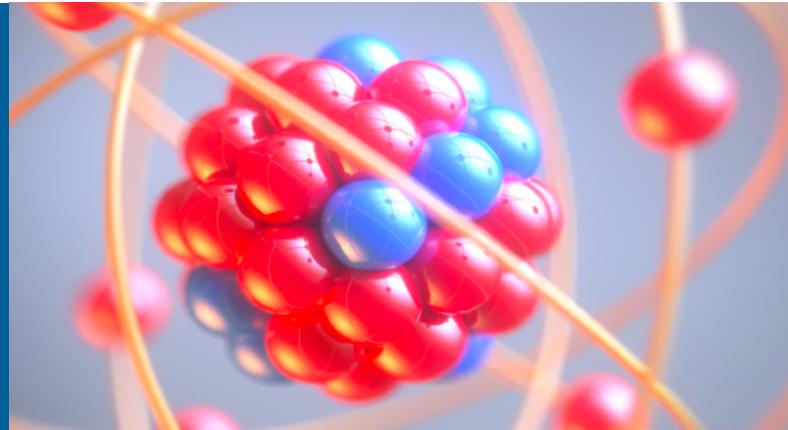


VARIATIONAL LEARNING NUCLEI AND HYPER NUCLEI



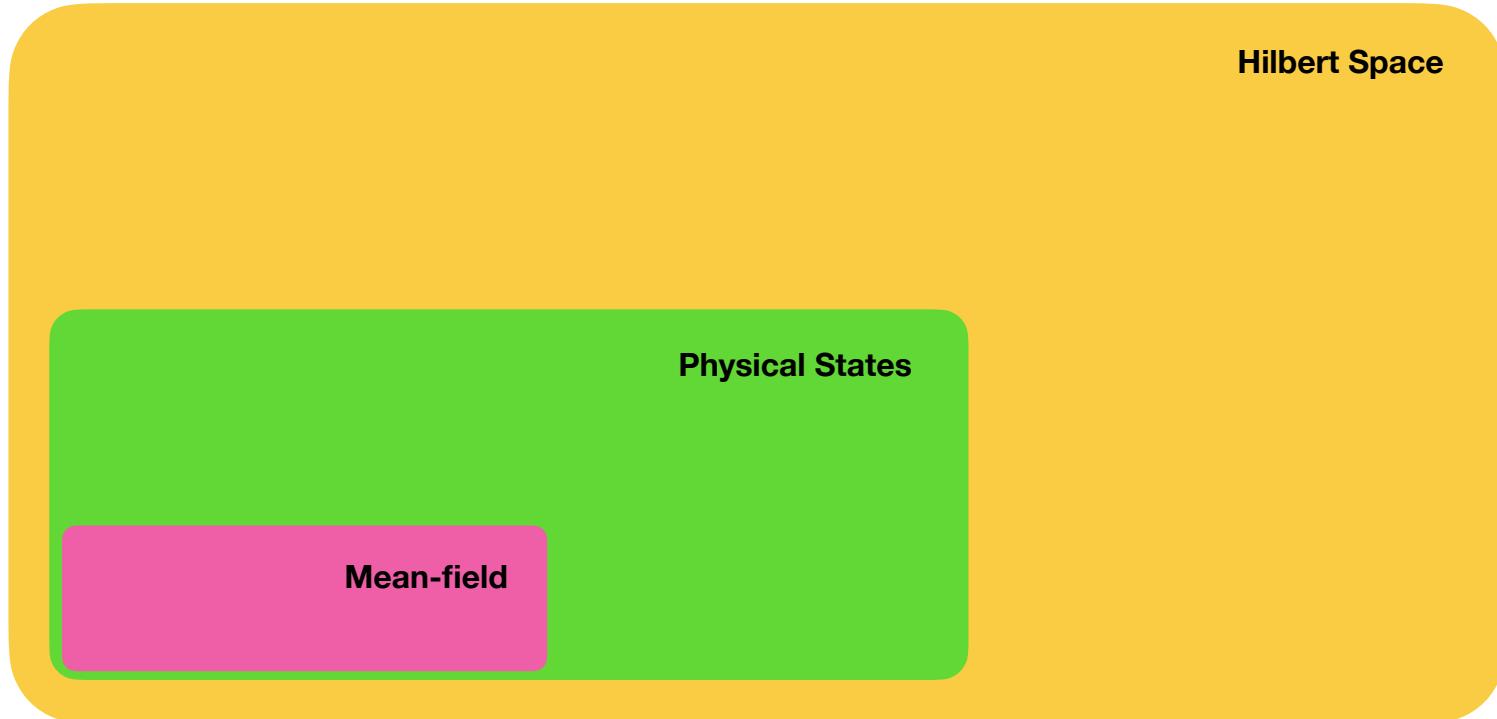
ALESSANDRO LOVATO

Argonne
NATIONAL LABORATORY

STREAMLINE Collaboration Symposium
Michigan State University & FRIB

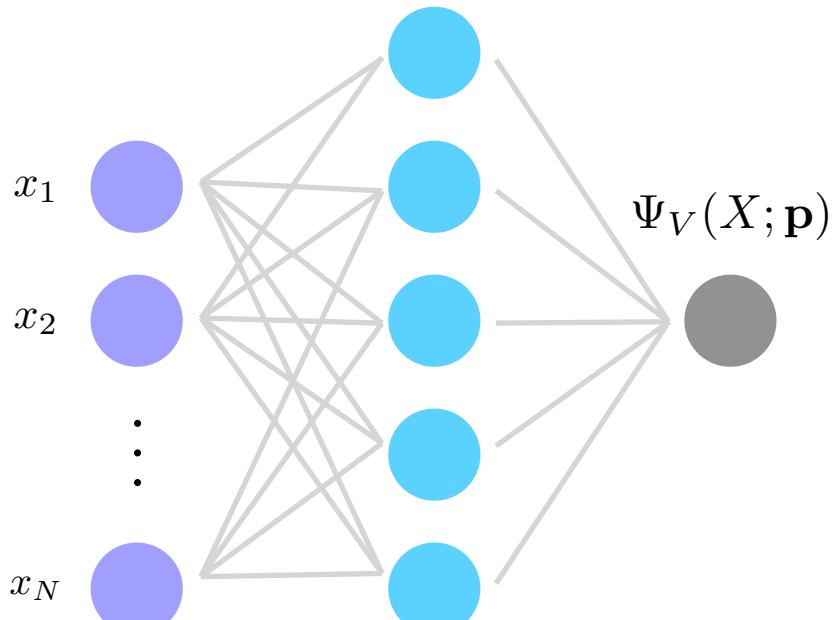
May 10, 2024

NEURAL-NETWORK QUANTUM STATES



NEURAL-NETWORK QUANTUM STATES

Originally introduced by Carleo and Troyer for spin systems, NQS are now widely and successfully applied to study condensed-matter systems



