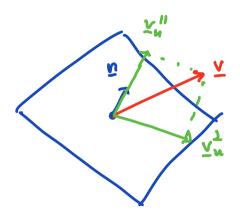
Lecture 6: Hore on stress , check points on cannon Logistics: - HW1 (5/7) please submit - Ituz will be posted =>mahe use of office hours if HW is not clear Last time: - infinitesimal force balance f-ya $\frac{1}{A} \oint \underline{t}_n dA = \underline{0}$ (Vol. terms vanish) - Canchy's postulate - 3rd law (Achien-Roachion) tu (-11,2) = -t(1,2) - Cauchy's theorem | == = (x) n $\underline{\underline{\underline{c}}} = Cauchy stress$ $\underline{\underline{c}}_{33} = \underline{\underline{t}}_{3}(\underline{\underline{c}}_{3},\underline{\underline{x}}) \\ \underline{\underline{c}}_{3} \uparrow$ <u>©</u> = &; ⊆; ⊗ ⊆; 6; = t; (e;, x) 6; = i-th component of trachion on jtk coordinate plans Today: Projections -> shear & normal elem Simple etress states Spherical + deviatoric stress

Projection Tensors



$$\overline{\Lambda}_{1}^{n} = \overline{L}_{1}^{n} \quad \overline{\Lambda}$$

$$\overline{\Lambda}_{1}^{n} = \overline{L}_{1}^{n} \quad \overline{\Lambda}$$

$$\bar{\Lambda} = \bar{\Lambda}_{\parallel}^{n} + \bar{\Lambda}_{\overline{\Lambda}}^{n}$$

dot product

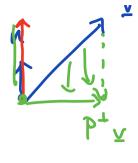
$$\vec{\lambda}_{\parallel}^{\mathsf{u}} = (\vec{\nu} \cdot \vec{\lambda}) \vec{\nu}$$

