

ET4340 Electronics for Quantum Computing

Homework 5

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Problem 1: Warmup exercises

1. Consider an operator M and one of its eigenstates $|\psi\rangle$ (with eigenvalue λ). Consider another operator A that anticommutes with M (i.e., $\{M, A\} \equiv MA + AM = 0$). Show that the state $A|\psi\rangle$ is an eigenstate of M with eigenvalue $-\lambda$.

We can write $M|\psi\rangle = \lambda|\psi\rangle$.

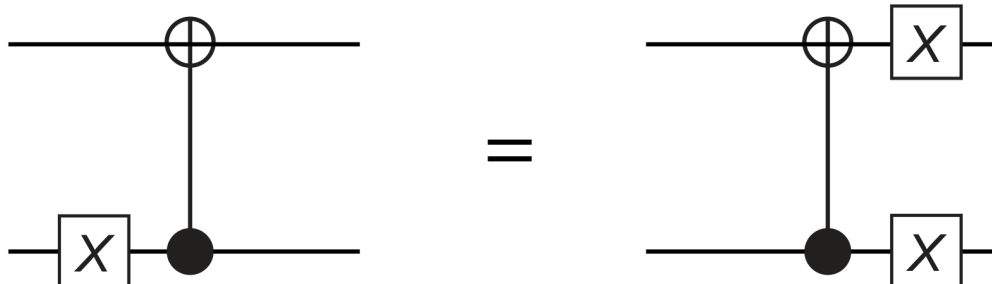
$$\begin{aligned}MA + AM &= 0 \\MA|\psi\rangle + AM|\psi\rangle &= 0 \\MA|\psi\rangle + A\lambda|\psi\rangle &= 0 \\M(A|\psi\rangle) &= -\lambda(A|\psi\rangle)\end{aligned}$$

2. Now consider an operator B that commutes with M (i.e., $[M, B] \equiv MB - BM = 0$). Show that the state $B|\psi\rangle$ is an eigenstate of M with eigenvalue λ .

Assuming M again has an eigenstate $|\psi\rangle$ with eigenvalue λ .

$$\begin{aligned}MB - BM &= 0 \\MB|\psi\rangle - BM|\psi\rangle &= 0 \\M(B|\psi\rangle) &= \lambda(B|\psi\rangle)\end{aligned}$$

3. Prove the identity:



Feel free to do this either by multiplying matrices or by manipulating circuit diagrams. From this we see that a single-qubit bit-flip error prior to CNOT proliferates into a double bit-flip error.

It is an intuitive identity in my opinion. Flipping the control bit flips the output. So to ‘simulate’ the flip before a CNOT you can flip both the output and the control after the CNOT.

By matrix multiplication (note that the order of operations is reversed with respect of the diagram because we compute the combined matrix $M = M_n \dots M_2 M_1$ in $|out\rangle = M |in\rangle$ where M_n is operation n):

$$\begin{aligned}
 M_a = \text{CNOT}_{01}(I \otimes X) &= \begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 1 \\ 0 & 0 & 1 & 0 \\ 0 & 1 & 0 & 0 \end{bmatrix} \begin{bmatrix} 0 & 1 & 0 & 0 \\ 1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 1 \\ 0 & 0 & 1 & 0 \end{bmatrix} = \begin{bmatrix} 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \\ 1 & 0 & 0 & 0 \end{bmatrix} \\
 M_b = (X \otimes X)\text{CNOT}_{01} &= \begin{bmatrix} 0 & 0 & 0 & 1 \\ 0 & 0 & 1 & 0 \\ 0 & 1 & 0 & 0 \\ 1 & 0 & 0 & 0 \end{bmatrix} \begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 1 \\ 0 & 0 & 1 & 0 \\ 0 & 1 & 0 & 0 \end{bmatrix} = \begin{bmatrix} 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \\ 1 & 0 & 0 & 0 \end{bmatrix} \\
 M_a &= M_b
 \end{aligned}$$

Problem 2: Three-qubit bit flip code

Consider the 3-qubit bit-flip code as covered in the lectures. In this code, a one-qubit state $|\psi\rangle = \alpha |0_2\rangle + \beta |1_2\rangle$ is encoded as $|\Psi\rangle = \alpha |0_3 0_2 0_1\rangle + \beta |1_3 1_2 1_1\rangle$.