$$E_V \equiv \frac{\langle \Psi_V | H | \Psi_V \rangle}{\langle \Psi_V | \Psi_V \rangle} > E_0$$

$$E_V \simeq \frac{1}{N} \sum_{X \in |\Psi_V(X)|^2} \frac{\langle X | H | \Psi_V \rangle}{\langle X | \Psi_V \rangle}$$

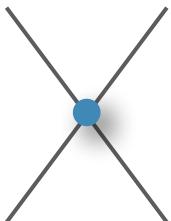
$$\mathbf{p}_{\tau+\delta\tau} = \mathbf{p}_\tau - \eta(S_\tau + \epsilon I)^{-1}\mathbf{g}_\tau$$

PIONLESS EFT HAMILTONIAN

Input: Hamiltonian inspired by a LO pionless-EFT expansion

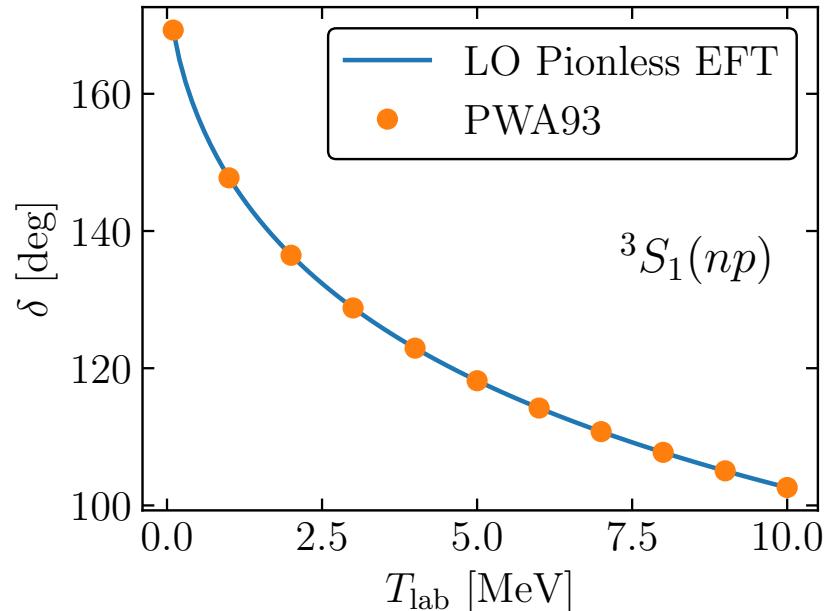
$$H_{LO} = - \sum_i \frac{\vec{\nabla}_i^2}{2m_N} + \sum_{i < j} v_{ij} + \sum_{i < j < k} V_{ijk}$$

- NN potential fit to s-wave np scattering lengths and effective ranges



$$v_{ij}^{\text{CI}} = \sum_{p=1}^4 v^p(r_{ij}) O_{ij}^p,$$

$$O_{ij}^{p=1,4} = (1, \tau_{ij}, \sigma_{ij}, \sigma_{ij}\tau_{ij})$$

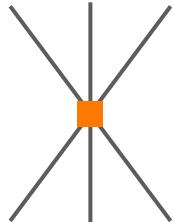


PIONLESS EFT HAMILTONIAN

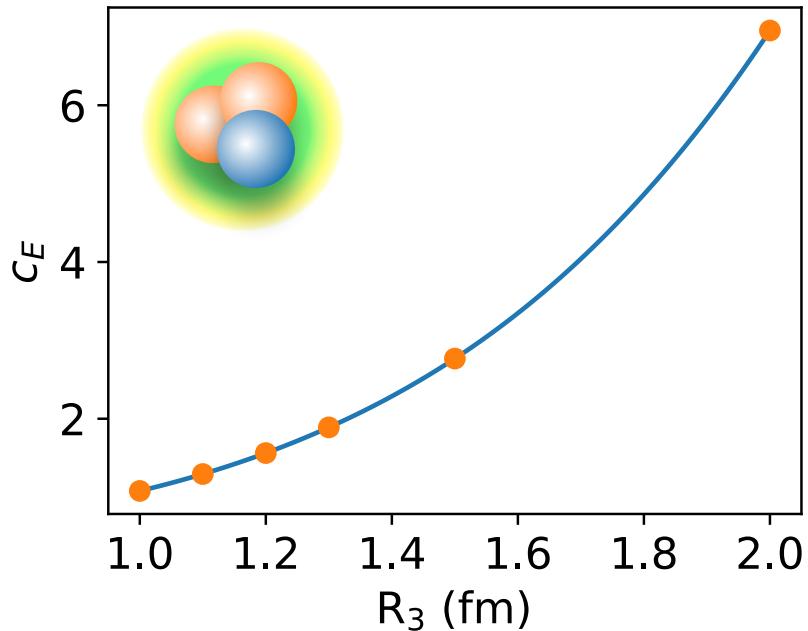
Input: Hamiltonian inspired by a LO pionless-EFT expansion

$$H_{LO} = - \sum_i \frac{\vec{\nabla}_i^2}{2m_N} + \sum_{i < j} v_{ij} + \sum_{i < j < k} V_{ijk}$$

- 3NF adjusted to reproduce the energy of ${}^3\text{H}$.



