

Quantum Gates

How to implement a $X \otimes b$?

$a \backslash b$	0	1
0	0	1
1	1	0

but $(a, b) \mapsto a \text{ XOR } b$ not unitary

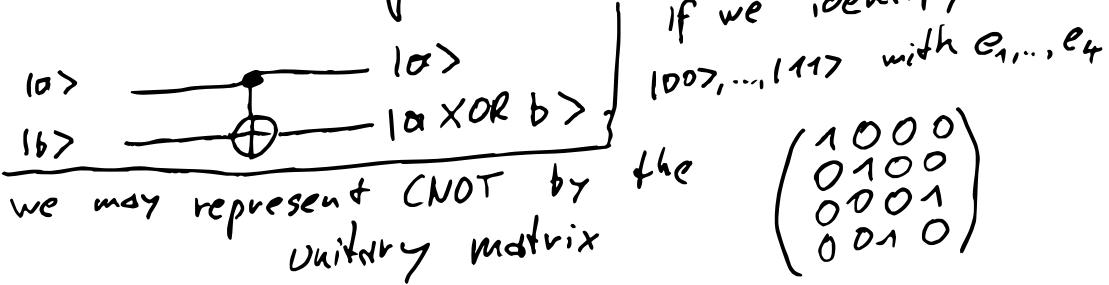
hence $(a, b) \mapsto (a, a \text{ XOR } b)$

(a, b)	$ 100\rangle$	$ 101\rangle$	$ 110\rangle$	$ 111\rangle$
$(a, a \text{ XOR } b)$	$ 100\rangle$	$ 101\rangle$	$ 111\rangle$	$ 110\rangle$

is permutation \Rightarrow unitary

This operation is called "Controlled Not", "CNOT"

In circuit diagrams, this is denoted by



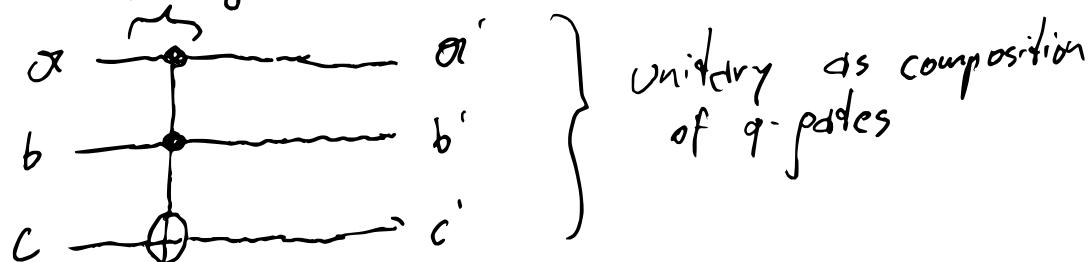
What about „AND“?

$a \text{ AND } b \Leftrightarrow$

$a \backslash b$	0	1
0	0	0
1	0	1

- $(a, b) \mapsto (a, a \text{ AND } b)$ not bijective
- requires three qubits $|a, b, c\rangle \mapsto |a, b, c \text{ XOR } (a \text{ AND } b)\rangle$

Toffoli gate



$ abc\rangle$	000	001	010	011	100	101	110	111
	000	001	010	011	100	101	111	110

↓
Input $c = 10\rangle$

If Input $c = 10\rangle \Rightarrow$ Output $c' = |a, b, a \text{ AND } b\rangle$

- extra qbit $|c\rangle$ is called „ancilla“ qbit
ancilla = Diener

Deutsch-Jozsa Algorithm

- "Problem": Given $\overline{f}(x_1, \dots, x_n) \in \{0, 1\}$ $x_i \in \{0, 1\}$
- "determine if
 - $\overline{f} = 0$
 - $\overline{f} = 1$
 - \overline{f} is unbalanced; i.e., exactly half of (x_1, \dots, x_n) lead to $\overline{f}(x_1, \dots, x_n) = 1$

- We assume the \overline{f} satisfies one of the three options
- Not useful in practice, but demonstrates that hard problems can be easier on quantum computers

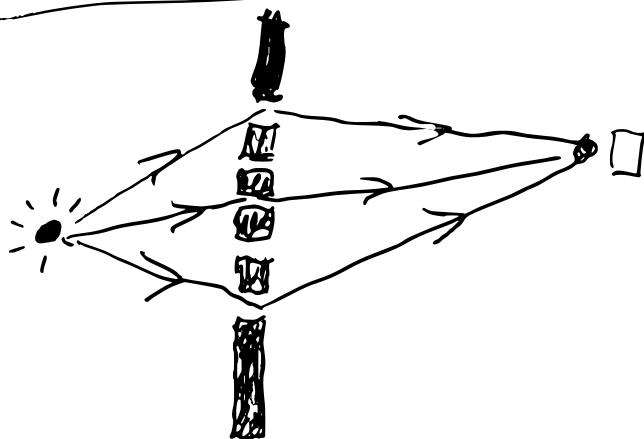
Classical algorithm:

To exclude the "balanced" case, we need to check at least $2^{n-1} + 1$ inputs

$$(x_1, \dots, x_n) \in \{0, 1\}^n$$

\Rightarrow Problem size $N = 2^n$, Runtime $N/2 + 1$

Physical inspiration



Intensity at detector is determined by "phase" difference of light-paths
 \rightarrow interference

Number the holes with }
 $x \in \{0, 1, 3\}$ } 2 " holes

Assume the holes' are located such that the light amplitude is $\approx (-1)^{F(y)}$ at detector if light travels through hole y .

Light intensity at detector is given by

$$\approx \sum_{y \in \{0, 1, 3\}} (-1)^{F(y)} = 0 \text{ if } F \text{ is balanced}$$

$= \pm 1$ if F is const

Quantum Algorithm

We may identify $(x_1, \dots, x_n) \in \{0,1\}^n$ with the basis element $|i\rangle$, where $i = x_1 + x_2 2 + x_3 4 + \dots + x_n 2^{n-1}$. In this sense, define $\overline{f}(|i\rangle) := |\overline{f}(x_1, \dots, x_n)\rangle$

An essential part of this (and many other algorithms) is the "Query":

Given 1-qbit state $|b\rangle$ and $\overline{f}(x_1, \dots, x_n) \in \{0,1\}$, define

$$Q_{\overline{f}}(|i\rangle \otimes |b\rangle) = |i\rangle \otimes (|b\rangle \oplus \overline{f}(|i\rangle))$$

Particularly for $|b\rangle = |-\rangle := \frac{1}{\sqrt{2}}(|0\rangle - |1\rangle)$,

we have

$$\begin{aligned} Q_{\overline{f}}(|i\rangle \otimes |-\rangle) &= |i\rangle \otimes \frac{1}{\sqrt{2}} [\overline{f}(|i\rangle) - |1\rangle \oplus \overline{f}(|i\rangle)] \\ &= \begin{cases} |i\rangle \otimes \frac{1}{\sqrt{2}}[|0\rangle - |1\rangle] & \overline{f}(|i\rangle) = |0\rangle \\ |i\rangle \otimes \frac{1}{\sqrt{2}}[|1\rangle - |0\rangle] & \overline{f}(|i\rangle) = |1\rangle \end{cases} \\ &= \overbrace{(-1)^{\overline{f}(|i\rangle)} |i\rangle}^{\overline{f}(|i\rangle)} \end{aligned}$$

Note that we identify $\overline{f}(|i\rangle)$ with $\overline{f}(x_1, \dots, x_n)$ instead of $|\overline{f}(x_1, \dots, x_n)\rangle$, i.e., the last equality only makes sense for basis element $|i\rangle$.

Note that Q_F is unitary since

$$|b\rangle \mapsto (|b\rangle \oplus F(|a\rangle))$$

is permutation for any value of $F(a)$

We make extensive use of the Hadamard gate \boxed{H} given by $\frac{1}{\sqrt{2}} \begin{pmatrix} 1 & 1 \\ 1 & -1 \end{pmatrix}$

we already know:

$$H|0\rangle = \frac{1}{\sqrt{2}}(|0\rangle + |1\rangle)$$

$$H|1\rangle = \frac{1}{\sqrt{2}}(|0\rangle - |1\rangle)$$

Lemma 1. There holds

$$\begin{aligned} H^{\otimes n}|0^n\rangle &:= \bigotimes_{i=1}^n H|0\rangle = \bigotimes_{i=1}^n \frac{1}{\sqrt{2}}(|0\rangle + |1\rangle) \\ &= \sum_{\substack{(x_1, \dots, x_n) \\ \in \{0,1\}^n}} \frac{1}{\sqrt{2^n}} |x_1 \dots x_n\rangle = \frac{1}{\sqrt{N}} \sum_{i=0}^{2^n-1} |i\rangle \end{aligned}$$

as well as for $i = x_1 + 2x_2 + \dots + 2^{n-1}x_n$

$$\begin{aligned} H^{\otimes n}|i\rangle &:= \bigotimes_{i=1}^n H|x_i\rangle = \bigotimes_{i=1}^n \frac{1}{\sqrt{2}}(|0\rangle + (-1)^{x_i}|1\rangle) \\ &= \sum_{\substack{(y_1 \dots y_n) \\ \in \{0,1\}^n}} \frac{1}{\sqrt{N}} (-1)^{x \cdot y} |y_1 \dots y_n\rangle \\ &= \frac{1}{\sqrt{N}} \sum_{j=0}^{2^n-1} (-1)^{i \cdot j} |j\rangle, \quad \text{where} \end{aligned}$$

$$i \cdot j := x \cdot y := \sum_{k=1}^n x_k y_k .$$

The Algorithm reads:

1) Initial state $|0^n\rangle := \underbrace{|0\rangle \otimes \dots \otimes |0\rangle}_{n\text{-times}}$

2) Apply \boxed{H} to each qbit to obtain

uniform state $\frac{1}{\sqrt{N}} \sum_{i=0}^{2^n-1} |i\rangle$

3) Tensorize result with $|-\rangle \Rightarrow \frac{1}{\sqrt{N}} \sum |i-\rangle$

