

Gravitational Waves from Cosmological Phase Transitions

Cyril Lager

2019 Sydney Particle Physics and Cosmology meeting

Introduction

What are Gravitational Waves?

- A prediction of General Relativity and Einstein Field Equations:

$$8\pi G T_{\mu\nu} = G_{\mu\nu}$$

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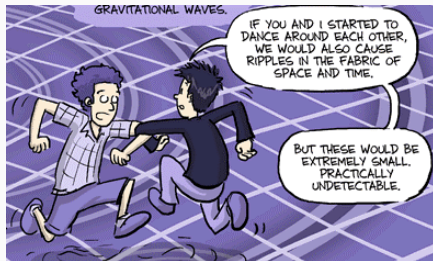
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- Gravitational wave source

Astrophysical?
Cosmological?
...?

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- Very weak coupling
- Gravitational wave source
- Gravitational wave properties

Frequency - Amplitude?
Transient or stochastic?

Detector design?

Gravitational Wave (GW) detections by LIGO/Virgo are promising for theoretical physics:

- confirm a prediction of General Relativity
- allow us to test GR (and its modifications) in a strong and dynamical regime
- suggest to look for other sources of GWs in relation to particle physics: phase transitions, inflation cosmic strings,...

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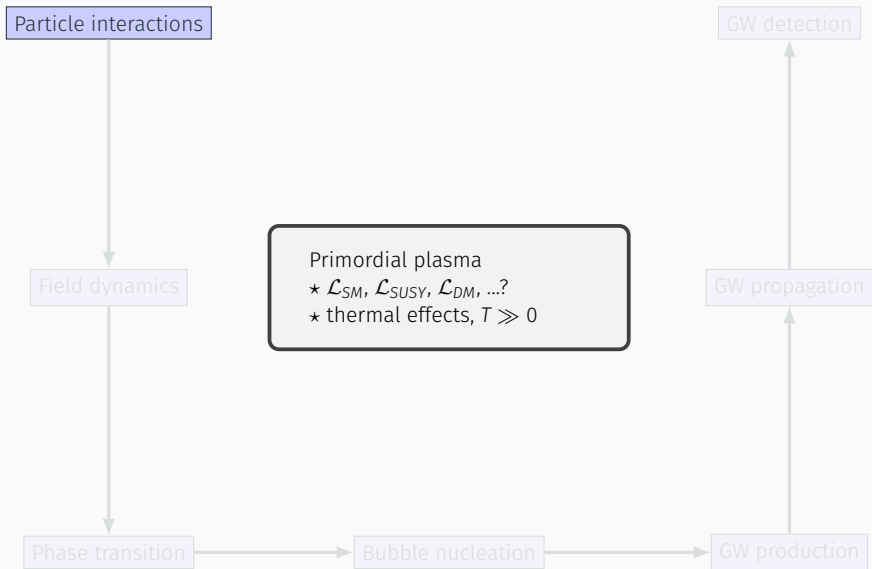
Main topic of this talk:

- exploring beyond the Standard Model physics with GWs from phase transitions

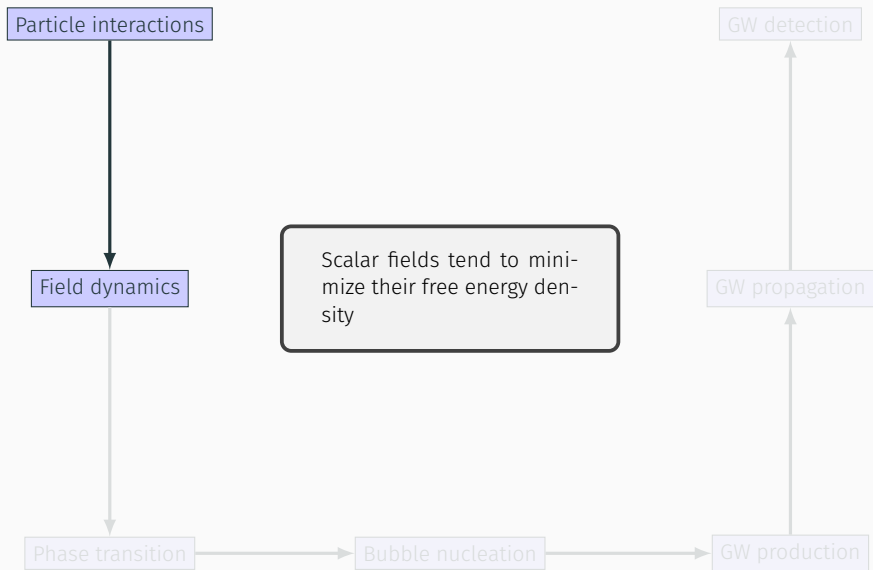
From quantum fields to Gravitational Waves!



