

PEU 323 Assignment 4

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1 Problem 1

(a) Yes, this wave function is physically acceptable if the normalization constant A is chosen such that $\int_{-a}^a |\psi(x)|^2 dx = 1$. The wave function is continuous and single-valued within $[-a, a]$, meeting the requirements for physical acceptability.

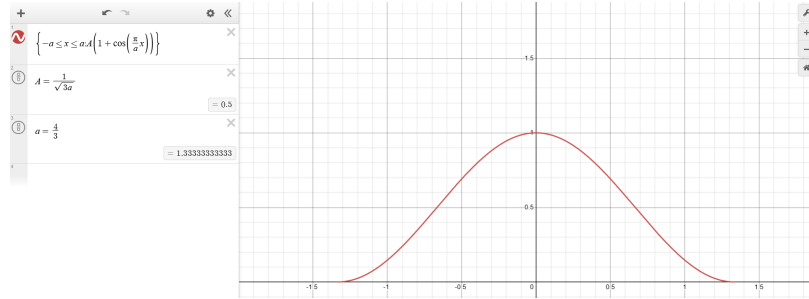


Figure 1: ψ

(b) The classically allowed region is:

$$[-a, a]$$

since the wave function $\psi(x) = 0$ outside this interval, meaning the classical particle cannot be found outside $[-a, a]$.

2 Problem 2

(a)

$$\psi(x, t) = \sin\left(\frac{n\pi}{a}x\right) e^{-i\omega t}$$

$$\frac{-\hbar^2}{2m} \frac{\partial^2 \psi}{\partial x^2} + V(x)\psi = i\hbar \frac{\partial \psi}{\partial t}$$

$$\begin{aligned} V(x)\psi &= \frac{\hbar^2}{2m} \frac{\partial^2 \psi}{\partial x^2} + i\hbar \frac{\partial \psi}{\partial t} \\ &= \frac{\hbar^2}{2m} \frac{\partial^2 (\sin(\frac{n\pi}{a}x))}{\partial x^2} e^{-i\omega t} + i\hbar \sin\left(\frac{n\pi}{a}x\right) \frac{\partial (e^{-i\omega t})}{\partial t} \\ &= -\frac{\hbar^2 n^2 \pi^2}{2ma^2} \sin\left(\frac{n\pi}{a}x\right) e^{-i\omega t} + \omega \hbar \sin\left(\frac{n\pi}{a}x\right) e^{-i\omega t} \\ &= \left(\omega \hbar - \frac{\hbar^2 n^2 \pi^2}{2ma^2}\right) \sin\left(\frac{n\pi}{a}x\right) e^{-i\omega t} \\ &= \left(\omega \hbar - \frac{\hbar^2 n^2 \pi^2}{2ma^2}\right) \psi \\ V &= \omega \hbar - \frac{\hbar^2 n^2 \pi^2}{2ma^2} \end{aligned}$$

- (b) To find the expectation value $\langle x \rangle$, we observe that the term $\sin^2\left(\frac{n\pi}{a}x\right)$ is symmetric around $x = \frac{a}{2}$. By shifting the domain of $\psi(x)$ as $\psi(x) \rightarrow \psi\left(x - \frac{a}{2}\right)$, we redefine the domain as $-\frac{a}{2} \leq x \leq \frac{a}{2}$.

Under this transformation, the shifted wave function becomes an even function, while x (as a variable) is odd with respect to $x = 0$. As a result, the integrand for $\langle x' \rangle = \int_{-\frac{a}{2}}^{\frac{a}{2}} x' |\psi(x')|^2 dx'$ is an odd function over a symmetric interval, leading to:

$$\langle x' \rangle = 0.$$

To obtain $\langle x \rangle$ in the original coordinates, we perform the inverse transformation $x' = x - \frac{a}{2}$, which implies $x = x' + \frac{a}{2}$. Thus:

$$\langle x \rangle = \langle x' + \frac{a}{2} \rangle = \langle x' \rangle + \frac{a}{2} = 0 + \frac{a}{2} = \frac{a}{2}.$$

Therefore, the average position of the particle is:

$$\langle x \rangle = \frac{a}{2}.$$

3 Problem 3

Fourier transform:

$$\phi(k) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} \psi(x) e^{-ikx} dx$$

