Magnetic vortex effects on first-order reversal curve (FORC) diagrams for greigite dispersions

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Abstract

First-order reversal curve (FORC) diagrams are used increasingly in geophysics for magnetic domain state identification. The domain state of a magnetic particle is highly sensitive to particle size, so FORC diagrams provide a measure of magnetic particles size distributions. However, the FORC signal of particles with nonuniform magnetisations, which are the main carrier of natural remanent magnetisations in many systems, is still poorly understood. In this study, the properties of non-interacting, randomly oriented dispersions of greigite (Fe_3S_4) in the uniform single-domain (SD) to non-uniform single-vortex (SV) size range are investigated via micromagnetic calculations. Signals for SD particles (< 50 nm) are found to be in excellent agreement with previous SD coherent-rotation studies. A transitional range from ~50 nm to ~70 nm is identified for which a mixture of SD and SV behaviour produces complex FORC diagrams. Particles $>$ ~70 nm have purely SV behaviour with the remanent state for all particles in the ensemble represented by the vortex state. It is found that for SV ensembles the FORC diagram provides a map of vortex nucleation and annihilation fields and that the FORC distribution peak should not be interpreted simply as the coercivity of the sample, but as a vortex annihilation field on the path to saturation.

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1. Introduction

First-order reversal curve (FORC) diagrams are a powerful tool in rock magnetic studies, which allow mineral and domain state identification as well as quantification of magnetostatic interactions among particles (Pike et al., 1999; Roberts et al., 2000, 2014; Dumas et al., 2007; Egli et al., 2010). As such, they have been the subject of numerical studies aimed at relating the behaviour of individual magnetic particles and small assemblages to experimental bulk properties (Pike et al., 1999; Carvallo et al., 2003, 2006; Muxworthy et al., 2004; Muxworthy and Williams, 2005; Newell, 2005; Harrison and Lascu, 2014; Valdez-Grijalva and Muxworthy, 2018; Roberts et al., 2017).

With the exceptions of Carvallo et al. (2003) and Roberts et al. (2017), all of these numerical studies have concentrated on FORC diagrams for ideal, uniformly magnetised single-domain (SD) particles. They have shown that uniaxial SD particles produce patterns in FORC diagrams (Muxworthy et al., 2004; Newell, 2005; Harrison and Lascu, 2014), that are distinct from those for SD materials with cubic anisotropy (Muxworthy et al., 2004; Harrison and Lascu, 2014; Valdez-Grijalva and Muxworthy, 2018). However, it is well-documented that most natural systems have magnetic signals dominated by larger grains with more complex magnetic domain states (Dunlop and Özdemir, 1997; Roberts et al., 2017). Grains just above the SD threshold size (e.g., ~64 nm for equidimensional magnetite, ~54 nm for greigite), are typically in a single-vortex (SV) state. The SV state dominates magnetic structures over an order of magnitude of size variations (Nagy et al., 2017; Valdez-Grijalva et al., 2018), which is much wider than the stable SD size range. SV grains have recently been found to be geologically meta-stable and retain relatively high remanences (Almeida et al., 2014; Nagy et al., 2017; Valdez-Grijalva et al., 2018).

Previous experimental studies on nano-patterned arrays of SV particles (Pike and Fernandez, 1999; Dumas et al., 2007) found that FORC diagrams are significatively more complex than for SD signals, with complex off-axis "butterfly" patterns that are related to vortex nucleation/annihilation processes. However, it is difficult to relate the behaviour of 2D nano-patterned arrays to the behaviour of natural particle systems found in geological samples. In natural samples, particles with varying size and orientation are dispersed in 3 dimensions. Thus, it is important to understand the contribution of dispersions of randomly aligned SV particles to FORC diagrams. Numerical modelling can aid the study of such systems. Carvallo et al. (2003) used a finite-difference model to calculate the FORC distributions of SV magnetite particles; however, that study primarily examined the effects of interactions between small clusters of cubic grains, and neither random particle distributions nor realistc grain morphologies were included.