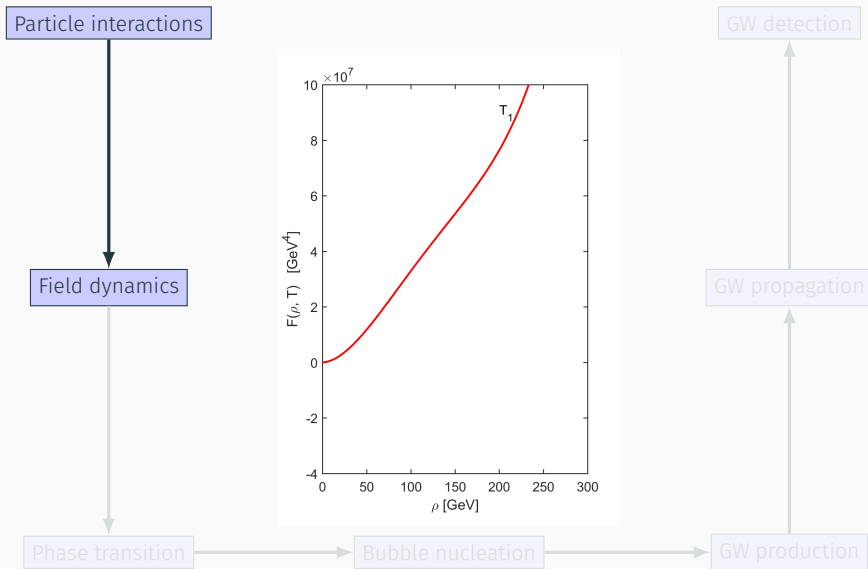
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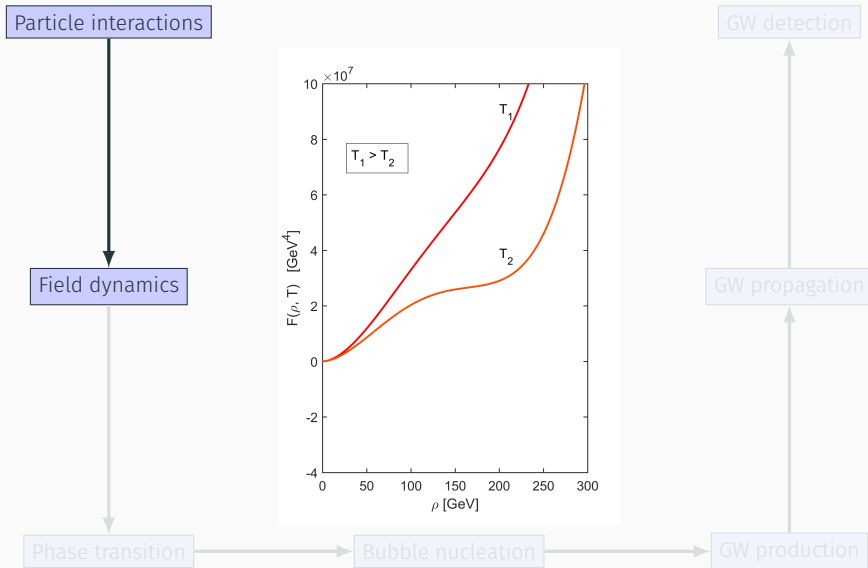
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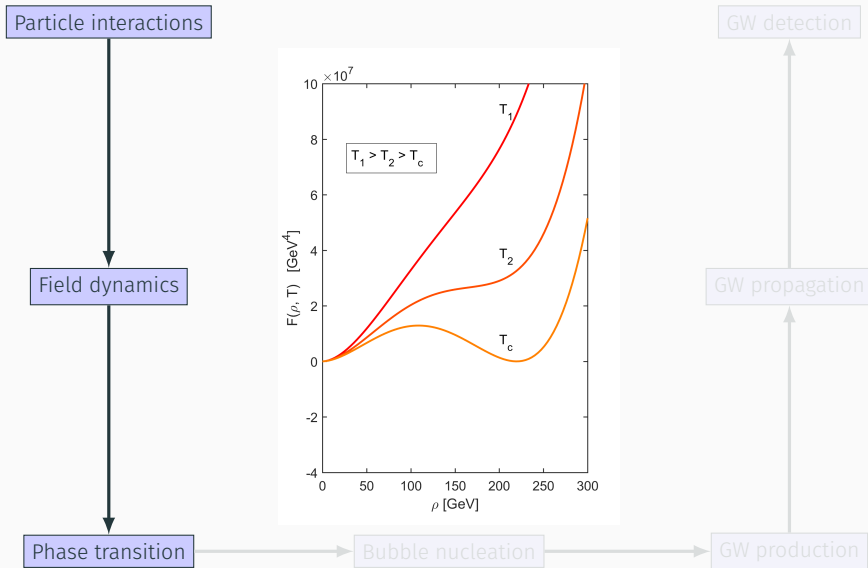
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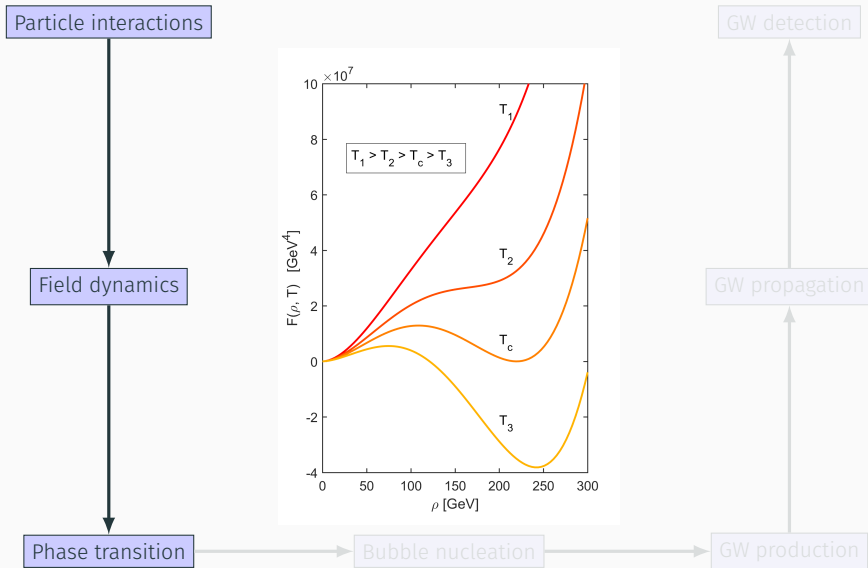
