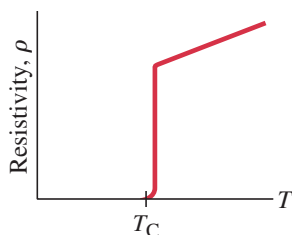


## \*18–9 Superconductivity

At very low temperatures, well below  $0^{\circ}\text{C}$ , the resistivity (Section 18–4) of certain metals and certain compounds or alloys becomes zero as measured by the highest-precision techniques. Materials in such a state are said to be **superconducting**. This phenomenon was first observed by H. K. Onnes (1853–1926) in 1911 when he cooled mercury below 4.2 K ( $-269^{\circ}\text{C}$ ) and found that the resistance of mercury suddenly dropped to zero. In general, superconductors become superconducting only below a certain *transition temperature* or *critical temperature*,  $T_C$ , which is usually within a few degrees of absolute zero. Current in a ring-shaped superconducting material has been observed to flow for years in the absence of a potential difference, with no measurable decrease. Measurements show that the resistivity  $\rho$  of superconductors is less than  $4 \times 10^{-25} \Omega \cdot \text{m}$ , which is over  $10^{16}$  times smaller than that for copper, and is considered to be zero in practice. See Fig. 18–26.



**FIGURE 18–26** A superconducting material has zero resistivity when its temperature is below  $T_C$ , its “critical temperature.” At temperatures above  $T_C$ , the resistivity jumps to a “normal” nonzero value and increases with temperature as most materials do (Eq. 18–4).

Before 1986 the highest temperature at which a material was found to superconduct was 23 K, which required liquid helium to keep the material cold. In 1987, a compound of yttrium, barium, copper, and oxygen (YBCO) was developed that can be superconducting at 90 K. Since this is above the boiling temperature of liquid nitrogen, 77 K, liquid nitrogen is sufficiently cold to keep the material superconducting. This was an important breakthrough because liquid nitrogen is much more easily and cheaply obtained than is the liquid helium needed for earlier superconductors. Superconductivity at temperatures as high as 160 K has been reported, though in fragile compounds.

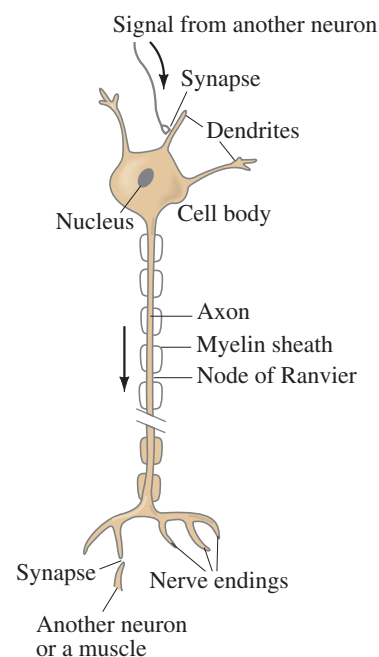
To develop high- $T_C$  superconductors for use as wires (such as for wires in “superconducting electromagnets”—Section 20–7), many applications today utilize a bismuth-strontium-calcium-copper oxide (BSCCO). A major challenge is how to make a useable, bendable wire out of the BSCCO, which is very brittle. (One solution is to embed tiny filaments of the high- $T_C$  superconductor in a metal alloy, which is not resistanceless but has resistance much less than a conventional copper cable.)

## \*18–10 Electrical Conduction in the Human Nervous System

An interesting example of the flow of electric charge is in the human nervous system, which provides us with the means for being aware of the world, for communication within the body, and for controlling the body’s muscles. Although the detailed functioning of the hugely complex nervous system still is not well understood, we do have a reasonable understanding of how messages are transmitted within the nervous system: they are electrical signals passing along the basic element of the nervous system, the **neuron**.

