

Thm (Sturm Liouville): There is an infinite sequence $(\lambda_j)_{j=1}^{\infty}$ of eigenvalues, λ_j real and $|\lambda_j| \xrightarrow{j \rightarrow \infty} \infty$. The corresponding $(e_j)_{j=1}^{\infty}$ form a complete orthonormal basis for $L^2[a, b]$.

We proved this under assumption that 0 is not an eigenvalue. We address this now.
 ↘ related to Wronskian being not 0.

To remove this assumption, we'll show:

Thm: not every real number is an eigenvalue of the Sturm-Liouville problem.

proof: let $(e_n)_{n=1}^{\infty}$ be complete orthonormal basis for $L^2(a, b)$. Assume that for any $\lambda \in \mathbb{R}$ is an eigenvalue with eigenfunction f_{λ} . We claim we prove later for any $\lambda \neq \mu$, $f_{\lambda} \perp f_{\mu}$.

Then let $\|f_{\lambda}\| = 1$. Consider not countable

$$E_n = \{\lambda \in \mathbb{R} : (e_n, f_{\lambda}) \neq 0\}$$

we'd like to consider $\bigcup_{n=1}^{\infty} E_n$. Then we write $E_n = \bigcup_{m=1}^{\infty} \{\lambda \in \mathbb{R} : |(e_n, f_{\lambda})| \geq \frac{1}{m}\}$.

Consider $\{\lambda \in \mathbb{R} : |(g, f_{\lambda})| \geq 1\}$ is finite for any $g \in L^2$.

This is a consequence of Bessel's inequality.

So E_n is countable b/c each one is countable.

So $\bigcup_{n=1}^{\infty} E_n$ is countable.

\Rightarrow This is a proper subset of \mathbb{R} (\mathbb{R} is not countable). Take $\lambda \in \mathbb{R} \setminus \bigcup_{n=1}^{\infty} E_n$.

Then $(e_n, f_{\lambda}) = 0 \forall n$. Then $f_{\lambda} = 0$. But $\|f_{\lambda}\| = 1$. \downarrow □

We need to prove claim $f_\lambda \perp f_\mu$.

Essentially we WTS $Lu = (pu')' + qu$ is "self adjoint" but L is not bounded. So we mean "self-adjoint" by

$$(Lu, v) = (u, Lv) \text{ for } u, v \in C^2 \text{ solutions of Sturm Liouville problem. (subset of } L^2).$$

$$\begin{aligned} \text{To see this, } (Lu, v) - (u, Lv) &= \int_a^b v Lu - u Lv \, dx = p(uv' - vu') \Big|_a^b \\ &= 0 \end{aligned}$$

B/c u, v are solutions to SL , so at say $x=a$

$$\alpha u(a) + \alpha' u'(a) = 0$$

$$\alpha v(a) + \alpha' v'(a) = 0$$

$$\Rightarrow \det \begin{pmatrix} u(a) & u'(a) \\ v(a) & v'(a) \end{pmatrix} \neq 0 \text{ otherwise } \alpha = \alpha' = 0$$

The same works for $x=b$

For f_λ, f_μ eigenfunctions of λ, μ , we have

$$0 = (Lf_\lambda, f_\mu) - (f_\lambda, Lf_\mu)$$

$$= (-\lambda f_\lambda, f_\mu) - (f_\lambda, -\mu f_\mu)$$

$$= (\mu - \lambda)(f_\lambda, f_\mu) \Rightarrow (f_\lambda, f_\mu) = 0 \text{ if } \lambda \neq \mu.$$

Final part of Sturm-Liouville

Take $\mu \in \mathbb{R}$ which is not eigenvalue of SL .

$$\text{Replace } (pu')' + qu = -\lambda u \quad (*)$$

$$\text{by } (pu')' + (q+\mu)u = -\lambda u \quad (**)$$

$\Rightarrow f$ is eigen function of (Φ) with eigen value λ

$\Leftrightarrow f$ is eigen function of (Φ) with eigenvalue $\lambda + \mu$.