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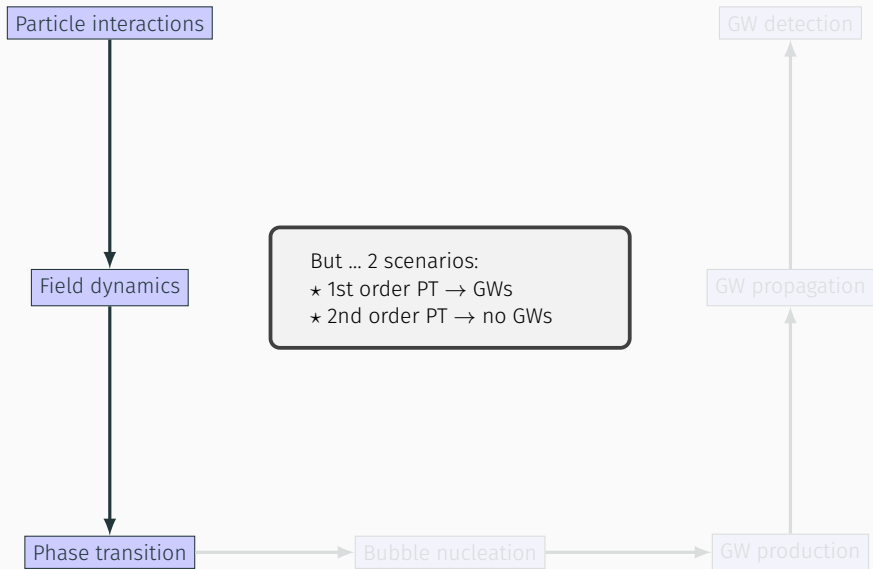
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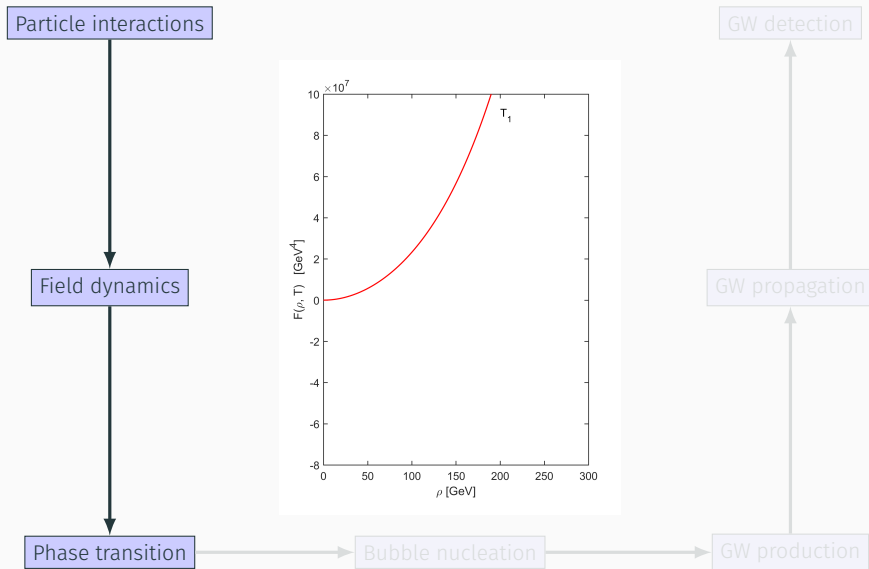
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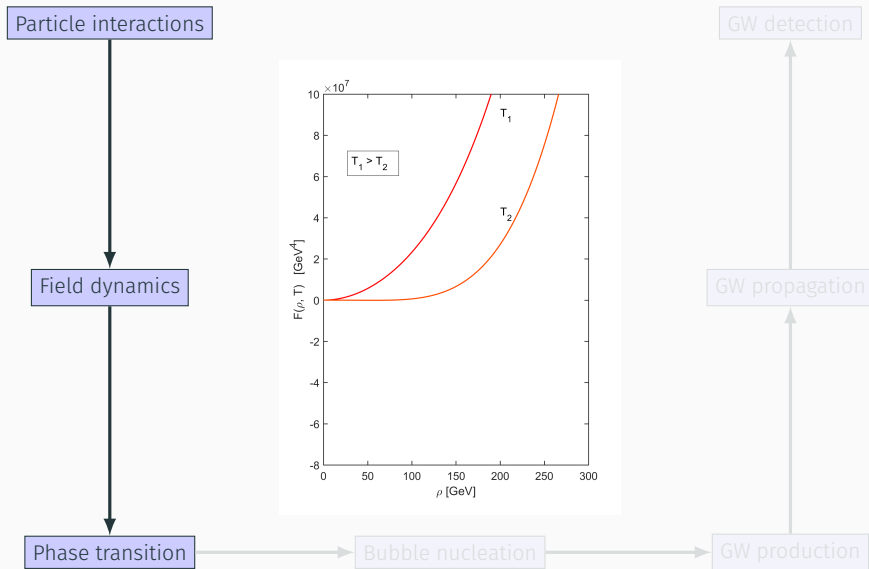
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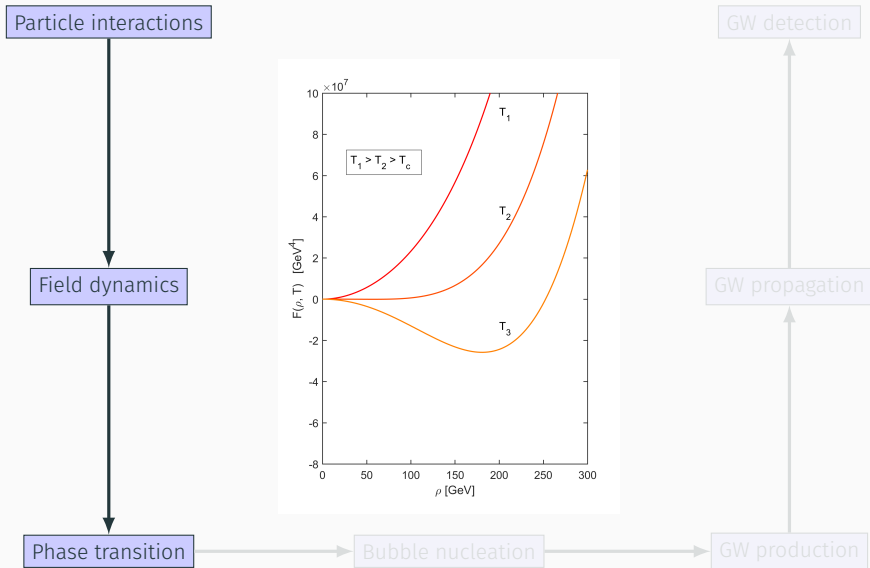