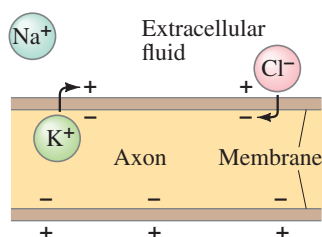
Neurons are living cells of unusual shape (Fig. 18–27). Attached to the main cell body are several small appendages known as *dendrites* and a long tail called the *axon*. Signals are received by the dendrites and are propagated along the axon. When a signal reaches the nerve endings, it is transmitted to the next neuron or to a muscle at a connection called a *synapse*.

**FIGURE 18–27** A simplified sketch of a typical neuron.



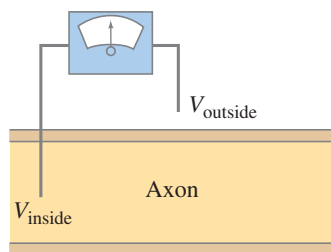
**TABLE 18–2**  
**Concentrations of Ions Inside and Outside a Typical Axon**

	Concentration inside axon (mol/m <sup>3</sup> )	Concentration outside axon (mol/m <sup>3</sup> )
K <sup>+</sup>	140	5
Na <sup>+</sup>	15	140
Cl <sup>−</sup>	9	125

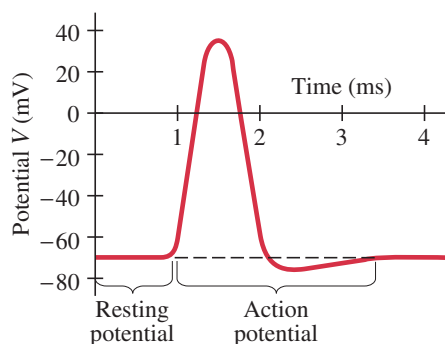


**FIGURE 18–28** How a dipole layer of charge forms on a cell membrane.

**FIGURE 18–29** Measuring the potential difference between the inside and outside of a nerve cell.



**FIGURE 18–30** Action potential.



A neuron, before transmitting an electrical signal, is in the so-called “resting state.” Like nearly all living cells, neurons have a net positive charge on the outer surface of the cell membrane and a negative charge on the inner surface. This difference in charge, or **dipole layer**, means that a potential difference exists across the cell membrane. When a neuron is not transmitting a signal, this **resting potential**, normally stated as

$$V_{\text{inside}} - V_{\text{outside}},$$

is typically  $-60$  mV to  $-90$  mV, depending on the type of organism. The most common ions in a cell are K<sup>+</sup>, Na<sup>+</sup>, and Cl<sup>−</sup>. There are large differences in the concentrations of these ions inside and outside an axon, as indicated by the typical values given in Table 18–2. Other ions are also present, so the fluids both inside and outside the axon are electrically neutral. Because of the differences in concentration, there is a tendency for ions to diffuse across the membrane (see Section 13–13 on diffusion). However, in the resting state the cell membrane prevents any net flow of Na<sup>+</sup> through a mechanism of active transport<sup>‡</sup> of Na<sup>+</sup> ions out of the cell by a particular protein to which Na<sup>+</sup> attach; energy needed comes from ATP. But it does allow the flow of Cl<sup>−</sup> ions, and less so of K<sup>+</sup> ions, and it is these two ions that produce the dipole charge layer on the membrane. Because there is a greater concentration of K<sup>+</sup> inside the cell than outside, more K<sup>+</sup> ions tend to diffuse outward across the membrane than diffuse inward. A K<sup>+</sup> ion that passes through the membrane becomes attached to the outer surface of the membrane, and leaves behind an equal negative charge that lies on the inner surface of the membrane (Fig. 18–28). The fluids themselves remain neutral. What keeps the ions on the membrane is their attraction for each other across the membrane. Independently, Cl<sup>−</sup> ions tend to diffuse *into* the cell since their concentration outside is higher. Both K<sup>+</sup> and Cl<sup>−</sup> diffusion tends to charge the interior surface of the membrane negative and the outside positive. As charge accumulates on the membrane surface, it becomes increasingly difficult for more ions to diffuse: K<sup>+</sup> ions trying to move outward, for example, are repelled by the positive charge already there. Equilibrium is reached when the tendency to diffuse because of the concentration difference is just balanced by the electrical potential difference across the membrane. The greater the concentration difference, the greater the potential difference across the membrane ( $-60$  mV to  $-90$  mV).