3) Apply query $\boxed{Q_F}$ with $|b\rangle = |-\rangle$

to obtain $\frac{1}{\sqrt{N}} \sum_{i=0}^{2^n-1} (-1)^{F(|i\rangle)} |i-\rangle$

4) Ignore last qbit $|-\rangle$ and apply \boxed{H}
to remaining state to obtain (Lemma 1)

$$\frac{1}{N} \sum_{i=0}^{2^n-1} \sum_{j=0}^{2^n-1} (-1)^{ij} (-1)^{F(|i\rangle)} |ij\rangle$$

5) The amplitude of the $j=0 \Leftrightarrow |0^n\rangle$ -state

$$is \quad \frac{1}{N} \sum_{i=0}^{2^n-1} \underbrace{(-1)^{i \cdot 0}}_{=1} (-1)^{F(|i\rangle)}$$

- If \bar{F} is balanced \Rightarrow

$$\frac{1}{N} \sum_{i=0}^{2^n-1} (-1)^{\bar{F}(1i)} = 0$$

$$\cdot \text{If } \bar{F}=1 \Rightarrow \frac{1}{N} \sum_{i=0}^{2^n-1} (-1)^{\bar{F}(1i)} = -1$$

$$\cdot \text{If } \bar{F}=0 \Rightarrow \frac{1}{N} \sum_{i=0}^{2^n-1} (-1)^{\bar{F}(1i)} = 1$$

Complexity of Quantum Alp:

$O(1)$ steps, $O(\log_2(N))$ quantum gates

$O(\log(N))$ qubits \Rightarrow exponential speedup

Example for \bar{F} :

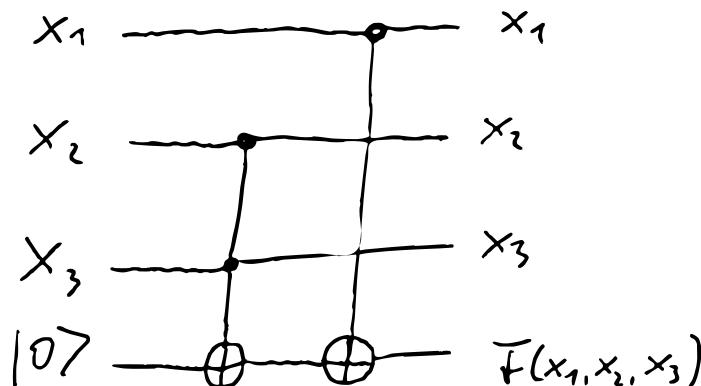
$$\cdot \bar{F}(x_1, x_2, x_3) = \text{mod}(x_1 + x_2 x_3, 2)$$

$\Rightarrow \bar{F}$ is balanced since

$$\bar{F}(0, x_2, x_3) = 1 - \bar{F}(1, x_2, x_3)$$

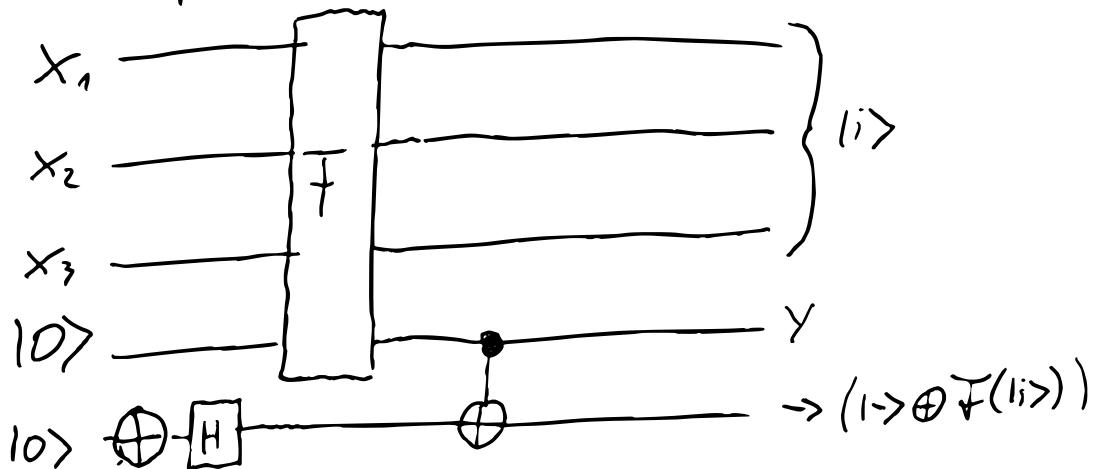
For the algorithm to work, \bar{F} needs to be a quantum circuit, i.e.

$$F(x_1, x_2, x_3) = x_1 \text{ XOR } (x_2 \text{ AND } x_3)$$

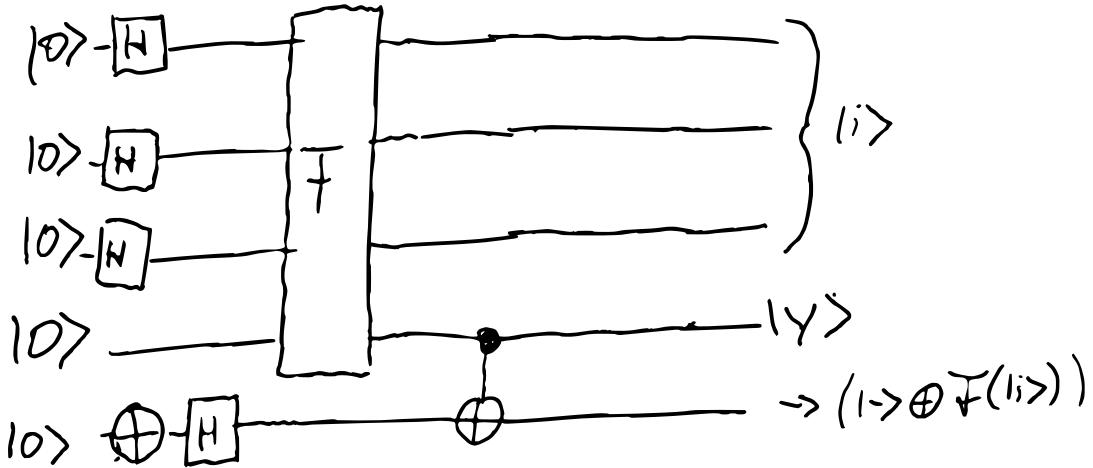


ancilla qubit

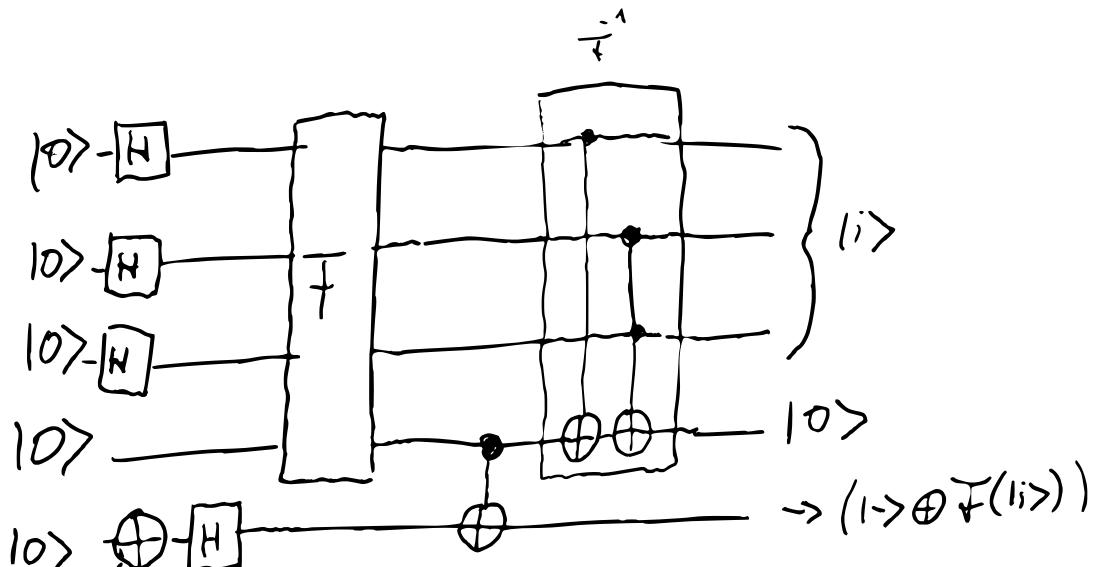
Next, implement Query $Q_{\bar{F}}(|i\rangle) = |i\rangle \otimes (\bar{I} \otimes \bar{F}(|i\rangle))$



Apply Steps 1-3 of ALP

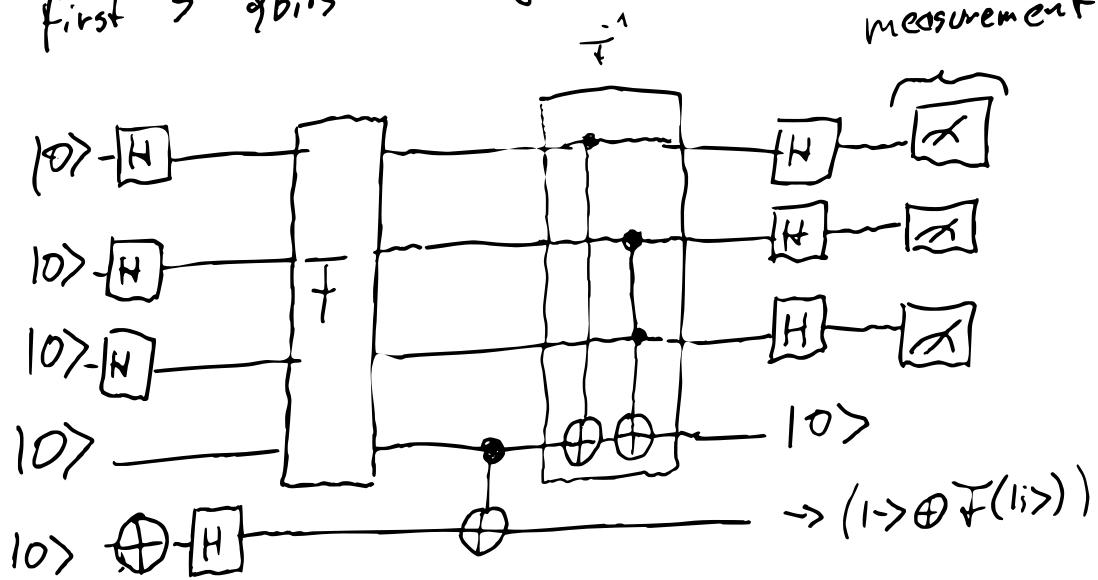


Problem: $|y\rangle$ and $|i\rangle$ are entangled
 Solution: reverse F such that $|y\rangle = |0\rangle$