Inverse Fourier transform:

$$\psi(x) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} \phi(k) e^{ikx} dk$$

$$\langle p \rangle = \int_{-\infty}^{\infty} \psi^*(x) \left(-i\hbar \frac{d}{dx} \right) \psi(x) dx$$

$$\langle p \rangle = \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \phi^*(k') e^{-ik'x} dk' \left(-i\hbar \frac{d}{dx} \right) \int_{-\infty}^{\infty} \phi(k) e^{ikx} dk dx$$

$$-i\hbar \frac{d}{dx} e^{ikx} = \hbar k e^{ikx}$$

$$\langle p \rangle = \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \phi^*(k') \phi(k) \hbar k \left(\frac{1}{2\pi} \int_{-\infty}^{\infty} e^{i(k-k')x} dx \right) dk' dk$$

$$\frac{1}{2\pi} \int_{-\infty}^{\infty} e^{i(k-k')x} dx = \delta(k - k')$$

$$\langle p \rangle = \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \phi^*(k') \phi(k) \hbar k \delta(k - k') dk' dk$$

$$\therefore \langle p \rangle = \int_{-\infty}^{\infty} \hbar k |\phi(k)|^2 dk$$

where $|\phi(k)|^2$ represents the probability density in k -space and $\hbar k$ is the momentum associated with each k .

4 Problem 4

(a) If a system is in a pure energy eigenstate, its time evolution is given by:

$$\psi(x, t) = \psi(x, 0)e^{-iEt/\hbar}$$

$$\psi(x, t + T) = \psi(x, t)$$

$$\psi(x, 0)e^{-iE(t+T)/\hbar} = \psi(x, 0)e^{-iEt/\hbar}$$

$$\therefore e^{-iET/\hbar} = 1$$

$$\therefore \frac{ET}{\hbar} = 2\pi n, \quad n \in \mathbb{N}$$

$$T = \frac{2\pi\hbar}{E}n$$

(b)

$$E = \frac{2\pi\hbar}{T}n$$

(c) For a particle in an infinite square well

$$E_n = \frac{n^2\pi^2\hbar^2}{2mL^2}$$

$$T = \frac{2\pi\hbar}{E_n}n = \frac{4m}{n\pi\hbar}L^2$$

5 Problem 5

The wave function stays undisturbed (the same) after the sudden expansion

$$\psi_0 = \sqrt{\frac{2}{L}} \cos\left(\frac{\pi}{L}x\right), \quad -\frac{L}{2} \leq x \leq \frac{L}{2}$$

The new ground state:

$$\psi'_0 = \sqrt{\frac{1}{L}} \cos\left(\frac{\pi}{2L}x\right), \quad -L \leq x \leq L$$

$$\begin{aligned} \langle \psi | \psi'_0 \rangle &= \frac{\sqrt{2}}{L} \int_{-\frac{L}{2}}^{\frac{L}{2}} \cos\left(\frac{\pi}{2L}x\right) \cos\left(\frac{\pi}{L}x\right) dx \\ &= \frac{\sqrt{2}}{L} \int_{-\frac{L}{2}}^{\frac{L}{2}} \cos\left(\frac{\pi}{2L}x\right) [1 - 2 \sin^2\left(\frac{\pi}{2L}x\right)] dx \\ &= \frac{\sqrt{2}}{L} \left[\frac{2L}{\pi} \sin\left(\frac{\pi}{2L}x\right) - \frac{4L}{3\pi} \sin^3\left(\frac{\pi}{2L}x\right) \right]_{-\frac{L}{2}}^{\frac{L}{2}} \\ &= \frac{2\sqrt{2}}{\pi} \left[\sin\left(\frac{\pi}{2L}x\right) - \frac{2}{3} \sin^3\left(\frac{\pi}{2L}x\right) \right]_{-\frac{L}{2}}^{\frac{L}{2}} \\ &= \frac{8}{3\pi} \end{aligned}$$

$$P = |\langle \psi | \psi'_0 \rangle|^2 = \left| \frac{8}{3\pi} \right|^2 = \frac{64}{9\pi^2}$$

6 Problem 6

References

- [1] M.H. El-Deeb. [PEU-323 Assignments](#).