

Second-order Tensors

Here we are interested in second-order tensors

Linear operators : $\underline{\underline{A}} \underline{u} = \underline{\underline{A}} \underline{u}$

maps vector $\underline{u} \in \mathcal{V}$ into vector $\underline{\underline{A}} \underline{u} \in \mathcal{V}$

Linearity requires that

$$1) \quad \underline{\underline{A}}(\underline{u} + \underline{v}) = \underline{\underline{A}}\underline{u} + \underline{\underline{A}}\underline{v} \quad \text{for all } \underline{u}, \underline{v} \in \mathcal{V}$$

$$2) \quad \underline{\underline{A}}(\alpha \underline{v}) = \alpha \underline{\underline{A}}\underline{v} \quad \text{for all } \alpha \in \mathbb{R} \text{ and } \underline{v} \in \mathcal{V}$$

Example: $\underline{\underline{A}}$ maps every $\underline{v} \in \mathcal{V}$ into $\underline{n} \neq \underline{0} \in \mathcal{V}$.

Is $\underline{\underline{A}}$ a tensor?

Consider $\underline{u}, \underline{v}, \underline{w} \in \mathcal{V}$

$$\underline{w} = \underline{u} + \underline{v}$$

$$\underline{\underline{A}}(\underline{u} + \underline{v}) \stackrel{?}{=} \underline{\underline{A}}\underline{u} + \underline{\underline{A}}\underline{v}$$

$$\underline{\underline{A}}\underline{w} \stackrel{?}{=} \underline{\underline{A}}\underline{u} + \underline{\underline{A}}\underline{v}$$

$$\underline{n} \neq \underline{n} + \underline{n}$$

$\Rightarrow \underline{\underline{A}}$ is not a tensor, because it is not linear

Tensor algebra

For all $\underline{v} \in \mathcal{V}$ we define

$$1) (\alpha \underline{A}) \underline{v} = \underline{A}(\alpha \underline{v}) \quad \text{scalar multiplication}$$

$$2) (\underline{A} + \underline{B}) \underline{v} = \underline{A}\underline{v} + \underline{B}\underline{v} \quad \text{tensor sum}$$

$$3) (\underline{A} \underline{B}) \underline{v} = \underline{A}(\underline{B}\underline{v}) \quad \text{tensor product}$$

Note there is also a scalar product introduced later.

The set of all 2nd-order tensors \mathcal{V}^2 is a vector space

$$1) \alpha \underline{A} \in \mathcal{V}^2 \quad \text{for all } \underline{A} \in \mathcal{V} \text{ and } \alpha \in \mathbb{R}$$

$$2) \underline{A} + \underline{B} \in \mathcal{V}^2 \quad \text{for all } \underline{A}, \underline{B} \in \mathcal{V}^2$$

$$3) \underline{A} \underline{B} \in \mathcal{V}^2 \quad \text{for all } \underline{A}, \underline{B} \in \mathcal{V}^2$$

Any of these operations will produce another second-order tensor.

Q: What is a basis for \mathcal{V}^2 ?

Two tensors $\underline{\underline{A}}$ and $\underline{\underline{B}}$ are equal if

$$\underline{\underline{A}} \underline{v} = \underline{\underline{B}} \underline{v} \quad \text{for all } \underline{v} \in V$$

Zero tensor: $\underline{\underline{0}} \underline{v} = \underline{0} \quad \text{for all } \underline{v} \in V$

Identity tensor: $\underline{\underline{I}} \underline{v} = \underline{v} \quad \text{for all } \underline{v} \in V$

Representation of a tensor

In a frame $\{\underline{e}_i\}$ a second order tensor

$\underline{\underline{S}}$ is represented by nine numbers

$$S_{ij} = \underline{e}_i \cdot \underline{\underline{S}} \underline{e}_j$$

Consider $\underline{v} = \underline{\underline{S}} \underline{u}$ where $\underline{v} = v_i \underline{e}_i$, $\underline{u} = u_j \underline{e}_j$

$$v_i \underline{e}_i = \underline{\underline{S}} (u_j \underline{e}_j) = \underline{\underline{S}} \underline{e}_j u_j$$

