# Searches for cLFV at Current and Future Colliders

Michael A. Schmidt

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@ CLFV2019, Fukuoka, Japan



#### The Standard Model is very successful...

... but incomplete

In particular neutrinos are massive

- ightarrow Flavour changing processes are a sensitive probe
  - ullet in SM $+m_
    u$  suppressed by unitarity,  ${\cal A}\sim G_F m_
    u^2\simeq 10^{-26}$
  - many neutrino mass models have large charged LFV
     due to non-unitarity or new contributions,
     a single-constant radiative mass models.
  - e.g. inverse seesaw, radiative mass models
  - could be completely unrelated to neutrino mass, e.g. SUSY

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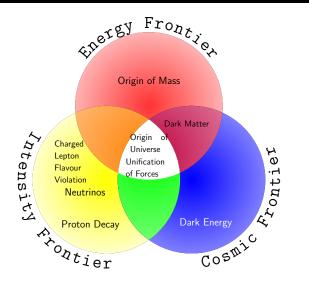
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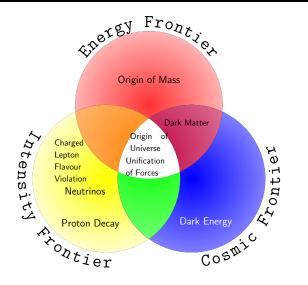
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Can high-energy colliders compete with the intensity frontier?



Can high-energy colliders compete with the intensity frontier?

#### **Overview**

Z boson decays

Higgs boson decay

Top-quark decay

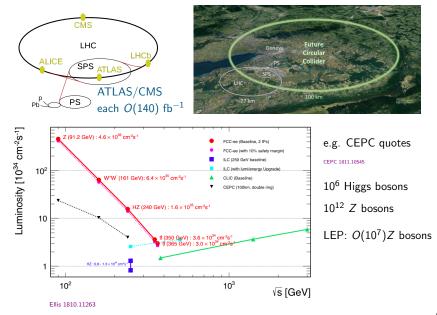
Heavy resonance decay

Scattering at the LHC

Scattering at future lepton colliders

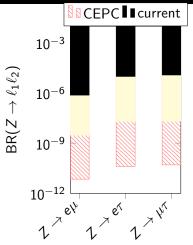
Conclusions

#### **Colliders**



### Z boson decays

#### cLFV Z boson decays



 $Z o e\mu$ : ATLAS 1408.5774, CMS EXO-13-005  $Z o \ell au$ : DELPHI  $(\mu au)$ , OPAL (e au) ATLAS, 13 TeV, 36.1 fb $^{-1}$  1804.09568 almost same sensitivity for  $\mu au$ 

No tree-level FCNC in SM induced at 1 loop in SM  $+m_{\nu}$ 

$$Z \sim V \qquad \ell_2^+ \propto \frac{G_F m_\nu^2}{16\pi^2} \simeq 10^{-28}$$

Observation clear sign of new physics e.g. due to a leptoquark

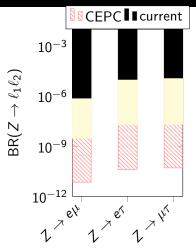
$$Z \longrightarrow \begin{cases} q' & \ell_2^+ \\ \phi & \ell_1^- \end{cases}$$

today typically less stringent as low-energy precision experiments

but will be more interesting with new Z boson factory

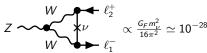
or if there is a signal to disentangle physics

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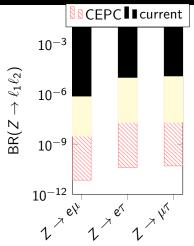
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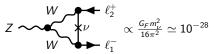
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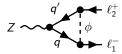


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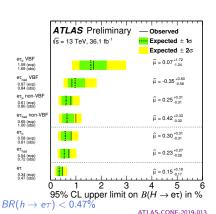
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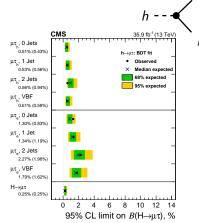
### Higgs boson decay

#### cLFV Higgs decay

Dimension-6 SMEFT operators Grzadkowski et al 1008.4884

$$\mathcal{L} = \left[ Y_{ij} + \frac{c_{ij}}{\Lambda^2} \left( H^{\dagger} H \right) \right] \ \bar{L}_i P_R \ell_j H + h.c. \rightarrow \left[ \frac{m_{ij}}{v} + \frac{c_{ij}}{\sqrt{2}} \frac{v^2}{\Lambda^2} \right] h \bar{\ell}_i P_R \ell_j + h.c.$$

