a} \left( \ddot{a} + \frac{4\pi}{3} \frac{G\rho_i}{a^2} \right) = 0$$

$$\frac{dE}{dt} = 0 \rightarrow E = \text{const}$$

$$\frac{\dot{a}^2}{2} - \frac{4\pi}{3} \frac{G\rho_i}{a} = E$$

$$\frac{\dot{a}^2}{2} - \frac{4\pi G\rho(t)a^2}{3} = E \equiv -\frac{K}{2}$$

$$\frac{\dot{a}^2}{2} - \frac{4\pi G\rho(t)a^3}{3a} = -\frac{K}{2}$$

$$\frac{\dot{a}^2}{2} - \frac{4\pi G\rho_i a_i^3}{3a} = -\frac{K}{2}$$

$$V = -\frac{4\pi G\rho_i a_i^3}{3a}$$

$$\frac{\dot{a}^2}{2} + V = -\frac{K}{2}$$

(c)

$$\begin{aligned}
\frac{\dot{a}^2}{2} + V &= 0 \\
\dot{a} &= \sqrt{-2V} \\
\sqrt{a}\dot{a} &= \sqrt{\frac{8\pi G\rho_i a_i^3}{3}} \\
\sqrt{a} \, da &= \sqrt{\frac{8\pi G\rho_i a_i^3}{3}} \, dt \\
\frac{2}{3}a^{\frac{3}{2}} &= \sqrt{\frac{8\pi G\rho_i a_i^3}{3}} t \\
a &= a_i \sqrt[3]{12\pi G\rho_i t^{\frac{2}{3}}}
\end{aligned}$$

(d)

$$\begin{aligned}
\frac{\dot{a}^2}{2} - \frac{4\pi G\rho(t)a^2}{3} &= -\frac{K}{2} \\
K &= \frac{8\pi G\rho(t)a^2}{3} - \dot{a}^2 \\
&= \frac{8\pi G\rho(t)a^2}{3} - a^2 \frac{\dot{a}^2}{a^2} \\
&= \frac{8\pi G\rho(t)a^2}{3} - a^2 H^2 \\
&= \frac{8\pi G(\rho_c \Omega)a^2}{3} - a^2 \left( \frac{8}{3}\pi G\rho_c \right) \\
&= \frac{8\pi G}{3} a^2 \rho_c (\Omega - 1)
\end{aligned}$$

### **$\Omega$ value indications**

Since by measuring  $\Omega$ , we have the  $K$  value

$$\begin{aligned}
\frac{\dot{a}^2}{2} - \frac{4\pi G\rho_i a_i^3}{3a} &= -\frac{K}{2} \\
\dot{a}^2 &= -K + \frac{8\pi G\rho_i a_i^3}{3a}
\end{aligned}$$

If  $\Omega < 1$  then  $K < 0$ ,  $\dot{a}^2 > 0$ : the universe expands forever.

If  $\Omega = 1$  then  $K = 0$ : the universe is flat

If  $\Omega > 1$  then  $K > 0$ ,  $a_{\max} \geq \dot{a} \geq 0$ : the universe approaches maximum size and then contracts to a Big Crunch.

## 2 Problem 2

$$\Phi = \phi + a\ddot{a}\frac{x^2}{2}$$

$$\nabla_x^2 \Phi = \nabla_x^2 \phi + \nabla_x^2 \left( a\ddot{a}\frac{x^2}{2} \right)$$

$$\nabla_r^2 \phi = 4\pi G\rho$$

$$\nabla_r^2 = \frac{1}{a^2} \nabla_x^2 \quad \rightarrow \quad \nabla_x^2 \phi = 4\pi G a^2 \rho$$

$$\nabla_x^2 \left( a\ddot{a}\frac{x^2}{2} \right) = \sum_{i=1,2,3} \frac{\partial^2}{\partial x_i^2} \left( a\ddot{a}\frac{x_i^2}{2} \right) = 3a\ddot{a} = -4\pi G \bar{\rho} a^2$$

$$\therefore \nabla_x^2 \Phi = 4\pi G a^2 (\rho - \bar{\rho}) = 4\pi G a^2 \bar{\rho} \delta$$