1. Suppose the encoded state is distorted by a rotation of 60° about the $+\hat{x}$ axis of qubit 3. What are the possible error syndromes you could measure (i.e., the measurement results m_a and m_b)? Show that the state $|\Psi\rangle$ is recovered after error correction, every time.

Since the distortion is not a full flip you can measure qubit 3 as either -1 or $+1$. The other two bits are not affected and will retain their (equal) states.

The Z (parity) measurement m_a of qubits q_1 and q_2 will always give $m = +1$ because $q_1 = q_2$. The same does not hold for qubits q_2 and q_1 , you will measure $m = \pm 1$ by chance.

The error syndromes thus are $m_a = +1, m_b = +1$ and $m_a = +1, m_b = -1$. In the first case you won't have to do anything. In the second case you will have to flip qubit q_3 back.

Thus the final state is always ‘correct’ or valid so to speak.

2. Suppose now that instead the encoded state is distorted by a rotation of 45° about the $+\hat{y}$ axis of qubit q_2 , but you don't know it and stick to using the bit flip code without modifications. What are the possible error syndromes you would measure? Can you recover the state $|\Psi\rangle$ every time? When do you succeed and when do you not? When you don't recover it, what is the erroneous final state of the logical qubit?

This time we apply a maybe- Y gate to qubit q_2 . Since $Y = iXZ$ and the measurements we perform are two Z measurement, we can make some conclusions about the different cases.

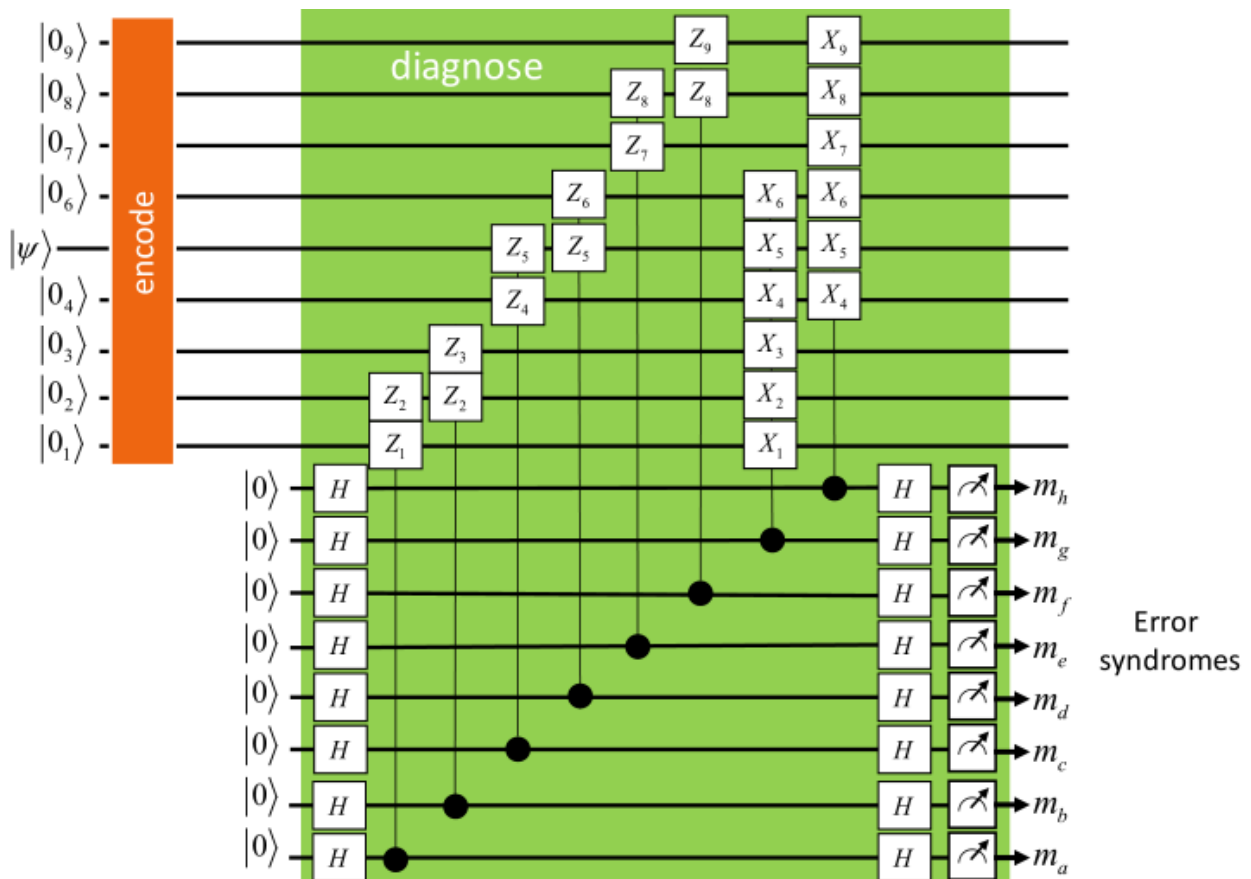
We can recover the error when the q_2 becomes Xq_2 because the syndrome is $m_a = -1, m_b = -1$. This syndrome induces a flip on q_2 which makes $XXq_2 = q_2$.