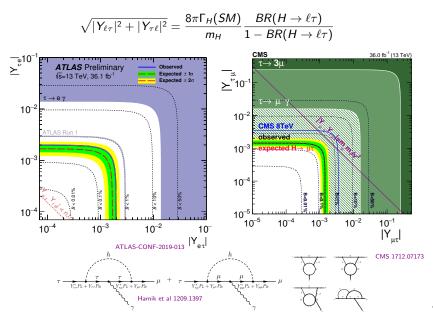




 $BR(h \rightarrow \mu \tau) < 0.25\%$ 

CMS 1712.07173

#### cLFV Higgs decay cont.



#### General (type-III) 2 Higgs doublet model

#### **EFT**

$$\mathcal{L} = \left[\frac{m_i}{v}\delta_{ij} + \frac{c_{ij}}{\sqrt{2}}\frac{v^2}{\Lambda^2}\right]h\bar{\ell}_i P_R \ell_j$$

two neutral CP even Higgs

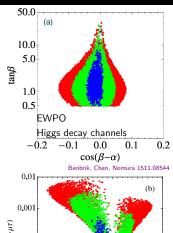
$$\Phi_i = (v_i + \phi_i)/\sqrt{2} \qquad \frac{v_2}{v_1} = t_\beta$$

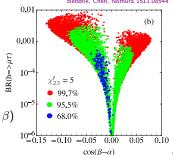
SM Higgs:  $h=-s_{\alpha}\phi_{1}+c_{\alpha}\phi_{2}$  with Yukawa couplings

$$Y_{ij} = -\frac{s_{\alpha}}{c_{\beta}} \frac{m_i}{v} \delta_{ij} + \frac{\cos(\beta - \alpha)}{c_{\beta}} \frac{\sqrt{m_i m_j}}{v} \chi_{ij}^{\ell}$$

Not suppressed by  $v^2/\Lambda^2 \rightarrow \text{large contribution}$ 

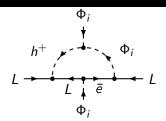
$$BR(h o \mu au) \propto \left(\left|\chi_{23}^\ell\right|^2 + \left|\chi_{32}^\ell\right|^2\right) \cos^2(eta - lpha) (1 + an^2eta)$$

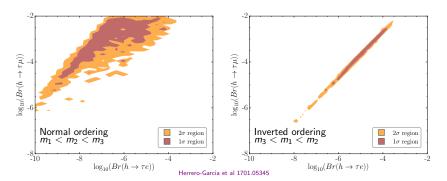




#### Example: Zee model

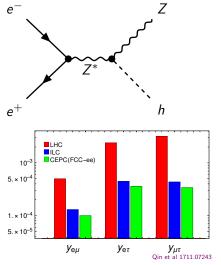
- Non-zero neutrino masses
- generated at loop level Zee 1980
- Simplest model with 2 Higgs doublets and charged singlet scalar h<sup>+</sup>





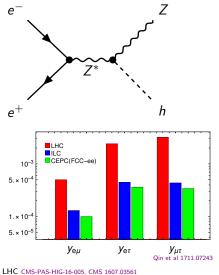
[see Herrero-Garcia et al 1605.06091 for Higgs cLFV in other neutrino mass models]

#### Future lepton collider

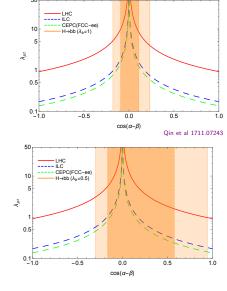


LHC cMs-PAS-HIG-16-005, CMS 1607.03561 ILC  $\sqrt{s}=250$  GeV, 4 polarizations,  $\mathcal{L}=2$  ab  $^{-1}$  CEPC  $\sqrt{s}=240$  GeV,  $\mathcal{L}=5$  ab  $^{-1}$ 

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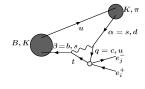
Top-quark decay