### 3 Problem 3

(a)

$$\begin{aligned}
\ddot{\delta}_m + 2H\dot{\delta}_m &= 4\pi G\bar{\rho}_m\delta_m \\
\ddot{\delta}_m + 2H\dot{\delta}_m &= 4\pi G\bar{\rho}_{m,0}a^{-3}\delta_m \\
\frac{d}{dt} &= \frac{da}{dt} \cdot \frac{d}{da} = \dot{a} \frac{d}{da} = Ha \frac{d}{da} \\
\dot{\delta}_m &= Ha \frac{d\delta_m}{da}, \quad \ddot{\delta}_m = \frac{d}{dt}(\dot{\delta}_m) = \frac{d}{dt} \left( Ha \frac{d\delta_m}{da} \right) = Ha \frac{d}{da} \left( Ha \frac{d\delta_m}{da} \right) \\
\therefore Ha \frac{d}{da} \left( Ha \frac{d\delta_m}{da} \right) + 2H^2 a \frac{d\delta_m}{da} &= 4\pi G\bar{\rho}_{m,0}a^{-3}\delta_m
\end{aligned}$$

multiply both sides by  $\frac{a}{H}$

$$\begin{aligned}
a^2 \frac{d}{da} \left( Ha \frac{d\delta_m}{da} \right) + 2Ha^2 \frac{d\delta_m}{da} &= 4\pi G\bar{\rho}_{m,0} \frac{\delta_m}{Ha^2} \\
\frac{d}{da} \left( a^3 H \frac{d\delta_m}{da} \right) &= 4\pi G\bar{\rho}_{m,0} \frac{\delta_m}{Ha^2}
\end{aligned}$$

(b)

$$\delta_m \propto \begin{cases} H \\ H \int \frac{da}{(aH)^3} \end{cases}$$

For  $\delta_m = H$

$$\begin{aligned}
\frac{d}{da} \left( a^3 H \frac{d\delta_m}{da} \right) &= 4\pi G\bar{\rho}_{m,0} \frac{\delta_m}{Ha^2} = \frac{3}{2}aH^2 \\
\frac{d}{da} \left( a^3 H \frac{dH}{da} \right) &= \frac{3}{2}aH^2 \\
\frac{1}{2} \frac{d}{da} \left( a^3 \frac{d}{da} (H^2) \right) &= \frac{3}{2}aH^2 \\
\frac{1}{2} \left( 3a^2 \dot{a} \frac{d}{da} (H^2) + a^3 \frac{d^2 H^2}{da^2} \right) &= \frac{3}{2}aH^2 \\
\frac{3}{2}a^3 \left( H \frac{d}{da} (H^2) + \frac{1}{3} \frac{d^2 H^2}{da^2} \right) &= \frac{3}{2}aH^2 \\
\frac{2}{3} \left( (3\dot{a} - 1)\dot{H} + \frac{\ddot{H}}{H} \right) &= 0
\end{aligned}$$