$$\bar{\Lambda}_{T}^{n} = \bar{\Lambda} - \bar{\Lambda}_{\parallel}^{n}$$

Use dyadic product: (asb) c = (b.e) q (neh) v n v h

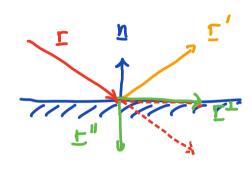
Projection keepers

$$\frac{1}{2} = \overline{N} \otimes \overline{N}$$





Reflections

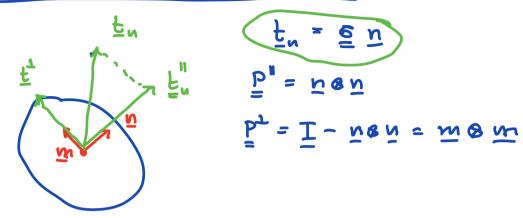




$$= \begin{bmatrix} -\frac{1}{2} - \frac{1}{2} & -\frac{1}{2} \\ -\frac{1}{2} - \frac{1}{2} & -\frac{1}{2} \end{bmatrix} = \begin{bmatrix} -\frac{1}{2} - \frac{1}{2} & -\frac{1}{2} \\ -\frac{1}{2} - \frac{1}{2} & -\frac{1}{2} \end{bmatrix} = \begin{bmatrix} -\frac{1}{2} - \frac{1}{2} & -\frac{1}{2} \\ -\frac{1}{2} - \frac{1}{2} & -\frac{1}{2} \end{bmatrix} = \begin{bmatrix} -\frac{1}{2} - \frac{1}{2} & -\frac{1}{2} \\ -\frac{1}{2} - \frac{1}{2} & -\frac{1}{2} \end{bmatrix} = \begin{bmatrix} -\frac{1}{2} - \frac{1}{2} & -\frac{1}{2} \\ -\frac{1}{2} & -\frac{1}{2} & -\frac{1}{2} \end{bmatrix} = \begin{bmatrix} -\frac{1}{2} - \frac{1}{2} & -\frac{1}{2} \\ -\frac{1}{2} & -\frac{1}{2} & -\frac{1}{2} \end{bmatrix} = \begin{bmatrix} -\frac{1}{2} - \frac{1}{2} & -\frac{1}{2} \\ -\frac{1}{2} & -\frac{1}{2} & -\frac{1}{2} \end{bmatrix} = \begin{bmatrix} -\frac{1}{2} - \frac{1}{2} & -\frac{1}{2} \\ -\frac{1}{2} & -\frac{1}{2} & -\frac{1}{2} \end{bmatrix} = \begin{bmatrix} -\frac{1}{2} - \frac{1}{2} & -\frac{1}{2} \\ -\frac{1}{2} & -\frac{1}{2} & -\frac{1}{2} \end{bmatrix} = \begin{bmatrix} -\frac{1}{2} - \frac{1}{2} & -\frac{1}{2} \\ -\frac{1}{2} & -\frac{1}{2} & -\frac{1}{2} \end{bmatrix} = \begin{bmatrix} -\frac{1}{2} & -\frac{1}{2} & -\frac{1}{2} \\ -\frac{1}{2} & -\frac{1}{2} & -\frac{1}{2} \end{bmatrix} = \begin{bmatrix} -\frac{1}{2} & -\frac{1}{2} & -\frac{1}{2} \\ -\frac{1}{2} & -\frac{1}{2} & -\frac{1}{2} \end{bmatrix} = \begin{bmatrix} -\frac{1}{2} & -\frac{1}{2} & -\frac{1}{2} \\ -\frac{1}{2} & -\frac{1}{2} & -\frac{1}{2} \end{bmatrix} = \begin{bmatrix} -\frac{1}{2} & -\frac{1}{2} & -\frac{1}{2} \\ -\frac{1}{2} & -\frac{1}{2} & -\frac{1}{2} \end{bmatrix} = \begin{bmatrix} -\frac{1}{2} & -\frac{1}{2} & -\frac{1}{2} \\ -\frac{1}{2} & -\frac{1}{2} & -\frac{1}{2} \end{bmatrix} = \begin{bmatrix} -\frac{1}{2} & -\frac{1}{2} & -\frac{1}{2} \\ -\frac{1}{2} & -\frac{1}{2} & -\frac{1}{2} \end{bmatrix} = \begin{bmatrix} -\frac{1}{2} & -\frac{1}{2} & -\frac{1}{2} \\ -\frac{1}{2} & -\frac{1}{2} & -\frac{1}{2} \end{bmatrix} = \begin{bmatrix} -\frac{1}{2} & -\frac{1}{2} & -\frac{1}{2} \\ -\frac{1}{2} & -\frac{1}{2} & -\frac{1}{2} \end{bmatrix} = \begin{bmatrix} -\frac{1}{2} & -\frac{1}{2} & -\frac{1}{2} \\ -\frac{1}{2} & -\frac{1}{2} & -\frac{1}{2} \end{bmatrix} = \begin{bmatrix} -\frac{1}{2} & -\frac{1}{2} & -\frac{1}{2} \\ -\frac{1}{2} & -\frac{1}{2} & -\frac{1}{2} \end{bmatrix} = \begin{bmatrix} -\frac{1}{2} & -\frac{1}{2} & -\frac{1}{2} \\ -\frac{1}{2} & -\frac{1}{2} & -\frac{1}{2} \end{bmatrix} = \begin{bmatrix} -\frac{1}{2} & -\frac{1}{2} & -\frac{1}{2} \\ -\frac{1}{2} & -\frac{1}{2} & -\frac{1}{2} \end{bmatrix} = \begin{bmatrix} -\frac{1}{2} & -\frac{1}{2} & -\frac{1}{2} \\ -\frac{1}{2} & -\frac{1}{2} & -\frac{1}{2} \end{bmatrix} = \begin{bmatrix} -\frac{1}{2} & -\frac{1}{2} & -\frac{1}{2} \\ -\frac{1}{2} & -\frac{1}{2} & -\frac{1}{2} \end{bmatrix} = \begin{bmatrix} -\frac{1}{2} & -\frac{1}{2} & -\frac{1}{2} \\ -\frac{1}{2} & -\frac{1}{2} \end{bmatrix} = \begin{bmatrix} -\frac{1}{2} & -\frac{1}{2} & -\frac{1}{2} \\ -\frac{1}{2} & -\frac{1}{2} & -\frac{1}{2} \end{bmatrix} = \begin{bmatrix} -\frac{1}{2} & -\frac{1}{2} & -\frac{1}{2} \\ -\frac{1}{2} & -\frac{1}{2} \end{bmatrix} = \begin{bmatrix} -\frac{1}{2} & -\frac{1}{2} & -\frac{1}{2} \end{bmatrix} = \begin{bmatrix} -\frac{1}{2} & -\frac{1}{2} & -\frac{1}{2} \end{bmatrix} = \begin{bmatrix} -$$