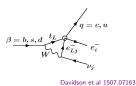
#### cLFV top-quark [Davidson et al 1507.07163]

described by D6 operators with 1 top quark and 2 charged leptons

$$\mathcal{L} = 2\sqrt{2} G_F \sum_i \epsilon_i \mathcal{O}_i$$

e.g. 
$$\mathcal{O}_{LL,RR,LR,RL}^{AV} = (\bar{\ell}_i \gamma^{\alpha} P_X \ell_j) (\bar{u}_q \gamma_{\alpha} P_Y t)$$





- HERA  $\sigma(e^{\pm}p \rightarrow e^{\pm}t + X) \leq 0.3pb$
- $\bullet \ \ \textit{K} \rightarrow \textit{e}\mu \textrm{, } \mu \rightarrow \textit{e}\gamma$
- radiative corrections

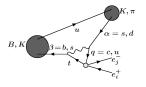
$$e\mu$$
 op's: most  $|\epsilon|\lesssim O(10^{-3}-10^{-2})$ , some  $O(1)$   $au\ell$  op's  $O(1-100)$   $|\epsilon^{ut}_{S+P,L}|\leq$  0.03

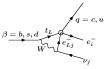
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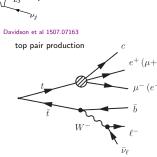
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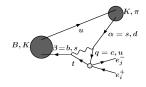


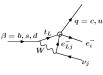
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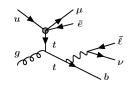


Davidson et al 1507.07163

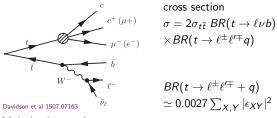
single top quark production (more diag's)

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#### cLFV top quark decay: top-quark pair production



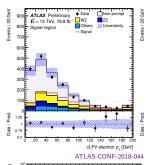
Main backgrounds:

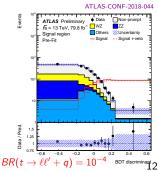
- ullet  $tar{t}$  with non-prompt lepton
- *Z*+ jets

Multi-variate analysis w/ 14 var's using BDT observed [expected] limit

$$BR(t \to \ell \ell' q) < 1.86[1.36^{+0.61}_{-0.37}] \times 10^{-5}$$
  
 $BR(t \to e\mu q) < 6.6[4.8^{+2.1}_{-1.4}] \times 10^{-5}$ 

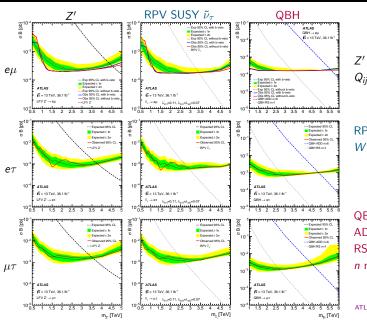
 $\rightarrow |\epsilon| \lesssim 0.1$ , more stringent for  $t \rightarrow \tau + X$  low-energy lim's stronger for most  $e\mu$  op's:  $\epsilon_{LL,RL}$ ,  $\epsilon_{S\pm P,R}$ ,  $\epsilon_{T,R}$ 





## Heavy resonance decay

### Heavy resonance: Z', RPV SUSY $\tilde{\nu}_{\tau}$ , quantum black hole



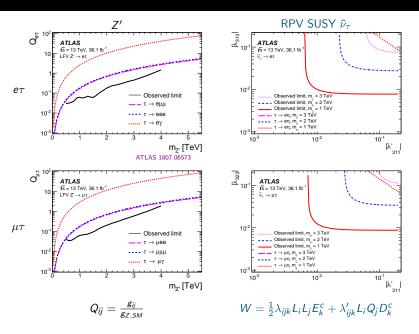
$$Z'$$
 $Q_{ij} = \frac{g_{ij}}{g_{Z,SM}}$ 

 $\begin{aligned} & \mathsf{RPV} \; \mathsf{SUSY} \; \tilde{\nu}_{\tau} \\ & W = \frac{1}{2} \lambda_{ijk} L_i L_j E_k^c \\ & + \lambda'_{ijk} L_i Q_j D_k^c \end{aligned}$ 

QBH
ADD (universal ED)
RS (warped ED)
n number of ED

ATLAS 1807.06573

#### Heavy resonance: Z', RPV SUSY $\tilde{\nu}_{\tau}$ cont.



### Scattering at the LHC

#### Relevant effective operators [Cai, MS 1510.02486]

#### D6 Operators with 2 Quarks and 2 Leptons

Buchmüller, Wyler NPB268(1986)621; Grzadkowski et al 1008.4884; Carpentier, Davidson 1008.0280; Petrov, Zhuridov 1308.6561