$$\begin{aligned}
& \frac{d}{da} \left( a^3 H \frac{dH}{da} \right) \\
&= \frac{1}{a^2} \left( 3\dot{a}a^4 H \frac{dH}{da} + a^5 \left( \frac{dH}{da} \right)^2 + a^5 H \frac{d^2 H}{da^2} \right) \\
&= \frac{1}{a^2} \left( 3\dot{a}^2 a^3 \frac{d(\frac{\dot{a}}{a})}{da} + a^5 \left( \frac{d(\frac{\dot{a}}{a})}{da} \right)^2 + \dot{a}a^4 \frac{d^2(\frac{\dot{a}}{a})}{da^2} \right) \\
&= \frac{1}{a^2} \left( 3\dot{a}^2 a^3 \frac{a\ddot{a} - \dot{a}^2}{a^2} + a^5 \left( \frac{a\ddot{a} - \dot{a}^2}{a^2} \right)^2 + \dot{a}a^4 \frac{d(\frac{a\ddot{a} - \dot{a}^2}{a^2})}{da} \right) \\
&= \frac{1}{a^2} \left( 3\dot{a}^2 a(a\ddot{a} - \dot{a}^2) + a(a\ddot{a} - \dot{a}^2)^2 + \dot{a}a^4 \left( \frac{a^2(\ddot{a}^2 + a\ddot{a} - 2\dot{a}\ddot{a}) - (a\ddot{a} - \dot{a}^2)(2a\dot{a})}{a^4} \right) \right) \\
&= \frac{1}{a^2} \left( 3\dot{a}^2 a(a\ddot{a} - \dot{a}^2) + a(a\ddot{a} - \dot{a}^2)^2 + \dot{a}(a^2(\ddot{a}^2 + a\ddot{a} - 2\dot{a}\ddot{a}) - 2a\dot{a}(a\ddot{a} - \dot{a}^2)) \right) \\
&= \frac{1}{a^2} \left( a\dot{a}^2(a\ddot{a} - \dot{a}^2) + a(a\ddot{a} - \dot{a}^2)^2 + \dot{a}a^2(\ddot{a}^2 + a\ddot{a} - 2\dot{a}\ddot{a}) \right) \\
&= \frac{1}{a^2} (a^2\dot{a}^2\ddot{a} - a\dot{a}^4 + a^3\ddot{a}^2 + a\dot{a}^4 - 2a^2\dot{a}^2\ddot{a} + \dot{a}a^2\ddot{a}^2 + \dot{a}a^3\ddot{a} - 2a^2\dot{a}^2\ddot{a}) \\
&= \frac{1}{a^2} (a^3\ddot{a}^2 - 3a^2\dot{a}^2\ddot{a} + a^2\dot{a}\ddot{a}^2 + a^3\dot{a}\ddot{a}) \\
&= a\ddot{a}^2 - 3\dot{a}^2\ddot{a} + \dot{a}\ddot{a}^2 + a\dot{a}\ddot{a}
\end{aligned}$$

$$H^2 = \frac{8\pi G}{3} \rho_m$$

$$H = \frac{\dot{a}}{a}$$

$$\dot{H} = \frac{a\ddot{a} - \dot{a}^2}{a^2}$$

$$\ddot{H} = \frac{a^2(\ddot{a}^2 + a\ddot{a} - 2\dot{a}\ddot{a}) - (a\ddot{a} - \dot{a}^2)(2a\dot{a})}{a^4}$$

$$\ddot{\delta}_m + 2H\dot{\delta}_m = 4\pi G\bar{\rho}_m\delta_m$$

$$\ddot{H} + 2H\dot{H} = 4\pi G\bar{\rho}_m H$$

$$\frac{a^2(\ddot{a}^2 + a\ddot{a} - 2\dot{a}\ddot{a}) - (a\ddot{a} - \dot{a}^2)(2a\dot{a})}{a^4} + 2\left(\frac{\dot{a}}{a}\right)\left(\frac{a\ddot{a} - \dot{a}^2}{a^2}\right) = 4\pi G\bar{\rho}_m\left(\frac{\dot{a}}{a}\right)$$

$$\frac{a^2(\ddot{a}^2 + a\ddot{a} - 2\dot{a}\ddot{a}) - (a\ddot{a} - \dot{a}^2)(2a\dot{a})}{a^3} + 2\frac{a\dot{a}\ddot{a} - \dot{a}^3}{a^2} = 4\pi G\bar{\rho}_m\dot{a}$$

$$a(\ddot{a}^2 + a\ddot{a} - 2\dot{a}\ddot{a}) - 2\dot{a}(a\ddot{a} - \dot{a}^2) + 2(a\dot{a}\ddot{a} - \dot{a}^3) = 4\pi G\bar{\rho}_m\dot{a}^2$$