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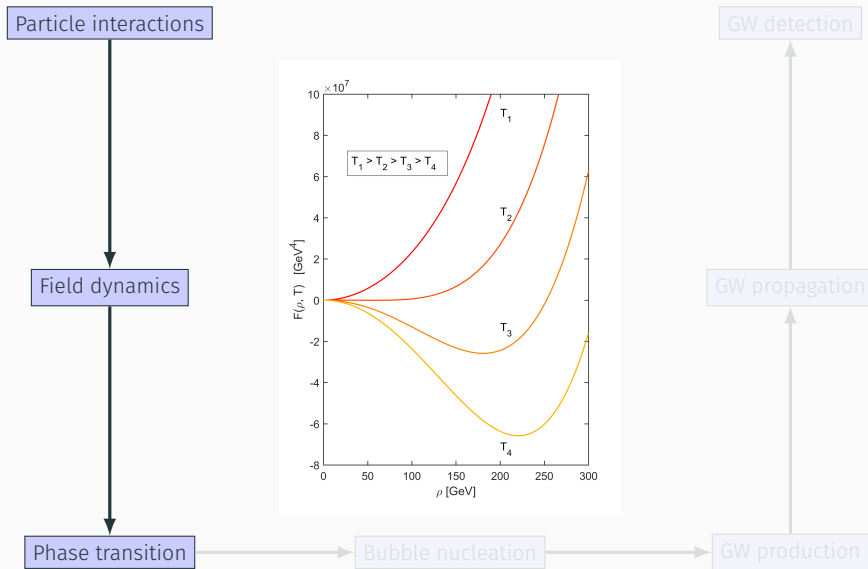
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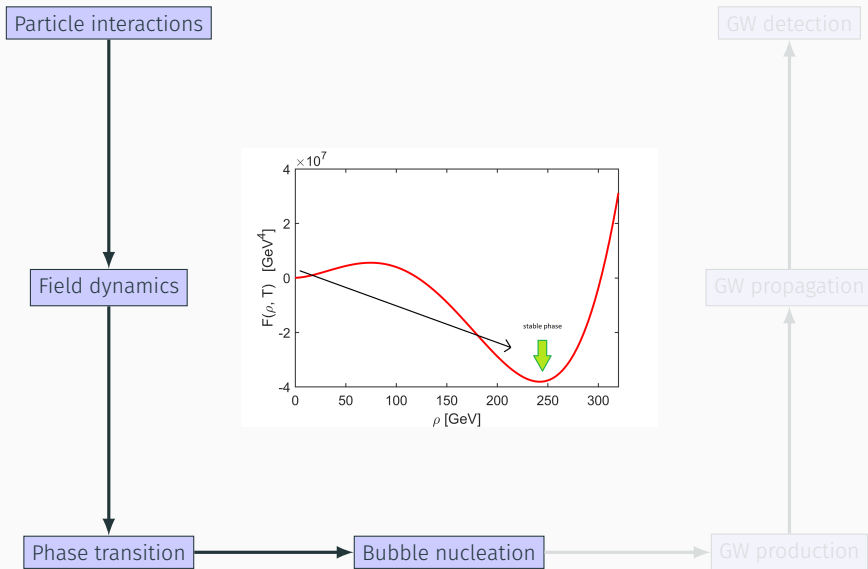
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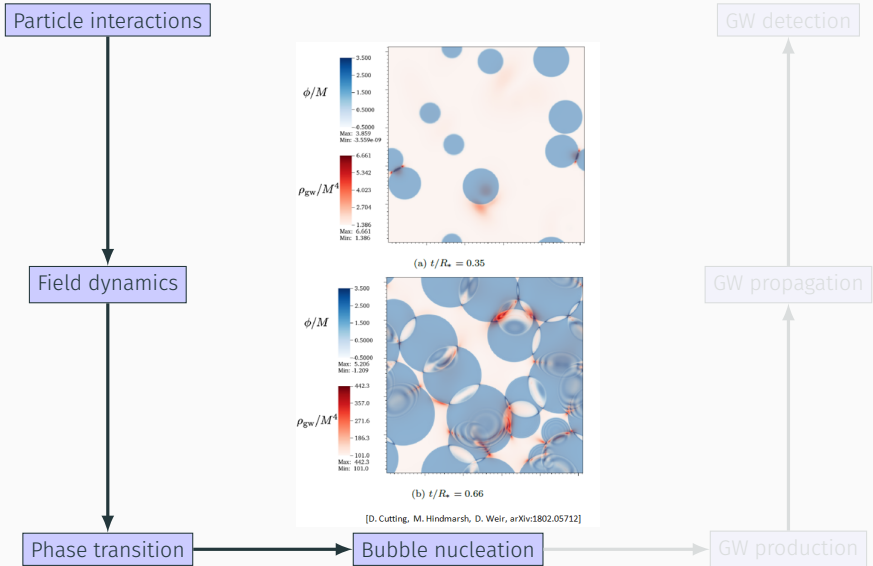