#### Vector

$$\begin{aligned} \mathcal{Q}_{lq}^{(1)} &= (\bar{L}\gamma_{\mu}L)(\bar{Q}\gamma^{\mu}Q) & \mathcal{Q}_{lq}^{(3)} &= (\bar{L}\gamma_{\mu}\tau^{I}L)(\bar{Q}\gamma^{\mu}\tau^{I}Q) \\ \mathcal{Q}_{eu} &= (\bar{\ell}\gamma_{\mu}\ell)(\bar{u}\gamma^{\mu}u) & \mathcal{Q}_{ed} &= (\bar{\ell}\gamma_{\mu}\ell)(\bar{d}\gamma^{\mu}d) \\ \mathcal{Q}_{lu} &= (\bar{L}\gamma_{\mu}L)(\bar{u}\gamma^{\mu}u) & \mathcal{Q}_{ld} &= (\bar{L}\gamma_{\mu}L)(\bar{d}\gamma^{\mu}d) \\ \mathcal{Q}_{qe} &= (\bar{Q}\gamma_{\mu}Q)(\bar{\ell}\gamma^{\mu}\ell) & \end{aligned}$$

$$\mathcal{Q}_{ledq} = (\bar{L}^{lpha}\ell)(\bar{d}Q^{lpha})$$

$$\mathcal{Q}_{lequ}^{(1)} = (\bar{L}^{\alpha}\ell)\epsilon_{\alpha\beta}(\bar{Q}^{\beta}u)$$

with same-flavour quark

$$\mathcal{Q}_{lequ}^{(3)} = (\bar{L}^{lpha} \sigma_{\mu
u} \ell) \epsilon_{lphaeta} (\bar{Q}^{eta} \sigma^{\mu
u} u)$$

#### D8 Operators with 2 Gluons and 2 Leptons

$$\mathcal{O}_{X}^{ij} = \alpha_{s} G_{\mu\nu}^{a} G^{a\mu\nu} \left( \bar{e}_{Ri} L_{j} \cdot \phi^{*} + h.c. \right) \qquad \mathcal{O}_{X}^{\prime ij} = i \alpha_{s} G_{\mu\nu}^{a} \tilde{G}^{a\mu\nu} \left( \bar{e}_{Ri} L_{j} \cdot \phi^{*} - h.c. \right)$$

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$$\mathcal{O}_{Y}^{ij} = i \alpha_{s} G_{\mu\nu}^{a} G_{\sigma\nu}^{a} \eta^{\rho\sigma} \bar{e}_{Ri} \gamma^{\mu} D^{\nu} L_{i} \qquad \mathcal{O}_{Z}^{ij} = i \alpha_{s} G_{\mu\nu}^{a} G_{\sigma\nu}^{a} \eta^{\rho\sigma} \bar{e}_{Ri} \gamma^{\mu} D^{\nu} e_{Ri}$$

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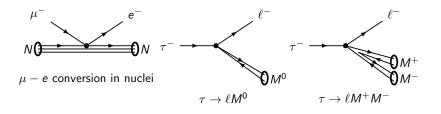
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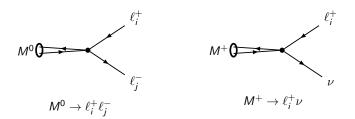
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#### D8 Operators with 2 Gluons and 2 Leptons