$$a\ddot{a}^2 + a^2\ddot{a} - 2a\dot{a}\ddot{a} = 4\pi G\bar{\rho}_ma^2\dot{a}$$

$$\frac{\ddot{a}^2}{a\dot{a}} + \frac{\ddot{a}}{a} - \frac{2\dot{a}\ddot{a}}{a} = 4\pi G\bar{\rho}_m$$

$$\frac{\ddot{a}^2 - 2\dot{a}\ddot{a} + \dot{a}\ddot{a}}{a\dot{a}} = 4\pi G\bar{\rho}_m$$

$$\frac{d(\dot{a}\ddot{a} - \dot{a}^2)}{dt} \frac{1}{a\dot{a}} = 4\pi G\bar{\rho}_m$$

For  $\delta_m = H \int \frac{da}{(aH)^3}$

(c)

$$H^2 = \frac{8\pi G}{3}\rho_{\text{total}} = \frac{8\pi G}{3}(\rho_r + \rho_m)$$

$$\rho_r \propto a^{-4}, \quad \rho_m \propto a^{-3}$$

$$\rho_r = \rho_{\text{eq}}\left(\frac{a_{\text{eq}}}{a}\right)^4, \quad \rho_m = \rho_{\text{eq}}\left(\frac{a_{\text{eq}}}{a}\right)^3 \quad \text{where } \rho_{\text{eq}} = \rho_{r, \text{eq}} = \rho_{m, \text{eq}}$$

$$y \equiv \frac{a}{a_{\text{eq}}} \rightarrow \rho_r = \rho_{\text{eq}}y^{-4}, \quad \rho_m = \rho_{\text{eq}}y^{-3}$$

$$H^2 = \frac{8\pi G}{3}\rho_{\text{eq}}(y^{-4} + y^{-3}) = \frac{8\pi G}{3}\frac{\rho_{\text{eq}}}{y^4}(1 + y)$$

$$\therefore H = \frac{A}{y^2}\sqrt{1 + y} \quad \text{where } A = \sqrt{\frac{8\pi G}{3}\rho_{\text{eq}}}$$

$$\frac{d}{da} \left( a^3 H \frac{d\delta_m}{da} \right) = 4\pi G \bar{\rho}_{m,0} \frac{\delta_m}{H a^2}$$

$$a = a_{\text{eq}} y, \quad \frac{d}{da} = \frac{1}{a_{\text{eq}}} \frac{d}{dy}$$

$$\frac{1}{a_{\text{eq}}} \frac{d}{dy} \left( (a_{\text{eq}} y)^3 H \frac{1}{a_{\text{eq}}} \frac{d\delta_m}{dy} \right) = 4\pi G \bar{\rho}_{m,0} \frac{\delta_m}{H (a_{\text{eq}} y)^2}$$

$$\frac{d}{dy} \left( y^3 H \frac{d\delta_m}{dy} \right) = 4\pi G \frac{\bar{\rho}_{m,0}}{a_{\text{eq}}^3} \frac{\delta_m}{H y^2}$$

$$\frac{d}{dy} \left( y^3 H \frac{d\delta_m}{dy} \right) = 4\pi G \rho_{\text{eq}} \frac{\delta_m}{H y^2}$$

Substituting with  $H = \frac{A}{y^2} \sqrt{1+y}$

$$\frac{d}{dy} \left( A y \sqrt{1+y} \frac{d\delta_m}{dy} \right) = 4\pi G \rho_{\text{eq}} \frac{\delta_m}{A \sqrt{1+y}}$$

$$\frac{d}{dy} \left( y \sqrt{1+y} \frac{d\delta_m}{dy} \right) = \frac{4\pi G \rho_{\text{eq}}}{A^2} \frac{\delta_m}{\sqrt{1+y}}$$

$$\frac{d}{dy} \left( y \sqrt{1+y} \frac{d\delta_m}{dy} \right) = \frac{3}{2} \frac{\delta_m}{\sqrt{1+y}}$$