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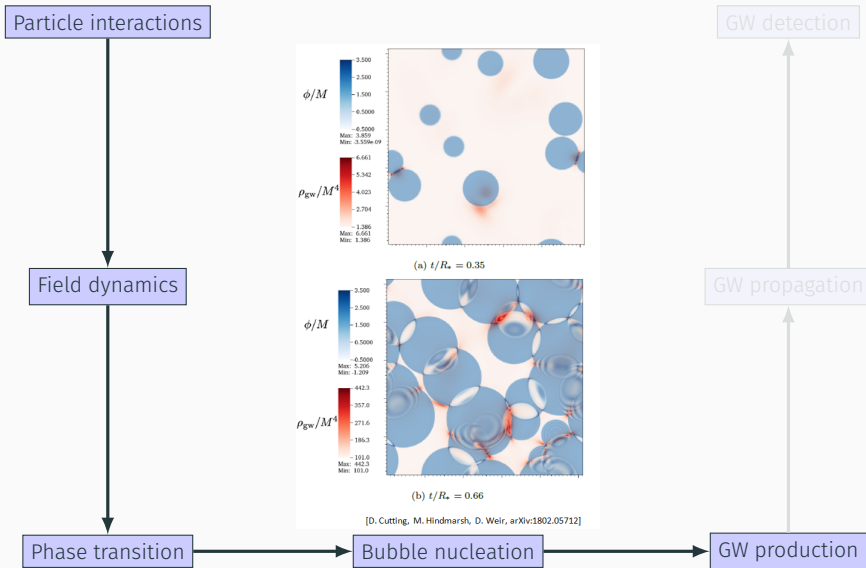
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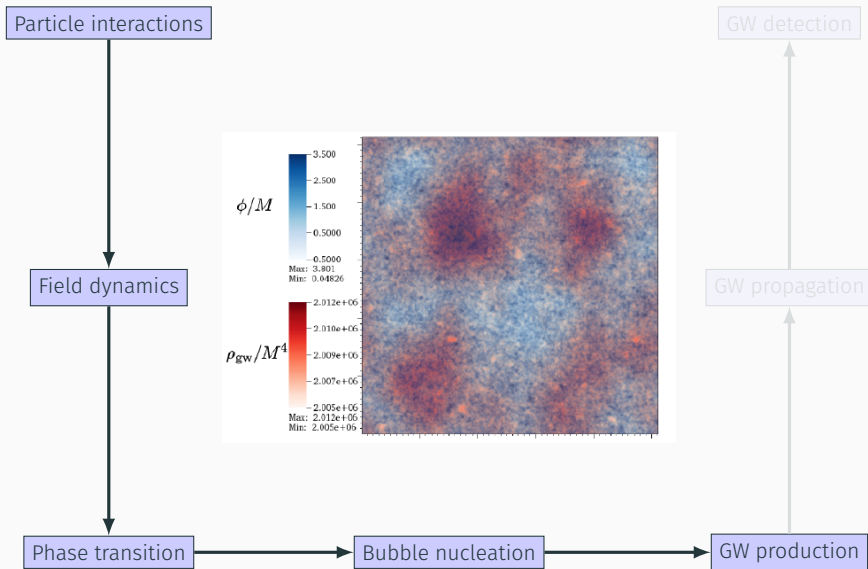
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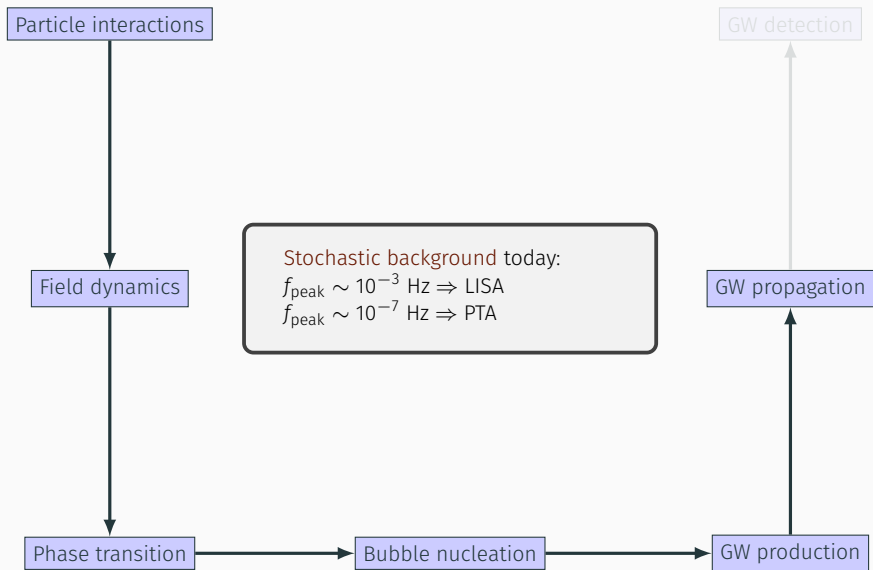
