CE 225: Dynamic of Structures

Fall 2025

Discussion 3: Arbitrary Forcing Functions

Instructor: Matthew DeJong GSI: Miguel A. Gomez

Plan for today

The plan for today's discussion includes:

- 1. Conceptual review of transmissibility (study summary).
- 2. Example of solving the EOM for an arbitrary forcing function.

Transmissibility 1

$$TR = \frac{\sqrt{1 + \left(2\zeta \frac{\omega}{\omega_n}\right)^2}}{\sqrt{\left(1 - \left(\frac{\omega}{\omega_n}\right)^2\right)^2 + \left(2\zeta \frac{\omega}{\omega_n}\right)^2}}$$

This can mean two different things, depending on the type of EOM that we have.

Harmonic Force Case

If our system is subjected to a harmonic force, then the EOM looks like:

 $m\ddot{u} + c\dot{u} + ku = p_0 \sin \omega t$

In this case:

Harmonic Ground Motion Case

If our system is subjected to a ground motion at its base, then our EOM for the relative displacement of the system w/r to the ground looks like:

$$\boxed{m\ddot{u} + c\dot{u} + ku = -m\ddot{u}_g} \implies \textbf{p(t)} = -m \text{ ig(t)} \tag{1.2}$$
 Where $\ddot{u}_g = \ddot{u}_{go} \sin{(\omega t)}$ = -m \text{ ign} \sin{(\omega t)} \text{ (ut)} \text{}.

In this case, we can use the Transmissibility equation for two purposes: (1) to find the ratio of the motion

transmitted from the ground to the system:

$$TR = \frac{u_o^t}{u_{go}} = \frac{\ddot{u}_o^t}{\ddot{u}_{go}} \Longrightarrow \begin{array}{c} \text{Plak total motion on our system.} \\ \text{peak motion of the ground.} \end{array}$$

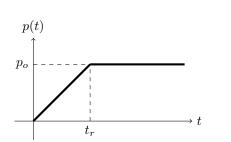
Or, if we compare our equations (1.1) and (1.2), we can define $p_o = -m\ddot{u}_{qo}$ and use the transmissibility equation to find a relation between the force transmitted to our system, and the inertial force generated by the ground motion:

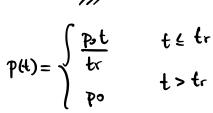
$$TR = \frac{f_{T,\text{max}}}{p_o} = \frac{f_{T,\text{max}}}{m\ddot{u}_{qo}}$$

2 Arbitrary Forcing Functions

Example

An undamped SDF system, starting from rest, is subjected to a step force with finite rise time, as the one shown in the Figure below.





The response of the system can be shown to be:

$$u(t) = (u_{st})_o \left(\frac{t}{t_r} - \frac{\sin \omega_n t}{\omega_n t_r}\right) \quad \text{for } t \le t_r$$
(2.3)

$$u(t) = (u_{st})_o \left\{ 1 + \frac{1}{\omega_n t_r} \left((1 - \cos(\omega_n t_r)) \sin(\omega_n (t - t_r)) - \sin(\omega_n t_r) \cos(\omega_n (t - t_r)) \right) \right\} \quad \text{for } t > t_r \quad (2.4)$$

Equation (2.4) can be simplified through trigonometric identities to:

$$u(t) = (u_{st})_o \left\{ 1 - \frac{1}{\omega_n t_r} \left[\sin(\omega_n t) - \sin(\omega_n (t - t_r)) \right] \right\}$$

Derive these results with:

- (a) Using the convolution integral. (Dhame's Integral).
- (b) Solving the EOM with analytical procedure.
- (c) Using superposition

(a) Convolution Integral

For undamped systems, starting from rest subjected to a force defined by p(t), we can find the response as:

$$u(t) = \frac{1}{m\omega_n} \int_0^t p(\tau) \sin{(\omega_n(t-\tau))} d\tau \qquad \text{and} \qquad \text{case} .$$

We also have the damped version of the convolution integral as:

$$P(\tau) = \begin{cases} \frac{1}{m\omega_{D}} \int_{0}^{t} p(\tau)e^{-t\omega_{n}(t-\tau)} \sin(\omega_{D}(t-\tau))d\tau \\ \frac{1}{t} \int_{0}^{t} \tau \geq t\tau \\ p_{D} \tau > t\tau \end{cases}$$

$$U(t) = \frac{1}{m\omega_{n}} \int_{0}^{t} \frac{p_{D}\tau}{t\tau} \cdot \sin(\omega_{n}(t-\tau)) d\tau \iff \text{integration by parts}$$

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Discussion 3: Arbitrary Forcing Functions