multiply by \underline{e}_k from left

$$v_i \underline{e}_k \cdot \underline{e}_i = \underline{e}_k \cdot \underline{\underline{S}} \underline{e}_j u_j$$

$$v_i \delta_{ki} = \underline{e}_k \cdot \underline{\underline{S}} \underline{e}_j u_j$$

$$v_i = (\underline{e}_i \cdot \underline{\underline{S}} \underline{e}_j) u_j$$

$$v_i = S_{ij} u_j$$

Matrix representation of tensor in $\{\underline{e}_i\}$

$$[\underline{\underline{S}}] = \begin{bmatrix} S_{11} & S_{12} & S_{13} \\ S_{21} & S_{22} & S_{23} \\ S_{31} & S_{32} & S_{33} \end{bmatrix} \in \mathbb{R}^3 \times \mathbb{R}^3$$

Note that $[\underline{\underline{S}}]_{ij} = S_{ij}$

Dyadic Product

The dyadic product of two vectors \underline{a} and \underline{b} is the 2nd-order tensor $\underline{a} \otimes \underline{b}$ defined by

$$(\underline{a} \otimes \underline{b}) \underline{v} = (\underline{b} \cdot \underline{v}) \underline{a} \quad \text{for all } \underline{v} \in V$$

This has the form: $\underline{\underline{A}} \underline{v} = \alpha \underline{a}$

in components: $A_{ij} v_j = \alpha a_i$

$$\alpha = \underline{b} \cdot \underline{v} = b_j v_j$$

$$A_{ij} = [\underline{a} \otimes \underline{b}]_{ij}$$

$$\Rightarrow [\underline{a} \otimes \underline{b}]_{ij} v_j = b_j v_j a_i$$

$$[\underline{a} \otimes \underline{b}]_{ij} v_j = (a_i b_j) v_j$$

$$\Rightarrow [\underline{a} \otimes \underline{b}]_{ij} = a_i b_j$$

So that

$$[\underline{a} \otimes \underline{b}] = \begin{bmatrix} a_1 b_1 & a_1 b_2 & a_1 b_3 \\ a_2 b_1 & a_2 b_2 & a_2 b_3 \\ a_3 b_1 & a_3 b_2 & a_3 b_3 \end{bmatrix} = \underline{a} \underline{b}^T$$

Linearity of dyadic product:

for scalars $\alpha, \beta \in \mathbb{R}$ and vectors $\underline{a}, \underline{b}, \underline{v}, \underline{w} \in V$

$$(\underline{a} \otimes \underline{b}) (\alpha \underline{v} + \beta \underline{w}) = \alpha (\underline{a} \otimes \underline{b}) \underline{v} + \beta (\underline{a} \otimes \underline{b}) \underline{w}$$

The product of two dyadic products

$$(\underline{a} \otimes \underline{b})(\underline{c} \otimes \underline{d}) = (\underline{b} \cdot \underline{c}) \underline{a} \otimes \underline{d} \Rightarrow \text{HW2}$$

needed for tensor product.

Basis for V^2

Given any frame $\{e_i\}$ the nine dyadic products $\{e_i \otimes e_j\}$ form a basis for V^2 .

Any second-order tensor $\underline{\underline{S}}$ can be written as linear combination

$$\underline{\underline{S}} = S_{ij} e_i \otimes e_j$$

where $S_{ij} = e_i \cdot \underline{\underline{S}} e_j$

Consider $\underline{v} = \underline{\underline{S}} \underline{u}$ with $\underline{v} = v_i e_i$, $\underline{u} = u_k e_k$

$$\begin{aligned} v_i e_i &= S_{ij} (e_i \otimes e_j) (u_k e_k) \\ &= S_{ij} u_k (e_i \otimes e_j) \cdot e_k \quad \text{apply def. of dyadic} \\ &= S_{ij} u_k (e_j \cdot e_k) e_i = S_{ij} u_k \delta_{kj} e_i \end{aligned}$$

$$\begin{aligned} v_i e_i &= S_{ij} u_j e_i \\ v_i &= S_{ij} u_j \end{aligned}$$

used often

Tensor algebra in components

Addition : $\underline{H} = \underline{S} + \underline{T}$

$$\begin{aligned} H_{ij} e_i \otimes e_j &= S_{ij} e_i \otimes e_j + T_{ij} e_i \otimes e_j \\ &= (S_{ij} + T_{ij}) e_i \otimes e_j \\ H_{ij} &= S_{ij} + T_{ij} \end{aligned}$$

