

# Backflow in Neural Quantum States

Slater–Jastrow–Backflow Wavefunctions and Links to Transformers

MHJ

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# Slater–Jastrow (SJ) variational ansatz

For  $N$  fermions with coordinates  $\mathbf{R} = (\mathbf{r}_1, \dots, \mathbf{r}_N)$ , a standard correlated ansatz is the Slater–Jastrow form

$$\Psi_{\text{SJ}}(\mathbf{R}) = \det[\phi_k(\mathbf{r}_i)] e^{J(\mathbf{R})}.$$

- $\det[\phi_k(\mathbf{r}_i)]$  enforces antisymmetry and defines the nodal surface.
- $e^{J(\mathbf{R})}$  (Jastrow factor) captures symmetric correlations (often pairwise).

A common two-body Jastrow form:

$$J(\mathbf{R}) = \sum_{i < j} u(r_{ij}), \quad r_{ij} = |\mathbf{r}_i - \mathbf{r}_j|.$$

# Limitations of Slater–Jastrow

- The Jastrow factor improves correlation energy but **does not change the nodes**:

$$\psi_{\text{SJ}}(\mathbf{R}) = 0 \iff \det[\phi_k(\mathbf{r}_i)] = 0.$$

- For strongly correlated fermions, the **nodal surface** is often the main source of error.
- Backflow is designed specifically to introduce flexible, configuration-dependent nodal structure.

# Backflow: configuration-dependent coordinates

Backflow replaces bare coordinates  $\mathbf{r}_i$  by *effective* coordinates

$$\tilde{\mathbf{r}}_i = \mathbf{r}_i + \boldsymbol{\xi}_i(\mathbf{R}),$$

where the backflow displacement  $\boldsymbol{\xi}_i$  depends on the full configuration.  
A classical (pairwise) backflow form:

$$\boldsymbol{\xi}_i(\mathbf{R}) = \sum_{j \neq i} \eta(r_{ij}) (\mathbf{r}_i - \mathbf{r}_j),$$

with a scalar backflow function  $\eta(\cdot)$  (often parameterized).

# Slater–Jastrow–Backflow (SJB) ansatz

The Slater determinant is evaluated at backflow coordinates:

$$\Psi_{\text{SJB}}(\mathbf{R}) = \det[\phi_k(\tilde{\mathbf{r}}_i)] e^{J(\mathbf{R})}.$$

- Backflow alters the nodes:

$$\Psi_{\text{SJB}}(\mathbf{R}) = 0 \iff \det[\phi_k(\tilde{\mathbf{r}}_i(\mathbf{R}))] = 0,$$

so the nodal surface becomes an implicit, many-body object.

- $J(\mathbf{R})$  still captures symmetric correlation and cusp conditions.

# Backflow as orbital deformation (equivalent view)

Instead of moving coordinates, one may view backflow as making orbitals configuration-dependent:

$$\phi_k(\mathbf{r}_i) \longrightarrow \phi_k(\mathbf{r}_i; \mathbf{R}) := \phi_k(\tilde{\mathbf{r}}_i(\mathbf{R})).$$

Thus, each particle experiences an *effective environment-dependent orbital* set by the positions of all other particles.

# Lattice analogue: backflow features for occupations

For lattice fermions with occupations  $\mathbf{n} = (n_1, \dots, n_L)$ , a backflow-like construction introduces *configuration-dependent features*

$$\tilde{h}_i(\mathbf{n}) = h_i + \sum_{j \neq i} f_{ij} \Phi(n_i, n_j, \mathbf{n}),$$

where  $\tilde{h}_i$  plays the role of an effective site feature/embedding.  
A determinant-based ansatz may then use effective orbitals

$$\Psi(\mathbf{n}) \propto \det[\varphi_k(\tilde{h}_i(\mathbf{n}))] \times e^{J(\mathbf{n})},$$

with a Jastrow-like factor  $J(\mathbf{n})$  capturing density–density correlations.



# Neural backflow: learned displacements or embeddings

In NQS practice, backflow is often **neural**:

$$\xi_i(\mathbf{R}) = \mathcal{N}_\theta(i; \mathbf{R}),$$

with  $\mathcal{N}_\theta$  a symmetry-respecting neural network (e.g. equivariant).  
Then

$$\tilde{\mathbf{r}}_i = \mathbf{r}_i + \xi_i(\mathbf{R}), \quad \Psi_{\text{SJB}}(\mathbf{R}) = \det[\phi_k(\tilde{\mathbf{r}}_i)] e^{J(\mathbf{R})}.$$

This generalizes pairwise backflow to *collective* many-body backflow.

# Transformer view: tokens, context, and backflow

A Transformer builds *context-dependent* representations:

$$\text{token embedding } \mathbf{e}_i \longrightarrow \text{contextual embedding } \mathbf{h}_i(\{\mathbf{e}_j\}_{j=1}^N).$$

Backflow is analogous:

$$\mathbf{r}_i \longrightarrow \tilde{\mathbf{r}}_i(\mathbf{R}) = \mathbf{r}_i + \boldsymbol{\xi}_i(\mathbf{R}),$$

i.e. each particle coordinate (or feature) becomes **contextual**, conditioned on the entire configuration.

# Self-attention as a backflow-like map (schematic)

In a Transformer layer, one often has

$$\mathbf{h}_i = \mathbf{e}_i + \sum_j \alpha_{ij} \mathbf{V} \mathbf{e}_j, \quad \alpha_{ij} = \text{softmax}_j \left( \frac{(\mathbf{Q} \mathbf{e}_i) \cdot (\mathbf{K} \mathbf{e}_j)}{\sqrt{d}} \right).$$

Compare with pairwise backflow:

$$\tilde{\mathbf{r}}_i = \mathbf{r}_i + \sum_{j \neq i} \eta(r_{ij}) (\mathbf{r}_i - \mathbf{r}_j).$$

## Analogy:

- weights  $\alpha_{ij}$  or  $\eta(r_{ij})$  quantify *influence* of  $j$  on  $i$ ,
- both produce nonlocal, permutation-symmetric updates of per-particle features.

# Why Transformer-style backflow is powerful for NQS

- Long-range correlations are captured naturally via attention over all particles.
- Backflow/attention directly produces **configuration-dependent nodes** (fermions).
- Parameter-efficient: global structure can be learned with fewer parameters than very deep MLPs.
- Symmetry handling: permutation symmetry and (with equivariance) rotational symmetry can be enforced.

**Takeaway:** backflow is the physics analogue of contextual embeddings; Transformers are a natural modern architecture to implement neural backflow maps.

# Summary

- Slater–Jastrow:

$$\psi_{\text{SJ}}(\mathbf{R}) = \det[\phi_k(\mathbf{r}_i)] e^{J(\mathbf{R})}$$

captures correlations but leaves nodes fixed.

- Backflow introduces effective coordinates

$$\tilde{\mathbf{r}}_i = \mathbf{r}_i + \boldsymbol{\xi}_i(\mathbf{R})$$

and yields Slater–Jastrow–Backflow:

$$\psi_{\text{SJB}}(\mathbf{R}) = \det[\phi_k(\tilde{\mathbf{r}}_i)] e^{J(\mathbf{R})}.$$

- Backflow  $\leftrightarrow$  contextualization: conceptually aligned with Transformer self-attention.