Now, we have the desired state
 $|i\rangle \otimes (I -> \oplus F(Ii>)) (\otimes |0\rangle)$

Apply step 4 of Algo and measure the first 3 qubits



The measurement must contain the projection onto $|000\rangle \otimes \text{span}\{|0>, |1>\}$

Note that this is not the most efficient implementation of Q_F



at this point we have $(-1)^{x_1} |x>$

Controlled Z-gate:
Z-gate only active if $x_2=1$

$Z = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$

$\Rightarrow (-1)^{x_2 x_3} |x>$

Remarks on Deutsch-Josza:

- First algorithm with exponential speedup over classical computer
- However, consider the following randomized classical algorithm $R(\bar{f})$:
 - 1) generate two random $x, y \in \{0,1\}^n$
 - 2) Evaluate $\bar{f}(x), \bar{f}(y)$
 - 3) Output: $R(\bar{f}) = 1$ if $\bar{f}(x) = \bar{f}(y) = 1$
 $R(\bar{f}) = 0$ if $\bar{f}(x) = \bar{f}(y) = 0$
 $R(\bar{f})_{\text{tolerant}}$ if $\bar{f}(x) \neq \bar{f}(y)$
- The algorithm is correct if \bar{f} is constant and answers correctly with prob. $\frac{1}{2}$ if \bar{f} is balanced.
- Apply the alg k -times with i.i.d random samples to get $(R_1(\bar{f}), \dots, R_k(\bar{f}))$. Return
 $1 \text{ if } R_1(\bar{f}) = \dots = R_k(\bar{f}) = 1$
 $0 \text{ if } R_1(\bar{f}) = \dots = R_k(\bar{f}) = 0$
 $"\text{balanced}" \text{ else}$
 $\Rightarrow \text{Probability of error } 2^{-k}, \text{ cost } O(k)$

Remark: Why do you have to uncompute \tilde{F} in order to remove entanglement?

After Step 3, we have

$$\frac{1}{\sqrt{N}} \sum_{i=0}^{2^n-1} (-1)^{\tilde{F}(1i>)} |i>$$

but in implementation, we actually have

$$\frac{1}{\sqrt{N}} \sum_{i=0}^{2^n-1} (-1)^{\tilde{F}(1i>)} |i> \otimes |0_i> \otimes |>$$

for some ancilla qbit $|0_i>$. We may ignore last qbit since not entangled.

Step 4 applies $H^{\otimes n}$ to obtain

$$\frac{1}{N} \sum_{i=0}^{N-1} \sum_{j=0}^{N-1} (-1)^{\tilde{F}(1i>)} (-1)^{i,j} |j> \otimes |0_i>$$

\Rightarrow Amplitude of $|j> = |0> \Rightarrow$

$$\frac{1}{N} \sum_{i=0}^{N-1} (-1)^{\tilde{F}(1i>)}$$

Mixed to non-zero even for balanced \tilde{F}

Simon's algorithm

First quantum alg. with exp speedup over
any (randomized) classical alg.

Given $i, j \in \{0, 1\}^n$, define $i \oplus j := (i_1 \oplus j_1, \dots, i_n \oplus j_n)$

Simon's problem: Let $\bar{F} : \{0, 1\}^n \rightarrow \{0, 1\}^n$ s.t.
there exists $s \in \{0, 1\}^n$

$$\bar{F}(x) = \bar{F}(y) \iff x = y \text{ or } x \oplus s = y$$

$$\forall x, y \in \{0, 1\}^n$$

Goal: Find s .

Quantum Algorithm

1) Start with $|0^n\rangle$ and apply \boxed{H} to

the first n qubits to obtain

$$\frac{1}{\sqrt{N}} \sum_{i=0}^{2^n-1} |i\rangle |0^n\rangle$$

2) Apply the query $|i\rangle \otimes |b^n\rangle \mapsto |i\rangle \otimes (|b\rangle \otimes \bar{F}(|i\rangle))$
with $|b^n\rangle = |0^n\rangle$ to obtain

$$\frac{1}{\sqrt{N}} \sum_{i=0}^{2^n-1} |i\rangle |\bar{F}(i)\rangle$$

3) Measure the second n qubits in computational basis, i.e. $\text{span}\{|i\rangle, i=0, \dots, 2^n-1\} \otimes \text{span}\{|j\rangle, j=0, \dots, 2^n-1\}$ to

obtain some output

$$|a\rangle \otimes |j\rangle = \frac{\sum_j \left(\frac{1}{\sqrt{n}} \sum_i |i\rangle F(i,j) \right)}{\|\sum_j P_j(\dots)\|}$$

Note that

$$|i\rangle F(i,j) \perp \text{ran } P_j \text{ for } F(i,j) \neq |j\rangle$$

by assumption, there exist exactly two inputs
 $|i_0\rangle$ and $|i_k\rangle = |i_0 \oplus s\rangle$ with $F(i_k, j) = |j\rangle$ $k=0,1$.

Hence $|a\rangle = \frac{1}{\sqrt{2}} (|i_0\rangle + |i_0 \oplus s\rangle)$

4) Ignore second n qubits and apply \boxed{H}
to the first n to obtain

$$\frac{1}{\sqrt{2^n}} \sum_{j=0}^{2^n-1} \left[(-1)^{b \cdot j} + (-1)^{(i_0 \oplus s) \cdot j} \right] |j\rangle$$

5) Measure in computational basis

- $|j\rangle$ has non-zero amplitude if

$$(i_0 \oplus s) \cdot j \equiv i_0 \cdot j \iff s \cdot j \equiv 0 \pmod{2}$$

\Rightarrow We obtain random element of the set

$$\{ j \mid s \cdot j \equiv 0 \pmod{2} \}$$

6) Repeat the procedure to obtain
 $n-1$ linearly ind. elements $j^{(1)}, \dots, j^{(n-1)}$ with
 $j^{(i)} \cdot s = 0 \pmod{2}$

or

$$\begin{pmatrix} j^{(1)} \\ \vdots \\ j^{(n-1)} \end{pmatrix} s = 0 \pmod{2}$$

7) Solve linear system mod 2 on
 classical computer in $O(n^3)$

Note that #spans $\{j^{(1)}, \dots, j^{(k)}\} \leq 2^k$

Hence, with probability $\frac{2^n - 2^k}{2^n} = 1 - 2^{k-n} \geq \frac{1}{2}$
 for $k \leq n-1$, we find

• Linearly independent vector j_{k+1}

Conclusion: number of qubits $O(n)$,
 number of gates $O(n)$, number of iterations
 $O(n)$ with high probability + $O(n^3)$

Classical algorithms for Simon's problem:

Lemma 2 Any randomized classical algorithm with less than $\delta 2^{\frac{n}{2}}$ queries to \tilde{f} fails with probability $\geq e^{-\frac{5}{4}\delta^2}$. (deterministic alg. fail certainly if less than $2^{\frac{n}{2}}$ queries).

Proof Every alg generates a sequence of queries x_1, \dots, x_n with $\tilde{f}(x_1), \dots, \tilde{f}(x_n)$. If all y_i are distinct, the alg can't distinguish between $s=0$ and $s \neq 0$.

Assume all y_1, \dots, y_k are distinct. Then the alg chooses x_{k+1} based on some prob. measure μ on $\{0,1\}^n$ [in the def. case, μ is a delta distribution].

$$\sum_{k=1}^c \sum_{s \in \{0,1\}^n} \sum_{i=1}^k \mu(x_i \oplus s) = \sum_{k=1}^c \underbrace{\sum_{i=1}^k}_{s} \underbrace{\sum_{i=1}^k \mu(x_i \oplus s)}_{=1}$$

$$\Rightarrow \exists s \in \{0,1\}^n: \sum_{k=1}^c \underbrace{\sum_{i=1}^k \mu(x_i \oplus s)}_{p_i} \leq \frac{c(c+1)}{2^{n+1}} \leq \frac{1}{2} \quad \text{for } c(c+1) \leq 2^n$$

Hence,

$$\begin{aligned} \overline{P}(y_1, \dots, y_{k+1} \text{ distinct}) &= \underbrace{\overline{P}(y_{k+1} \notin \{y_1, \dots, y_k\} | y_1, \dots, y_k \text{ distinct})}_{\cdot \overline{P}(y_1, \dots, y_k \text{ distinct})} \\ &\geq 1 - p_i \end{aligned}$$

Iterate the argument to obtain

$$P(\gamma_1, \dots, \gamma_c \text{ distinct}) \geq \prod_{i=1}^c (1-p_i)$$

Taking logarithms shows

$$\log\left(\prod_{i=1}^c (1-p_i)\right) = \sum_{i=1}^c \log(1-p_i)$$

$$\left(p_i \leq \frac{1}{2}\right) \rightarrow \geq -\frac{5}{4} \sum_{i=1}^c p_i \geq -\frac{5}{4} \frac{c(c+1)}{2^{n+1}}$$

$$\Rightarrow P(\gamma_1, \dots, \gamma_c \text{ distinct}) \geq e^{-\frac{5}{4} \frac{c(c+1)}{2^{n+1}}}$$