In this study, we employ a micromagnetic finite element method (FEM) to obtain FORC diagrams for non-interacting ensembles of SD and SV greigite (Fe_3S_4). Greigite is the iron-sulphide counterpart to magnetite. Recent interest in greigite comes from both its promising properties for material science (Li et al., 2014) and the abundance of this mineral in sedimentary rocks for Earth science (Roberts et al., 2011). FORC diagrams are often used to help identify greigite. The relatively high anisotropy of greigite means that the behaviour of this mineral is representative of cubic-anisotropic ferri- and ferro- magnets like magnetite and iron. We calculate FORC diagrams for simulations of non-interacting dispersions of randomly oriented greigite with sizes 30–80 nm; this size range covers the SD–SV threshold (Valdez-Grijalva et al., 2018). Simulations are carried out on an ensemble of 500 particles with random orientations. The unstructured discretisation of FEMs allows us to study realistic greigite particle shapes as observed in nature. We determine the onset of SV behaviour and its consequences for FORC diagram interpretation.

2. Methods

2.1. The micromagnetic algorithm

A ferromagnetic material—neglecting thermal and magnetostrictive effects—has a Gibbs free-energy functional given by (Brown, 1963):

$$E_{\mathrm{G}} = \int_{\Omega} (\phi_{\mathrm{exchange}} + \phi_{\mathrm{anisotropy}} + \phi_{\mathrm{stray}} + \phi_{\mathrm{external}})\, \mathrm{d}^3\boldsymbol{r}, \tag{1}$$

where Ω is the ferromagnetic volume. Here,

$$\phi_{\mathrm{exchange}} = A|\nabla \boldsymbol{m}|^2, \tag{2}$$

where $\boldsymbol{m}$ is the reduced magnetisation vector and A the exchange stiffness constant, provides an expression for the energy density due to quantum-mechanical exchange forces (Landau and Lifshitz, 1935).

$$\phi_{\mathrm{anisotropy}} = \frac{K_1}{2} \sum_{i \neq j} \gamma_i^2 \gamma_j^2 + K_2 \prod_i \gamma_i^2, \tag{3}$$

where γ_i represent the direction cosines and K_1 and K_2 the first and second magnetocrystalline anisotropy (MCA) constants, is the MCA energy density in the cubic anisotropy system. In terms of the reduced magnetisation vector components, this becomes:

$$\phi_{\mathrm{anisotropy}} = K_1(m_x^2 m_y^2 + m_y^2 m_z^2 + m_z^2 m_x^2), \tag{4}$$

where K_2 is neglected because K_1 is the dominant term at room temperature. The magnetostatic self-energy density is given by:

$$\phi_{\mathrm{stray}} = -\frac{\mu_0 M_{\mathrm{S}}}{2} \boldsymbol{m} \cdot \boldsymbol{H}_{\mathrm{stray}}, \tag{5}$$

where $\boldsymbol{H}_{\mathrm{stray}}$ is the stray field produced by the ferromagnetic body and M_{S} is the saturation magnetisation. Finally, the energy density due to an external magnetic field $\boldsymbol{H}_{\mathrm{external}}$ is:

$$\phi_{\mathrm{external}} = -\mu_0 M_{\mathrm{S}} \boldsymbol{m} \cdot \boldsymbol{H}_{\mathrm{external}}. \tag{6}$$

Such magnetic particle systems will be driven spontaneously toward an equilibrium state with a locally minimal magnetic Gibbs free-energy (Brown, 1963). In this study we utilise a modified gradient descent method to find the equilibrium magnetisation (Ó Conbhuí et al., 2017).

Discretisation of the spatial domain is achieved by decomposing the volume into tetrahedral elements. This allows modelling of particles with arbitrary geometries. To model accurately nonuniform magnetisations, spatial discretisation in the model should be smaller than the exchange length $l_{\text{exch.}} = \sqrt{2A/\mu_0 M_\text{S}^2}$ (Rave et al., 1998), which for greigite is $l_{\text{exch.}} \approx 6.6\,\text{nm}$; a maximum element size of 5 nm was used for all simulations. The non-local problem of calculating the stray field is resolved by a hybrid finite-element/boundary-element method (BEM) formulation (Fredkin and Koehler, 1990).

The magnetic parameters of greigite used in this investigation are: the saturation magnetisation $M_\text{S} = 2.7\times10^5$ A/m (Li et al., 2014); the exchange stiffness constant $A = 2 \times 10^{-12}$ J/m (Chang et al., 2008); and the first MCA constant $K_1 = -1.7 \times 10^4\,\text{J/m}^3$ (Winklhofer et al., 2014). This set of parameters has been used in recent numerical studies of greigite (Valdez-Grijalva et al., 2018; Valdez-Grijalva and Muxworthy, 2018).