$$V_{ijk} \propto c_E \sum_{\text{cyc}} e^{-(r_{ij}^2 + r_{jk}^2)/R_3^2}$$



HIDDEN NUCLEONS ANSATZ

$$\Psi_{\text{HN}}(X) = \det \begin{bmatrix} \phi_1(x_1) & \phi_1(x_2) & \phi_1(x_3) & \phi_1(x_4) \\ \phi_2(x_1) & \phi_2(x_2) & \phi_2(x_3) & \phi_2(x_4) \\ \phi_3(x_1) & \phi_3(x_2) & \phi_3(x_3) & \phi_3(x_4) \\ \phi_4(x_1) & \phi_4(x_2) & \phi_4(x_3) & \phi_4(x_4) \\ \hline \chi_1(x_1) & \chi_1(x_2) & \chi_1(x_3) & \chi_1(x_4) \\ \chi_2(x_1) & \chi_2(x_2) & \chi_2(x_3) & \chi_2(x_4) \end{bmatrix} \begin{bmatrix} \phi_1(y_1) & \phi_1(y_2) \\ \phi_2(y_1) & \phi_1(y_2) \\ \phi_3(y_1) & \phi_1(y_2) \\ \phi_4(y_1) & \phi_1(y_2) \\ \hline \chi_1(y_1) & \chi_2(y_2) \\ \chi_2(y_1) & \chi_2(y_2) \end{bmatrix}$$

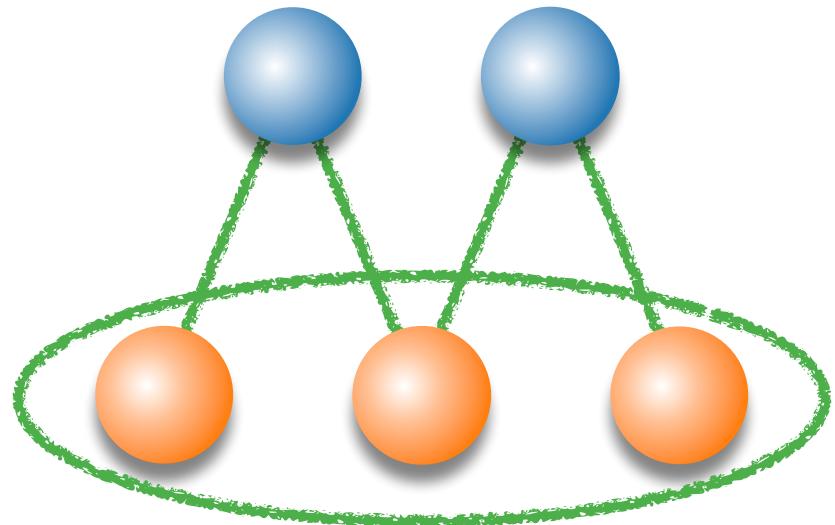
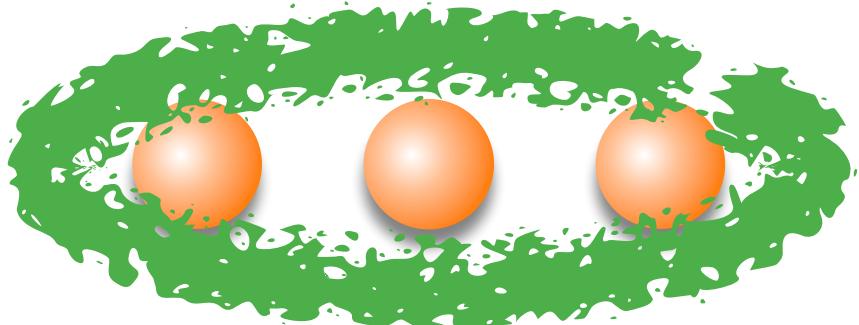
Visible orbitals on visible coordinates

Visible orbitals on hidden coordinates

Hidden orbitals on visible coordinates

Hidden orbitals on hidden coordinates

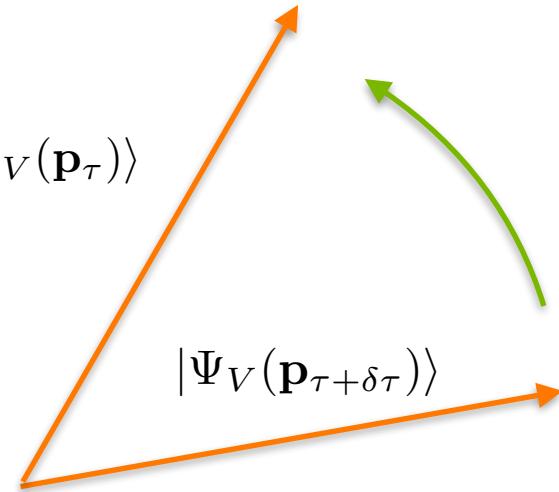
HIDDEN NUCLEONS ANSATZ



WAVE FUNCTION OPTIMIZATION

ANN trained by performing an imaginary-time evolution in the variational manifold

$$\left\{ \begin{array}{l} |\bar{\Psi}_V(\mathbf{p}_\tau)\rangle \equiv (1 - H\delta\tau)|\Psi_V(\mathbf{p}_\tau)\rangle \\ \mathbf{p}_{\tau+\delta\tau} = \arg \max_{\mathbf{p} \in R^d} \left(|\langle \bar{\Psi}_V(\mathbf{p}_\tau) | \Psi_V(\mathbf{p}_{\tau+\delta\tau}) \rangle|^2 \right) \end{array} \right.$$