Choosing $c+1 = 5 \cdot 2^{\frac{n}{2}}$ satisfies $c(c+1) \leq 2^n$

and $P(\gamma_1, \dots, \gamma_c \text{ distinct}) \geq e^{-\frac{5}{4} \delta^2}$

□

$$\log(1-p) = 0 + \frac{1}{1}(-p) - \frac{1}{2} \frac{(-p)^2}{2} \text{ for some } 1-p \leq \frac{p}{2} \leq p$$

$$\geq -p - \frac{p^2}{2} = -p\left(1 + \frac{p}{2}\right)$$

$$\geq -\frac{5}{4}p$$

The quantum Fourier Transform

Discrete FT: For x_0, \dots, x_{N-1} define

$$\hat{X}_k := \frac{1}{\sqrt{N}} \sum_{j=0}^{N-1} x_j e^{-\frac{2\pi i}{N} k j}$$

$$X_k = \frac{1}{\sqrt{N}} \sum_{j=0}^{N-1} \hat{X}_j e^{\frac{2\pi i}{N} k j}$$

in matrix notation

$$\hat{X} = F_N X \quad \text{with} \quad F_N \in \mathbb{C}^{N \times N}$$

$$(F_N)_{kj} := \frac{1}{\sqrt{N}} e^{-\frac{2\pi i}{N} k j}$$

F_N is unitary matrix ✓

Fast FT: Assume $N = 2^n$

$$\hat{X}_k = \frac{1}{\sqrt{2}} \left(\frac{1}{\sqrt{N/2}} \sum_{j=0}^{N-1} x_j e^{-\frac{2\pi i}{N/2} k \frac{j}{2}} \right.$$

$$\left. + e^{-\frac{2\pi i}{N} k \frac{1}{2}} \sum_{j \text{ even}} e^{-\frac{2\pi i}{N/2} k \frac{j+1}{2}} \right)$$

⇒ split FT into 2 FTs of half size

leads to $O(N \log N)$ complexity.

The QFT maps a state

$$|x\rangle = \sum_{j=0}^{N-1} \hat{x}_j |j\rangle \text{ to } |y\rangle = \sum_{j=0}^{N-1} x_j |j\rangle$$

where x_j is given by the classical FFT

If $|x\rangle = |j_0\rangle$, then $x_j = \delta_{jj_0}$ and hence

$$\hat{x}_k = \frac{1}{\sqrt{N}} e^{\frac{2\pi i k j_0}{N}}$$

$$\Rightarrow |j\rangle \xrightarrow{\text{QFT}} \frac{1}{\sqrt{N}} \sum_{k=0}^{N-1} e^{\frac{2\pi i k j}{N}} |k\rangle$$
$$= \overline{f}_N |j\rangle$$

This is a convention.
Everything works with FT

To implement \overline{f}_N efficiently, we observe

$$\frac{k}{N} = \frac{k}{2^n} = \sum_{e=1}^n k_e \frac{1}{2^e} \quad \text{if } k = k_1 2^{n-1} + k_2 2^{n-2} + \dots + k_n$$

$$\overline{f}_N |j\rangle = \frac{1}{2^n} \sum_{k \in \{0,1\}^n} e^{2\pi i j \sum_{e=1}^n k_e \frac{1}{2^e}} |k_1 \dots k_n\rangle$$
$$= \bigotimes_{e=1}^n e^{-2\pi i j k_e \frac{1}{2^e}} |k_e\rangle$$

$$= \bigotimes_{e=1}^n \left(|0\rangle + e^{2\pi i j / 2^e} |1\rangle \right) \frac{1}{\sqrt{2}}$$

Furthermore note that

$$e^{2\pi i j / 2^e} = e^{2\pi i \sum_{m=1}^{n-e} j_m 2^{n-m-e}} = e^{\sum_{m=n-e+1}^n j_m 2^{n-m-e}}$$

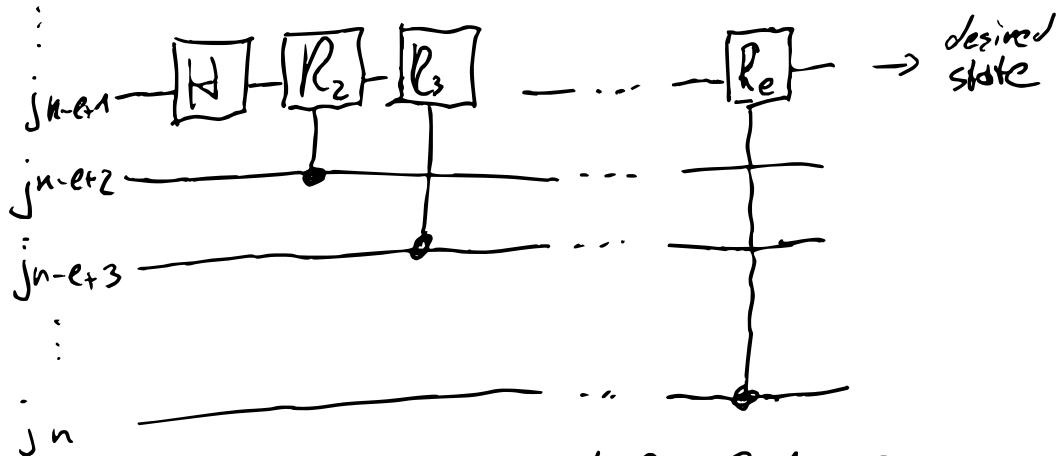
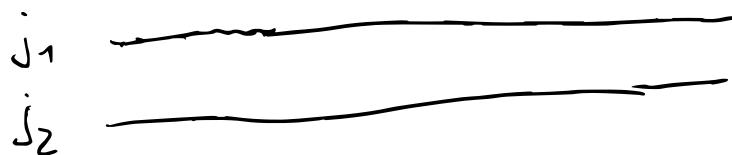
$\in \mathbb{N}$

\Rightarrow the first $n-e$ significant bits of j do not matter.

To implement

$$\frac{1}{2^e} (|0\rangle + e^{2\pi i \sum_{m=n-e+1}^n j_m 2^{n-m-e}} |1\rangle)$$

we use the $[R_s]$ gate given by the matrix $\begin{pmatrix} 1 & 0 \\ 0 & e^{2\pi i / 2^e} \end{pmatrix}$ or $[P]$ in IBM q-Composer



need to repeat that for $e=1, \dots, n$

Cost of QFT:

We need n qubits and $O(n)$ gates per qubit

$\Rightarrow O(n^2)$ gates

\Rightarrow exponential speedup over FFT with $O(n 2^n)$ operations.

Remark Strictly speaking, QFT does something different than FT. The state

$$QFT(|x\rangle) = \sum_{j=0}^{N-1} x_j |j\rangle$$

can only be accessed via measurement and hence will collapse to some $|j'\rangle$ with a certain probability. We will see that this is still very useful.

Remark Rs gates don't do very much for large s . One can show that $O(n \log n)$ gates suffice if one accepts a small error probability

Remark reversing the order of the
passes and using the adjoint gates
gives an efficient implementation of
the inverse QFT F_N^{-1}

Application Phase estimation

Suppose we have unitary operator $U: V \rightarrow V$
 $\dim V = 2^n$
 with eigenvector ψ

$$\Rightarrow U\psi = e^{2\pi i \phi} \psi \quad \text{for some } \phi \in [0, 1)$$

$$(|U\psi|^2 = |\lambda|^2 |\psi|^2 = |\psi|^2 \Rightarrow |\lambda| = 1)$$

Assume that $\phi = \sum_{j=1}^n \phi_j 2^{-j}$ can be
 written with n bits

Since U is operator on 2^n -dim Hilbert space
 classical computation of $U\psi \cdot \psi$ costs
 at least $O(2^n)$

Quantum algorithm

1) Start with $|0^n\rangle |\psi\rangle$

2) Apply $H^{\otimes n}$ to the first n qubits to obtain
 $\frac{1}{\sqrt{N}} \sum_{j=0}^{2^n-1} |j\rangle |\psi\rangle$ (Applying QFT_N would
 do the same)

3) Apply $|j\rangle |\psi\rangle \mapsto |j\rangle |U^j|\psi\rangle$ to obtain

$$\frac{1}{\sqrt{N}} \sum_{j=0}^{2^n-1} e^{2\pi i j \phi} |j\rangle |\psi\rangle$$

Note that the first n qubits satisfy

$$\frac{1}{\sqrt{N}} \sum_{j=0}^{N-1} e^{2\pi i j \phi} |j\rangle = \frac{1}{\sqrt{N}} \sum_j e^{2\pi i j \frac{N\phi}{N}} |j\rangle = F_N(|N\phi\rangle)$$

→ This might be a bit confusing since suddenly $\phi \in \mathbb{R}$ is interpreted as an element of the vector space in which ψ is contained

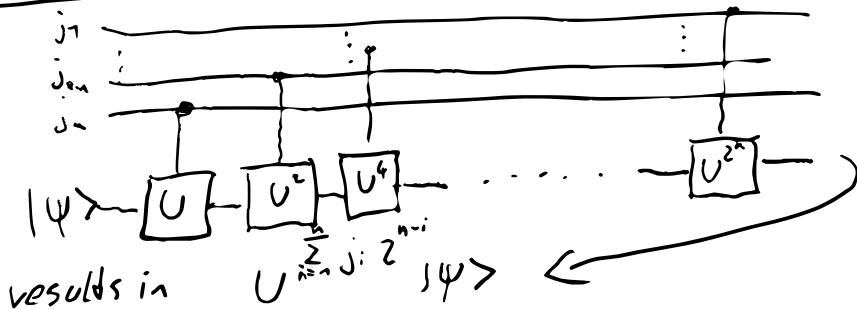
But $N\phi = \sum_{j=1}^n \phi_j 2^{n-j}$ is a basis element and hence this makes sense

4) Apply IQFT to the first n qubits to obtain

$$F_N^{-1}\left(\frac{1}{\sqrt{N}} \sum_j e^{2\pi i j \phi} |j\rangle\right) = |N\phi\rangle$$

5) Measure in computational basis to obtain $N\phi$ and hence ϕ

Remark How to implement step 3?