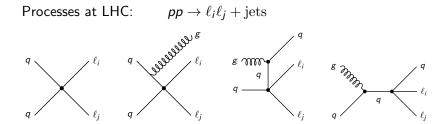
$$\begin{split} \mathcal{O}_{X}^{ij} &= \alpha_{s} G_{\mu\nu}^{a} G^{a\mu\nu} \left( \bar{e}_{Ri} L_{j} \cdot \phi^{*} + h.c. \right) & \mathcal{O}_{X}^{\prime ij} &= i \, \alpha_{s} G_{\mu\nu}^{a} \, \tilde{G}^{a\mu\nu} \left( \bar{e}_{Ri} L_{j} \cdot \phi^{*} - h.c. \right) \\ \bar{\mathcal{O}}_{X}^{ij} &= i \, \alpha_{s} G_{\mu\nu}^{a} G^{a\mu\nu} \left( \bar{e}_{Ri} L_{j} \cdot \phi^{*} - h.c. \right) & \bar{\mathcal{O}}_{X}^{\prime ij} &= \alpha_{s} G_{\mu\nu}^{a} \, \tilde{G}^{a\mu\nu} \left( \bar{e}_{Ri} L_{j} \cdot \phi^{*} + h.c. \right) \\ \mathcal{O}_{Y}^{ij} &= i \, \alpha_{s} G_{\mu\rho}^{a} G_{\sigma\nu}^{a} \eta^{\rho\sigma} \bar{L}_{i} \gamma^{\mu} D^{\nu} L_{j} & \mathcal{O}_{Z}^{ij} &= i \, \alpha_{s} G_{\mu\rho}^{a} G_{\sigma\nu}^{a} \eta^{\rho\sigma} \bar{e}_{Ri} \gamma^{\mu} D^{\nu} e_{Rj} \end{split}$$

#### Precision Experiments [Cai, MS 1510.02486]





### cLFV at the Large Hadron Collider (LHC) [Cai, MS 1510.02486]



### Signal: opposite-sign different flavour pair of leptons

#### Several existing searches:

- ullet ATLAS 7 TeV: LFV heavy neutral particle decay to  $e\mu$  ATLAS 1103.5559
- ullet CMS 8 TeV: LFV heavy neutral particle decay to  $e\mu$  CMS-PAS-EXO-13-002
- ATLAS 7 TeV: LFV in eμ continuum in R SUSY<sub>ATLAS</sub> 1205.0725
- ATLAS 8 TeV: LFV heavy neutral particle decayatlas 1503.04430
- ullet CMS 8 TeV: LFV heavy neutral particle decay to  $e\mu$  cms 1604.05239
- ATLAS 13 TeV, 3.2 fb<sup>-1</sup>: LFV heavy neutral particle decay ATLAS 1607.08079
- ATLAS 13 TeV, 36.1 fb<sup>-1</sup> atlas 1807.06573

#### Interesting ATLAS searches [Cai, MS 1510.02486]

#### Recast limits of most sensitive previous searches

ATLAS 1503.04430	ATLAS 1205.0725
8 TeV	7 TeV
$20.3~{ m fb}^{-1}$	$2.1 \; { m fb}^{-1}$
$e\mu$ , $e au$ , $\mu au$	e $\mu$
inclusive	exclusive
including arbitrary number of jets	separated by number of jets

#### Projection to 14 TeV

- Assuming 300 fb<sup>-1</sup>
- Follow searching strategy of exclusive 7 TeV search

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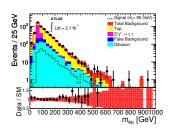
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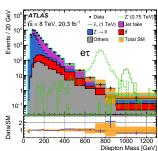
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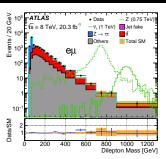
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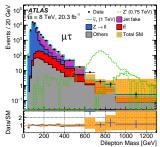
### ATLAS Searches [Cai, MS 1510.02486]



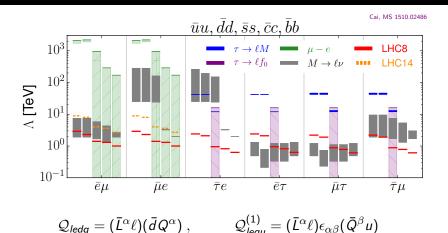
#### ATLAS 7TeV 1205.0725







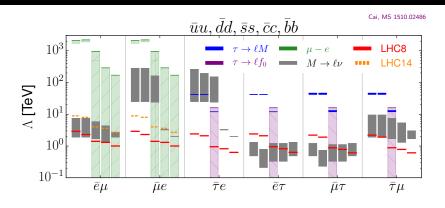
## cLFV at hadron colliders: quarks



LHC more interesting for vector operators with right-handed quark currents due to weaker constraints from intensity frontier

$$[\bar{q}\gamma_{\mu}P_{R}q][\bar{\ell}\gamma_{\mu}P_{R,L}\ell$$

## cLFV at hadron colliders: quarks



$$\mathcal{Q}_{ledq} = (\bar{L}^{\alpha}\ell)(\bar{d}Q^{\alpha}) \;, \qquad \quad \mathcal{Q}_{lequ}^{(1)} = (\bar{L}^{\alpha}\ell)\epsilon_{\alpha\beta}(\bar{Q}^{\beta}u)$$