$$y \sqrt{1+y} \frac{d^2 \delta_m}{dy^2} + \left( \sqrt{1+y} + \frac{y}{2\sqrt{1+y}} \right) \frac{d\delta_m}{dy} = \frac{3}{2} \frac{\delta_m}{\sqrt{1+y}}$$

$$\frac{d^2 \delta_m}{dy^2} + \left( \frac{1}{y} + \frac{1}{2(1+y)} \right) \frac{d\delta_m}{dy} = \frac{3}{2} \frac{\delta_m}{y(1+y)}$$

$$\frac{d^2 \delta_m}{dy^2} + \frac{2+3y}{2y(1+y)} \frac{d\delta_m}{dy} = \frac{3}{2} \frac{\delta_m}{y(1+y)}$$

### Verifying the solutions

for  $\delta_{m1} = 1 + \frac{3}{2}y$ :

$$\frac{d\delta_m}{dy} = \frac{3}{2}, \quad \frac{d^2 \delta_m}{dy^2} = 0$$

$$0 + \frac{2+3y}{2y(1+y)} \frac{3}{2} = \frac{3}{2} \frac{1 + \frac{3}{2}y}{y(1+y)}$$

$$\frac{2+3y}{2y(1+y)} \frac{3}{2} = \frac{3}{2} \frac{2+3y}{2y(1+y)} \quad (\text{verified})$$

for  $\delta_{m2} = (1 + \frac{3}{2}y) \ln \left( \frac{\sqrt{1+y}+1}{\sqrt{1+y}-1} \right) - 3\sqrt{1+y}$ :



$$\begin{aligned}
\frac{d\delta_m}{dy} &= \frac{3}{2} \ln\left(\frac{\sqrt{1+y}+1}{\sqrt{1+y}-1}\right) - \left(1 + \frac{3}{2}y\right) \frac{1}{y\sqrt{y+1}} - \frac{3}{2\sqrt{1+y}} \\
&= \frac{3}{2} \ln\left(\frac{\sqrt{1+y}+1}{\sqrt{1+y}-1}\right) - \frac{1}{y\sqrt{y+1}} - \frac{3}{\sqrt{1+y}} \\
&= \frac{3}{2} \ln\left(\frac{\sqrt{1+y}+1}{\sqrt{1+y}-1}\right) - \left(\frac{1}{y} + 3\right) \frac{1}{\sqrt{y+1}} \\
\frac{d^2\delta_m}{dy^2} &= -\frac{3}{2} \frac{1}{y\sqrt{y+1}} + \frac{1}{y^2\sqrt{y+1}} + \frac{1}{2} \left(\frac{1}{y} + 3\right) \frac{1}{(y+1)^{\frac{3}{2}}} \\
&= \left(-\frac{3}{2} + \frac{1}{y}\right) \frac{1}{y\sqrt{y+1}} + \frac{1}{2} \left(\frac{1}{y} + 3\right) \frac{1}{(y+1)^{\frac{3}{2}}} \\
&= \frac{2-3y}{2y} \frac{1}{y\sqrt{y+1}} + \frac{3y+1}{2y} \frac{1}{(y+1)^{\frac{3}{2}}}
\end{aligned}$$