But we will need to assume that

$\boxed{U^{2^k}}$ can be implemented effectively
 (This might not be true in general, but
 is in the applications below)

Note the the input doesn't need to be a single eigenvector. Let $|1\phi\rangle, |1\phi'\rangle$ denote two EVs with phase ϕ, ϕ' . Linearity implies that input $\frac{1}{\sqrt{2}}(|1\phi\rangle + |1\phi'\rangle)$ produces output $\frac{1}{\sqrt{2}}(|N\phi\rangle + |N\phi'\rangle)$

A measurement will produce ϕ, ϕ' with equal probability

if ϕ requires more than n bits, the final state is a small perturbation of $|N\phi\rangle$ with

$$|\delta| = \left| \frac{1}{\sqrt{N}} \sum_{j=0}^{N-1} \left[e^{2\pi i j \phi} - e^{2\pi i j \hat{\phi}} \right] |1j\rangle \right| / \left| \frac{1}{\sqrt{N}} \sum_{j=0}^{N-1} e^{2\pi i j \phi} |1j\rangle \right| \simeq O(2^{-n})$$

Shor's integer factorization

Factoring problem: given $n \in \mathbb{N} \setminus \{\text{primes}\}$
find $1 < k < n$ with $\frac{n}{k} \in \mathbb{N}$.

- $n \in \mathbb{N}$ can be defined using $\log_2(n)$ bits \Rightarrow Pdy. complexity of tps means $O(\log_2(n)^P)$ operations
- There exist efficient algorithms to check whether n is prime (see, e.g. AKS-test or power of prime (or BPSW-test))
- Security of many crypto-systems is based on the fact that integer factorization is hard
- Best known classical algorithm is the general number field sieve with runtime $O(e^{(\log n)^{\frac{1}{3}} (\log \log n)^{\frac{2}{3}}})$

Note that the routine is a conjecture!

- It is not known that classical algs can not be faster. Latest paper which (falsely) claimed an $O(\log^p)$ -alg for factorization was by Schnorr (Pom) in 2021

Reduction to period finding

Want to factor $N \in \mathbb{N}$.

1) Choose random $x \in \{2, \dots, N-1\}$ with $\gcd(x, N) = 1$

$\Rightarrow x \in (\mathbb{Z}/N\mathbb{Z})^\times$ multiplicative group
 $\mod n$

(Lemma 3)
 $\Rightarrow x$ has period r with $x^r \mod N = 1$

\Rightarrow with prob. $\frac{1}{2}$ r is even and

$$x^{\frac{r}{2}} \pm 1 \not\equiv 0 \mod N$$

$$\Rightarrow (x^{\frac{r}{2}} - 1)(x^{\frac{r}{2}} + 1) \equiv x^r - 1 \equiv 0 \mod N$$

$\Rightarrow \gcd(x^{\frac{r}{2}} - 1, N), \gcd(x^{\frac{r}{2}} + 1, N)$
are non-trivial factors of N

Lemma 3 Every $x \in (\mathbb{Z}/N\mathbb{Z})^*$ has period v which is minimal with $x^v \equiv 1 \pmod{N}$.

Proof Consider the set $S = \{1, x, \dots\} \subseteq (\mathbb{Z}/N\mathbb{Z})^*$. Since $(\mathbb{Z}/N\mathbb{Z})^*$ is finite, there exist $j, k \in \mathbb{N}$ with $x^j = x^k \pmod{N}$. W.L.O.G assume $j < k$

$$\Rightarrow x^{k-j} = 1 \pmod{N}.$$

□

Euler totient func: $\varphi(n) := \#\{1 \leq k \leq n : \gcd(k, n) = 1\}$

$$\varphi(n) = n \prod_{p|n} \left(1 - \frac{1}{p}\right) \Rightarrow \varphi(p^\alpha) = p^\alpha \left(1 - \frac{1}{p}\right) = p^{\alpha-1}(p-1)$$

Lemma [Chinese remainder Thm]

Let $m_1, \dots, m_n \in \mathbb{N}$ with $\gcd(m_i, m_j) = 1$ if $i \neq j$. Then

$$\begin{bmatrix} \text{Find } \\ x \text{ s.t. } \\ \begin{array}{l} x \equiv a_1 \pmod{m_1} \\ \vdots \\ x \equiv a_n \pmod{m_n} \end{array} \end{bmatrix}$$

is solvable and any two solutions are equal mod $M = m_1 m_2 \dots m_n$

Proof Define $M := \frac{M}{m_i} \Rightarrow \gcd(M_i, m_i) = 1$

$\Rightarrow M_i$ has mult. inverse mod m_i denoted by N_i

$\Rightarrow x = \sum_i a_i M_i N_i$ is solution since

$$M_i N_i \equiv 1 \pmod{m_i} \text{ and } M_i N_i \equiv 0 \pmod{m_j \text{ if } i \neq j}$$

Two solutions x, x' satisfy $x - x' \equiv 0 \pmod{m_i}$

$\Rightarrow M = m_1 m_2 \dots m_n$ divides $x - x'$ (since m_i are coprime)

□

Lemma 4 Let $p > 2$ prime. Let 2^d maximal power of 2 dividing $\varphi(p^\alpha)$. Let x denote randomly chosen element in $(\frac{\mathbb{Z}}{p^\alpha \mathbb{Z}})^\times$.

$$\Rightarrow P(2^d \text{ divides order of } x) = \frac{1}{2}$$

Proof: $\varphi(p^\alpha) = p^{d-1}(p-1)$ is even $\Rightarrow d \geq 1$.

Let g denote generator of $(\frac{\mathbb{Z}}{p^\alpha \mathbb{Z}})^\times \Rightarrow x = g^k \pmod{p^\alpha}$ for some $k \in \{1, \dots, \varphi(p^\alpha)\}$. Let r denote order of x .

1) k is odd: $g^{kr} \equiv 1 \pmod{p^\alpha} \Rightarrow \varphi(p^\alpha) \text{ divides } kr$
 since k is odd $\Rightarrow 2^d \text{ divides } r$ since $\varphi(p^\alpha)$ is minimal with $g^{\varphi(p^\alpha)} = 1$

2) k is even: $g^{k\varphi(p^\alpha)/2} \equiv 1 \pmod{p^\alpha}$
 $\Rightarrow r \text{ divides } \varphi(p^\alpha)/2$ since $x^r = g^{kr} \equiv 1 \pmod{p^\alpha}$
 r is minimal with

$\Rightarrow 2^d \text{ does not divide } r$

Hence, for exactly half of $x \in (\frac{\mathbb{Z}}{p^\alpha \mathbb{Z}})^\times$ we have $x = g^k$ with k even/odd □

Lemma 5 Let $N = p_1^{\alpha_1} \cdots p_m^{\alpha_m}$ prime factorization of odd N . Let x be random in $\left(\frac{N}{\mathbb{Z}}\right)^*$ with order $r \bmod N$.

$$\Rightarrow P(r \text{ is even and } x^{\frac{r}{2}} \not\equiv -1 \bmod N) \geq 1 - 2^{-m+1}$$

Proof Chinese remainder theorem \Rightarrow choose x random is equivalent to choose x_j random in $\left(\frac{N}{p_j^{\alpha_j}}\right)^*$ with $x \equiv x_j \pmod{p_j^{\alpha_j}}$ if $j=1, \dots, m$

If this is because $x \leftrightarrow (x_1, \dots, x_m)$ is one-to-one

Let r_j be order $x_j \pmod{p_j^{\alpha_j}}$ and let 2^{d_j} max power of 2 dividing r_j . r is odd or $x^{\frac{r}{2}} \not\equiv -1 \bmod N \Rightarrow$
We will show $d_1 = d_2 = \dots = d_m$. If this is the case, the following argument concludes the proof:

Let d_j' maximal s.t. $2^{d_j'}$ divides $\varphi(p_j^{\alpha_j})$

$$\text{Lemma 4} \Rightarrow P(d_j = d_j') = \frac{1}{2} \Rightarrow P(d_j = k) \leq \frac{1}{2}$$

$\forall k \in \mathbb{N}$. Since d_j independent

$$\Rightarrow P(d_1 = \dots = d_m) \leq 2^{-m+1}$$

It remains to show $d_1 = \dots = d_m$.

Case 1: r is odd. $x^r \equiv 1 \pmod{N} \Rightarrow x^{r_j} \equiv 1 \pmod{p_j^{d_j}}$
 $\Rightarrow r_j$ divides $r \Rightarrow r_j$ is odd $\Rightarrow d_j = 0$.

Case 2: r is even and $x^{\frac{r}{2}} \equiv -1 \pmod{N}$
 $\Rightarrow x^{\frac{r}{2}} \equiv -1 \pmod{p_j^{d_j}}$. If r_j would divide $\frac{r}{2}$, we would
 have $x^{\frac{r}{2}} = \underbrace{x^{kr_j}}_{= x_j^{kr_j} \pmod{p_j^{d_j}}} = -1 \pmod{p_j^{d_j}}$ } $\Rightarrow r_j$ does not divide $\frac{r}{2}$.

but r_j divides r . Hence largest power
 of 2 dividing r must be equal to d_j . □

Note that if r is period of $x \pmod{N}$
 and r even, we have $x^{\frac{r}{2}} \not\equiv 1 \pmod{N}$
 $\Rightarrow P(r \text{ even } x^{\frac{r}{2}} \not\equiv \pm 1 \pmod{N}) \geq 1 - 2^{-m+1} \geq \frac{1}{2}$
 if $m \geq 2$, i.e., if N is not a prime power.

Define $|Uy\rangle := |xy \bmod N\rangle$

Note that $|j\rangle \mapsto |xj \bmod N\rangle$ is bijective

since $xj = xj' \bmod N \quad x(j-j') = 0 \bmod N$
 $\Rightarrow \gcd(x, N) \neq 1$.