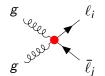
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#### cLFV at the Large Hadron Collider (LHC): gluons [Cai, MS, Valencia 1802.09822]

Processes at LHC:

$$pp \to \ell_i \ell_j$$



## Signal:

opposite-sign different flavour pair of leptons

#### Most sensitive searches

ATLAS 1607.08079 CMS-PAS-EXO-16-058 1802.01122

13 TeV 13 TeV 3.2 fb<sup>-1</sup> 35.9 fb<sup>-1</sup>

 $e\mu$ ,  $e\tau$ ,  $\mu\tau$   $e\mu$  inclusive inclusive

newer ATLAS search: 13 TeV, 36.1 fb<sup>-1</sup> 1807.06573

EFT scattering amplitudes

$$\mathcal{A}(s) \simeq \frac{s}{\Lambda^2} \stackrel{s \to \infty}{\longrightarrow} \infty$$

 $\Rightarrow$   $\mathsf{Violation}$  of  $\mathsf{perturbative}$   $\mathsf{unitarity}$ 

#### Solutions

- UV-complete models/simplified models
- apply unitarization procedure, e.g.

Wigner 1964; Wigner, Eisenbud 1947; Gupta 1950

Recent application to monojets: Bell, Busoni, Kobakhidze, Long, MS 1606.027

ullet couplings o form factor

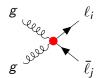
Baur, Zeppenfeld hep-ph/9309227

$$C 
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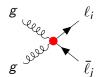
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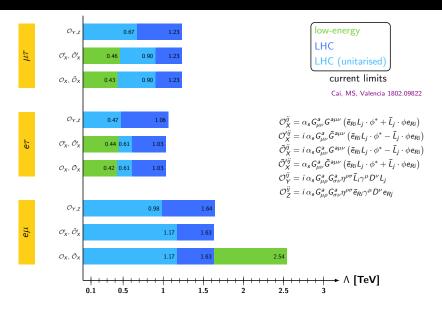
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### cLFV at hadron colliders: gluons



See also Bhattacharya et al 1802.06082 for a related analysis

Scattering at future lepton colliders

## Bileptons - seven simplified models [Li,MS 1809.07924]

$$\Delta L = 0$$
 complex scalar  $H_2 \sim (2, \frac{1}{2})$ 

$$\mathcal{L} = y_2^{ij} \frac{\mathsf{H}_2}{\mathsf{L}_i} P_R \ell_j + h.c.$$

LH singlet vector  $H_1 \sim (1,0)$ 

$$\mathcal{L} = y_1^{ij} \frac{\mathbf{H}_{1\mu}}{\mathbf{L}_i} \bar{\mathbf{L}}_i \gamma^{\mu} P_L \mathbf{L}_j$$

LH triplet vector  $H_3 \sim (3,0)$ 

$$\mathcal{L} = y_3^{ij} \bar{L}_i \gamma^\mu \vec{\sigma} \cdot {\color{black} H_{3\mu}} P_L L_j$$

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 right-handed scalar  $\Delta_1 \sim (1,2)$ 

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assumption: real and symmetric Yukawa coupling matrices

related work: Dev, Mohapatra, Zhang 1711.08430, also 1712.03642, 1803.11167

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vector  $\Delta_2 \sim \left(2,\frac{3}{2}\right)$ 

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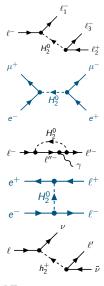
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## Existing (low-energy) precision constraints [LI,MS 1809.07924]

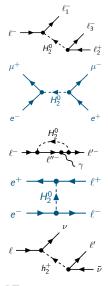
- $\bullet$  LFV trilepton decays,  $\ell \to \ell_1 \bar{\ell}_2 \bar{\ell}_3$
- Muonium antimuonium conversion,  $\mu^+e^-\to\mu^-e^+$
- ullet anomalous magnetic (and electric) dipole moments,  $a_\ell$
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- ullet lepton flavour non-universality,  $\ell o \ell' 
  u ar{
  u}$



Future sensitivity improvements at e.g. Belle 2, Mu3E, ...