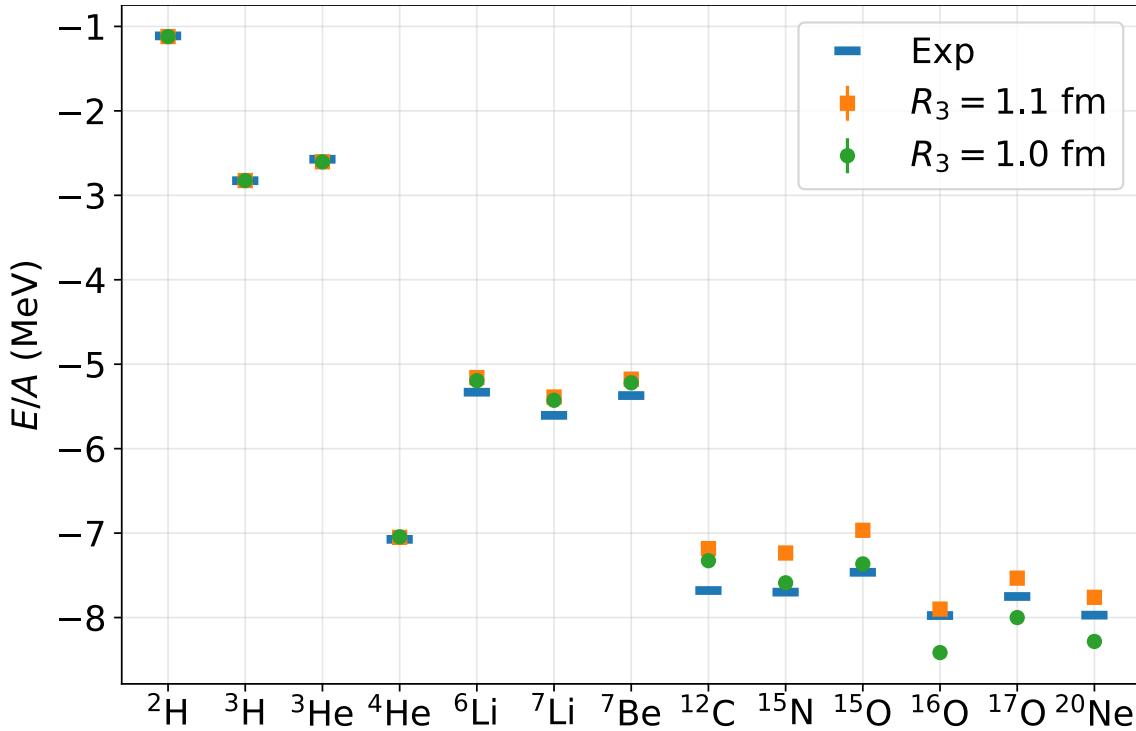
The parameters are updated as

$$\mathbf{p}_{\tau+\delta\tau} = \mathbf{p}_\tau - \delta\tau S^{-1} \mathbf{g}_\tau$$

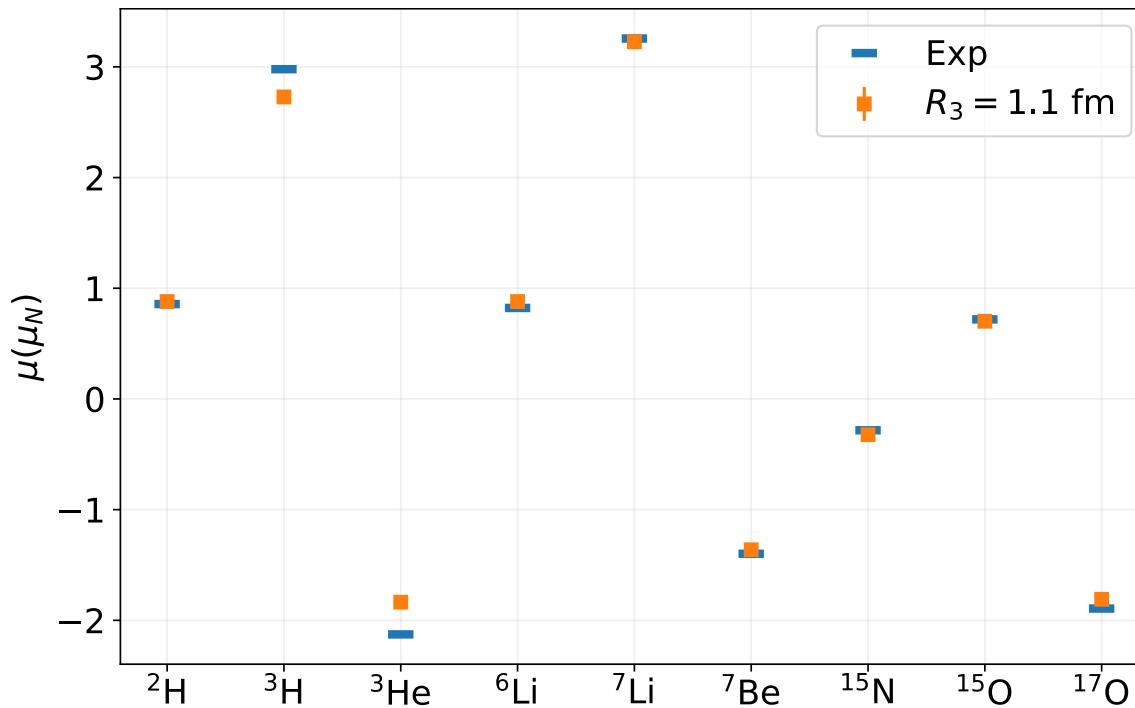
J. Stokes, *et al.*, *Quantum* 4, 269 (2020).