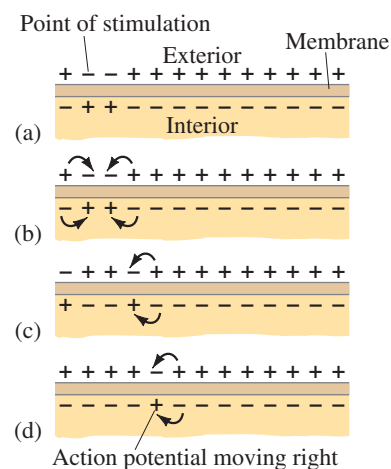
The most important aspect of a neuron is not that it has a resting potential (most cells do), but rather that it can respond to a stimulus and conduct an electrical signal along its length. The stimulus could be thermal (when you touch a hot stove) or chemical (as in taste buds); it could be pressure (as on the skin or at the eardrum), or light (as in the eye); or it could be the electric stimulus of a signal coming from the brain or another neuron. In the laboratory, the stimulus is usually electrical and is applied by a tiny probe at some point on the neuron. If the stimulus exceeds some threshold, a voltage pulse will travel down the axon. This voltage pulse can be detected at a point on the axon using a voltmeter or an oscilloscope connected as in Fig. 18–29. This voltage pulse has the shape shown in Fig. 18–30, and is called an **action potential**. As can be seen, the potential increases from a resting potential of about  $-70$  mV and becomes a positive  $30$  mV or  $40$  mV. The action potential lasts for about  $1$  ms and travels down an axon with a speed of  $30$  m/s to  $150$  m/s. When an action potential is stimulated, the nerve is said to have “fired.”

What causes the action potential? At the point where the stimulus occurs, the membrane suddenly alters its permeability, becoming much more permeable to Na<sup>+</sup> than to K<sup>+</sup> and Cl<sup>−</sup> ions. Thus, Na<sup>+</sup> ions rush into the cell and the inner surface of the wall becomes positively charged, and the potential difference quickly swings positive ( $\approx +30$  mV in Fig. 18–30). Just as suddenly, the membrane returns to its original characteristics; it becomes impermeable to Na<sup>+</sup> and in fact pumps out Na<sup>+</sup> ions. The diffusion of Cl<sup>−</sup> and K<sup>+</sup> ions again predominates and the original resting potential is restored ( $-70$  mV in Fig. 18–30).

<sup>‡</sup>This transport mechanism is sometimes referred to as the “sodium pump.”

What causes the action potential to travel along the axon? The action potential occurs at the point of stimulation, as shown in Fig. 18–31a. The membrane momentarily is positive on the inside and negative on the outside at this point. Nearby charges are attracted toward this region, as shown in Fig. 18–31b. The potential in these adjacent regions then drops, causing an action potential there. Thus, as the membrane returns to normal at the original point, nearby it experiences an action potential, so the action potential moves down the axon (Figs. 18–31c and d).

You may wonder if the number of ions that pass through the membrane would significantly alter the concentrations. The answer is no; and we can show why (and again show the power and usefulness of physics) by treating the axon as a capacitor as we do in Search and Learn Problem 6 (the concentration changes by less than 1 part in  $10^4$ ).



**FIGURE 18–31** Propagation of action potential along axon membrane.

## Summary

An electric **battery** serves as a source of nearly constant potential difference by transforming chemical energy into electric energy. A simple battery consists of two electrodes made of different metals immersed in a solution or paste known as an electrolyte.