Reflection tensor: R = I - 2 nou

Normal and Shear Stresses



normal stan:
$$\underline{\underline{t}}_{n}^{"} = \underline{\underline{P}}^{"}\underline{\underline{t}}_{n} = (\underline{\underline{n}}\underline{\underline{s}}\underline{\underline{n}})\underline{\underline{t}}_{n}$$

$$= (\underline{\underline{n}}\underline{\underline{t}}\underline{\underline{n}})\underline{\underline{n}} = \underline{\underline{s}}\underline{\underline{n}}\underline{\underline{s}}\underline{\underline{n}}$$

on is may. of normal steers

shear stres: $\underline{t}_{n} = \underline{P}^{t}\underline{t}_{n} = (\underline{m} \cdot \underline{u}) \underline{t}_{n} = (\underline{m} \cdot \underline{t}_{n}) \underline{u}$

mag. of shew shew: I = motu = mogn = mogn

En >0 tousile shers

en <0 combusins ston

From geometry: $|\underline{t}_n|^2 = \sigma_n^2 + T^2$

Simple states of stress I) Hydrostatic sters

$$F_{\parallel} = F^{u} \implies F_{T} = 0$$

$$= (\overline{N}, \overline{N}) \overline{N} (-b) = -b \overline{N}$$

$$F_{\parallel} = \overline{b}_{\parallel} \overline{p}^{n} = (\overline{N} \otimes \overline{N}) (-b \overline{N}) = -b (\overline{N} \otimes \overline{N}) \overline{N}$$

$$\overline{F}_{\parallel} = \overline{Q} \overline{N} = (-b \overline{I}) \overline{N} = -b \overline{N}$$

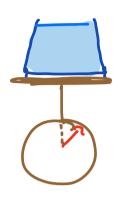
normal stress: $a_n = -p$ on all planes show shows: $a_n = -p$

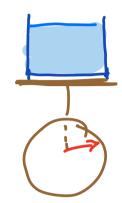
Paseal's law:

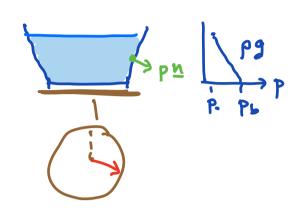
The pressure in a fluid at rest is independent of the direction of a surface.

Hydrostatic peradox:

force ou bouse is some







II Uniaxial stress

$$\bar{\chi}$$
 $\bar{\Gamma}$ $\bar{\Gamma}^{N} = g(\underline{x}, \bar{n}) (\bar{x}, \underline{x})$
 $\bar{\Gamma}^{N} = \bar{g} \bar{\nu} = g(\underline{x} \circ x) \bar{n} = g(\underline{x} \cdot \bar{\nu}) \bar{\lambda}$

$$\underline{t}_{n} = \underline{t} \underbrace{s} \underbrace{M}$$

$$\underline{t}_{n} = s \underbrace{(\underline{T}_{n})}_{n} \underbrace{M}_{n} = 0$$

traction is always pasallel to x

III Pure shear stren

two relications & and & X. 7 = 0

$$\begin{aligned}
\bar{n} &= \underline{L} : & = \underline{r} = \underline{r} \\
\bar{n} &= \underline{L} : & = \underline{r} \\
\bar{n} &= \underline{L} : & = \underline{r} \\
\bar{n} &= \underline{L} : & = \underline{L} \\
\bar{n} &= \underline{$$

IV Plane stress

If there exist a pair of orthog. vectors

The stand of the standard of orthog. vectors

If there exist a pair of orthog. vectors

If there exist a pair of orthog. vectors

I a frame of o

a: Is uniaxial stema plane streno?

{I, e, e, e,}

I'ez=0 I'ez=0

dou't need to specify

other basis vectors

(not unique)

$$\delta_{ij} = \delta(e_1 \cdot e_j)(e_i \cdot e_1)$$

$$c_{11} = c \quad (e_1 \cdot e_1) \quad (e_1 \cdot e_1) = c$$

$$c_{12} = c \quad (e_1 \cdot e_1) \quad (e_1 \cdot e_1) = c$$

$$c_{13} = c \quad (e_1 \cdot e_1) \quad (e_1 \cdot e_1) = c$$

Spherical & Deviatoric Stress tensers

sphrical alread tensor:
$$\underline{\underline{\sigma}}_s = -p \underline{\underline{I}} \quad p = -\frac{1}{3} \text{ tr}(\underline{\underline{\sigma}})$$
deviatoric structures: $\underline{\underline{\sigma}}_b = \underline{\underline{\sigma}} - \underline{\underline{\sigma}}_s$

$$= \underline{\underline{\sigma}} + p \underline{\underline{I}}$$

The pressure $p = -\frac{1}{3} \operatorname{tr}(\bar{z})$ is mean normal traction.

Sperical shers -> volumetric changes
Deviatoric slæss -> changes shape of booky