$\Rightarrow U$ is permutation on basis elements \Rightarrow unitary

Lemma Let r denote the period of x in $(\frac{Z}{N\mathbb{Z}})^*$.

Then, $|U_s\rangle := \frac{1}{\sqrt{r}} \sum_{k=0}^{r-1} \exp\left(\frac{-2\pi i sk}{r}\right) |x^k \bmod N\rangle$

are EVs of U with eigenvalues

$$\lambda_s := \exp\left(\frac{2\pi i s}{r}\right), \quad s = 0, \dots, r-1$$

Proof

$$\begin{aligned} U|U_s\rangle &= \frac{1}{\sqrt{r}} \sum_{k=0}^{r-1} \exp\left(\frac{-2\pi i sk}{r}\right) |x^{k+1} \bmod N\rangle \\ &= \lambda_s \frac{1}{\sqrt{r}} \sum_{k=0}^{r-1} \exp\left(\frac{-2\pi i s(k+1)}{r}\right) |x^{k+1} \bmod N\rangle \end{aligned}$$

$x^r = x^0 \bmod N$
 $e^{-2\pi i s} = e^0$

$\Rightarrow \lambda_s |U_s\rangle$

□

Note that there holds

$$\frac{1}{r} \sum_{s=0}^{r-1} |us\rangle = \frac{1}{r} \sum_{s=0}^{r-1} \sum_{k=0}^{r-1} e^{\frac{-2\pi iks}{r}} |x^k \bmod N\rangle$$

$$= \frac{1}{r} \sum_{k=0}^{r-1} \underbrace{\sum_{s=0}^{r-1} e^{\frac{-2\pi iks}{r}}}_{\begin{cases} 0 & k \neq 0 \\ r & k=0 \end{cases}} |x^k \bmod N\rangle$$

$$= |x^0 \bmod N\rangle = |1\rangle$$

\Rightarrow Quantum phase estimation for f with input $|1\rangle$ produces the state

$$\frac{1}{r} \sum_{s=0}^{r-1} |2^{\frac{s}{r}}\rangle$$

Note that $r \leq N$ hence $\frac{s}{r}$ can be exactly represented with n bits

- To efficiently implement $U^{2^k}|y\rangle = |xy^{2^k} \bmod N\rangle$

We use the fact k -times

$$y^{2^k} = (((y^2)^2)^2). \text{ Hence } |xy^{2^k} \bmod N\rangle$$

requires 1 multiplication, k squares mod N

\Rightarrow can be implemented as in classical circuits with XOR & AND gates.

Shor's period finding algorithm

Note that this is the only quantum part of Shor's alg. $N \leq 2^n$

1) Prepare $| \psi \rangle = |\tilde{1}\rangle = |\overbrace{0 \dots 0}^n 1\rangle$ and U as before

2) Use quantum phase estimation with U and $|\psi\rangle$ to obtain the state

$$\frac{1}{\sqrt{r}} \sum_{s=0}^{r-1} |2^s \frac{s}{r}\rangle$$

3) Measurement gives a random number r from the set $\left\{ \frac{1}{r}, \dots, \frac{r-1}{r} \right\} \subseteq \left\{ \frac{i}{2^n} : i=0, \dots, 2^n - 1 \right\}$

4) Write $\frac{s}{r} = \frac{s_0}{r_0}$ with $\text{gcd}(s_0, r_0) = 1$. If s, r were coprime already, we have $r=r_0$ and found the period. Otherwise, repeat.

Remark: There are $O(\frac{r}{\text{ggr}})$ — Euler's totient function numbers $1 \leq k \leq r$ with $\text{gcd}(s, r)=1$.

\Rightarrow Success prob of step 4 is $O(\frac{1}{\text{ggr}})$

Since $\frac{1}{\text{ggr}} \geq \frac{1}{\text{ggr}N} = \frac{1}{\text{gpr}^n}$, step 4 requires on average gpr^n runs.

Remark

Implicitly, we used QFT to find the frequency of

$x^k \bmod N$, i.e. the period of x

Grover's algorithm

Problem: Given $F : \{0,1\}^n \rightarrow \{0,1\}$, find $x \in \{0,1\}^n$ with $F(x) = 1$ or determine $F = 0$.

Recall the query

$$Q_F(|i\rangle) = (-1)^{F(i)} |i\rangle$$

(again we identify $x \in \{0,1\}^n$ with $i = 0, \dots, 2^{n-1}$)

Define the Grover diffusion operator

$$U_S := 2 \underbrace{|S\rangle\langle S|}_{\text{Projection onto } |S\rangle} - I$$

where $|S\rangle$ denotes the uniform state

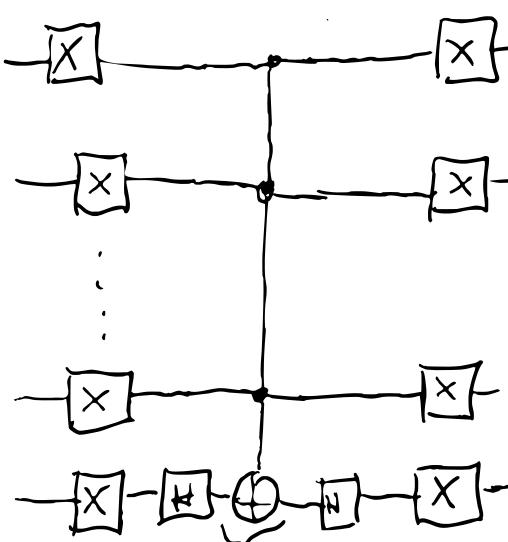
$$|S\rangle = \frac{1}{\sqrt{N}} \sum_{j=0}^{N-1} |j\rangle$$

Note that $|S\rangle = H^{\otimes n} |0^n\rangle$ and hence

$$U_S = H^{\otimes n} \underbrace{\left(2|0^n\rangle\langle 0^n| - I \right)}_{\text{corresponds to the matrix}} H^{\otimes n}$$

corresponds to the matrix

$$\begin{pmatrix} +1 & & & \\ -1 & 0 & & \\ & \ddots & \ddots & \\ 0 & & & -1 \end{pmatrix} \in \mathbb{C}^{2^n \times 2^n}$$



$$-(2|0><01 - I)$$

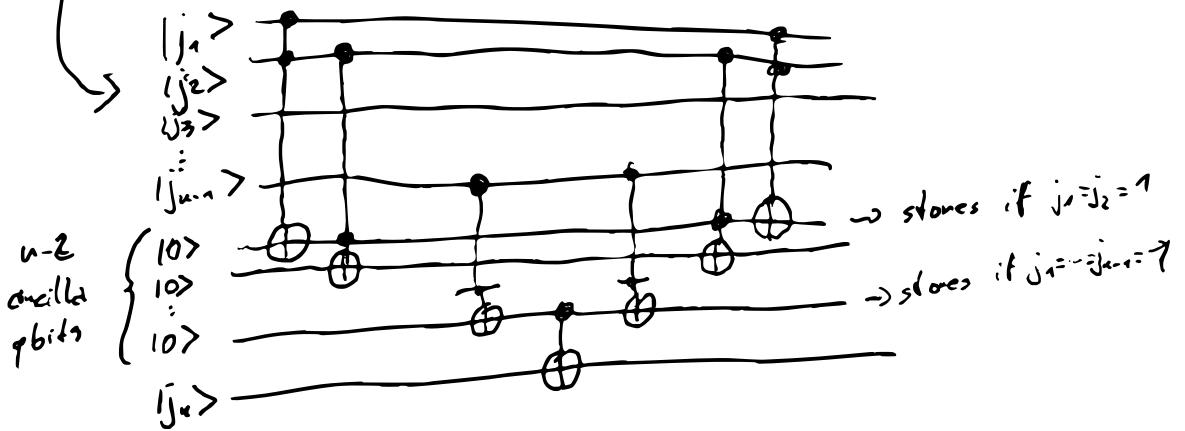
since we can only measure amplitudes,
 a global phase change
 (multiplication with
 $\lambda \in \mathbb{C}, |\lambda|=1$)
 is not noticeable
 alternatively, one

can apply

$$[I|Y][X]-[Y|X] \quad \text{do obtain}$$

$$\begin{pmatrix} 0 & 1 \\ 0 & 0 \end{pmatrix} \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \begin{pmatrix} 0 & 1 \\ 0 & 0 \end{pmatrix} \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} = \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}$$

generalized Toffoli gate



Algorithm

1) Apply $H^{\otimes n}$ to obtain $|S\rangle = \frac{1}{\sqrt{N}} \sum_{j=0}^{N-1} |j\rangle$

2) For $k=1, \dots, r(N)$ do

2a) Apply Q_F to current state $\sum_{j=0}^{N-1} \alpha_j |j\rangle$
to obtain $\sum_{j=0}^{N-1} (-1)^{F(j)} \alpha_j |j\rangle \rightarrow$

2b) Apply U_S to first n qubits

3) Measure in computational basis finds state

Then: With $r(N) = \frac{\arccos(\sqrt{\frac{1}{N}})}{2\arcsin(\sqrt{\frac{1}{N}})}$, Grover's alg. finds state

$|x\rangle$ with $F(x)=1$ with probability $\geq 1 - \frac{t}{N}$

Proof:

Define $|B\rangle := \frac{1}{\sqrt{N-t}} \sum_{F(j)=0} |j\rangle$, where $t = \#\bar{F}(|1\rangle)$

Note that the first n qubits of $|x\rangle \mapsto Q_S^{-1}(x\rangle)$
satisfy $|x\rangle \mapsto (2|B\rangle\langle B| - I)|x\rangle = U_B(|x\rangle)$