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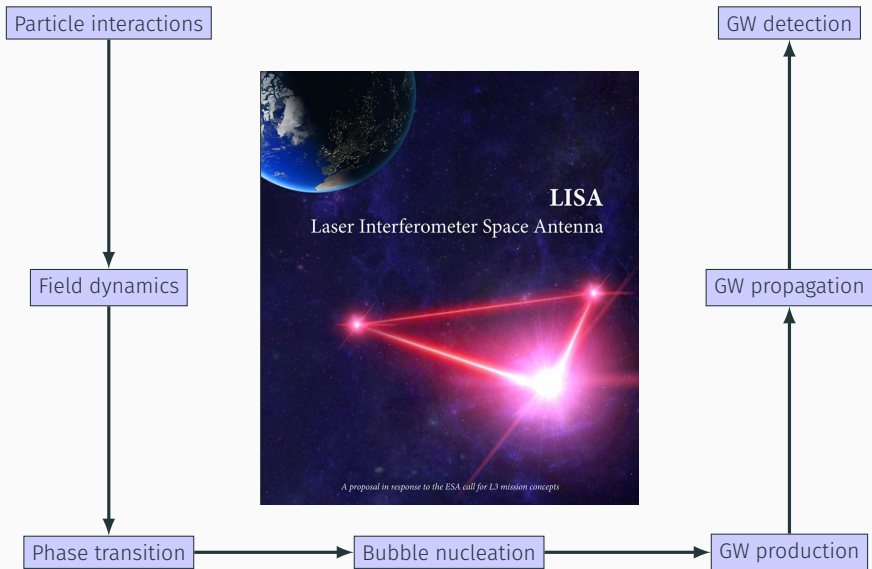
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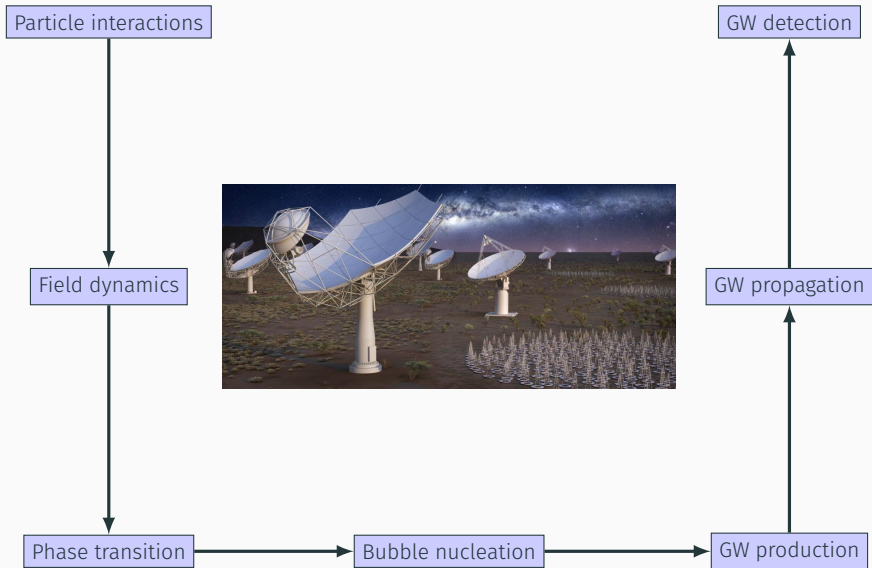
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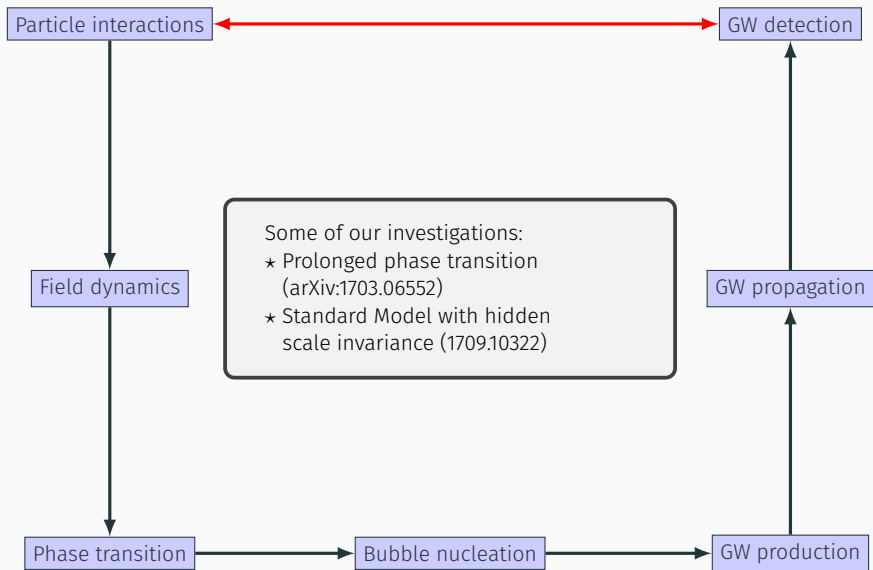
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GW background from bubble collisions