When q_2 becomes Zq_2 , the syndrome is $m_a = +1, m_b = +1$. The final state in this case is $|\Psi\rangle = \alpha |0_3 +_2 1_1\rangle + \beta |0_3 -_2 1_1\rangle$.

I think my answer is in the right direction but I'm having trouble writing it down and proving it mathmatically.

Shor's 9-qubit code

Consider Shor's 9-qubit code as covered in the lecture.



1. Suppose a phase flip occurred on qubit q_4 . What error syndromes (m_a, \dots, m_h) will you measure?

A phase flip means a Z transition so the error syndrome for Z_4 is m_g, m_h . I will write down only the abnormal measurements ($m = -1$).

2. Now suppose a phase flip occurred on qubit q_5 .

The syndrome for Z_5 is m_g, m_h as well.

3. Now suppose a phase flip occurred on qubit q_6 .

The syndrome for Z_6 is m_g, m_h as well.

4. Explain the results from the previous three assignments.

Apparently we get the same syndrome when a phase flip occurs on qubits q_4, q_5 and q_6 . Judging from Shor's previous work there has to be something clever going on behind the scenes here.

The fact is that we can fix a phase flip on a qubit q_z that is in any of the sets of three qubits $q_1 \cdots q_3, q_4 \cdots q_6$ and $q_7 \cdots q_9$ by performing a phase flip in any one (or all three) of the qubits in the set of qubits that q_z is in. So Z_4 can be fixed by doing a Z on qubit q_4, q_5 or q_6 (or all three).

This property follows from the fact that if you do an even number of Z transformations, the state of the logical qubit will stay the same.

Lets see what Z_2 does:

$$(I \otimes Z \otimes I)(\alpha|000\rangle + \beta|111\rangle) = \begin{bmatrix} 1 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & -1 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & -1 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & -1 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & -1 \end{bmatrix} \begin{bmatrix} \alpha \\ 0 \\ 0 \\ 0 \\ 0 \\ 0 \\ 0 \\ \beta \end{bmatrix} = \begin{bmatrix} \alpha \\ 0 \\ 0 \\ 0 \\ 0 \\ 0 \\ 0 \\ -\beta \end{bmatrix}$$

This just demonstrates that a Z transformation indeed does not affect $\alpha|000\rangle$ and flips the phase of $\beta|111\rangle$.