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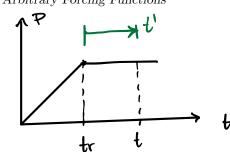
$$\frac{Po}{m \cdot w_n^2} \left[\cos \left(w_n(t-tr) \right) - \frac{1}{w_n} \cdot \sin \left(w_n(t-tr) \right) + \frac{Po}{m \cdot w_n^2} \left[\cos \left(w_n(t-t) \right) - \cos \left(w_n(t-tr) \right) \right]$$

$$= \frac{Po}{k} \left[1 - \frac{1}{w_n} \left(\sin \left(w_n \left(t - tr \right) \right) \right) \right]$$

目

(b) Solving the EOM

(i)
$$t \in tr$$
 $P(t) = \frac{Bot}{tr}$
 $min + kn = \frac{Bot}{tr}$
 $= A \cdot \omega s \cdot (unt) + B \cdot sin \cdot (unt) + up(t)$
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Eou: mitku = Po

$$t'(0)$$
 @ $t = tr$
 $t' = t - tr$

$$u(t') = A \cdot \cos(\omega_n t') + B \cdot \sin(\omega_n t') + up.$$

$$u(t) = A \cdot (os (unt))$$

$$u_{\rho} = C \Rightarrow u_{\rho}(o) = 0 \Rightarrow 0 + kC = P_{0} \Rightarrow C = \frac{P_{0}}{k} = (ust)_{0}.$$

$$\Rightarrow u(t') = A \cdot \cos(\omega_n t') + B \cdot \sin(\omega_n t') + (u_s t')_0$$

Naw, need initial conditions
$$(t') = 0$$
 (or $t = t_r$) $\rightarrow \frac{u_{camp}(t_r)}{u_{ramp}(t_r)}$

$$u_{ramp}(t) = (u_{st}), \left(\frac{t}{tr} - \frac{\sin(\omega_{n}t)}{tr \cdot \omega_{n}}\right) \Rightarrow u_{ramp}(tr) = (u_{st}), \left[1 - \frac{\sin(\omega_{n}t)}{tr \cdot \omega_{n}}\right]$$

$$\dot{u}$$
 ramp (t) = (Nut) $\left(\frac{1}{tr} - \frac{\cos(u_n t)}{u_n}\right) \Rightarrow \dot{u}$ ramp (tr) = (Nust) $\left[\frac{1}{tr} - \frac{\cos(u_n t)}{tr}\right]$

$$\Rightarrow A \cdot (1) + B \cdot (0) + (uct)_0 = (vet)_0 - (ut)_0 \cdot \frac{\sin(untx)}{tx \cdot un}$$

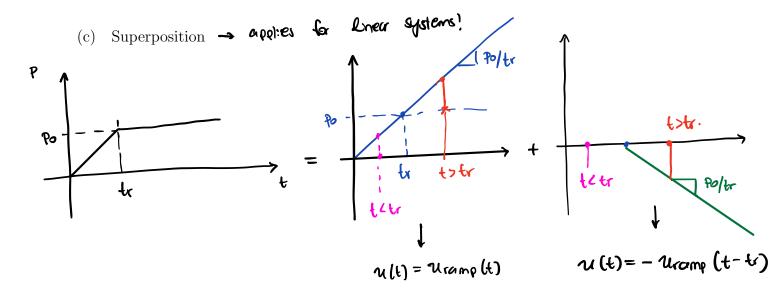
$$\Rightarrow A = -(ust)_0 \cdot \frac{s_0(u_0 t)}{u_0 tr}$$

$$\Rightarrow B = (nst)_0 \left[\frac{1}{w_n t_r} - \frac{\cos(w_n t_r)}{w_n t_r} \right]$$

So:
$$u(t') = (ust)_0 \left[-sin \frac{(unt)}{wnt} cos(unt') + \frac{1}{unt} (1-cos(unt)) \cdot sin(unt') + 1 \right]$$

with
$$t' = t - tr$$

$$u(t) = (ust)_0 \left[-\frac{\sin(untr)}{untr}, \cos(un(t-tr)) + \frac{1}{untr} (1 - \cos(untr)) \sin(un(t-tr)) + 1 \right]$$



for
$$t > tr$$

$$u(t) = u_{(comp)}(t) - u_{(comp)}(t-tr)$$

$$= (ust)_0 \left[\frac{t}{tr} - \frac{s_{(in)}(u_n t)}{tr} - \frac{t-tr}{tr} + \frac{s_{(in)}(u_n (t-tr))}{u_n tr} \right]$$

$$= (ust)_0 \left[1 - \frac{1}{u_n tr} \left(\frac{s_{(in)}(u_n t)}{u_n tr} - \frac{s_{(in)}(u_n (t-tr))}{u_n tr} \right) \right].$$