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Abstract: This is some abstract about this thesis.

Dedicated to

Bárbara, Humberto and Luciana.

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Declaration

The work in this thesis is based on research carried out in the Department of Physics at Durham University. No part of this thesis has been submitted elsewhere for any degree or qualification.

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*"It seems very pretty, [...] somehow it seems to fill my head with
ideas—only I don't exactly know what they are!"*

— from *Through the Looking-Glass, and What Alice Found There* by
Lewis Carroll

Chapter 1

Introduction

1.1 The Standard Model at 51

1.1.1 Fields and symmetries

1.1.2 Spontaneous symmetry breaking

1.1.3 Shortcomings

1.2 Portals to beyond the Standard Model

Chapter 2

Aspects of neutrino physics

2.1 Oscillations

2.2 Sources

2.3 Scattering

2.4 Mass mechanisms

For a review on low-scale models see [1].

2.4.1 Conventional seesaws

2.4.2 Low scale seesaw variants

2.4.3 Radiative masses

The most straightforward extension of the SM which can generate neutrino masses at loop level is perhaps the so-called Zee-Babu model.

Chapter 3

New forces in Neutrino Scattering

Chapter 4

A dark neutrino sector

Chapter 5

The MiniBooNE experimental anomaly

Chapter 6

Future precision neutrino physics

Chapter 7

Conclusions

Particle physics is at a very important moment of its history. The Standard Model has surprised us with its unprecedent accuracy in describing particle data, but a few mysteries remain. Perhaps neutrino masses and dark matter are to the Standard Model what blackbody radiation was to classical physics in the beginning of the 20th century, a scientific revolution on the wait.

Appendix A

Phase Space

In this appendix we show explicitly the factorization of N -final state phase space factors into $N - 2$ 2-body ones. To study arbitrary observables that depend on the final state momenta, one must also be able to write down the four momenta in terms of the phase space variables. We carry out this exercise for 3 and 4 body phase space below.

A.1 Factorization into 2PS

The final decomposition of the phase space factors then reads:

$$dPS^2 = dm^2 dPS^2(p_1, p_2) dPS^2(p_2, p_3). \quad (\text{A.1.1})$$

A.2 Three-body phase space

A.3 Four-body phase space

Appendix B

One loop ν masses in Type-I seesaw

B.1 Type-I seesaw neutrino masses in the SM

In this appendix, we compute the one-loop corrections to the light neutrino masses in the SM, following [2–4].

We will also be making use of the On-Shell (OS) renormalization scheme, known . This is ensured by requiring that the off-diagonal elements of the self-energy be diagonal when the external particles are on their mass shell, and that the residue of the renormalized propagator are equal to one.

Assuming Majorana neutrino fields, one can write the self-energy tensor in its most general form:

$$\Sigma_{ij}(\not{q}) = \not{q}P_L\Sigma_{ij}^L(q^2) + \not{q}P_R\Sigma_{ij}^R(q^2) + P_L\Sigma_{ij}^M(q^2) + P_R\Sigma_{ij}^{M*}(q^2), \quad (\text{B.1.1})$$

where by virtue of the Majorana nature the previous terms obey

$$\Sigma_{ij}^L(q^2) = \Sigma_{ij}^{R*}(q^2), \quad \Sigma_{ij}^M(q^2) = \Sigma_{ji}^M(q^2).$$

$$\delta m_L = \quad (\text{B.1.2})$$

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