S. Sorella, *Phys. Rev. B* 64, 024512 (2001)

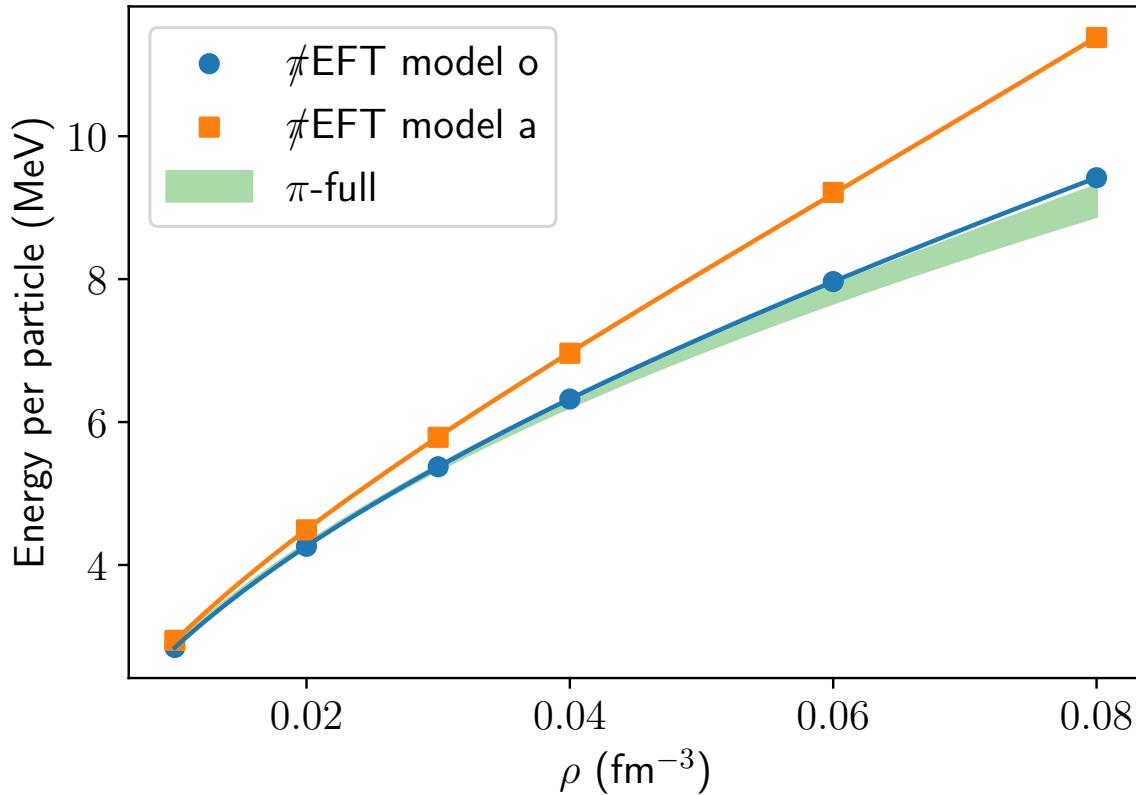
BINDING ENERGIES



MAGNETIC MOMENTS

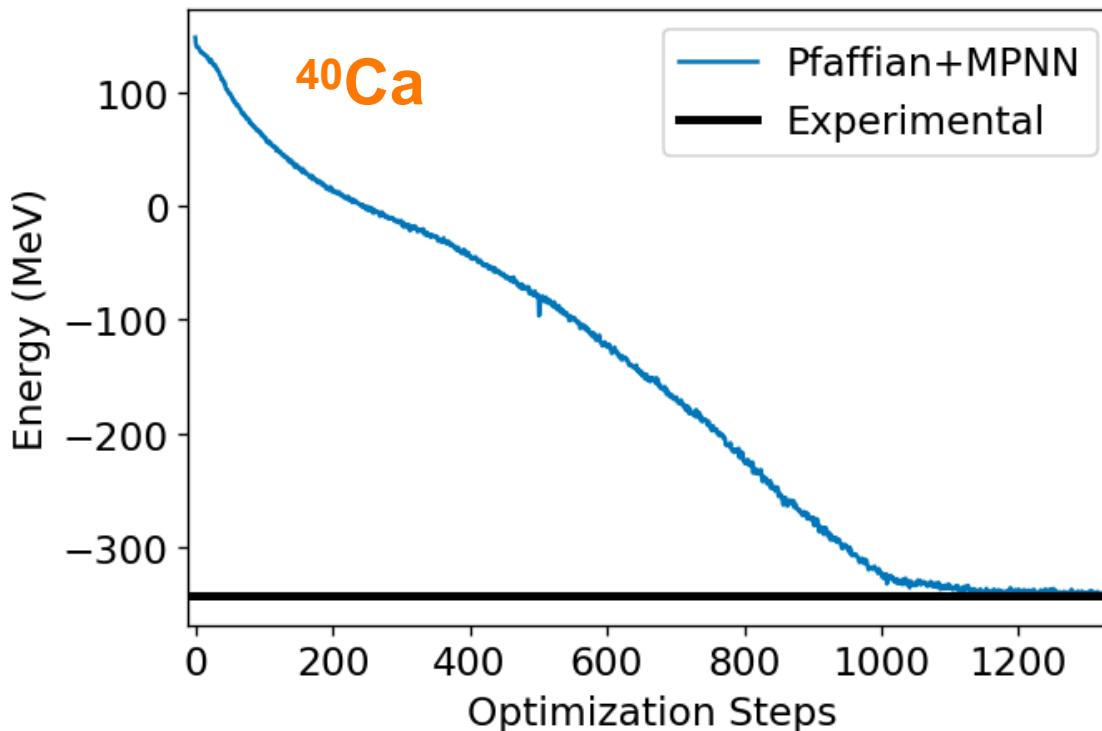


DILUTE NEUTRON MATTER



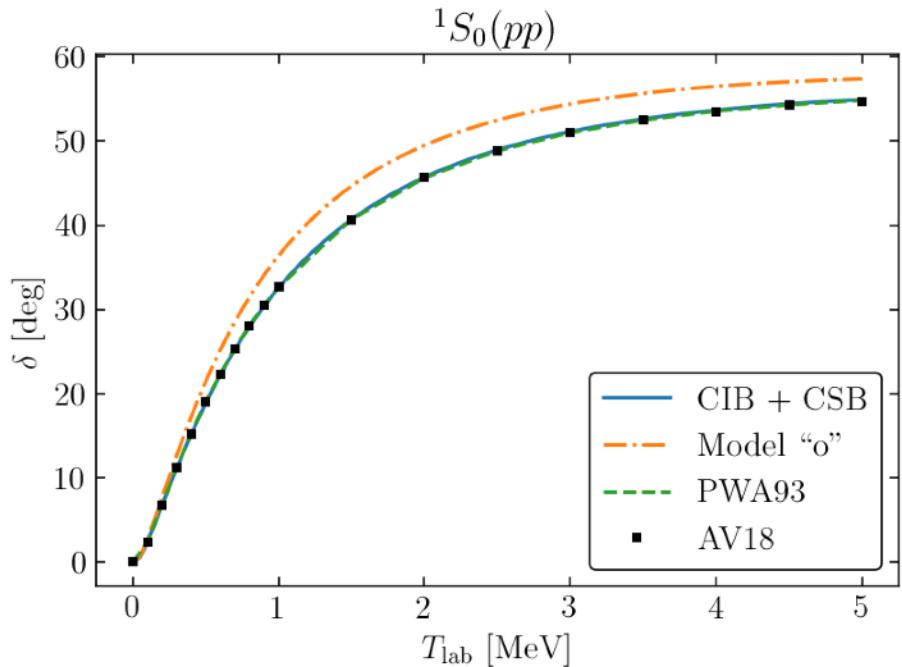
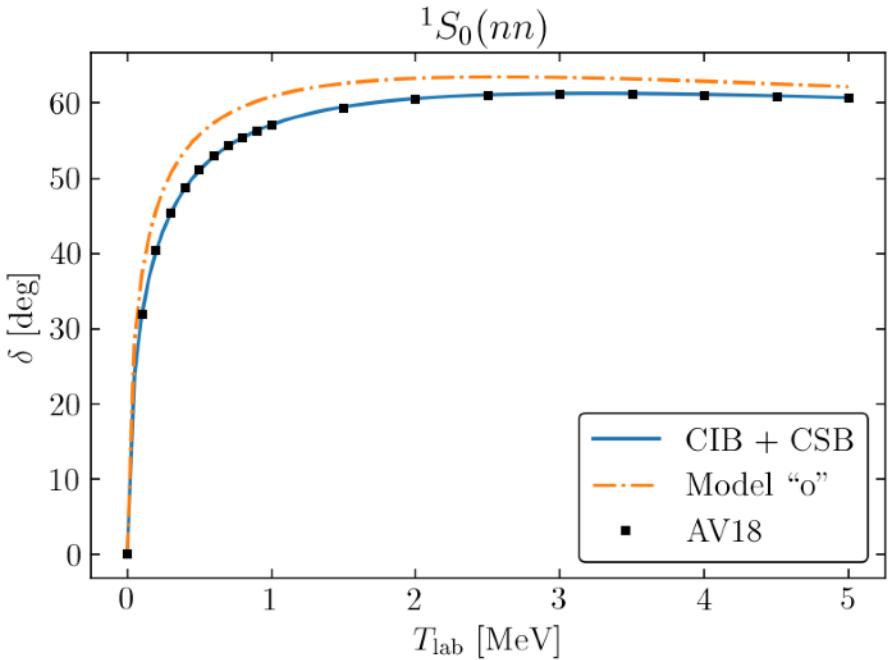
FIRST GOAL: EXTENSION TO LARGER NUCLEI

Bryce Fore has significantly improved the scalability of the code to multiple GPU nodes



SECOND GOAL: MORE REALISTIC POTENTIALS

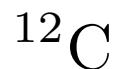
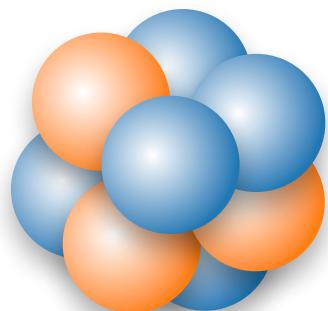
Anthony Tropiano has included CIB and CSB terms in the Hamiltonian



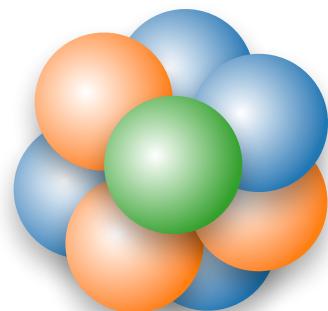
ADDITIONAL GOAL: TACKLE STRANGENESS

Hypernuclei: bound state between an ordinary nucleus with one (or more) hyperons

Our work: consider single- Λ hypernuclei to begin with