Scalar multiplication : $\underline{H} = \alpha \underline{S} \Rightarrow H_{ij} = \alpha S_{ij}$

Product : $\underline{H} = \underline{S} \underline{T}$

$$\begin{aligned} \underline{H} &= S_{ij} (e_i \otimes e_j) T_{kl} (e_k \otimes e_l) \\ &= S_{ij} T_{kl} \underbrace{(e_i \otimes e_j)(e_k \otimes e_l)}_{\text{product of two dyads}} \\ &= S_{ij} T_{kl} (e_j \otimes e_k) e_i \otimes e_l \\ &\quad \delta_{jk} \\ &= S_{ij} T_{jl} e_i \otimes e_l \end{aligned}$$

$$H_{il} e_i \otimes e_l = S_{ij} T_{jl} e_i \otimes e_l$$

\Rightarrow

$$H_{il} = S_{ij} T_{jl}$$

note the dummy j !

Transpose of a tensor

To any $\underline{\underline{S}} \in V^2$ we associate a transpose $\underline{\underline{S}}^T \in V^2$ the unique tensor such that

$$\underline{\underline{S}}\underline{u} \cdot \underline{v} = \underline{u} \cdot \underline{\underline{S}}^T \underline{v} \quad \text{for all } \underline{u}, \underline{v} \in V$$

This implies that $S_{ij}^T = S_{ji}$ as follows

$$(S_{ij} u_j e_i) \cdot (v_l e_l) = (u_k e_k) \cdot (S_{ij}^T v_j e_i)$$

$$S_{ij} u_j v_l (e_i \cdot e_l) = S_{ij}^T v_j u_k (e_k \cdot e_i)$$

$$S_{ij} u_j v_l \delta_{il} = S_{ij}^T v_j u_k \delta_{ki}$$

$$S_{ij} u_j v_i = S_{ij}^T v_j u_i \quad \begin{matrix} \text{rename indices} \\ i \rightarrow j \text{ on rhs} \end{matrix}$$

$$S_{ij} u_j v_i = S_{ji}^T u_j v_i$$

$$\Rightarrow S_{ij} = S_{ji}^T \quad \checkmark$$

Properties of transpose!

$$(\underline{\underline{A}}^T)^T = \underline{\underline{A}}$$

$$(\underline{\underline{A}} \underline{\underline{B}})^T = \underline{\underline{B}}^T \underline{\underline{A}}^T$$

$$(\underline{u} \otimes \underline{v})^T = \underline{v} \otimes \underline{u}$$

$\underline{\underline{S}}$ is symmetric if $\underline{\underline{S}} = \underline{\underline{S}}^T$ $S_{ij} = S_{ji}$

$\underline{\underline{S}}$ is skew-symmetric if $\underline{\underline{S}} = -\underline{\underline{S}}^T$ $S_{ij} = -S_{ji}$

Symmetric-Skew decomposition:

Any tensor $\underline{\underline{S}} \in V^2$ can be written as

$$\underline{\underline{S}} = \underline{\underline{E}} + \underline{\underline{W}}$$

$$\underline{\underline{E}} = \frac{1}{2} (\underline{\underline{S}} + \underline{\underline{S}}^T)$$

$$\underline{\underline{W}} = \frac{1}{2} (\underline{\underline{S}} - \underline{\underline{S}}^T)$$

$$\underline{\underline{E}} = \underline{\underline{E}}^T$$

$$\underline{\underline{W}} = -\underline{\underline{W}}^T$$

Note: $\underline{\underline{W}} = \begin{bmatrix} 0 & W_{12} & W_{13} \\ -W_{12} & 0 & W_{23} \\ -W_{13} & -W_{23} & 0 \end{bmatrix}$ only 3 indep. comp.