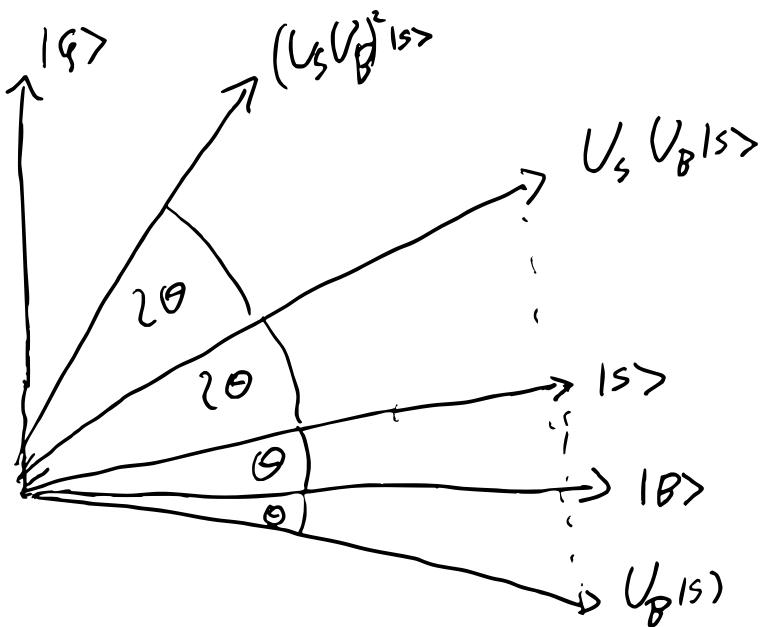
(check for basis elements $|x\rangle = |j\rangle$)

Define $|g\rangle := \frac{1}{\sqrt{t}} \sum |j\rangle$ and note

$|S\rangle \in \text{span}\{|g\rangle, |B\rangle\}$; $U_S, U_B : \text{span}\{|g\rangle, |B\rangle\} \hookrightarrow$

hence, the Grover iteration never leaves the plane

$\text{span}\{|g\rangle, |B\rangle\}$



Each application of $U_s U_B$ rotates a state $|x\rangle$ towards $|g\rangle$ by an angle Θ given by $\text{span}\{|g\rangle, |x\rangle\}$

$$\cos \Theta = \frac{\langle s | B \rangle}{|s| |B|} = \frac{1}{\sqrt{N}} \frac{1}{\sqrt{N-t}} \sum_{f(j)=0} \langle j | j \rangle = \frac{N-t}{\sqrt{N(N-t)}}$$

$$= \sqrt{1 - \frac{t}{N}}$$

$$\Rightarrow \sin^2 \Theta = 1 - \cos^2 \Theta = 1 - 1 + \frac{t}{N} \Rightarrow \sin \Theta = \sqrt{\frac{t}{N}}$$

$$\Rightarrow \Theta \approx \sqrt{\frac{t}{N}}$$

for large N .

Since the angle between $|g\rangle$ and $|s\rangle$ is α

$$\cos \alpha = \frac{1}{\sqrt{N}} \frac{1}{\sqrt{t}} \sum_{f(j)=1} \langle j | j \rangle = \frac{t}{\sqrt{N}} \Rightarrow r(N) = \text{round} \left(\frac{\arccos \left(\frac{t}{\sqrt{N}} \right)}{2 \arcsin \frac{t}{\sqrt{N}}} \right)$$

to obtain a state $|x\rangle$ with $\langle g|x\rangle \geq \cos(\theta)$

$$\Rightarrow |\langle g|x\rangle|^2 \geq \cos^2 \theta = 1 - \frac{t}{N}$$



Remark in practise, we don't know when $|x\rangle$ gets close to $|g\rangle$. But we can measure and check whether $F(x) = 1$. If not, restart the algorithm. If $t < N$

\Rightarrow

$$r(N) = \frac{\arccos \sqrt{\frac{t}{N}}}{\arcsin \sqrt{\frac{t}{N}}} \approx \frac{\pi}{4} \sqrt{\frac{N}{t}}$$

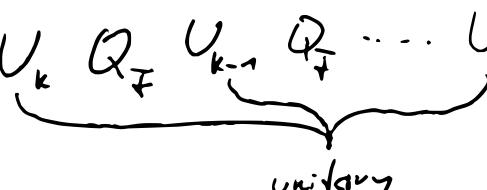
Remark If $t=1$, a classical algorithm requires at least N evaluations of F . Grover's algorithm only needs $O(\sqrt{N})$ iterations with $O(n)$ gates

Optimality of prover's Algo

Lemma Any Quantum Algo based on the Query Q_F requires at least $O(\delta \sqrt{N})$ applications of Q_F to succeed with prob. $\geq \delta^2$.

Proof Any Q-Algo starts with some state $|\psi\rangle$ and applies Unitary transformations as well as Q_F . i.e. we may write the state after k applications of Q_F as

$$|\psi_k^F\rangle := U_k Q_F U_{k-1} Q_F \dots U_1 Q_F |\psi\rangle$$



Additionally, we will consider

$$|\psi_k\rangle := U_k U_{k-1} \dots U_1 |\psi\rangle$$

Let $\overline{f}_j : \{0,1\}^n \rightarrow \{0,1\}$ with $\overline{f}_j(|i\rangle) = d_{ij}$

and define

$$D_k = \sum_{j=0}^{N-1} \|\psi_k^F - \psi_k\|^2$$

Idea: If D_k is small, the evolution of \tilde{F} doesn't make a big difference and it will be hard to find $\tilde{F}(|i\rangle) = 1$.

Step 1: Show that $D_k \leq 4k^2$ by induction.

$$k=0: D_0 = 0 \quad \checkmark$$

$$\begin{aligned} k \mapsto k+1: D_{k+1} &= \sum_{j=0}^{N-1} \| U_{k+j} Q_{\tilde{F}_j} \psi_k^{\tilde{F}_j} - U_{k+j} \psi_k \| ^2 \\ &= \sum_{j=0}^{N-1} \| Q_{\tilde{F}_j} \psi_k^{\tilde{F}_j} - \psi_k \| ^2 \\ &= \sum \| Q_{\tilde{F}_j} (\psi_k^{\tilde{F}_j} - \psi_k) + (Q_{\tilde{F}_j} - I) \psi_k \| ^2 \end{aligned}$$

$$\begin{aligned} \text{Note that } Q_{\tilde{F}_j} &= I - 2|j\rangle\langle j| \Rightarrow (Q_{\tilde{F}_j} - I)|\psi_k\rangle \\ &= -2|j\rangle(\langle j|\psi_k\rangle) \end{aligned}$$

$$\begin{aligned} \Rightarrow D_{k+1} &\leq \sum_{j=0}^{N-1} \| \psi_k^{\tilde{F}_j} - \psi_k \| ^2 + 4 \| \psi_k^{\tilde{F}_j} - \psi_k \| |\langle j|\psi_k\rangle| \\ &\quad + 4 |\langle j|\psi_k\rangle|^2 \end{aligned}$$

Cauchy - Schwartz shows

$$D_{k+1} \leq D_k + 4 \left(\sum_j \| \psi_k^{\pi_j} - \psi_k \|^2 \right)^{\frac{1}{2}} \left(\sum_j |\langle j | \psi_k \rangle|^2 \right)^{\frac{1}{2}}$$
$$= 4 \underbrace{\langle \psi_k | \psi_k \rangle}_{=1}^2$$
$$\leq D_k + 4\sqrt{D_k} + 4$$

Induction Hyp. $\rightarrow \leq 4k^2 + 8k + 4 = \underline{4(k+1)^2}$ ✓

This concludes the induction and shows

$$\underline{D_k \leq 4k^2}$$

Step 2: Assume $|\langle j | \psi_k^{\pi_j} \rangle|^2 \geq \sum_{j=0}^N \text{if } j=0, \dots, k, \text{ i.e.}$
the Alp. works with $\kappa^{j, \theta}$ prob. for each input.

Replacing $|j\rangle$ with $e^{i\theta}|j\rangle$ does not change
success prob. \Rightarrow we may assume $|\langle j | \psi_k^{\pi_j} \rangle| = |\langle j | \psi_k \rangle|$.

$$\Rightarrow \| \psi_k^{\pi_j} - j \|^2 = 2 - 2|\langle j | \psi_k^{\pi_j} \rangle| \leq \underline{2(1-\delta)}$$

$$\text{Define } E_k := \sum_{j=0}^{k-1} \| \psi_k^{\pi_j} - j \|^2 \Rightarrow E_k \leq 2N(1-\delta)$$

$$\overline{\pi}_k := \sum_{j=0}^{k-1} \| j - \psi_k \|^2$$

$$\begin{aligned}
 \Rightarrow D_k &= \sum_{j=0}^{N-1} \|(\psi_k^{\top j} - j) + (j - \psi_k)\|^2 \\
 &\geq \sum_j \| \psi_k^{\top j} - j \|^2 - 2 \| \psi_k^{\top j} - j \| \| j - \psi_k \| \\
 &\quad - \| j - \psi_k \|^2 \\
 &= E_k + F_k - 2 \left(\sum_j \| \psi_k^{\top j} - j \|^2 \right)^{\frac{1}{2}} \left(\sum_j \| j - \psi_k \|^2 \right)^{\frac{1}{2}} \\
 &= E_k + F_k - 2 \sqrt{E_k F_k} = \underline{(\sqrt{E_k} - \sqrt{F_k})^2}
 \end{aligned}$$

Note that any state $|\phi\rangle$ satisfies

$$\begin{aligned}
 \sum_{j=0}^{N-1} \| \phi - j \|^2 &= \sum_j \| \phi \|^2 - 2 \langle \phi | j \rangle + \sum_j \| j \|^2 \\
 &\geq 2N - 2 \sqrt{\sum_j 1} \underbrace{\sqrt{\sum_j |\langle \phi | j \rangle|^2}}_{=1} \\
 &= 2(N - \sqrt{N})
 \end{aligned}$$