5. Suppose you measure the syndromes $m_a = m_b = m_c = m_d = m_g = m_h = -1, m_e = m_f = +1$. What error does this syndrome detect? Hint: it is not a single-qubit error. Interestingly, this shows that Shor's code can correct at least some two-qubit errors!

Table 1 table describes what cubits parity measurements are affected by which transformations (X, Y or Z errors) on each qubit.

We are looking for the syndrome in the first row:

?	•	•	•	•	•	•	5
Y_5			•	•	•	•	5
X_2	•	•			○		2

As you can see from the table, this syndrome can be produced by an Y on q_5 and an X on q_2 . This is the only 2-qubit error that will generate this syndrome.

	<div style="display: flex; justify-content: space-around; align-items: center;"> <div style="text-align: center;">Y</div> <div style="text-align: center;">X</div> <div style="text-align: center;">Z</div> </div>								
	a	b	c	d	e	f	g	h	
9						•		•	9
8					•	•		•	8
7					•			•	7
6				•			•	•	6
5			•	•			•	•	5
4			•				•	•	4
3		•					•		3
2	•	•					•		2
1	•						•		1
	a	b	c	d	e	f	g	h	

Table 1: Pauli operations affecting qubits in Shor's 9-qubit code

Problem 4: Operations on logical qubits One aspect that makes certain error-correction codes more practical than others is the ability to perform logical operations directly on the encoded qubits. This is the reason why the most 'economical' 5-qubit code has never been very popular. Consider again the Shor 9-qubit code.

1. Show that the operation $X_9 \otimes X_8 \dots X_1$ performs the logical Z operation (i.e., $|0_{Shor}\rangle \rightarrow |0_{Shor}\rangle$, $|1_{Shor}\rangle \rightarrow -|1_{Shor}\rangle$).

$$|0_{Shor}\rangle = \frac{1}{\sqrt{8}} (|000\rangle + |111\rangle)^{\otimes 3}$$

$$|1_{Shor}\rangle = \frac{1}{\sqrt{8}} (|000\rangle - |111\rangle)^{\otimes 3}$$

By flipping all the qubits in $|0_{Shor}\rangle$ we get:

$$(X_9 \otimes X_8 \dots X_1) |0_{Shor}\rangle = \frac{1}{\sqrt{8}} (|111\rangle + |000\rangle)^{\otimes 3} = |0_{Shor}\rangle$$

By doing the same for $|1_{Shor}\rangle$ we obtain:

$$\begin{aligned} (X_9 \otimes X_8 \dots X_1) |1_{Shor}\rangle &= \frac{1}{\sqrt{8}} (|111\rangle - |000\rangle)^{\otimes 3} \\ &= (-1)^3 \frac{1}{\sqrt{8}} (|000\rangle - |111\rangle)^{\otimes 3} = -|1_{Shor}\rangle \end{aligned}$$

2. Show that the operation $Z_9 \otimes Z_8 \dots Z_1$ performs the logical X operation (i.e., $|0_{Shor}\rangle \rightarrow |1_{Shor}\rangle$, $|1_{Shor}\rangle \rightarrow |0_{Shor}\rangle$).

Z flips the phase so $\frac{1}{\sqrt{2}}(|0\rangle + |1\rangle)$ becomes $\frac{1}{\sqrt{2}}(|0\rangle - |1\rangle)$ and vice versa. Evidently, $\frac{1}{\sqrt{8}}(|000\rangle + |111\rangle)^{\otimes 3}$ becomes $\frac{1}{\sqrt{8}}(|000\rangle - |111\rangle)^{\otimes 3}$ and vice versa because of the 9 phase flips. So the operation $Z_9 \otimes Z_8 \dots Z_1$ performs the logical X operation because $|0_{Shor}\rangle \rightarrow |1_{Shor}\rangle$ and $|1_{Shor}\rangle \rightarrow |0_{Shor}\rangle$.