2.2. The FORC model

FORC diagrams are constructed from a class of partial hysteresis curves called first-order reversal curves (Mayergoyz, 1986), each starting at a value B_a of the applied field along the main hysteresis branch and tracing the magnetisation as the field B_b is increased to saturation. A magnetisation function on two variables $M = M(B_a, B_b)$ is thus obtained. The FORC distribution ρ is then defined as (Roberts et al., 2000):

$$\rho = -\frac{1}{2}\frac{\partial^2 M}{\partial H_a \partial H_b} = -\frac{\mu_0^2}{2}\frac{\partial^2 M}{\partial B_a \partial B_b}, \tag{7}$$

where μ_0 is the magnetic constant (or vacuum permeability) and $H = B/\mu_0$.

Once $M(B_a, B_b)$ is obtained, calculation of $\rho(B_a, B_b)$ is done by least-squares fitting of a degree 2 polynomial surface $a_0 + a_1 B_a + a_2 B_b + a_3 B_a B_b + a_4 B_a^2 +$

$a_5 B_b^2 + \text{error} = M(B_a, B_b)$ on a subgrid of $M(B_a, B_b)$ centered around (B_a, B_b) as determined by the so-called smoothing factor (SF) and including $(2\times\text{SF} + 1)^2$ points; the value of ρ is then simply $-\mu_0^2 a_3/2$ (Pike et al., 1999). FORC diagrams are usually presented with rotated axes $B_c = (B_b - B_a)/2$, $B_u = (B_b + B_a)/2$.

Distributions with random orientation of magnetic particles with respect to the applied field were determined by taking 500 field orientations from a sector of the unit sphere (Fig. 1). We use 500 field orientations as a workable compromise between accuracy and calculation speed. Also, for each particle/field-orientation, the hysteresis curve consists mostly of reversible motion of the magnetisation; thus, we only need to calculate the main branch of the hysteresis loop and the few reversal curves starting at the different switching fields along the main branch (Valdez-Grijalva and Muxworthy, 2018). These simplifications reduce vastly the number of calculations needed without loss of important information. The external-field rate of change for all models was 1 mT with a saturation field of 250 mT, so that 501 reversal curves were calculated for each particle/field-orientation.

Scanning electron and transmission electron micrographs of naturally occurring greigite samples (Snowball, 1997; Vasiliev et al., 2008; Roberts, 2015) reveal that greigite tends to grow authigenically as well-defined regular truncated octahedral particles. Micromagnetic calculations for truncated octahedral greigite particles indicate that the SD–SV threshold occurs at ∼54 nm (Valdez-Grijalva et al., 2018). In this study we model FORC diagrams for non-interacting ensembles of truncated octahedral greigite particles sized 30–80 nm (where size is normalised to the volume of a cube) at 2 nm size intervales. This range is chosen because it spans the transition from SD to SV behaviour.

3. Results

For ensembles with SD particles < 50 nm, hysteresis behaviour is dominated by coherent rotation (Fig. 2). This is seen by comparing FORC diagrams for these ensembles (Fig. 2b) with those of idealised SD (effectively a single magnetic dipole), coherently rotating greigite particles (Fig. 2a) determined using the method outlined in Valdez-Grijalva and Muxworthy (2018). Diagrams for particles $< 50\,\mathrm{nm}$ obtained with the micromagnetic algorithm (Fig. 2b) are offset $\sim 3\,\mathrm{mT}$ to the left compared to the dipole model (Fig. 2a); lower coercivities due to the micromagnetic algorithm, which includes flowering (small deviations from a perfect SD structure) as a result of magnetostatic self-interaction effects, account for this effect.