→ can be related to an axial vector \underline{w}

$$\underline{\underline{W}} \underline{v} = \underline{w} \times \underline{v} \quad \text{for all } \underline{v} \in V$$

Relation:

$$W_{ij} = -\epsilon_{ijk} w_k$$

$$w_k = -\frac{1}{2} \epsilon_{ijk} W_{ij}$$

$$\underline{\underline{W}} = \begin{bmatrix} 0 & -w_3 & w_2 \\ w_3 & 0 & -w_1 \\ -w_2 & w_1 & 0 \end{bmatrix}$$

Trace of a tensor

We define the trace of a dyad as

$$\text{tr}(\underline{a} \otimes \underline{b}) = \underline{a} \cdot \underline{b} = a_i b_i$$

this implies that

$$\boxed{\text{tr}(\underline{\underline{A}}) = A_{ii} = A_{11} + A_{22} + A_{33}}$$

as follows $\text{tr}(A_{ij} e_i \otimes e_j) = A_{ij} \text{tr}(e_i \otimes e_j)$
 $= A_{ij} S_{ij} = A_{ii}$

Properties: $\text{tr}(\underline{\underline{A}}^T) = \text{tr}(\underline{\underline{A}})$

$$\text{tr}(\underline{\underline{A}} \underline{\underline{B}}) = \text{tr}(\underline{\underline{B}} \underline{\underline{A}})$$

$$\text{tr}(\underline{\underline{A}} + \underline{\underline{B}}) = \text{tr}(\underline{\underline{A}}) + \text{tr}(\underline{\underline{B}})$$

$$\text{tr}(\alpha \underline{\underline{A}}) = \alpha \text{tr}(\underline{\underline{A}})$$

Decomposition: $\boxed{\underline{\underline{A}} = \alpha \underline{\underline{I}} + \text{dev } \underline{\underline{A}}}$

Spherical tensor: $\alpha \underline{\underline{I}}$ where $\alpha = \frac{1}{3} \text{tr}(\underline{\underline{A}})$

Deviatoric tensor: $\text{dev } \underline{\underline{A}} = \underline{\underline{A}} - \alpha \underline{\underline{I}}$

$$\text{tr}(\text{dev } \underline{\underline{A}}) = 0$$

Tensor scalar product (Contraction)

analogous to scalar product of vectors

$$\underline{\underline{A}} : \underline{\underline{B}} = \text{tr}(\underline{\underline{A}}^T \underline{\underline{B}}) = A_{ij} B_{ij}$$

scalar ∇

explicitly:

$$\underline{\underline{A}} : \underline{\underline{B}} = \sum_{i=1}^3 \sum_{j=1}^3 A_{ij} B_{ij} = A_{11} B_{11} + A_{12} B_{12} + A_{13} B_{13} + \dots$$

$$A_{21} B_{21} + A_{22} B_{22} + A_{23} B_{23} + \dots$$

$$A_{31} B_{31} + A_{32} B_{32} + A_{33} B_{33}$$

The index expression is derived as follows

$$\underline{\underline{A}} : \underline{\underline{B}} = \text{tr}(\underline{\underline{A}}^T \underline{\underline{B}}) =$$

$$\underline{\underline{A}}^T \underline{\underline{B}} = (A_{ji} e_i \otimes e_j)(B_{kl} e_k \otimes e_l)$$

$$= A_{ji} B_{kl} (e_i \otimes e_j)(e_k \otimes e_l)$$

$$= A_{ji} B_{kl} (\delta_{jk} \cdot e_i \otimes e_k)(e_i \otimes e_l) = A_{ji} B_{kl} \delta_{jk} e_i \otimes e_l$$

$$= A_{ji} B_{jl} e_i \otimes e_l$$

$$\text{tr}(\underline{\underline{A}}^T \underline{\underline{B}}) = A_{ji} B_{jl} \delta_{il} = A_{ji} B_{ji} = A_{ij} B_{ij} \quad \checkmark$$

Properties: 1) $\underline{\underline{A}} : \underline{\underline{B}} = \underline{\underline{B}} : \underline{\underline{A}}$

$$2) (\underline{a} \otimes \underline{b}) : (\underline{c} \otimes \underline{d}) = (\underline{a} \cdot \underline{c})(\underline{b} \cdot \underline{d})$$