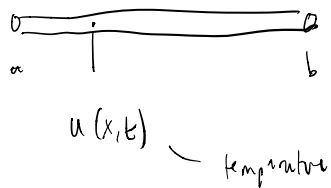
If $\lambda = 0$ is eigen value of $(\Phi) \Rightarrow 0$ is not eigen value of $(\Phi + \mu)$.

(Shift of spectrum: even if $\lambda = 0$, we shift so λ cannot be 0. Apply theorem, then shift eigenvalue back).

(v) Application to PDE:

The Sturm-Liouville theory is useful in the method called separation of variables for solving PDE.

Consider the heat eqn on a finite interval $[a, b] \subseteq \mathbb{R}$.



Let $u(x,t)$ be temperature of bar at position x and time t . The PDE for u is: $u_t = \alpha^2 u_{xx}$, $\alpha > 0$. α^2 is the diffusivity (take $\alpha^2 = 1$).

↙
derivatives

We need to prescribe boundary/initial conditions.

(BC) Boundary condition: $\alpha u(a,t) + \alpha' u_x(a,t) = 0$ $0 < t < \infty$
 $\beta u(b,t) + \beta' u_x(b,t) = 0$

(IC) Initial condition $u(x,0) = \psi(x)$ $a \leq x \leq b$.

(Remark: can also consider wave equation $u_{tt} = c^2 u_{xx}$).

Solve using separation of variables. Look for solution

$$u(x, t) = f(x) g(t)$$

$$\Rightarrow \text{PDE becomes } f(x) g'(t) = f''(x) g(t)$$

$$\Rightarrow \frac{g'(t)}{g(t)} = \frac{f''(x)}{f(x)} = \text{constant} = -\lambda$$

↙ b/c fix t then RHS must be constant and vice versa.

\Rightarrow we get 2 ODEs...

$$g'(t) + \lambda g(t) = 0 \quad \text{and} \quad f''(x) + \lambda f(x) = 0 \quad \textcircled{1}$$

$$\text{BC} \Rightarrow \alpha f(a) g(t) + \alpha' f'(a) g(t) = 0$$

we don't want $g(t) = 0$ so

$$\alpha f(a) + \alpha' f'(a) = 0 \quad \textcircled{2}$$

$$\text{Similarly } \beta f(b) + \beta' f'(b) = 0 \quad \textcircled{3}$$

Then $\textcircled{1}, \textcircled{2}, \textcircled{3}$ is a Sturm Liouville problem! so there is $(\lambda_j)_{j=1}^{\infty}$ and $(\varphi_j)_{j=1}^{\infty}$ eigenfunction and $(\varphi_j)_{j=1}^{\infty}$ complete orthonormal basis for L^2 .

we got λ , f
↙ will become φ

Then solve $g'(t) + \lambda_j g(t) = 0$ (linear 1st order ODE).

$$\Rightarrow g_j(t) = c_j e^{-\lambda_j t}$$

$$\text{Then we get } u_j(x, t) = f_j(x) g_j(t) = c_j e^{-\lambda_j t} \varphi_j(x)$$

Finally we get $u(x,t) = \sum_{j=1}^{\infty} c_j e^{-\lambda_j t} \varphi_j(x)$ (superposition principle).

To find c_j , use IC:

$$u(x,0) = \varphi(x).$$

$$\Rightarrow \varphi(x) = \sum_{j=1}^{\infty} c_j \varphi_j(x), \quad c_j = (\varphi, \varphi_j) \quad (\text{Hilbert complete}) \text{ for } L_2 \quad \star$$

$$\Rightarrow u(x,t) = \sum_{j=1}^{\infty} (\varphi, \varphi_j) e^{-\lambda_j t} \varphi_j(x).$$

Ex: Consider Heat equation on $[0, \pi]$. With $u(0)=0$, $u(\pi)=0$
the eigen values are $\lambda_j = j^2$ and $\varphi_j(x) = \sin jx$ so

$$u(x,t) = \sum_{j=1}^{\infty} c_j e^{-j^2 t} \sin jx.$$

we need to show $\sum_{j=1}^{\infty} c_j e^{-j^2 t} \varphi_j(x)$ converges.

For heat equation in example, you can show that $u(x,t) \in C^2([a,b]; \mathbb{R})$
why Weierstrass M-test. \wedge converges in