Kecture 23: Contitutive Theory notes on

raks dou't work

Logistics: - ItW6 complete yey?

- ItU7 shill out standing

- HW8 due Th

Last time: Enogg balance

Fret law of Thronodynamics

 $du = dQ + dU \qquad \int_{\Gamma}^{\Gamma} U [D] = Q[D] + W[D]$

not heating net working

Heating:

9=- K VT

Q[r] = SprdV - Jage q·ndV

Working

M[26] = D[26] - 9 K[66] = 2 0: 9 an

Strans bona: \$:0 9 = = (\(\tau \) + \(\tau \)

Today: Constitutive laws

relation between stress and strain

Constitutive Theory

Common constitutive laus:

New tou iau fluid:
$$\underline{\sigma} = p \underline{I} + y (\nabla y + \nabla x^{T})$$

(in compressible) $p = -\frac{1}{3} tr(\underline{\sigma}) y = viscosily \underline{v} = velocity$

Both derive from the same functional form
$$\underline{G}(\underline{A}) = \underline{C}\underline{A} = \lambda \operatorname{tr}(\underline{A}) + 2 \mu \operatorname{sym}(\underline{A})$$

4th a solur fecusos

liner electic solid: $\Delta = \nabla u$ Newboulen fluid: $\Delta = \nabla v$ $tr(\nabla u) = \nabla v$

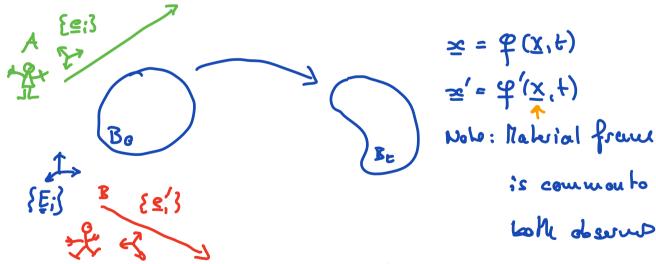
⇒ for livear dastic solid just substitute for Newtonian Sluid thre is a complication du to the incompassibility constraint. Why do constitution laws have this form?

Change lu obsolut

Constitutue lans cannot depend on the observer -> lecture ou Change of basis

⇒ parsive drauge in frame.

Change in observer -> active change of frame



Note: both ob sours

Change in obsessue cannot induce defermation

>> two reference frames must be related by

rigid body motion

$$x' = Q(t)x + c(t)$$

$$Z' = Q(t)q(x,+) + c(t)$$
Eulerian transformation

Q= robatien c= translation => constitution laws cannot depend on obsume

Effect on kinematic quantities:

$$\nabla \varphi = F \qquad \nabla \varphi' = F' = GF \\
 x' = \varphi(x,t) + c = G(t) \varphi(x,t) + c(t)$$
Right (anchy-Green Shain:

$$C' = F' F' = (FG') (GF) = F' G' GF = F' F = C$$
what about epatial tensors?

What about spatial tensors?

Axion of frame in difference

Fields are called frame indifferent or objective if for any superimposed rigid body motion z' = Q(t) z + c we have

scalar field:
$$\phi'(x',t) = \phi(x,t)$$

L'=QLGT+QQT => no L=√z is pobjective

that's why √no is not use

in outshifutive laws.

The non-objective term $\mathcal{Q} = \mathcal{Q} \mathcal{Q}^T$ \Rightarrow rigid body anguler velocity between the observers

Show $\mathcal{Q} = -\mathcal{Q}^T$ is show symatric $\mathcal{Q} = \mathcal{Q}^T = \mathcal{Z}^T =$

Neu-objechine part of $\underline{l} = \nabla_{xv}$ is skew-sym \Rightarrow talu sym (∇_{xv}) to be objective

$$\Rightarrow \underline{d} = \operatorname{sym}(\nabla_{\underline{x}}\underline{v}) = \frac{1}{2}(\nabla_{\underline{x}}\underline{v} + \nabla_{\underline{x}}\underline{v}^{T})$$

rate of strain tensor is objective

> wed in constitutive laws.

Objective functions

Fields: $\phi(z,t)$ eculor

veches veches

<u>S</u> (2,+)

cwor

field -> depends on z

Conshitutive functions are not fields but they take fields as in put inhual ewgg: μ(z,t) = û(ρ(z,t), θ(z,t)) output fiela Timput fields
conshibehing fuuch'n

word $f(cw): q(x,t) = \hat{q}(\Theta(x,t))$ q=-KQB

Couchy shess: $g(z,t) = \hat{g}(d(z,t))$

Isotropic functions

Functions that are frame invariant are called isofropic.

φ = scales fauc.
$$\hat{\omega}$$
 = vect. fauc. $\hat{\underline{\varsigma}}$ = tenser fauct.
θ = scales field \underline{v} = vect. field $\underline{\underline{\varsigma}}$ = tenser field

For two frames related by rigid body rotation & we have the following isotropic functions:

$$\frac{1}{2}(0|\mathbf{x}_{i}) = \frac{1}{6} \mathbf{s}(0|\mathbf{x}_{i}) = \frac{1}{6} \mathbf{s}(0|\mathbf{x}_{i})$$

Examples!

1)
$$\hat{\phi}(\underline{s}) = dut(\underline{s})$$

$$\hat{\phi}(\underline{s}') = \hat{\phi}(\underline{Q} \underline{s} \underline{G}^{T}) = dut(\underline{Q} \underline{s} \underline{G}^{T}) = dut(\underline{S})$$

$$= dut(\underline{G}) dut(\underline{S}) dut(\underline{G}^{T}) = dut(\underline{S})$$

$$\hat{\mathcal{L}}(\underline{\mathcal{L}},\underline{A}) = \underline{A} \times \\ \hat{\mathcal{L}}(\underline{\mathcal{L}},\underline{A}) = \underline{A} \times$$

Representation theorem for isotropic functions. Rivlin-Erichson Rep. Thu

$$\underline{G}(\underline{A}) = \alpha_o(\underline{T}_A)\underline{I} + \alpha_i(\underline{T}_A)\underline{A} + \alpha_z(\underline{T}_A)\underline{A}^z$$

where X_{a} , X_{c} , X_{c} are scaled functions of the set of invariants of A $T_{A}=\{T_{i}(A), T_{i}(A)\}$

· G is sque if A is sque.

- & is isolopic

$$G(GAG^{\dagger}) = \alpha_{0}I + \alpha_{1}GAG^{\dagger} + \alpha_{2}GAG^{\dagger}GAG^{\dagger}$$

$$= \alpha_{0}QQ^{\dagger} + \alpha_{1}GAQ^{\dagger} + \alpha_{2}GAG^{\dagger}GAG^{\dagger}$$

$$= G(\alpha_{0}I + \alpha_{1}A^{\dagger} + \alpha_{2}A^{\dagger})GT$$

$$= G(A)QT$$

only true if do, d, , d, depend only on invariants of A