First follows from prop. of trace

$$\underline{\underline{A}} : \underline{\underline{B}} = \text{tr}(\underline{\underline{A}}^T \underline{\underline{B}}) = \text{tr}((\underline{\underline{A}}^T \underline{\underline{B}})^T) = \text{tr}(\underline{\underline{B}}^T \underline{\underline{A}}) = \underline{\underline{B}} : \underline{\underline{A}}$$

Second property $[\underline{a} \otimes \underline{b}]_{ij} = a_i b_j$

$$\begin{aligned} (\underline{a} \otimes \underline{b}) : (\underline{c} \otimes \underline{d}) &= [\underline{a} \otimes \underline{b}]_{ij} [\underline{c} \otimes \underline{d}]_{ij} = a_i b_j c_i d_j \\ &= a_i c_i b_j d_j \\ &= (\underline{a} \cdot \underline{c})(\underline{b} \cdot \underline{d}) \end{aligned}$$

A common norm for tensors is

$$|\underline{\underline{A}}| = \sqrt{\underline{\underline{A}} : \underline{\underline{A}}^T} = \sqrt{A_{ij} A_{ij}^T} \geq 0$$

Note: Tensor scalar product will be important to express the work done during deformation.

For example the shear heating in glaciology.

Determinant and Inverse

The determinant of $\underline{\underline{A}} \in \mathbb{V}^2$ is the scalar

$$\det(\underline{\underline{A}}) = \det \begin{vmatrix} A_{11} & A_{12} & A_{13} \\ A_{21} & A_{22} & A_{23} \\ A_{31} & A_{32} & A_{33} \end{vmatrix} = \epsilon_{ijk} [\underline{\underline{A}}]_{i1} [\underline{\underline{A}}]_{j2} [\underline{\underline{A}}]_{k3}$$

where $[\underline{\underline{A}}]_{i1}, [\underline{\underline{A}}]_{j2}, [\underline{\underline{A}}]_{k3}$ are the columns of $[\underline{\underline{A}}]$

Properties: $\det(\underline{\underline{AB}}) = \det(\underline{\underline{A}}) \det(\underline{\underline{B}})$

$$\det(\underline{\underline{A}}^T) = \det(\underline{\underline{A}})$$

$$\det(\alpha \underline{\underline{A}}) = \alpha^n \det(\underline{\underline{A}}) \quad (\underline{\underline{A}} \text{ is } n \times n)$$

$\underline{\underline{A}}$ is singular if $\det \underline{\underline{A}} = 0$.

If $\det \underline{\underline{A}} \neq 0$ then the inverse $\underline{\underline{A}}^{-1}$ exists

$$\underline{\underline{A}}^{-1} \underline{\underline{A}} = \underline{\underline{A}} \underline{\underline{A}}^{-1} = \underline{\underline{I}}$$

$$\text{Properties: } (\underline{\underline{A}} \underline{\underline{B}})^{-1} = \underline{\underline{B}}^{-1} \underline{\underline{A}}^{-1}$$

$$(\underline{\underline{A}}^{-1})^{-1} = \underline{\underline{A}}$$

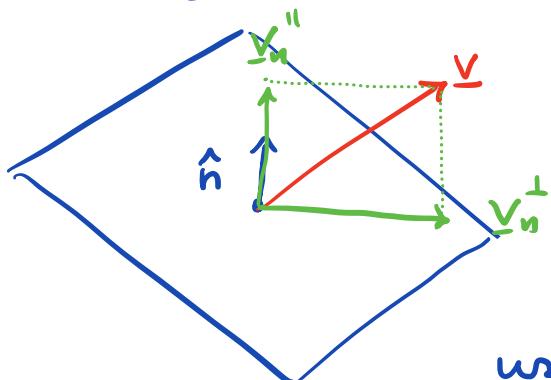
$$(\underline{\underline{A}}^{-1})^T = (\underline{\underline{A}}^T)^{-1}$$

$$(\alpha \underline{\underline{A}})^{-1} = \frac{1}{\alpha} \underline{\underline{A}}^{-1}$$

$$\det(\underline{\underline{A}}^{-1}) = \det(\underline{\underline{A}})^{-1} = \frac{1}{\det(\underline{\underline{A}})}$$

Projection & Reflection tensors

commonly used to partition forces on a surface.