Stochastic background from three sources [C. Caprini et al., JCAP 1604 (2016) no.04 001]:

$$h^2\Omega_{\text{GW}}(f) \simeq h^2\Omega_{\text{col}} + h^2\Omega_{\text{sw}} + h^2\Omega_{\text{MHD}}$$

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Peak frequency and amplitude of the background mainly depend on the mean bubble size \bar{R} at collision and kinetic energy ρ_{kin} stored in the bubbles:

- $f_{\text{peak}} \sim (\bar{R})^{-1}$
- $\Omega_{\text{col}} \sim (\bar{R}H_p)^2 \frac{\rho_{\text{kin}}^2}{(\rho_{\text{kin}} + \rho_{\text{rad}})^2}$

Going beyond dimensional analysis with simulations (and redshift)

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Bubble collision simulations

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Amplitude:

$$h^2\Omega_{\text{col}}(f) = 1.67 \times 10^{-5} \left(\frac{100}{g_*}\right)^{1/3} \left(\frac{\beta}{H_p}\right)^{-2} \kappa_v^2 \left(\frac{\alpha}{1+\alpha}\right)^2 \left(\frac{0.11v^3}{0.42+v^2}\right) S(f)$$

$$S(f) = \frac{3.8(f/f_0)^{2.8}}{1 + 2.8(f/f_0)^{3.8}}$$

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Peak frequency:

$$f_0 = 1.65 \times 10^{-7} \left(\frac{T_p}{1 \text{ GeV}}\right) \left(\frac{g_*}{100}\right)^{1/6} H_p^{-1} \beta \left(\frac{0.62}{1.8 - 0.1v + v^2}\right) \text{ Hz}$$

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The set of parameters (\bar{R} , ρ_{kin} , v , κ_ν) is determined by the underlying particle physics model.

Phase transitions beyond the Standard Model

Various scenarios of phase transition

In the Standard Model, both electroweak and QCD PTs are crossover

[K. Kajantie et al., Phys. Rev. Lett. 77 (1996) 2887] [Y. Aoki et al, Nature 443 (2006) 675]

⇒ no stochastic GW background predicted in the SM

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Typical BSM electroweak PT:

- $\mathcal{O}(1)$ bubbles produced per Hubble volume at $T_n \lesssim T_{EW}$
- they rapidly collide ⇒ percolation temperature $T_p \sim T_n$
- time scale of the process much shorter than Hubble time
- $f_{\text{peak}} \sim \text{milliHertz}$ ⇒ range of LISA [C. Caprini et al., JCAP 1604 (2016) no.04 001]

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Prolonged and supercooled PT [A. Kobakhidze, CL, A. Manning, J. Yue, arXiv:1703.06552]:

- weaker nucleation probability
- less bubbles produced ⇒ more time needed for them to collide
- ⇒ $T_p \ll T_n \lesssim T_{EW}$
- $f_{\text{peak}} \sim 10^{-8}$ Hertz ⇒ range of Pulsar Timing Arrays

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QCD-induced PT [S. Arunasalam, A. Kobakhidze, CL, S. Liang, A. Zhou, arXiv:1709.10322] :

- no EWPT before $T \sim T_{\text{QCD}}$
- at T_{QCD} : 1st order chiral phase transition with **6** massless quarks
- quark condensates trigger EWPT
- $f_{\text{peak}} \sim 10^{-8}$ **Hertz** ⇒ range of Pulsar Timing Arrays

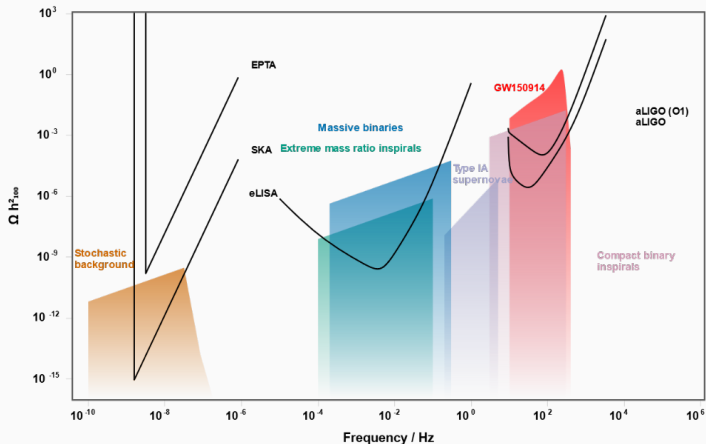
See also: [E. Witten Nucl.Pys.B177 (1981) 477] [W. Buchmuller, D. Wyler, PLB 249 (1990) 281] [S. Iso et al., PRL 119 (2017) 141301] [B. von Harling, G. Servant, JHEP 1801 (2018) 159]

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Prolonged electroweak phase transition

Non-linear realization of the electroweak gauge group

[A. Kobakhidze, A. Manning, J. Yue, Int.J.Mod.Phys. D26 (2017) no.10, 1750114]

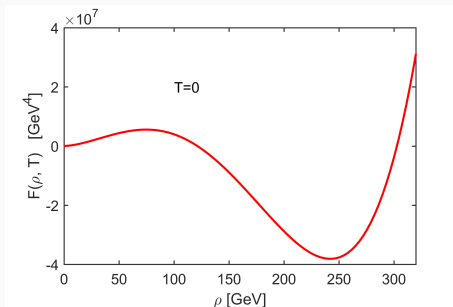
- $\mathcal{G}_{\text{coset}} = SU(2)_L \times U(1)_Y/U(1)_Q$ is gauged
- broken generators $T^i = \sigma^i - \delta^{i3}\mathbb{I}$ and Goldstone bosons $\pi^i(x)$
- physical Higgs as a singlet $\rho(x) \sim (1, 1)_0$
- SM Higgs doublet identified as $H(x) = \frac{\rho(x)}{\sqrt{2}} e^{i\frac{1}{2}\pi^i(x)T^i} \begin{pmatrix} 0 \\ 1 \end{pmatrix}, \quad i \in \{1, 2, 3\}$

SM particle content but BSM interactions

Minimal setup (usual SM configurations except Higgs potential):

$$V^{(0)}(\rho) = -\frac{\mu^2}{2}\rho^2 + \frac{\kappa}{3}\rho^3 + \frac{\lambda}{4}\rho^4.$$

Model specified by **one** parameter: $\kappa = \bar{\kappa} \cdot \frac{m_h^2}{v} \sim 63.5 \cdot \bar{\kappa} \text{ GeV}$.