Particles with cubic anisotropy have hysteresis behaviour that departs from that seen in the simple hysteron with one plus and one minus magnetisation states. There exist intermediate easy axis states along hysteresis curves for the SD state (Valdez-Grijalva and Muxworthy, 2018). The tilted, elongated, negative-valued ridge (Fig. 2) is a consequence of the cubic anisotropy and is produced by the fraction of particles with a hard axis aligned closely with the applied field. These particles have the lowest switching fields: from the plus-state to an intermediate state at $B = B_*^+$ and from the intermediate state to the minus-state at $B = B_-^*$. Reversal curves with $B_-^* < B_a < B_*^+$ experience a sharp upward discontinuity at $B_b = B_+^* < |B_*^+|$ when hard-aligned particles return to the plus-state from their intermediate states. The combination of this type of irreversible event in hard-aligned particles causes the local peak at $B_c \approx 15\,\mathrm{mT}, B_u \approx -3\,\mathrm{mT}$ (Fig. 2b). For reversal curves with $B_a < B_-^*$, hard-aligned particles are initially in the minus-state and undergo irreversible rotation to an intermediate state on the path to positive saturation at $B = B_*^- = |B_*^+|$ due to the symmetry of the particles and the lack of magnetostatic interactions. The combination of these irreversible events causes a negative FORC distribution response at $B_a = B_-^*, B_b = B_+^*$. The sum effect of this type of response for many particles with a distribution of switching fields produces the elongated negative contribution observed in all SD ensembles, roughly along the line segment connecting ($B_c = 18\,\mathrm{mT}, B_u = -6\,\mathrm{mT}$) with ($B_c = 42\,\mathrm{mT}, B_u = -30\,\mathrm{mT}$).

The fraction of particles with easy axis alignment close to the applied field orientation exhibits hysteron-like behaviour, i.e., just two switching fields: from the plus-state to the minus-state B_-^+ and *vice versa* B_+^-. The lack of interactions and the symmetry of particles in our simulations ensure that $|B_-^+| = B_-^+$. Thus, this fraction of particles produces FORC distribution responses at $B_a = B_-^+, B_b = B_+^-$. These types of irreversible responses accumulate on the line $B_a = -B_b$; they account for the most drastic changes in the magnetisation of the ensemble and, thus, account for the high slopes around the coercive field of the sample. This makes the position of the FORC diagram peak coincide with the coercivity of $B_\mathrm{C} \approx 24\,\mathrm{mT}$ for SD ensembles.

Particles with size $d \geq 50\,\mathrm{nm}$ switch incoherently; that is, the FORC diagrams depart from coherent rotation behaviour associated with SD particles as the tight boomerang-shaped FORC diagram pattern exhibited by the SD greigite (Fig. 2) becomes more fragmented (Fig. 3). This change is driven initially by particles with hard axes close to the applied field nucleating hard-aligned vortices (Valdez-Grijalva et al., 2018) as intermediate meta-stable states during hysteresis. Even though nucleation of hard-aligned vortices occurs in particles below the zero-field SD–SV threshold $d_0 \approx 54\,\mathrm{nm}$ (Valdez-Grijalva et al., 2018), this is expected because vortex nucleation greatly reduces the magnetic free-energy. A corollary of this is that a fraction of particles (with easy axis alignment close to the applied field) above the zero-field SD–SV threshold can remain in a SD state throughout hysteresis. These effects are due to distortion of the zero-field energy landscape by the applied field.

An appreciable positive source in the FORC distribution appears along the $B_u = 0$ axis at $B_c \approx 52\,\mathrm{mT}$ (the B_c axis is not to be confused with the coercivity B_C) for ensembles with particles $\geq 54\,\mathrm{nm}$ (Fig. 4, region 5); this contribution represents the annihilation of vortex states on the return to positive saturation. The elongated, negative ridge due to SD particles with cubic MCA and its corresponding symmetric positive response move to lower (B_c, B_u) values (Fig. 4, regions 1, 3) and the first responses for $B_u > 0$ begin to form (Fig. 4, region 2); these are elongated features at 45° to the $B_u = 0$ axis, which are different to the vertical widening usually attributed to magnetostatic interactions (Pike et al., 1999; Muxworthy et al., 2004; Muxworthy and Williams, 2005).