$$\begin{aligned}
\text{LHS} &= \left[ \frac{2-3y}{2y} \frac{1}{y\sqrt{y+1}} + \frac{3y+1}{2y} \frac{1}{(y+1)^{\frac{3}{2}}} \right] + \frac{2+3y}{2y(1+y)} \left[ \frac{3}{2} \ln\left(\frac{\sqrt{1+y}+1}{\sqrt{1+y}-1}\right) - \frac{3y+1}{y} \frac{1}{\sqrt{y+1}} \right] \\
&= \left[ \frac{2-3y}{2y} \frac{1}{y\sqrt{y+1}} + \frac{3y+1}{2y} \frac{1}{(y+1)^{\frac{3}{2}}} \right] + \frac{3}{2} \frac{2+3y}{2y(1+y)} \ln\left(\frac{\sqrt{1+y}+1}{\sqrt{1+y}-1}\right) - \frac{(2+3y)(3y+1)}{2y^2} \frac{1}{(y+1)^{\frac{3}{2}}} \\
&= \left[ \frac{2-3y}{2y} \frac{1}{y\sqrt{y+1}} + \frac{3y+1}{2y} \frac{1}{(y+1)^{\frac{3}{2}}} \right] + \frac{3}{2} \frac{2+3y}{2y(1+y)} \ln\left(\frac{\sqrt{1+y}+1}{\sqrt{1+y}-1}\right) - \frac{9y^2+9y+2}{2y^2} \frac{1}{(y+1)^{\frac{3}{2}}} \\
&= \left[ \frac{2-3y}{2y} \frac{1}{y\sqrt{y+1}} + \cancel{\frac{3y+1}{2y} \frac{1}{(y+1)^{\frac{3}{2}}}} \right] + \frac{3}{2} \frac{2+3y}{2y(1+y)} \ln\left(\frac{\sqrt{1+y}+1}{\sqrt{1+y}-1}\right) - \left[ \cancel{\frac{3y^2+y}{2y^2}} + \frac{6y^2+8y+2}{2y^2} \right] \frac{1}{(y+1)^{\frac{3}{2}}} \\
&= \left[ \frac{2-3y}{2y} \frac{1}{y\sqrt{y+1}} \right] + \frac{3}{2} \frac{2+3y}{2y(1+y)} \ln\left(\frac{\sqrt{1+y}+1}{\sqrt{1+y}-1}\right) - \left[ \frac{2(3y+1)(y+1)}{2y^2} \right] \frac{1}{(y+1)^{\frac{3}{2}}} \\
&= \left[ \frac{2-3y}{2y} \frac{1}{y\sqrt{y+1}} \right] + \frac{3}{2} \frac{2+3y}{2y(1+y)} \ln\left(\frac{\sqrt{1+y}+1}{\sqrt{1+y}-1}\right) - \left[ \frac{2(3y+1)}{2y} \right] \frac{1}{y\sqrt{y+1}} \\
&= \frac{3}{2} \frac{2+3y}{2y(1+y)} \ln\left(\frac{\sqrt{1+y}+1}{\sqrt{1+y}-1}\right) - \frac{9}{2} \frac{1}{y\sqrt{1+y}} \\
\text{RHS} &= \frac{3}{2} \frac{1}{y(1+y)} \left[ \left(1 + \frac{3}{2}y\right) \ln\left(\frac{\sqrt{1+y}+1}{\sqrt{1+y}-1}\right) - 3\sqrt{1+y} \right] \\
&= \frac{3}{2} \frac{2+3y}{2y(1+y)} \ln\left(\frac{\sqrt{1+y}+1}{\sqrt{1+y}-1}\right) - \frac{9}{2} \frac{1}{y\sqrt{1+y}} \quad (\text{verified})
\end{aligned}$$

### The growing and decaying modes

at early times,  $a \ll a_{\text{eq}}$ ,  $y \ll 1$ ,

$$\delta_{m1} \rightarrow 1, \quad \delta_{m2} \rightarrow \infty$$

The growing mode:  $\delta_{m2}$

The decaying mode:  $\delta_{m1}$

at late times,  $a \gg a_{\text{eq}}$ ,  $y \gg 1$ ,

$$\delta_{m1} \rightarrow \frac{3}{2}y, \quad \delta_{m2} \rightarrow -3\sqrt{y}$$