This implies $F_k \geq 2(N - \sqrt{N})$ and hence

$$\begin{aligned}
 F_k &\geq E_k \text{ for sufficiently large } N \text{ and} \\
 \sqrt{F_k} - \sqrt{E_k} &\geq \sqrt{2(N - \sqrt{N})} - \sqrt{2(1-\delta)N} = \frac{2\sqrt{N} - 2\sqrt{N}}{\sqrt{2(N - \sqrt{N})} + \sqrt{2(1-\delta)N}}
 \end{aligned}$$

$$\geq \frac{\epsilon \delta N - \epsilon \sqrt{N}}{\epsilon \sqrt{2N}} = \frac{\delta}{\sqrt{2}} \sqrt{N} - \frac{1}{\sqrt{2}}$$

$$\Rightarrow D_k \geq (\widehat{F}_k - \widehat{E}_k)^2 \simeq \delta^2 N$$

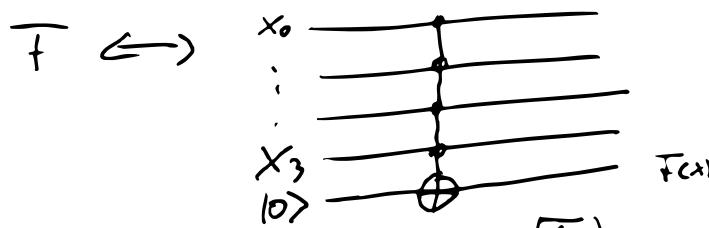
$$\text{Since } D_k \leq 4k^2 \Rightarrow k \simeq \delta \sqrt{N}$$

□

Example in Quantum Computer

Link on webpage, $N=16$, $n=4$

$$\overline{f}(x) = \overline{f}(x_0, \dots, x_3) = x_0 \text{ AND } x_1 \text{ AND } x_2 \text{ AND } x_3 \\ = \sum_{x=(1,1,1,1)} f = 1$$



There holds: $\frac{\arccos(\sqrt{\frac{1}{16}})}{2 \arcsin(\sqrt{\frac{1}{16}})} \approx 2.6083$

$$\Rightarrow \text{optimal } r(N) = 3.$$

$$\text{Probability of success} \geq 1 - \frac{1}{N} = \frac{15}{16} \approx 0.9375$$

Numerical Quadrature:

Problem: Given $f: \{0,1\}^n \rightarrow [-1,1]$, compute $\frac{1}{N} \sum_{i=0}^{N-1} f^{(i)}$

Quantum Super Sampling

Assumption: Oracle $Q_f: |0\rangle \otimes |i\rangle \mapsto \left[\sqrt{1-p_{(i)}} |0\rangle + \sqrt{p_{(i)}} |1\rangle \right] \otimes |i\rangle$

1) Start with $|0\rangle \otimes |0\rangle \otimes |0\rangle$

2) Apply $H^{\otimes p} \otimes I \otimes H^{\otimes n}$ \rightarrow

$$\frac{1}{\sqrt{PN}} \sum_{i=0}^{N-1} \sum_{m=0}^{p-1} |m\rangle \otimes |0\rangle \otimes |i\rangle$$

3) Apply Q_f \rightarrow

$$\frac{1}{\sqrt{PN}} \sum_i \sum_m |m\rangle \otimes \left[\sqrt{1-p_{(i)}} |0\rangle + \sqrt{p_{(i)}} |1\rangle \right] \otimes |i\rangle$$

4) Define $|G\rangle := \sqrt{\frac{\sum p_{(i)}}{N}} \sum_i \sqrt{p_{(i)}} |1\rangle \otimes |i\rangle$

$$|B\rangle := \sqrt{\frac{\sum 1-p_{(i)}}{N}} \sum_i \sqrt{1-p_{(i)}} |0\rangle \otimes |i\rangle$$

$$\bar{P} := \frac{1}{N} \sum p_{(i)}$$

Note

$$U_s := \left(2|0\rangle\langle 0| - I \right) \otimes I$$

$$U_s |B\rangle = |B\rangle$$

$$U_s |G\rangle = -|G\rangle$$

$$U_\psi = \left(2|\psi\rangle\langle\psi| - I \right)$$

$$\text{with } |\psi\rangle := \sqrt{1-p} |B\rangle + \sqrt{p} |G\rangle$$

Steps 1-3 provide quantum circuit to implement
the map $U: |0\rangle \otimes |0^n\rangle \mapsto |\psi\rangle$

$$\Rightarrow U_\psi = \tilde{U} \left(2|0^{n+1}\rangle\langle 0^{n+1}| - I \right) U$$

as in Grover's algorithm.

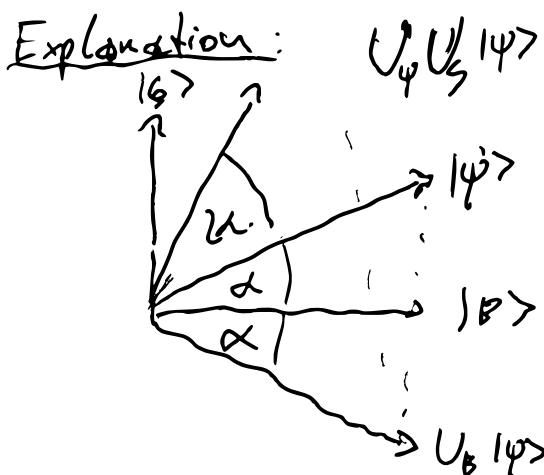
5) State reads

$$\frac{1}{\sqrt{P}} \sum_{m=0}^{P-1} |m\rangle \otimes [\sqrt{1-p}|B\rangle + \sqrt{p}|G\rangle]$$

for $\bar{P} := \frac{1}{n} \sum f(i)$

6) Apply grover iteration $U_\psi U_g$ m -times to second $n+1$ qubits to obtain

$$\frac{1}{\sqrt{P}} \sum_{m=0}^{P-1} |m\rangle \otimes (U_\psi U_g)^m [\sqrt{1-p}|B\rangle + \sqrt{p}|G\rangle]$$



Let $\sin \theta = \sqrt{p}$
 $\Rightarrow |ψ\rangle = \cos \theta |B\rangle + \sin \theta |G\rangle$

$$(U_\psi U_g)^m |ψ\rangle = \cos((2m+1)\theta) |B\rangle + \sin((2m+1)\theta) |G\rangle$$

U_g acts like reflection across $|B\rangle$

Note: m -times application of $U_\psi U_g$ is in "phase estimation"-alg. But: Cost in general $O(P_m)$

7) $|\psi\rangle$ rotates with rate 2θ in $|B\rangle - |g\rangle$ -plane. Use QFT to obtain rate.

Measure basis $\frac{1}{\sqrt{P}} \sum_{m=0}^{P-1} |m\rangle \otimes \left[\cos((2m+1)\theta)|B\rangle + \sin((2m+1)\theta)|g\rangle \right]$
and qbits in $\{|g\rangle, |B\rangle\}$ -basis to obtain

$$\frac{1}{C} \sum_{m=0}^{P-1} \underbrace{\sin[(2m+1)\theta]}_{x_m} |m\rangle \quad (\text{or } \cos(2m+1)\theta)$$

$$\text{where } C = \frac{1}{P} \sum_{m=0}^{P-1} \sin[(2m+1)\theta]^2.$$

→ Apply QFT to obtain [assume $\Theta = \frac{\pi\theta_0}{P}, \theta_0 \in \mathbb{R}$]

$$\frac{1}{C} \sum_{m=0}^{P-1} x_m |m\rangle, \text{ where } x_m \text{ is given}$$

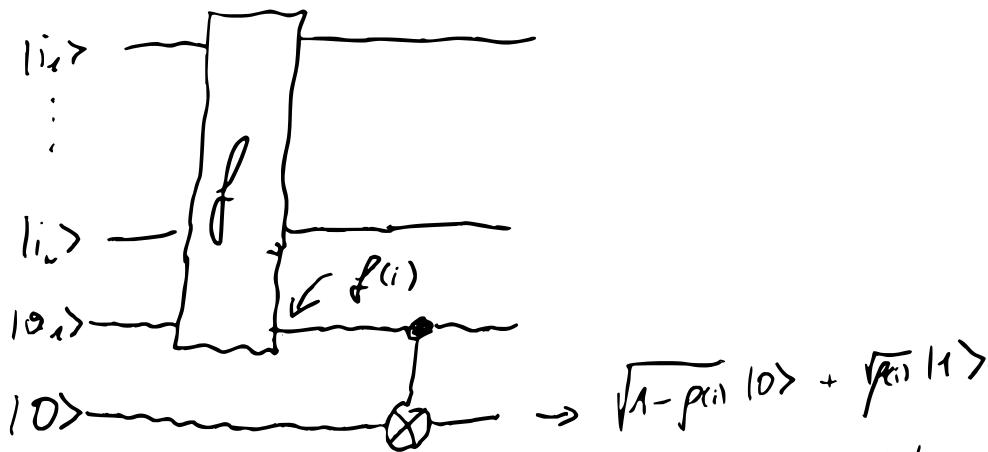
by FFT of \hat{x}_m ; i.e.

$$\begin{aligned} x_m &= \frac{1}{\sqrt{P}} \sum_{k=0}^{P-1} \hat{x}_k e^{\frac{2k\pi i m}{P}} \\ &= \frac{1}{\sqrt{P}} \sum_{k=0}^{P-1} \frac{1}{2i} \begin{pmatrix} e^{i(2k+1)\frac{\theta_0}{P}} & -e^{-i(2k+1)\frac{\theta_0}{P}} \end{pmatrix} e^{\frac{2k\pi i m}{P}} \\ &= e^{i\frac{\theta_0}{P}} \frac{1}{2i} \begin{cases} \sqrt{P} & m = \theta_0 \\ 0 & \text{else} \end{cases} \end{aligned}$$

8) Final measurement produces $|m\rangle = |\theta_0\rangle$
 and hence $\bar{f} = \sin\left(\frac{\pi\theta_0}{p}\right)^2$

Remark: implementation of f unclear.

If $f: \{0,1\}^n \rightarrow \{0,1\}$, we can use
 controlled NOT-gate to implement



general integrals can always be split into

$$\frac{1}{N} \sum_{i=0}^{N-1} f(i) = \frac{1}{N} \sum_{i=0}^{N-1} \sum_{k=1}^R 2^{-k} f_k(i) \in \{0,1\}$$

i.e., R integration problems for precision 2^{-R} .