For particles slightly below and above the SD–SV threshold d_0, vortex nucleation occurs only for negative applied field values, thus noticeable changes in the FORC diagrams (Fig. 3a–c) are not evident in changes in the saturation remanence M_{RS} to saturation magnetisation M_{S} ratio up to 72 nm, whereas coercivity decreases sharply above 48 nm (Fig. 5b). The monotonically-decreasing coercivity trend is preserved up to 62 nm when it rises from $B_{\mathrm{C}} \approx 15\,\mathrm{mT}$ to ~20 mT for $d = 68\,\mathrm{nm}$. With increasing size, coercivitiy decreases further, accompanied by a sharp decrease in M_{RS} (Fig. 5b). The drop in M_{RS} is driven by particles nucleating vortices at $B_a > 0$ for $d \geq 68\,\mathrm{nm}$. For $d \geq 76\,\mathrm{nm}$, all particles nucleate vortices so that the vortex state becomes the remanent magnetic domain state; this is reflected in the Day plot (Day et al., 1977), a scatter plot of the $M_{\mathrm{RS}}/M_{\mathrm{S}}$ ratio against the coercivity of remanence B_{CR} (the field necessary to reduce the remanence to zero) to B_{C} ratio, by particles 76 nm and larger (Fig. 5a), associated here with the SV state. Particles sized 62–72 nm move away from the top left of the Day plot (Fig. 5a) to a region with high remanence but larger $B_{\mathrm{CR}}/B_{\mathrm{C}}$ values. These sizes coincide with the anomalous coercivity increase for these sizes (Fig. 5b). The increased coercivities can be explained by vortex nucleation, which causes hysteresis loops to become increasingly wasp-waisted (Fig. 6) so that they cross the zero-magnetisation axis at increasing (absolute) values of the applied field strength. FORC diagrams for these sizes are the most complex of all those simulated here, and have a variety of features (Figs. 3c, 6) caused by the complex interplay of SV and SD effects. The elongated, negative ridge becomes more faint with increasing particle size, whereas the positive responses for $B_u > 0$ become larger and move toward the $B_c = 0$ axis with increasing size. Large, positive FORC responses for $B_u > 0$ along the $B_c = 0$ axis are expected for larger multi-domain (MD) grains (Pike et al., 2001; Roberts et al., 2006). The non-interacting nature of these ensembles means that the SD and SV FORC signals are linearly additive. Therefore, it is possible to discern the FORC responses due to SD (Fig. 6, regions 3, 6) and SV particles (Fig. 6, regions 1, 2, 4, 5, 7, 8).

The elongated, negative ridge typical of SD particles with cubic MCA (Valdez-

Grijalva and Muxworthy, 2018) disappears for particles $\geq 76\,\mathrm{nm}$ (Fig. 3d). A circular, negative feature centered roughly at ($B_c = 8\,\mathrm{mT}, B_u = -8\,\mathrm{mT}$) becomes larger and of a magnitude comparable to the largest positive reponses. For $d = 76$ and $78\,\mathrm{nm}$ the negative response has a larger absolute value than the distribution peak (Fig. 3d). For the $80\,\mathrm{nm}$ particle model, a faint negative response appears centered roughly at ($B_c = 40\,\mathrm{mT}, B_u = -12\,\mathrm{mT}$) (Fig. 7, region 6). Fig. 7 represents the contribution of purely SV particles, that is, ensembles of particles that are all in a SV remanent state. It is logical that this FORC diagram is somewhat less complex than those for ensembles with a fraction of particles still in the SD state as well as some in the SV state; the difference is due to the field angle relative to particle orientation, as has also been shown by Roberts et al. (2017) for magnetite.

Particles with hard axes aligned closely with the applied field nucleate hard-aligned vortices at high applied field values (Fig. 8); as the field decreases below $\sim$12 mT these vortices rotate irreversibly to an easy axis alignment. As the field is increased on reversal curves with $\sim 0\,\mathrm{mT} \leq B_a \leq \sim 12\,\mathrm{mT}$ these vortices switch irreversibly back to a hard alignment at $B_b \approx 28\,\mathrm{mT}$ to create a local peak at $B_c \approx 12\,\mathrm{mT}$, $B_u \approx 16\,\mathrm{mT}$ (Fig. 7, region 1); this is manifested in the raw hysteresis data by the smoothed discontinuity at $B \approx 28\,\mathrm{mT}$ whereas the reversible motion traced by the reversal curves around this region accounts for the tilted, elongated response surrounding the local peak.