$$\underline{\underline{v}} = \underline{\underline{v}}_n^{\parallel} + \underline{\underline{v}}_n^{\perp}$$

$$\underline{\underline{v}}_n^{\parallel} = (\underline{\underline{v}} \cdot \hat{\underline{\underline{n}}}) \hat{\underline{\underline{n}}}$$

$$\underline{\underline{v}}_n^{\perp} = \underline{\underline{v}} - \underline{\underline{v}}_n^{\parallel}$$

use dyadic property

$$\underline{\underline{v}}_n^{\parallel} = (\underline{\underline{v}} \cdot \hat{\underline{\underline{n}}}) \hat{\underline{\underline{n}}} = (\hat{\underline{\underline{n}}} \otimes \hat{\underline{\underline{n}}}) \underline{\underline{v}} = \underline{\underline{P}}_n^{\parallel} \underline{\underline{v}}$$

$$\underline{\underline{v}}_n^{\perp} = \underline{\underline{v}} - (\hat{\underline{\underline{n}}} \otimes \hat{\underline{\underline{n}}}) \underline{\underline{v}} = (\underline{\underline{I}} - \hat{\underline{\underline{n}}} \otimes \hat{\underline{\underline{n}}}) \underline{\underline{v}} = \underline{\underline{P}}_n^{\perp} \underline{\underline{v}}$$

$$\underline{\underline{P}}_n^{\parallel} = \hat{\underline{\underline{n}}} \otimes \hat{\underline{\underline{n}}}$$

$$\underline{\underline{P}}_n^{\perp} = \underline{\underline{I}} - \hat{\underline{\underline{n}}} \otimes \hat{\underline{\underline{n}}}$$

Properties:

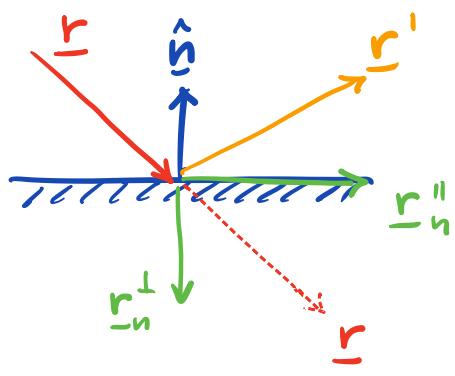
$$\underline{\underline{P}} = \underline{\underline{P}}^T$$

$$\underline{\underline{P}}^2 = \underline{\underline{P}}$$

$$\underline{\underline{P}}'' + \underline{\underline{P}}^\perp = \underline{\underline{I}}$$

$$\underline{\underline{P}}'' \underline{\underline{P}}^\perp = \underline{\underline{0}}$$
symmetric

Reflections



incoming: $\underline{\underline{r}} = \underline{\underline{r}}_n'' + \underline{\underline{r}}_n^\perp$

reflected: $\underline{\underline{r}}' = \underline{\underline{r}}_n'' - \underline{\underline{r}}_n^\perp$

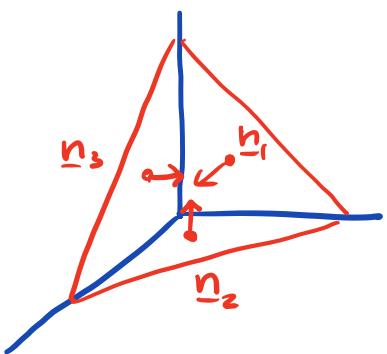
$$\underline{\underline{r}}' = (\underline{\underline{P}}_n'' - \underline{\underline{P}}_n^\perp) \underline{\underline{r}}$$

$$\underline{\underline{r}}' = (\underline{\underline{I}} - 2 \hat{n} \otimes \hat{n}) \underline{\underline{r}}$$

$$\underline{\underline{r}}' = \underline{\underline{R}}_n \underline{\underline{r}}$$

Reflection tensor: $\underline{\underline{R}}_n = \underline{\underline{I}} - 2 \hat{n} \otimes \hat{n}$

Example : Corner reflector



Inverts the direction of any ray that reflects off all three surfaces.

Direction of triply reflected ray:

$$\underline{r}''' = \underline{R}_{n_1} \underline{R}_{n_2} \underline{R}_{n_3} \underline{r}$$

Show that $\underline{r}''' = -\underline{r}$!