

(5)

Inductively, we constructed e_1, \dots, e_n of K with eigenvalue $\lambda_1, \dots, \lambda_n$ such that $|\lambda_j| = \|K_j\|$, $1 \leq j \leq n$. This can be continued if $K_n \neq 0$. If $K_n = 0$, the construction stops and we get

$$y = x - \sum_{j=1}^{n-1} (x, e_j) e_j \in H_n$$

$$\Rightarrow 0 = K_n y = K y = K x - \sum_{j=1}^{n-1} (x, e_j) K e_j$$

$$\Rightarrow Kx = \sum_{j=1}^{n-1} \lambda_j (x, e_j) e_j. \quad \text{We are done with finite sum}$$

If $K_n \neq 0 \quad \forall n$, choose $x \in X$ and $y_n = x - \sum_{j=1}^{n-1} (x, e_j) e_j$

$$\text{Then } y_n \in X_n, \text{ write } x = y_n + \sum_{j=1}^{n-1} (x, e_j) e_j$$

$$\Rightarrow \|Kx - \sum_{j=1}^{n-1} \lambda_j (x, e_j) e_j\| = \|K y_n\| = \|K_n y_n\| \leq \|K_n\| \|y_n\|$$

$$\leq \|\lambda_n\| \|y_n\|. \quad \text{But } \|x\|^2 = \|y_n\|^2 + \sum_{j=1}^{n-1} |(x, e_j)|^2 \Rightarrow \|y_n\| \leq \|x\|$$

$$\leq \|\lambda_n\| \|x\| \Rightarrow \text{limit} \rightarrow 0. \quad \text{as } \lambda_n \rightarrow 0. \quad ! \text{ (lemma)}$$

$$\text{So } Kx = \sum_{j=1}^{\infty} \lambda_j (x, e_j) e_j.$$

(ii) We need to show the completeness when K is separable.

$$\text{We know that } \forall x \in X, \quad Kx = \sum_{n=1}^{\infty} \lambda_n (x, e_n) e_n.$$

We can assume that $\lambda_n \neq 0$. Here e_n is a finite or infinite

orthonormal sequence. Each e_n is eigen function

②

So we observe that $x = \sum_{n=1}^{\infty} (x, e_n) e_n \in \ker K$.

Then let $\{f_m\}$ be a complete orthonormal basis for $(\ker K)^\perp$, which is a Hilbert space. Here we have $x = \underbrace{(\ker K)}_{f_m} \oplus \underbrace{(\ker K)^\perp}_{e_n}$

~~let e_n and~~ Because f_m are eigenfunctions of 0, so $f_m \perp e_n$

Thus $\{f_m\} \cup \{e_n\}$ is a complete orthonormal set in X .

For any $x \in X$, $x = \sum_n (x, e_n) e_n = \sum (x, f_m) f_m$

~~So $x =$~~ $\{f_m\} \cup \{e_n\}$ is a complete orthonormal set in X . \square

Remark: ① + ②

Ex: Let K be a ^{Hermitian} self-adjoint operator on a Hilbert space X . Let K be ~~positive~~ ^{positive}

operator namely $(Kx, x) \geq 0 \quad \forall x \in X$. Then we can define a

square root $A = K^{1/2}$ such that $A^2 = K$.

Proof: $Kx = \sum \lambda_n (x, e_n) e_n, \quad \lambda_n \geq 0$

Define $Ax = \sum \sqrt{\lambda_n} (x, e_n) e_n$.

Ex: (Canonical form of self adjoints) Let A be self adjoint on X .

Then there is orthonormal set $\{e_n\}_{n=1}^{\infty}$, $\{\lambda_n\}_{n=1}^{\infty}$ real positive real numbers λ_n with $\lambda_n \rightarrow 0$ so that

$$Ax = \sum_{n=1}^{\infty} \lambda_n (x, e_n) e_n.$$

Then λ_n are called eigen values of A .

①

Proof: Since A is cpt, so is A^*A . and A^*A is cpt, Hermitian

so there is orthonormal set $\{e_n\}_{n=1}^N$ w/ that $A^*A e_n = \mu_n e_n$, $\mu_n \neq 0$

and A^*A is zero on the subspace orthogonal to $\{e_n\}_{n=1}^N$. (w/ will be so)

Since A^*A is positive, each $\mu_n > 0$. Let $\lambda_n = \sqrt{\mu_n}$, let $f_n = A e_n / \lambda_n$

We can check f_n is ~~or~~ orthonormal and

~~App~~ $Ax = \sum \lambda_n (\overset{e_n}{\langle x, e_n \rangle}) \overset{f_n}{\text{the } f_n}$

Chap 9: Sturm-Liouville systems.

We apply the Hilbert space theories to ODEs. then we study PDE later.

1) Sturm-Liouville eigenvalue problem.