During hysteresis, as the remanent state is approached, all particles $\geq 76\,\mathrm{nm}$ have nucleated vortices: particles with easy axis alignment close to the applied field directly nucleate an easy-aligned vortex while the rest nucleate vortices initially oriented along hard <100> or <110> directions (Fig. 8b), which rotate irreversibly to an easy axis alignment as the field approaches zero. The latter fraction of particles then undergo irreversible rotations to intermediate positions for $\sim -10\,\mathrm{mT} \leq B \leq \sim -20\,\mathrm{mT}$. For FORCs with $\sim -10\,\mathrm{mT} \leq B_a \leq \sim -20\,\mathrm{mT}$, these vortices rotate back to the initial easy axis alignment at $B_b \approx 4\,\mathrm{mT}$. A combination of these irreversible events creates the lowest negative FORC response (Fig. 7, region 2). A further applied field increase to $\sim$30 mT causes these vortices to switch to the initial hard position from which they nucleated. These events cause the tilted, elongated FORC response (Fig. 7, region 3).

As the applied field decreases past $\sim-52\,\text{mT}$, the vortices of particles with easy axis alignment close to the applied field annihilate (Fig. 8). Reversal curves with $\sim-80\,\text{mT} \leq B_a \leq \sim-52\,\text{mT}$ trace lower slopes with decreasing B_a due to the combined reversible motion of vortices and single domains; this is the source of the faint negative contribution for $B_u < \sim45\,\text{mT}$ (Fig. 7, region 4). On increasing B_b on these curves, nucleation of easy-aligned vortices occurs at $\sim-5\,\text{mT}$ creating the boomerang-shaped response (Fig. 7, region 5) that limits the faint negative response in region 4; this corresponds with the smoothed discontinuity in hysteresis curves as the field approaches zero from the left. Increasing the applied field to positive values causes the easy-aligned vortices of particles with hard axes close to the applied field to switch to hard alignments at $\sim28\,\text{mT}$, creating a negative FORC region (Fig. 7, region 6). The distribution peak at region 7 (Fig. 7) corresponds to the average annihilation field of the vortices on the reversal paths to positive saturation.

There is a large spread in the vortex nucleation and annihilation fields (Fig. 8). Particles with hard axis alignment close to the applied field nucleate hard-aligned vortices for fields as high as $\sim200\,\text{mT}$ and annihilate on the opposite side of the particle for equally high (absolute) values. However, these nucleation and annihilation events make a negligible contribution to the FORC diagram because the change in magnetisation of a particle nucleating/annihilating a hard-aligned vortex from/to a SD state can be as low as 1%.

4. Discussion

Comparison of results for micromagnetic simulations presented here with the coherently rotating dipole model of Valdez-Grijalva and Muxworthy (2018) indicates excellent agreement (Fig. 2). This confirms the accuracy of our model using only 500 random field orientations instead of field orientations on a regular grid, which requires a high density of field orientations near the poles of the sphere. A FORC diagram for SD coherently rotating particles has the same general features as those obtained for weakly interacting SD particles with cubic MCA by Harrison and Lascu (2014), i.e., a positive ridge along the B_c axis, slightly offset toward $B_u < 0$ values and a tilted, negative ridge on the lower half of the FORC plane. For these ensembles, the horizontal spread along the B_c axis corresponds to the density of switching fields of the differently oriented particles and the FORC distribution peak position corresponds directly to the ensemble coercivity. The negative ridge is indicative of intermediate states along the hysteresis curve and, therefore, of SD particles with cubic MCA (Valdez-Grijalva and Muxworthy, 2018); this FORC response has been identified in simulations for magnetite (Harrison and Lascu, 2014), and is potentially unique to non-interacting to weakly interacting SD particles with cubic MCA.

Whereas the pure SD signal produces a tight, boomerang-shaped FORC distribution (Fig. 2), increasing particle size introduces SV structures that fragment this pattern. The FORC distribution peak is moved toward higher B_c values along the $B_u = 0$ axis. Paradoxically, as this occurs, the bulk coercivity of the ensembles decreases (Fig. 5). This paradox has been observed previously in synthetic size-controlled samples of sub-100 nm Fe dots (Dumas et al., 2007).

Fragmentation of the FORC diagram for non-uniformly magnetised particles has been observed in experimental studies (Pike and Fernandez, 1999; Dumas et al., 2007; Roberts et al., 2017; Zhao et al., 2017) and in numerical models (Carvallo et al., 2003; Roberts et al., 2017); however, these studies did not include random field orientation distributions. The trend is, nevertheless, clear and is representative of the complex self-interactions brought about by nonuniform structures and multiple vortex nucleation/annihilation fields (Pike and Fernandez, 1999). It is difficult to compare our results to the FORC signals measured by Muxworthy et al. (2006) and Krása et al. (2011) for synthetic patterned magnetite because many of their FORC diagrams appear to have smoothed the subtle features observed here, which raises questions about the integrity of these samples (e.g., crystallinity) or the adequateness of the FORC measurement density for these samples.