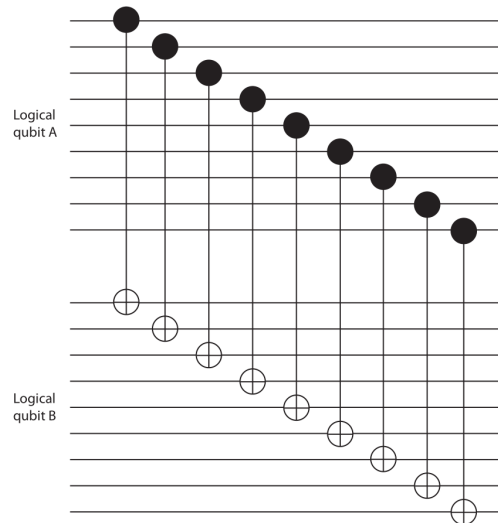
3. What logical operation does $Y_9 \otimes Y_8 \dots Y_1$ do?

Since $Y = iXZ$ and taking the previous results into account, the operation will implement the logical Y .

4. Can you think of a simpler way to realize this logical X operation?

We saw in problem 2.4 we can get away with a single Z transformation in each set of qubits $q_9 \dots q_7$, $q_6 \dots q_4$ and $q_3 \dots q_1$.

5. Finally, consider two logical qubits, A and B , each encoded using Shor's 9-qubit code, and the transversal quantum circuit below:



What operation does this circuit perform on the two logical qubits? (Hint: It's not quite what you think of at first glance!).

There are four cases to be considered because both A and B can be in either $|0_{Shor}\rangle$ or $|1_{Shor}\rangle$.

The Shor states consist of three parts as seen in previous exercises and we will look at them one at a time.

The set is in the state $\frac{1}{\sqrt{2}}(|000\rangle \pm |111\rangle)$. Since we do a **CNOT** on single qubits from A and B we can look at the problem from the perspective of qubits $q_{A,n}$ and $q_{B,n}$. But, I want to write them shorter so these qubits are now a and b respectively.

Using the following definitions:

$$|+\rangle = \frac{1}{\sqrt{2}}(|0\rangle + |1\rangle)$$

$$|-\rangle = \frac{1}{\sqrt{2}}(|0\rangle - |1\rangle)$$

We can calculate the result of applying CNOT_{ab} to the 4 possible combinations of the states of a and b :

$a \otimes b$	CNOT_{ab}
$ +\rangle \otimes +\rangle$	$\rightarrow +\rangle \otimes +\rangle$
$ +\rangle \otimes -\rangle$	$\rightarrow -\rangle \otimes -\rangle$
$ -\rangle \otimes +\rangle$	$\rightarrow -\rangle \otimes +\rangle$
$ -\rangle \otimes -\rangle$	$\rightarrow +\rangle \otimes -\rangle$

The strange thing is that a actually gets phase flipped if $b = |-\rangle$. Now we can extend this to the level of logical qubits and we get that:

$$|0_{Shor}\rangle \otimes |0_{Shor}\rangle \rightarrow |0_{Shor}\rangle \otimes |0_{Shor}\rangle$$

$$|0_{Shor}\rangle \otimes |1_{Shor}\rangle \rightarrow |1_{Shor}\rangle \otimes |1_{Shor}\rangle$$

$$|1_{Shor}\rangle \otimes |0_{Shor}\rangle \rightarrow |1_{Shor}\rangle \otimes |0_{Shor}\rangle$$

$$|1_{Shor}\rangle \otimes |1_{Shor}\rangle \rightarrow |0_{Shor}\rangle \otimes |1_{Shor}\rangle$$

So, the CNOT_{ab} 's actually implement a logical CNOT_{BA} .

These were mostly fun and valuable exercises :), thank you for your efforts so far I really appreciate it!

Happy holidays!