For $u \in C^2([a, b])$ (twice ~~continuous~~ ^{continuous} differentiable functions on $[a, b]$).

Consider a differential operator $Lu = (pu')' + qu = p^2 u'' + p'u' + qu$

where $p \in C^1([a, b], \mathbb{R})$; $p > 0$ and $q \in C^0([a, b], \mathbb{R})$.

Def: Sturm-Liouville eigenvalue problem:

$$\begin{cases} -Lu = \lambda u & \text{on } [a, b] \\ \alpha u(a) + \beta u'(a) = 0 \\ \gamma u(b) + \delta u'(b) = 0 \end{cases}$$

where, $\alpha, \alpha', \beta, \beta'$ are real constants and $\alpha = \alpha' = 0$, $\beta = \beta' = 0$ is excluded.

An eigenfunction to a scalar λ is a non-zero C^2 function u .

λ is called eigenvalue if such function exists.

Ex: ~~For~~ let $[a, b] = [0, \pi]$, ~~$p=0$, $p'=1$ on $[0, \pi]$~~ . Consider

$$\begin{cases} -u'' = \lambda u & \text{on } [a, b] \\ u(0) = 0 \\ u(\pi) = 0 \end{cases}$$

Here $q=0$, $p=1$, $\alpha=1$, $\alpha'=0$, $\beta=1$, $\beta'=0$.

(so $p'=0$)

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Find all eigenvalue and eigen functions.

Sol: We seek for all possible values

- $\lambda > 0 \Rightarrow$ general solution is $u(x) = A \cos \sqrt{\lambda} x + B \sin \sqrt{\lambda} x$.
(how do we know this?)

We ~~know~~ ^{boundary} condition, $u(0) = 0 \Rightarrow A = 0$, $u(\pi) = 0 \Rightarrow \cancel{B \cos \pi}$

$\sqrt{\lambda} B \cos \sqrt{\lambda} \pi = 0$. If $B = 0$, $u = 0$ trivial so $B \neq 0 \Rightarrow$

$\cos \sqrt{\lambda} \pi = 0 \Rightarrow \lambda = j^2, j = 1, 2, \dots$ So correspondingly eigenfunction is
(from prop of cosine)

$u_j = B \sin jx$. (not unique!)

• $\lambda = 0 \Rightarrow u'' = 0 \Rightarrow u = Ax + B$. Use ^(why do we mention form?) B.C. [?] $\Rightarrow A = B = 0$.

• $\lambda < 0 \Rightarrow u'' - \lambda u = 0 \Rightarrow u = A e^{+\sqrt{\lambda} x} + B e^{-\sqrt{\lambda} x} \Rightarrow A = B = 0$.

We find that eigenvalues $\lambda_j = j^2$, $j = 1, 2, \dots$ eigenfunctions
_{index}

$u_j = \sin jx$. Remember: ① Notice that u_j form a complete orthonormal

basis for $L^2([0, \pi]; \mathbb{R})$.

② the eigen value $\lambda_j \rightarrow \infty$ as $j \rightarrow \infty$.

③ ~~is~~ λ_j the "eigenvalue" of the differential operator?

(Unbounded operator).



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Q: How to find these for general Sturm-Liouville problems?

The idea is to look at the "inverse" L^{-1} of L . We show that L^{-1} is an integral operator (Green's function). It is self-adjoint, compact

hence we can apply spectral theory to get eigenvalues of L^{-1} .

Finally, we connect them to L .

I) Green's function

Wacky thought this might be useful

Consider the solution of the ODE problem

$$\begin{cases} Lw = g \\ \alpha w(a) + \alpha' w'(a) = 0 \\ \beta w(b) + \beta' w'(b) = 0 \end{cases} \quad (*)$$

The miracle of L is that given g we find w . There is a standard way to ~~find~~ solve this ODE called variation of parameters.

Let u, v be two linearly independent solutions of the homogeneous equation $Lu=0, Lv=0$. The general solution is $w = \alpha u + \beta v$.
 (is this something we know? or take this for granted)

Then we look for a solution of (*) of the form $w = \varphi(x)u(x) + \chi(x)v(x)$.

Where φ, χ to be determined (tbd)

Take derivative $\Rightarrow w' = \varphi u' + \chi v' + \varphi' u + \chi' v$

We choose $\varphi' u + \chi' v = 0$ and get $w' = \varphi u' + \chi v'$.

(Lw)

(how do we get this?)

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$$\text{Then } (pw')' + qw = \cancel{p'w} + p'u' + v'p'v' = q$$

$$\text{So } \begin{cases} p(u'u' + v'v') = q \end{cases}$$

$$\text{and } \begin{cases} u'u + v'v = 0 \end{cases}$$

$$\text{As } \begin{pmatrix} pu' & pv' \\ u & v \end{pmatrix} \Rightarrow \det = \frac{p(u'v - uv')}{\text{what this say?}}. \quad \left(\begin{array}{l} \text{this is value} \\ \text{Wronskian} \end{array} \right)$$

Prop: Suppose that $u, v \in C^2([a, b])$ are not identically zero.