**Electric current**,  $I$ , refers to the rate of flow of electric charge and is measured in **amperes** (A): 1 A equals a flow of 1 C/s past a given point.

The direction of **conventional current** is that of positive charge flow. In a wire, it is actually negatively charged electrons that move, so they flow in a direction opposite to the conventional current. A positive charge flow in one direction is almost always equivalent to a negative charge flow in the opposite direction. Positive conventional current always flows from a high potential to a low potential.

The **resistance**  $R$  of a device is defined by the relation

$$V = IR, \quad (18-2)$$

where  $I$  is the current in the device when a potential difference  $V$  is applied across it. For materials such as metals,  $R$  is a constant independent of  $V$  (thus  $I \propto V$ ), a result known as **Ohm's law**. Thus, the current  $I$  coming from a battery of voltage  $V$  depends on the resistance  $R$  of the circuit connected to it.

Voltage is applied *across* a device or between the ends of a wire. Current passes *through* a wire or device. Resistance is a property of the wire or device.

The unit of resistance is the **ohm** ( $\Omega$ ), where  $1 \Omega = 1 \text{ V/A}$ . See Table 18–3.

**TABLE 18–3 Summary of Units**

Current	1 A = 1 C/s
Potential difference	1 V = 1 J/C
Power	1 W = 1 J/s
Resistance	1 $\Omega$ = 1 V/A

The resistance  $R$  of a wire is inversely proportional to its cross-sectional area  $A$ , and directly proportional to its length  $\ell$  and to a property of the material called its resistivity:

$$R = \frac{\rho \ell}{A}. \quad (18-3)$$

The **resistivity**,  $\rho$ , increases with temperature for metals, but for semiconductors it may decrease.

The rate at which energy is transformed in a resistance  $R$  from electric to other forms of energy (such as heat and light)

is equal to the product of current and voltage. That is, the **power** transformed, measured in watts, is given by

$$P = IV, \quad (18-5)$$

which for resistors can be written as

$$P = I^2 R = \frac{V^2}{R}. \quad (18-6)$$

The SI unit of power is the **watt** ( $1 \text{ W} = 1 \text{ J/s}$ ).

The total electric energy transformed in any device equals the product of the power and the time during which the device is operated. In SI units, energy is given in joules ( $1 \text{ J} = 1 \text{ W} \cdot \text{s}$ ), but electric companies use a larger unit, the **kilowatt-hour** ( $1 \text{ kWh} = 3.6 \times 10^6 \text{ J}$ ).

Electric current can be **direct current** (**dc**), in which the current is steady in one direction; or it can be **alternating current** (**ac**), in which the current reverses direction at a particular frequency  $f$ , typically 60 Hz. Alternating currents are typically sinusoidal in time,

$$I = I_0 \sin \omega t, \quad (18-7b)$$

where  $\omega = 2\pi f$ , and are produced by an alternating voltage.

The **rms** values of sinusoidally alternating currents and voltages are given by

$$I_{\text{rms}} = \frac{I_0}{\sqrt{2}} \quad \text{and} \quad V_{\text{rms}} = \frac{V_0}{\sqrt{2}}, \quad (18-8)$$

respectively, where  $I_0$  and  $V_0$  are the **peak** values. The power relationship,  $P = IV = I^2 R = V^2/R$ , is valid for the average power in alternating currents when the rms values of  $V$  and  $I$  are used.

[\*The current in a wire, at the microscopic level, is considered to be a slow **drift velocity** of electrons,  $\vec{v}_d$ . The current  $I$  is given by

$$I = neA v_d, \quad (18-10)$$

where  $n$  is the number of free electrons per unit volume,  $e$  is the magnitude of the charge on an electron, and  $A$  is the cross-sectional area of the wire.]

[\*At very low temperatures certain materials become **superconducting**, which means their electrical resistance becomes zero.]

[\*The human nervous system operates via electrical conduction: when a nerve “fires,” an electrical signal travels as a voltage pulse known as an **action potential**.]