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Percolation temperature (\sim collision) [L. Leita et al., JCAP 1210 (2012) 024]: $p(t_p) \approx 0.7$

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Number density of produced bubbles:

$$\frac{dN}{dR}(t, t_R) = \Gamma(t_R) \left(\frac{a(t_R)}{a(t)} \right)^4 \frac{p(t_R)}{v(t_R)}$$

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Nucleation temperature T_n : maximum of $\frac{dN}{dR}(t_p, t_R)$

Bubbles properties at collision

By definition:

- most bubbles collide at t_p
- majority of them produced at t_n

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Kinetic energy stored in bubble-walls:

$$E_{\text{kin}} = \kappa_\nu \cdot 4\pi \int_{t_n}^{t_p} dt \frac{dR}{dt}(t, t_n) R^2(t, t_n) \epsilon(t)$$

- $\epsilon(t)$: latent heat (\sim vacuum energy)
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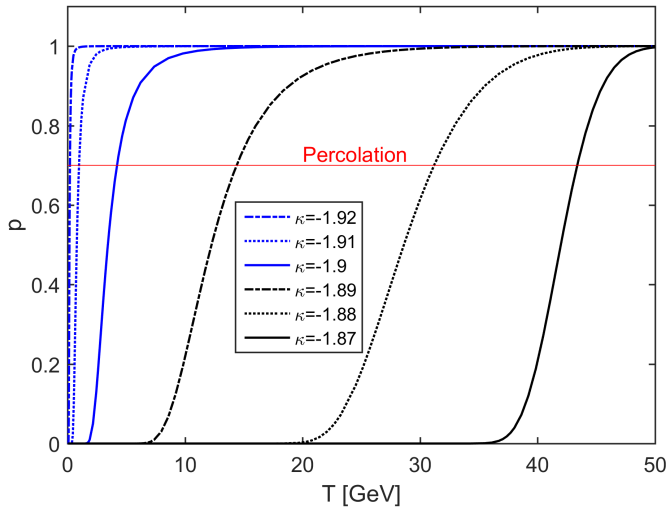
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\bar{R} and E_{kin} : key parameters to deduce the GW spectrum

Numerical results

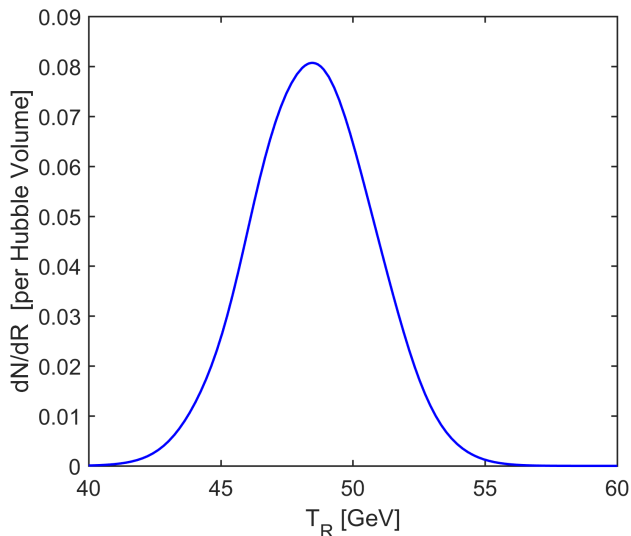
Probability $p(T)$:



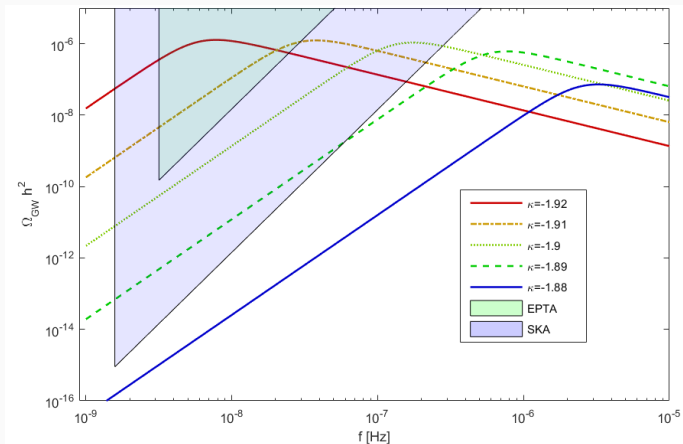
Numerical results

Number density distribution for $|\bar{\kappa}| = 1.9$:

$\Rightarrow T_n \sim 49 \text{ GeV}$



Numerical results



- Current constraints: EPTA, PPTA, NANOGrav
- Possible detection: Square Kilometre Array

Conclusion

- Gravitational Waves are allowing us to explore the fundamental laws of the Universe!
- Future detectors and theoretical investigations are likely to bring new discoveries!