Suppose that $Lu = 0, Lv = 0$. Then either $u \equiv cv$ for some constant c or $p(u'v - uv') \neq 0$ for all $t \in [a, b]$.

Moreover $p(u'v - uv') = \text{constant}$

~~Proof~~ (For this we need a existence result for ODE).

Proof: We consider

$$\begin{aligned} \frac{d}{dt}(p(u'v - v'u')) &= \frac{d}{dt}((pv')u - (pu')v) = (pv')'u + pv'u' \\ &\quad - pu'u' - (pw')v = +quv - quv = 0 \end{aligned}$$

So the quantity is a constant.

Now we consider the problem at $t_0 \in [a, b]$. If $p(u'v - uv')(t_0) \neq 0$ \Rightarrow nonvanishing for all t . If $p(u'v - uv')(t_0) = 0 \Rightarrow \vec{u} = \begin{pmatrix} u(t_0) \\ u'(t_0) \end{pmatrix}$ and $\vec{v} = \begin{pmatrix} v(t_0) \\ v'(t_0) \end{pmatrix}$

are linearly dependent $\Rightarrow \vec{u} = c\vec{v}$ for some constant c . Then

$$Lu = 0 \text{ satisfy } u(t_0) = cv(t_0), \quad u'(t_0) = cv'(t_0)$$

$$Lv = 0 \quad \dots \quad v(t_0) = v(t_0) \quad v'(t_0) = v'(t_0)$$

$\Rightarrow w = u - cv$ satisfy $Lw = 0, w(t_0) = 0, w'(t_0) = 0 \Rightarrow w = 0$ by uniqueness \square

(5)

Then we can solve for φ' , χ' :

$$\varphi' = -\frac{v g}{c}, \quad \chi' = \frac{u g}{c}$$

$$\text{So } \varphi(t) = \frac{1}{c} \left(\int_t^b v(z) g(z) dz + A \right)$$

$$\chi(t) = \frac{1}{c} \left(\int_a^t u(z) g(z) dz + B \right)$$

$$\text{Then the solution } w(t) = \varphi(t) u(t) + \chi(t) v(t)$$

Finally, we use Boundary condition to find A and B. (Actually, they are zero). If $A = B = 0$, we get

$$w(t) = \frac{1}{c} \int_t^b v(z) g(z) dz u(t) + \frac{1}{c} \int_a^t u(z) g(z) dz v(t)$$

We check BC:

$$\bullet \text{ at } t=a: w(a) = \frac{u(a)}{c} \int_a^b v(z) g(z) dz, \quad w'(a) = \frac{u'(a)}{c} \int_a^b g(z) v(z) dz$$

$$\text{Then } \alpha w(a) + \alpha' w'(a) = (\alpha u(a) + \alpha' u'(a)) \frac{1}{c} \int_a^b g(z) v(z) dz = 0$$

if u satisfies the BC at a .

$$\bullet \text{ at } t=b, w(b) = \frac{v(b)}{c} \int_a^b g(z) u(z) dz, \quad w'(b) = \frac{v'(b)}{c} \int_a^b g(z) u(z) dz$$

$$\Rightarrow \beta w(b) + \beta' w'(b) = (\beta v(b) + \beta' v'(b)) \int_a^b g(z) u(z) dz = 0$$

if v satisfies the BC at b .

⑧

If there are subyfund, we get

$$w(t) = \frac{1}{c} \int_a^b g(s) (H(s-t)u(t)v(s) - H(t-s)u(s)v(t)) ds$$

$$= \int_a^b k(s,t) g(s) ds$$

$$k(s,t) = \begin{cases} - & s < t \\ - & s > t \end{cases}$$

where $H(s)$ is the ~~Heaviside~~ Heaviside function

$$H(s) = \begin{cases} 0 & s < 0 \\ 1 & s > 0 \end{cases}$$

Here $k(s,t)$ is the integral kernel. It's a symmetric function, real valued, continuous on $[a,b] \times [a,b]$.

Now we need to check that the desired u, v can be found.

Thm: (Existence and uniqueness of ODE)

Let $g: [a,b] \rightarrow \mathbb{R}$ be a given C^0 function and suppose that

$u(t_0) = c_1, u'(t_0) = c_2$ at some $t_0 \in [a,b]$. Then there is

a unique solution $C^2([a,b])$ ~~with~~ of $Lu = g$ with $u(t_0) = c_1,$

~~u'(t_0) = c_2.~~

Using this result, we can find u with

$$\begin{cases} Lu = 0 \\ u(a) = c_1 \\ u'(a) = c_2 \end{cases}$$

such that $\alpha c_1 + \beta c_2 = 0!$ Similarly for v . Thus

shows that the inverse operator is found.