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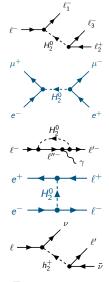
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  u ar{
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## Off-shell production $H_{1\mu}$ : $e^+e^- o e^\pm\mu^\mp(e^\pm au^\mp)$ [Li,MS 1809.07924]

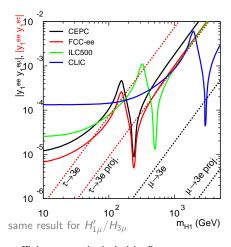
$$\mathcal{L} = y_1^{ij} H_{1\mu} \bar{L}_i \gamma^{\mu} P_L L_j$$

$$e^+ \qquad \qquad e^- \qquad e^+ \qquad \qquad \ell^+$$

$$e^- \qquad \qquad \ell^+ \qquad \qquad e^-$$

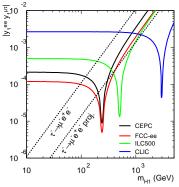
Basic cuts:  $p_T > 10$  GeV and  $|\eta| < 2.5$ 

Four collider configurations: CEPC:  $5 \text{ ab}^{-1}$  at 240 GeV FCC-ee:  $16 \text{ ab}^{-1}$  at 240 GeV ILC500:  $4 \text{ ab}^{-1}$  at 500 GeV CLIC:  $5 \text{ ab}^{-1}$  at 3 TeV

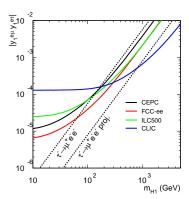


 $\tau$  efficiency not included in figure 60%  $\tau$  eff.  $\Rightarrow$  77% sensitivity reduction for 1  $\tau$ 

## $H_{1\mu}\colon e^+e^- o \mu^\pm au^\mp$ [Li,MS 1809.07924]



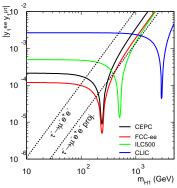
rel. couplings  $|y_1^{ee}y_1^{\mu\tau}|$   $e^+$   $\mu^\pm$ 



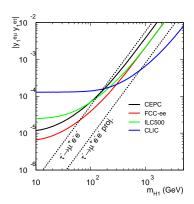
rel. couplings 
$$|y^{e\mu}y^{e\tau}|$$

$$e^+ \biguplus \mu^+ \qquad e^+ \biguplus \tau^+ \\ e^- \biguplus \tau^- \qquad e^- \biguplus \mu^-$$

## $extcolor{H}_{1\mu}$ : $e^+e^ightarrow\mu^\pm au^\mp$ [Li,MS 1809.07924]



rel. couplings  $|y_1^{ee}y_1^{\mu\tau}|$   $e^+$   $\mu^{\pm}$ 

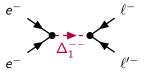


rel. couplings 
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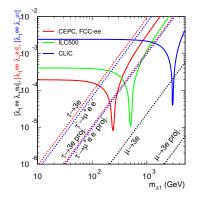
$$e^+ \xrightarrow[H_1]{} \mu^+ \qquad e^+ \xrightarrow[H_1]{} \tau^-$$

$$e^- \xrightarrow[]{} \tau^- \qquad e^- \xrightarrow[]{} \mu^-$$

## Same-sign lepton collider - $\Delta_1$ : $e^-e^- o \ell^-\ell'^-$ [Li,MS 1809.07924]



relevant couplings  $|\lambda_1^{ee}\lambda_1^{e\ell}| \text{ and } |\lambda_1^{ee}\lambda_1^{\mu\tau}|$ 



smaller integrated luminosity  $\mathcal{L} = 500\,\mathrm{fb}^{-1}$ 

## On-shell production $H_{1\mu}$ : $e^+e^- o e^\pm\mu^\mp(e^\pm au^\mp)+H_1$ [Li,MS in preparation]

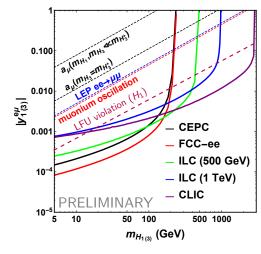
$$\begin{split} \mathcal{L} = & y_1^{ij} \mathbf{H}_{1\mu} \bar{L}_i \gamma^{\mu} P_L L_j \\ &+ y_3^{ij} \bar{L}_i \gamma^{\mu} \vec{\sigma} \cdot \mathbf{H}_{3\mu} P_L L_j \end{split}$$