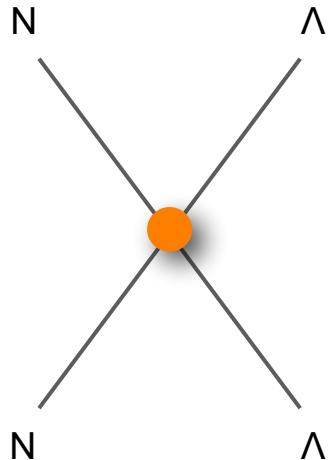
6 protons, 6 neutrons



6 protons, 5 neutrons, 1 lambda

PIONLESS EFT FOR HYPER-NUCLEI

Input: Hamiltonian inspired by a LO pionless-EFT expansion



$$V_{\Lambda N} = \sum_{S,T} v_{ST}(r_{ij}) \hat{P}_{ST}$$

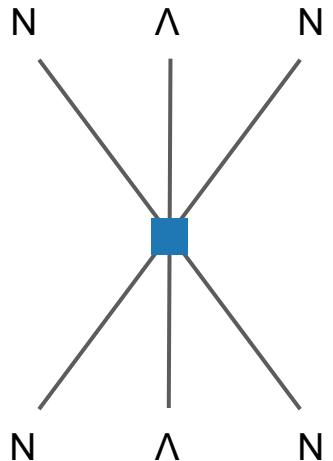
→ T=1/2, S=1

→ T=1/2, S=0

Parameters are determined by fitting proton- Λ scattering length and effective range

PIONLESS EFT FOR HYPER-NUCLEI

Input: Hamiltonian inspired by a LO pionless-EFT expansion



$$V_{\Lambda NN} = \sum_{S,T} D_{S,T} v_{ST}(r_{i\Lambda}) v_{ST}(r_{j\Lambda}) \hat{Q}_{ST}$$

→ T=0, S=1/2

→ T=0, S=3/2

→ T=1, S=1/2

Parameters are determined by fitting $^3_\Lambda\text{H}$, $^4_\Lambda\text{H}(S=0)$, $^4_\Lambda\text{H}(S=1)$, and $^5_\Lambda\text{He}$.