The growing mode:  $\delta_{m1}$

The decaying mode:  $\delta_{m2}$

## 4 Problem 4

(a)

$$c_s^2 = \frac{\partial P}{\partial \rho} = \frac{4}{3} \alpha \rho^{\frac{1}{3}}$$

$$t \propto a^{\frac{3}{2}}, \quad \rho \propto a^{-3} \quad \text{for a matter-dominated universe}$$

$$dt = a d\eta$$

$$d\eta = \frac{1}{a} dt \rightarrow \eta \propto a^{\frac{1}{2}} \rightarrow a \propto \eta^2$$

$$\therefore \rho \propto \eta^{-6} \rightarrow c_s^2 \propto \eta^{-2}$$

$$\therefore c_s^2 = c_0^2 \left( \frac{\eta_0}{\eta} \right)^2$$

(b) **Small Scales (large  $k$ , where  $c_s^2 k^2 \gg \frac{3}{2} H^2$ ):**

the term  $c_s^2 k^2$  dominates the equation because  $k^2$  (related to the wavelength  $\lambda = \frac{2\pi}{k}$ ) is large, and the sound speed  $C_s$  allows pressure to resist gravitational collapse. The equation approximates to  $\delta'' + c_s^2 k^2 \delta \approx 0$ , which resembles an oscillatory solution due to pressure waves. This indicates that density perturbations oscillate as acoustic waves rather than growing, as pressure prevents collapse on scales smaller than the Jeans length (where  $C_s k \sim H$ ).

**Large Scales (small  $k$ , where  $c_s^2 k^2 \ll \frac{3}{2} H^2$ ):**

the term  $\frac{3}{2} H^2$  dominates, as  $k^2$  is small. The equation simplifies to  $\delta'' + H\delta' - \frac{3}{2} H^2 \delta \approx 0$ . In a matter-dominated universe,  $H \propto \eta^{-1}$ , leading to a growing mode  $\delta \propto \eta^2 \propto a$ . This reflects gravitational instability, where perturbations grow over time due to self-gravity, unopposed by pressure on scales larger than the Jeans length.

(c)

$$a \propto \eta^2 \rightarrow H = \frac{2}{\eta}$$

$$\delta'' + \frac{2}{\eta} \delta' + \left( c_0^2 \left( \frac{\eta_0}{\eta} \right)^2 k^2 - \frac{6}{\eta^2} \right) \delta = 0$$

$$\eta^2 \delta'' + 2\eta \delta' + (c_0^2 \eta_0^2 k^2 - 6) \delta = 0$$

Let  $\delta = C\eta^\beta$ ,

$$\delta' = C\beta\eta^{\beta-1}, \quad \delta'' = C\beta(\beta-1)\eta^{\beta-2}$$

$$C\beta(\beta-1)\eta^\beta + 2C\beta\eta^\beta + (c_0^2 \eta_0^2 k^2 - 6)C\eta^\beta = 0$$

$$\beta(\beta-1) + 2\beta + (c_0^2 \eta_0^2 k^2 - 6) = 0$$

$$\beta^2 + \beta + (c_0^2 \eta_0^2 k^2 - 6) = 0$$

$$\begin{aligned}\beta &= \frac{1}{2} \left( -1 \pm \sqrt{1 - 4(c_0^2 \eta_0^2 k^2 - 6)} \right) \\ &= \frac{1}{2} \left( -1 \pm \sqrt{25 - 4c_0^2 \eta_0^2 k^2} \right)\end{aligned}$$

The transition occurs at

$$\begin{aligned}25 - 4c_0^2 \eta_0^2 k^2 &= 0 \\ k &= \frac{5}{2c_0 \eta_0}\end{aligned}$$

The growing mode is  $\eta^{\beta_+}$ , where  $\beta_+ = \frac{1}{2}(-1 + \sqrt{25 - 4c_0^2 \eta_0^2 k^2})$  it is constant when  $\beta_+ = 0$ , or

$$\begin{aligned}-1 + \sqrt{25 - 4c_0^2 \eta_0^2 k^2} &= 0 \\ 25 - 4c_0^2 \eta_0^2 k^2 &= 1 \\ 24 - 4c_0^2 \eta_0^2 k^2 &= 0 \\ k &= \frac{\sqrt{6}}{c_0 \eta_0}\end{aligned}$$

## References

- [1] M. El-Deeb, “PEU-405 Assignments.” [Online]. Available: <https://github.com/mhdeeb/peu-assignments/tree/main/peu-405>