Cuts:  $p_T > 10$  GeV and  $|\eta| < 2.5$ 

Five collider configurations: CEPC:  $5 \text{ ab}^{-1}$  at 240 GeV FCC-ee:  $16 \text{ ab}^{-1}$  at 240 GeV ILC (500 GeV):  $4 \text{ ab}^{-1}$  at 500 GeV ILC (1TeV):  $1 \text{ ab}^{-1}$  at 1 TeV

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## Conclusions

#### **Conclusions**

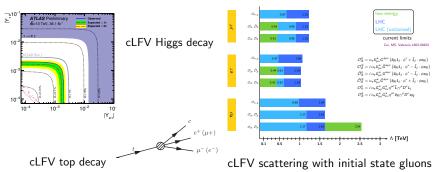
#### colliders complementary way to search for charged LFV

 $\mu \leftrightarrow e$  flavour: stringent limits from low-energy precision exp.

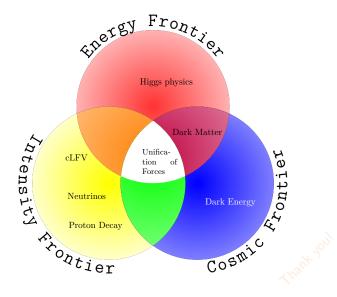
 $au \leftrightarrow \ell$  flavour complementary sensitivity at colliders

#### colliders test more Lorentz structures

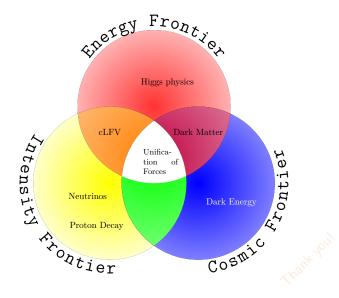
best for operators which are difficult to constrain at low energy



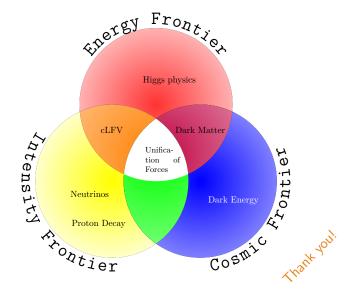
#### Conclusions cont.



#### Conclusions cont.



#### Conclusions cont.



# Backup slides

## cLFV D8 operator with 2 gluons and 2 leptons

process	exp. limit	operator	Λ [TeV]	
eμ				
$Br(\mu^{-\frac{48}{22}Ti} \to e^{-\frac{48}{22}Ti})$	$< 4.3 \times 10^{-12}$	$O_X, \bar{O}_X$	2.11	
${\sf Br}(\mu^{-197}_{79}{ m Au} o e^{-197}_{79}{ m Au})$	$<7\times10^{-13}$	$\mathcal{O}_X$ , $\bar{\mathcal{O}}_X$	2.54	
eτ				
$Br( au^+ o e^+\pi^+\pi^-)$	$< 2.3 \times 10^{-8}$	$O_X, \bar{O}_X$	0.42	
${\sf Br}( au^- o e^-{\sf K}^+{\sf K}^-)$	$< 3.4  imes 10^{-8}$	$\mathcal{O}_X$ , $\bar{\mathcal{O}}_X$	0.37	
$Br( au^- o e^-\eta)$	$<9.2\times10^{-8}$	$\mathcal{O}_X',\bar{\mathcal{O}}_X'$	0.40	
$Br(\tau^- \to e^- \eta')$	$<1.6\times10^{-7}$	$O'_X, \bar{O}'_X$	0.44	
μτ				
$Br(\tau^-  o \mu^- \pi^+ \pi^-)$	$<2.1\times10^{-8}$	$O_X, \bar{O}_X$	0.43	
$\mathrm{Br}( au^-  o \mu^- K^+ K^-)$	$<4.4\times10^{-8}$	$\mathcal{O}_X$ , $\bar{\mathcal{O}}_X$	0.36	
$Br( au^-  o \mu^- \eta)$	$<6.5\times10^{-8}$	$\mathcal{O}_X',\bar{\mathcal{O}}_X'$	0.42	
$Br(\tau^- \to \mu^- \eta')$	$< 1.3 \times 10^{-7}$	$O_X', \bar{O}_X'$	0.46	