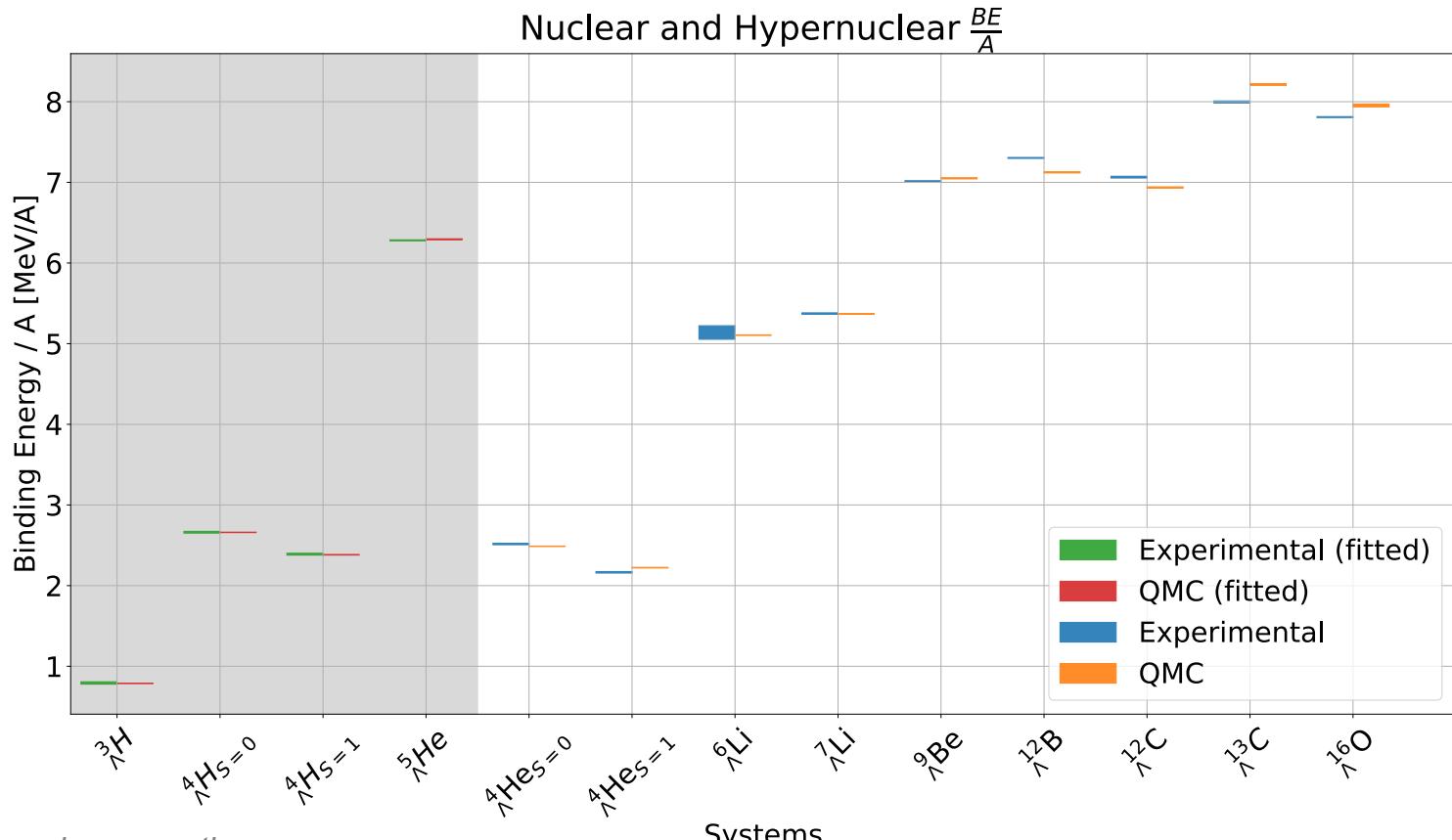
NQS FOR HYPER-NUCLEI

No terms in the Hamiltonian mix lambda and nucleons: **distinguishable**

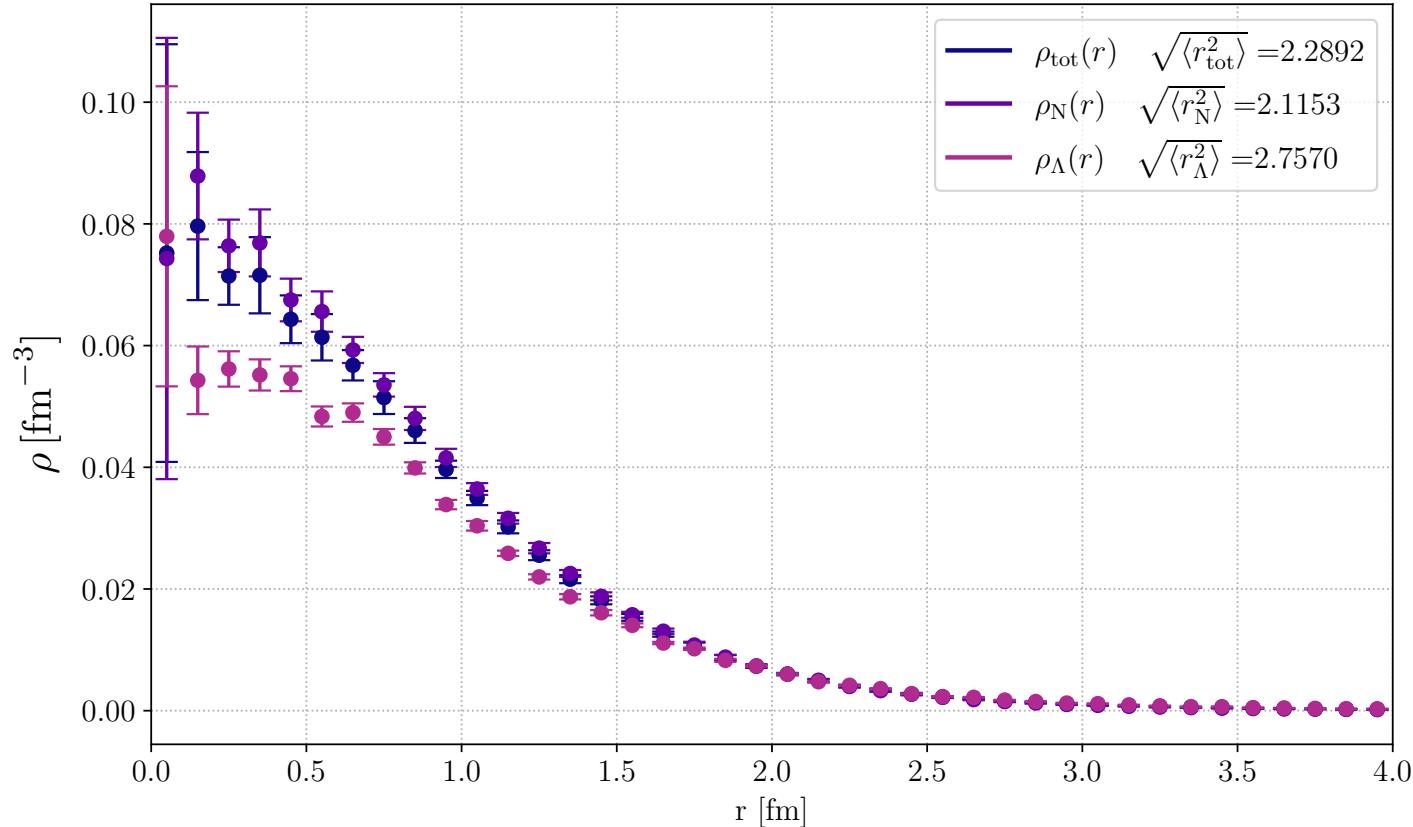
Andrea Di Donna proposed the following ansatz