Pike and Fernandez (1999) obtained asymmetric nucleation and annihilation fields of magnetic vortices in nano-patterned Co dots; our models agree with this finding (Fig. 8). However, Pike and Fernandez (1999) studied elongated disc-like particles where the vortex cores were always perpendicular to the particle plane that mostly underwent reversible motion from nucleation to annihilation as they traversed the particle. In this study, we demonstrate that different features on SV FORC diagrams are due to a variety of vortex nucleation and annihilation events, which depend on particle alignment with respect to the applied field and on the presence of distinctly different vortex states, i.e., the vortex energies and stabilities depend on their alignment within the crystalline structure (Valdez-Grijalva et al., 2018).

FORC diagrams were averaged for simulations between 30 and 80 nm (Fig. 9a) and between 60 and 80 nm (Fig. 9b). That is, the particle size distributions are uniform (flat) for these averaged FORC diagrams. The FORC diagram in Fig. 9a has a boomerang-shaped distribution surrounded by a variety of more complex responses. This pattern shows some similarities to the patterns observed by Dumas et al. (2007) for samples that included both SD and SV particles. The FORC distribution peak position coincides with the ensemble coercivity, while still having a response corresponding to the annihilation field of easy-aligned vortices.

Both diagrams in Fig. 9 have a significative spread in the positive B_u region. This effect is purely due to domain state, not magnetostatic interactions. The main peak for the averaged SV-dominant diagram (Fig. 9b) occurs along the $B_u = 0$ axis at $B_c \approx 52\,\mathrm{mT}$, which indicates a disconnect with the bulk coercivity of the ensemble ($B_\mathrm{C} \approx 16\,\mathrm{mT}$). This is a departure from the usual interpretation of FORC diagrams, i.e., that the FORC diagram provides a map of the coercivity distribution. This interpretation holds for SD coherently rotating grains, where the peak response coincides with the value of the ensemble coercivity. It does not hold, however, for SV grains because their coercivity decreases with size while the position of the maximum moves toward higher B_c values. Instead, for SV grains the FORC distribution peak, and most FORC features, should be interpreted as due to vortex nucleation/annihilation fields and their irreversible motions.

5. Conclusion

A micromagnetic FEM/BEM was employed to calculate FORC distributions for non-interacting ensembles of greigite across a size range that spans the SD to SV threshold. 500 random orientations from a uniform distribution over a sector of the unit sphere were used for each particle size. This choice was found to be in excellent agreement with previous calculations for SD greigite (Valdez-Grijalva and Muxworthy, 2018).

FORC diagrams are found to be extremely sensitive to the domain state of the simulated particles. When even a small fraction of particles starts to nucleate vortices, e.g., $d \approx 50\,\mathrm{nm}$, this is reflected in the FORC diagram (Fig. 3a compared to Fig. 2). The same cannot be said of the Day plot (Fig. 5a). Anomalous behaviour for particles sized 62 to 72 nm, with coercivity increasing with size was found; these particles plot in an unexpected region of the Day plot. The anomaly disappears for particles $> 72\,\mathrm{nm}$, and when $d \geq 76\,\mathrm{nm}$ they have much lower $M_{\mathrm{RS}}/M_{\mathrm{S}}$ and higher $B_{\mathrm{CR}}/B_{\mathrm{C}}$ values.

Detailed FORC analysis and micromagnetic solutions for $d = 80\,\mathrm{nm}$ particles reveals the meaning of the FORC diagram for SV ensembles as a map of vortex nucleation/annihilation fields. Interpretation of FORC diagrams as a coercivity distribution does not apply to SV systems (see Pike and Fernandez (1999); Roberts et al. (2017)). Recognition that the remanence in palaeomagnetic studies is often carried by vortex state particles should help users of FORC diagrams to avoid misinterpretation of vertical spread in FORC diagrams, just as it is recognised that vertical spread in MD particles is due to domain wall interactions within particles (Pike et al., 2001). For SD particles, the typical interpretation of the peak position coinciding with the coercivity of the sample holds; however, for SV-dominated samples, the position of the peak occurs at a value much higher than the bulk coercivity of the sample.

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