$$\Psi(x_\Lambda, x_1, \dots, x_A) = \mathcal{U}(x_\Lambda; x_1, \dots, x_A) \times \Psi_{HN}(x_1, \dots, x_A)$$

GROUND-STATE ENERGIES



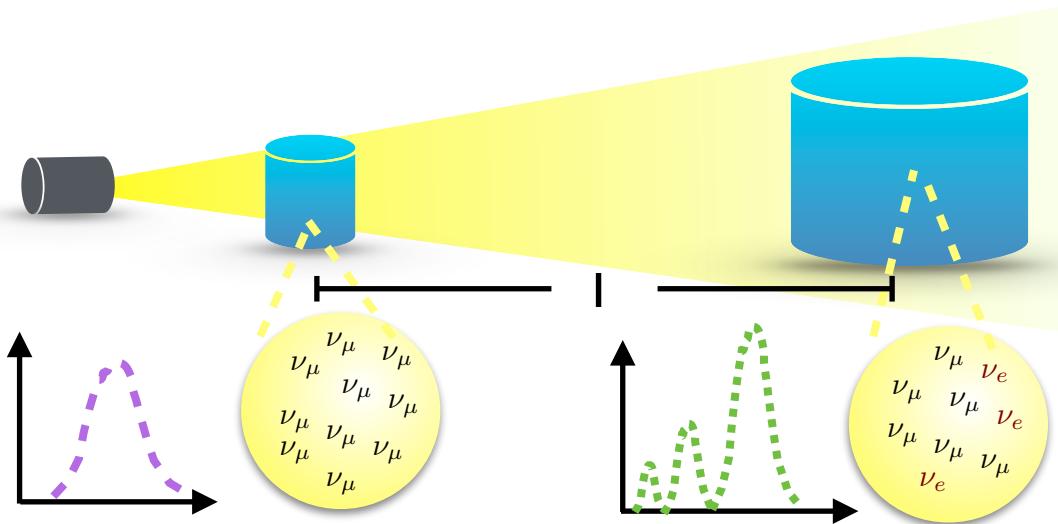
SINGLE-PARTICLE DENSITIES



REAL-TIME QUANTUM DYNAMICS

Real-time quantum dynamics is the prototypal exponentially-hard problem

- Relevant for: fission, fusion lepton- and hadron-nucleus scattering, collective neutrino oscillations



REAL-TIME QUANTUM DYNAMICS

Initial goal: compute the linear response function:

$$R(\omega) = \sum_n \langle \Psi_0 | D^\dagger | \Psi_n \rangle \langle \Psi_n | D | \Psi_0 \rangle \delta(\omega - E_f + E_0)$$

$$D = \frac{NZ}{A} \left(\frac{1}{Z} \sum_{p=1}^A \vec{r}_p - \frac{1}{N} \sum_{n=1}^N \vec{r}_n \right) = \frac{NZ}{A} (\vec{R}_p - \vec{R}_n)$$

With N. Rocco and K. Godbye, we re-derived off the expressions for the linear response

$$R(\omega) = \text{Im} \left[\int dt e^{i\omega t} (\langle \Psi(t) | D | \Psi(t) \rangle - \langle \Psi_0 | D | \Psi_0 \rangle) \right]$$

REAL-TIME QUANTUM DYNAMICS

The time-evolved state is given by

$$|\Psi(t)\rangle = e^{-iHt} e^{-ibD} |\Psi_0\rangle$$

Two options:

- Use the TDVP to evolve $|\Psi_0\rangle$ with $H' = D$ for a time b .
- Maximize the fidelity between $e^{-ibD} |\Psi_0\rangle$ and $|\hat{\Psi}(\mathbf{p})\rangle_b$

Real time evolution with the nuclear Hamiltonian

$$\mathcal{D}(|\Psi(\mathbf{p}_{t+\delta t})\rangle, e^{-iHt}|\Psi(\mathbf{p}_t)\rangle)^2 = \arccos \left(\sqrt{\frac{\langle\Psi(\mathbf{p}_{t+\delta t})|e^{-iHt}|\Psi(\mathbf{p}_t)\rangle\langle\Psi(\mathbf{p}_t)|e^{iHt}|\Psi(\mathbf{p}_{t+\delta t})\rangle}{\langle\Psi(\mathbf{p}_{t+\delta t})|\Psi(\mathbf{p}_{t+\delta t})\rangle\langle\Psi(\mathbf{p}_t)|\Psi(\mathbf{p}_{t+\delta t})\rangle}} \right)^2$$

THANK YOU